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Correspondence and requests for materials should be addressed to X.J.H. (xingjiuhuang@ iim.ac.cn) or Q.X.L. (liqun@ustc.edu.cn)

Facet-dependent electrochemical properties of $Co₃O₄$ nanocrystals toward heavy metal ions

Xin-Yao Yu', Qiang-Qiang Meng 2 , Tao Luo', Yong Jia', Bai Sun', Qun-Xiang Li 2 , Jin-Huai Liu' & Xing-Jiu Huang¹

¹ Research Center for Biomimetic Functional Materials and Sensing Devices, Institute of Intelligent Machines, Chinese Academy of Sciences, Hefei 230031, PR China, ²Hefei National Laboratory for Physical Science at the Microscale, University of Science and Technology of China, Hefei, Anhui 230026, PR China.

We revealed an interesting facet-dependent electrochemical behavior toward heavy metal ions (HMIs) based on their adsorption behaviors. The (111) facet of $Co₃O₄$ nanoplates has better electrochemical sensing performance than that of the (001) facet of $Co₃O₄$ nanocubes. Adsorption measurements and density-functional theory (DFT) calculations reveals that adsorption of HMIs is responsible for the difference of electrochemical properties. Our combined experimental and theoretical studies provide a solid hint to explain the mechanism of electrochemical detection of HMIs using nanoscale metal oxides. Furthermore, this study not only suggests a promising new strategy for designing high performance electrochemical sensing interface through the selective synthesis of nanoscale materials exposed with different well-defined facets, but also provides a deep understanding for a more sensitive and selective electroanalysis at nanomaterials modified electrodes.

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trodes were very often explored for electrochemical detection of very trace levels of toxic heavy metal ions
(HMIs). V trodes were very often explored for electrochemical detection of very trace levels of toxic heavy metal ions (HMIs). Very recently, nanoscale metal oxides as novel modifiers have been reported in electrochemical detection of HMIs¹⁻⁵. Compared with traditional modifiers such as noble metals and biomolecules, the electrodes fabricated with nanoscale metal oxides are easy to synthesize with low $cost^{1,3,5}$. However, as is often claimed, increased currents and increased analytical sensitivity are simply reflections of an increased microscopic surface area and not electrocatalytic activity or any other nano-effects. That is, the mechanism of using nanoscale metal oxides in electrochemical sensing of HMIs has not been proposed by considering their insulating property. Therefore, the design and implementation of new experimental approach combined with theoretical studies are extremely needed.

For a crystalline nanoscale material, different facets may have different geometric and electronic structures, and exhibit different physical and chemical properties^{6–8}. Much more attention has been given to investigate the facet effects on catalysts, photocatalysts, electrocatalysts, Li-ion battery, supercapacitors and so on^{9–24}. Take Co₃O₄ nanoscale materials as an example; Liet al. investigated the facet effect of $Co₃O₄$ on catalytic property for methane combustion¹⁰. Xiao *et al.* reported that the exposed facets of $Co₃O₄$ nanocrystals are very important for Li⁺ transport in Li-ion battery¹². Although oriented $Co₃O₄$ nanostructures have attracted wide attention, the direct experimental detection of HMIs with controlled crystalline morphology and orientation remains a significant challenge- that is, the facet effect in electrochemical sensor especially detection of HMIs is not proposed. Further, till now nanoscale $Co₃O₄$ has not been used to examine the electrochemical sensing of HMIs. It should be also pointed out that plenary theoretical investigations are expected to help to get deeper insight into the crystal facet effect, but most reports on the crystal facet effect did not adequately combined first-principles theoretical studies at atomic level with the experimental results^{13,14,17,19-21,24}.

In this work, we report the facile synthesis of $Co₃O₄$ with two different shapes. Using this non-conductive nanomaterial-modified electrode, we try to demonstrate the facet-dependent electrochemical behaviour of $Co₃O₄$ nanocrystals toward HMIs by combining the adsorption measurements and the density-functional theory (DFT) calculations. It has been recognized that the (111) facet of $Co₃O₄$ nanoplates has better electrochemical sensing performance than that of the (001) facet of $Co₃O₄$ nanocubes. Adsorption measurements and DFT calculations reveal that adsorption of Pb(II) is responsible for the difference of electrochemical

properties. To the best of our knowledge, this is the first study to investigate the effect of facet on electrochemical sensing behaviour toward HMIs. Our combined experimental and theoretical studies provide a solid hint to explain the mechanism of electrochemical detection of HMIs using nanoscale metal oxides. This strategy may be extended to other electrochemical sensors based on nanoscale metal oxides.

Results

Two simple routes have been designed to fabricate $Co₃O₄$ nanocubes and nanoplates (Supporting Information, Figure S1). $Co₃O₄$ nanocubes were synthesized by a one-step hydrothermal method with $Co(CH_3COO)_2$ as the only reactant. The Co_3O_4 nanoplates were obtained by a solvothermal method in ethylene glycol (EG) followed by a calcination treatment in air. Figure 1 shows the representative scanning electron microscopy (SEM), transmission electron microscopy (TEM) and high resolution TEM (HRTEM) images of $Co₃O₄$ nanoplates and nanocubes. The TEM and HRTEM images of $Co₃O₄$ nanocubes are presented in Figure 1a and b, respectively. As seen, the particle size of $Co₃O₄$ nanocubes is about 20 \sim 40 nm. HRTEM (Figure 1b) and corresponding selected area electron diffraction (SAED) pattern (Inset in Figure 1b) indicate that the lattice fringe is 0.28 nm and the nanocube is exposed with six (001) facets. A typical SEM image taken for $Co₃O₄$ nanoplates is shown in Figure 1c. The low magnification SEM and TEM images in Figure S2 demonstrated the homogeneity of $Co₃O₄$ nanoplates. It is seen that most of the nanoplates display well-defined hexagonal shape. The width and the thickness of the plates is determined to be 700 \sim 900 nm and 50 \sim 60 nm, respectively. The HRTEM image in Figure 1d clearly shows that the spacing between lattice fringes with an angle of 60° is 0.28 nm, which is consistent with the (202) and

($0\overline{2}2$) planes of cubic Co₃O₄. And it is also confirmed by the SAED pattern (Inset in Figure 1d) revealing a single crystal structure. Thus, the resulting dominant exposed plates of $Co₃O₄$ nanoplates are (111) facets. Powder X-ray diffraction patterns (XRD) (Supporting Information, Figure S3) of both nanocubes and nanoplates have identical peaks, which can be perfectly indexed to that of cubic spinel Co3O4 (Joint Committee on Powder Diffraction Standards (JCPDS) card no. 42-1467). No impurities have been detected, indicating the formation of pure cobalt oxides. The crystal sizes were calculated according to Sherrer equation based on the (311) diffraction peak of XRD data. The crystal size of $Co₃O₄$ nanocubes was 32.5 nm, which was consistent with the TEM observation. However, the crystal size of $Co₃O₄$ nanoplates was calculated to be 21.3 nm, which may be due to its porous structure. The specific surface areas of $Co₃O₄$ nanocubes and nanoplates have been measured by the Brunauer– Emmett–Teller (BET) method (Supporting Information, Figure S4). The measured specific surface areas for $Co₃O₄$ nanocubes and nanoplates are 20 and 13 $m^2 g^{-1}$, respectively.

In general, the obtaining of metal oxides requires a basic medium, and usually NaOH is employed for this purpose. However, in the fabrication of $Co₃O₄$ nanocubes, the hydrolysis of acetate anions from $Co(CH_3COO)_2$ becomes a good source of hydroxyl anions (Equation (1)). The O_2 dissolved in the deionized water is supposed to act as the main oxidant. Partial Co(II) is transformed into Co(III) and the coexistence of Co(II) and Co(III) under the ambient environment prefers the formation of $Co₃O₄$ (Equation (2)). The low amount of hydroxyl anions in the solution may promote the formation of $Co₃O₄$ nanocubes with (001) planes. Similar result has also been reported by Xiao et al¹².

$$
CH_3COO^- + H_2O \rightarrow CH_3COOH + OH^-
$$
 (1)

Figure 1 | Co₃O₄ nanocrystals. (a–b), TEM and HRTEM image of typical Co₃O₄ nanocubes. Inset in panel (b): SAED pattern of nanocubes. (c–d), SEM and HRTEM image of typical $Co₃O₄$ nanoplates. Inset in panel (d): SAED pattern of nanoplates.

$$
\left\langle \right\rangle
$$

$$
6Co^{2+} + 12OH^{-} + O_2 \rightarrow 2Co_3O_4 + 6H_2O \qquad (2)
$$

From the SEM and TEM images of the $Co₃O₄$ nanoplates precursors before calcination (Data not shown), we can see that the platelike structures of $Co₃O₄$ are inherited from these precursors. The typical XRD pattern of the precursors is shown (Supporting Information, Figure S3, black line). A strong diffraction peak around 10° in the XRD pattern is the characteristic of metal glycolates^{25–28}. It has been assumed that EG would lose its two protons and the dianion complex with metal center²⁵. Xia et al. has discussed the oligomerization process of metal glycolates. Longer chains of cobalt glycolate oligomers could self-assemble into ordered bundles (i.e., nanoplates) through van der Waals interactions and then precipitate out from the reaction medium²⁵. After calcination, $Co₃O₄$ exposed mainly with (111) facets are obtained. However, the reason for the exposed (111) facets after calcination is not clear.

 $Co₃O₄$ nanocubes and nanoplates modified glassy carbon electrodes (GCEs) are first electrochemically characterized by cyclic voltammetric (CV) and electrochemical impedance spectrum (EIS) (Supporting Information, Figure S5). As compared with the bare GCE, the anodic and cathodic peaks decrease at the $Co₃O₄$ modified electrode, which demonstrates that $Co₃O₄$ nanomaterials have been modified onto the surface of the GCE (Supporting Information, Figure S5a). It also indicates that the rate of electron transfer at the electrode surface is hindered with the attachment $Co₃O₄$ onto the GCE surface. In a typical Nyquist plot, the semicircle proton corresponding to the electron-transfer resistance (R_{et}) at higher frequency range while a linear part at lower frequency range represents the diffusion limited process³. The value corresponding to the bare GCE is about 27 Ω (Supporting Information, Figure S5b). And the value of Co₃O₄ nanocubes modified GCE (48 Ω) is comparable to that of Co₃O₄ nanoplates (50 Ω) (Supporting Information, Figure S5b). In this case, the surface area of the modified electrodes were calculated to be 0.0534 (nanoplates), 0.0519 (nanocubes), and 0.07065 (bare GCE) cm^2 . This may be due to the insulating property of $Co₃O₄$ nanostructures, as such, the direct contribution of microscopic surface area could be avoided, which is very helpful to understand that the difference of electrochemical behavior is indeed from the effect of crystal facet (as will be discussed in the following).

Among all of the HMIs, lead has been identified as one of the most toxic heavy metals because of its detrimental effects on the human nervous system, blood circulation system, kidneys, and reproductive system²⁹. Subsequently, the performances of these two types of Co3O4 with different exposed facets in electrochemical detection of heavy metal ions are comprehensively investigated. Pb(II) is used as a probe heavy metal ion. We first examined the square wave stripping voltammograms (SWASV) responses to Pb(II) in 0.1 M NaAc-HAc ($pH = 5.0$) at $Co₃O₄$ nanoplates and nanocubes modified electrodes (Figure S6). There is almost no obvious response at bare GCE, and at the meanwhile, a weak stripping response for nanocubes/GCE is observed. However, a strong and well-defined peak at -0.584 V is clearly seen for nanoplates modified GCE. And the peak current obtained is about 8 times and 13 times that of nanocubes/GCE and bare GCE, respectively. Moreover, the stripping peak shifts toward more positive potential and become progressively less symmetrical. This is the consequence of having more ions on the electrode (this is fully consistent with that more ions are adsorbed by nanoplates and release to the bare GCE, see Discussion section), and therefore requiring a longer sweep to remove the metal from the surface entirely. Figure 2 shows that the current densities increase linearly versus the Pb(II) concentrations. For nanocubes/GCE, the linearization equation is $j/(\mu A \text{ cm}^{-2}) = -122.9 + 174.4 \text{ c}/\mu M$. While for nanoplates/GCE, the linearization equation is j/(μ A cm⁻²) = -131.1 $+ 372.7$ c/ μ M. The results indicate that Co₃O₄ nanoplates with (111) facet exhibit better electrochemical detection performance than

Figure 2 | Electrochemical properties of $Co₃O₄$ nanocrystals. Calibration plots of $Co₃O₄$ nanoplates and $Co₃O₄$ nanocubes modified electrode toward Pb(II) at different concentrations in 0.1 M NaAc-HAc solution $(pH = 5.0)$. Considering the weak response at nanocubes modified electrode, the initial concentration is from 0.7 μ M.

 $Co₃O₄$ nanocubes with (001) facet. The sensitivity of $Co₃O₄$ nanoplates/GCE (372.7 μ A cm⁻²/ μ M) is over 2 times that of Co₃O₄ nanocubes/GCE (174.4 μ A cm⁻²/ μ M). The SWASV responses toward Pb(II) at various concentrations on $Co₃O₄$ nanoplates and nanocubes modified electrodes were seen in Figure S7. The limit of detection (LOD) as low as 0.12 nM and 0.16 nM (3 σ method) for Co₃O₄ nanoplates and nanocubes, respectively was achieved. This meets the requirements of the World Health Organization (WHO) maximum permissible limit for lead concentration in drinking water of 10 µg L^{-1} . The LODs are better than existing methods based on noble metal and ion-imprinted polymers, such as gold nanofilm $(0.1 \mu M)^{30}$ and nano-sized Pb^{2+} imprinted polymer (0.6 nM)³¹, etc. The LODs are also superior to some oxides, such as tube-in-tube $SnO₂ (1.6 nM)³²$ and $SnO₂/graphene$ nanocomposites (0.18 nM)³³.

Discussion

We suggest that the difference in electrochemical behavior may have a relationship with the adsorption capacities of different nanoscale $Co₃O₄$ toward Pb(II). To confirm this contribution, the adsorption measurement is conducted in 0.1 M NaAc-HAc ($pH = 5.0$) in order to be consistent with electrochemical sensing conditions. $Co₃O₄$ nanocubes and nanoplates are incubated in Pb(II) solutions with different initial concentrations for 24 h to reach the adsorption equilibrium. The adsorption isotherms are shown in Figure 3a and b. Both the adsorption isotherms fit the Langmuir isotherm very well (Supporting Information, Figure S8a and b). The maximum adsorption capacity (q_{max}) of Co₃O₄ nanoplates (33.8 mg g⁻¹) is about 19 times that of $Co₃O₄$ nanocubes (1.8 mg g⁻¹). In previous result we have shown that $Co₃O₄$ is non-conductive material (Supporting Information, Figure S5). We believe that the higher adsorption performance of $Co₃O₄$ nanoplates than that of $Co₃O₄$ nanocubes lead to its better performance in electrochemical detection. It is important to emphasize that as the adsorption experiments were done in solution with dispersed $Co₃O₄$, while the $Co₃O₄$ were just coated onto the surface of the GCE, the electrochemical sensitivities for two types of $Co₃O₄$ would be not proportional to the adsorption capacities of them.

The above results and discussions experimentally show that $Co₃O₄$ nanocubes and nanoplates exhibit obviously different electrochemical sensing and adsorption performances. $Co₃O₄$ nanoplates with (111) planes exhibit larger sensitivity and adsorption capacity for heavy metal ions. Besides, we suggest that Pb(II) on the (111) facet are easier to diffuse onto GCE for reduction and

Figure 3 | Adsorption measurements. Adsorption isotherms of Pb(II) onto (a), Co₃O₄ nanoplates and (b), Co₃O₄ nanocubes.

stripping reaction than those on the (001) facet. Regarding the above surface area analysis of two kinds of $Co₃O₄$ (nanocubes > nanoplates), this further confirms that the differences in their electrochemical sensing performances are due to the facet effect.

To better understand the effect of facet on electrochemical behavior of Pb(II) on $Co₃O₄$ nanoplates and nanocubes, we perform DFT calculations to explore the adsorption and diffusion behaviors of Pb on $Co₃O₄(001)$ and (111) surfaces. The adopted computational models are similar to the recent report on Li storage capability of $Ti₃C₂$ and $Ti_3C_2X_2$ (X = F, OH) monolayer³⁴. The adsorption configurations of Pb atom on $Co₃O₄(001)$ and (111) surfaces are optimized without any symmetry constraint. In the most energetically stable configuration of $Pb/Co_3O_4(001)$ system, the Pb atom is four-coordinated by two oxygen and two Co surface atoms of $Co₃O₄(001)$ surface (Supporting Information, Figure S9), and the Pb-O and Pb-Co bond lengths are 2.40 and 2.27 A, respectively. For $Pb/Co₃O₄(111)$ adsorption system, the Pb atom locates at the three-fold hollow site (Supporting Information, Figure S10), and the Pb-O bond length is 2.27 Å. Two corresponding adsorption configurations of Pb atom on $Co₃O₄(001)$ and (111) surfaces are illustrated in top panel of Figure 4 labeled with i and ii, respectively. The calculated adsorption energy of Pb atom on $Co₃O₄(111)$ surface is -4.11 eV, which is larger than that on $Co₃O₄(001)$ surface (-3.70 eV). The relative large

adsorption energy for $Pb/Co₃O₄(111)$ system mainly originates from the relative short Pb-O bond length. The Bader charge analysis³⁵ (Supporting Information, Figure S11) shows that the adsorbed Pb atom carries positive $+0.36$ and $+0.10$ |e| (e, the electron charge) on $Co₃O₄(111)$ and (001) surfaces, respectively, which gives the quantitative proof of the relative strong Pb-substrate interaction for Pb/ $Co₃O₄(111)$ system. Note that there are four equivalent stable adsorption sites in a $Co₃O₄(001)$ $\sqrt{2} \times \sqrt{2}$ cell and $Co₃O₄(111) - (1 \times 1)$ cell (Supporting Information, Figure S9 and S10), but the area of the former cell is larger than that of the latter one. In addition, the adsorption energy for Pb atom locating at the second stable adsorption site (labeled with II' symbol within twelve cyan circles in Supporting Information, Figure S10c) on $Co₃O₄(111)$ surface is predicted to be -3.81 eV, which is also larger than that of the most stable adsorption configuration of Pb/ $Co₃O₄(001)$ system. On the other hand, we have examined the full-coverage of Pb on the stable adsorption sites. The adsorption energy of each Pb on $Co₃O₄(001)$ and (111) surfaces is predicted to be -3.38 and -3.69 eV, respectively. It is clear that they are consistent with the calculated results of low-coverage of Pb on $Co₃O₄$ facets. Thus, the Pb atom adsorption ability on $Co₃O₄(111)$ surface is significantly larger than that on $Co₃O₄(001)$ surface. Moreover, the high adsorption energies and the large charge transfer suggest that

Figure 4 | DFT calculations. Top panel: side view of the optimized stable adsorption and transition-state (TS) structures for Pb on Co₃O₄(001) and (111) surfaces. Bottom panel: the TS barriers of Pb on $Co₃O₄(001)$ and (111) surfaces, and top view of these optimized stable adsorption and TS structures are also shown. After overcoming the TS barriers, Pb reaches the nearest neighboring stable adsorption site.

the adsorbed Pb can form a strong Coulomb interaction with the $Co₃O₄$ surface.

Now we turn to calculate the energy barriers of Pb diffusing on $Co₃O₄(001)$ and (111) surfaces using the climbing image nudged elastic band (CI-NEB) method³⁶. The Pb has two diffusion paths from the most stable adsorption site to its nearest neighboring stable site on $Co₃O₄(001)$ surface, while there are six diffusion directions for $Pb/Co₃O₄(111)$ adsorption system (Supporting Information, Figure S12). The geometric structures of the transition-states for Pb on $Co₃O₄(001)$ and (111) surfaces are shown in Figure 4 labeled with TS-i and TS-ii, respectively. The calculated diffusion energy barrier of Pb on $Co₃O₄(001)$ and (111) surfaces are predicted to be 1.12 and 0.74 eV, respectively. The relative low transition-state barrier results in the Pb fast diffusing on $Co₃O₄(111)$ surface. These DFT results confirm that $Co₃O₄(111)$ surface can capture more Pb than $Co₃O₄(001)$ surface, and the adsorbed Pb diffuse more easily on the $Co₃O₄(111)$ surface, consistent well with experimental observations.

Based on the above experimental and DFT calculated results, a schematic illustration of how adsorptive nanoscale materials with different exposed crystal facets could be designed for electrochemical sensing interface is depicted in Figure 5. As for bare GCE, there are not nanoscale materials which can effectively capture HMIs and the weaker stripping peak is obtained. Large amount of HMIs could be adsorbed onto the surface of nanoscale materials and then diffuse to the GCE surface. The more target HMIs are adsorbed onto the surface of nanoscale material, the more HMIs diffuse to the GCE and the stronger the stripping peak current obtains as a result. As $Co₃O₄$ nanoplates exposed with (111) facet can adsorb more metal ions than $Co₃O₄$ nanocubes with (001) facet and the metal ions on the surface of (111) facet can diffuse more easily onto the GCE, $Co₃O₄$ nanoplates modified GCE obtain larger stripping peak current. In contrast to previous studies $37-41$ focusing on elevating the specific surface area of nanoscale materials in electrochemical sensing; our findings demonstrate that the construction of nanoscale materials onto electrode with well-defined exposed facets is crucial when considering the design of high-performance electrochemical sensor.

In summary, $Co₃O₄$ nanocubes and nanoplates have been successfully synthesized by facile methods by using different solvents. The predominantly exposed facets are (001) in the $Co₃O₄$ nanocubes and (111) in the Co₃O₄ nanoplates, respectively. Studies of their sensing properties revealed that the as-prepared $Co₃O₄$ nanocrystals exhibit interesting facet-dependent electrochemical behaviors toward HMIs based on their adsorption behaviors. The (111)-bound $Co₃O₄$ nanoplates are superior to (001) -bound Co₃O₄ nanocubes. Adsorption measurements revealed that $Co₃O₄$ (111) facets can adsorb more metal ions than the (001) facets. DFT calculations suggest that $Co₃O₄(111)$ facets exhibit a relative larger adsorption energy, more adsorption sites, and a relative lower transition-state barrier than $Co₃O₄(001)$ facets, which may be the predominant reason accounting for the facet-dependent electrochemical sensing behavior toward HMIs. Both adsorption experiments and DFT results agree well with the electrochemical sensing results. This study not only suggests a promising new strategy for designing high performance electrochemical sensing interface through the selective synthesis of nanoscale materials exposed with different well-defined facets, but also provides a deep understanding for a more sensitive and selective electroanalysis at nanomaterials modified electrodes.

Methods

Materials. All reagents were commercially available from Sinopharm Chemical Reagent Co., Ltd (China) with analytical grade. Stock solution used in electrochemical measurement of Pb(II) was prepared by dissolving $Pb(NO₃)₂$ in deionized water. 0.1 M acetate buffer (NaAc-HAc) solution of $pH = 5.0$ was prepared with NaAc and HAc. Ultrapure fresh water obtained from a Millipore water purification system (MilliQ, specific resistivity > 18 M Ω cm, S.A., Molsheim, France) are used in all runs.

Synthesis of Co_3O_4 nanocubes. 0.25 g of $Co(CH_3COO)_2$ ^{-4H₂O was loaded into a} 23 mL poly(tetrafluoroethylene) (PTFE)-lined stainless steel autoclave, which was then filled with 18 mL water. The autoclave was sealed and maintained at 200° C for 12 h, and then cooled down to room temperature. The final products were centrifuged, rinsed with distilled water and ethanol for several times to remove any impurities.

Synthesis of Co_3O_4 nanoplates. 0.25 g of $Co(CH_3COO)_2$ · $4H_2O$ was loaded into a 23 mL poly(tetrafluoroethylene) (PTFE)-lined stainless steel autoclave, which was then filled with 18 mL ethylene glycol. The autoclave was sealed and maintained at 200° C for 12 h, and then cooled down to room temperature. The final products were centrifuged, rinsed with distilled water and ethanol several times to remove any impurities. The as-prepared precursors were finally calcined at 350° C in air for 3 h.

Characterization. The SEM images were taken by a FEI Quanta 200 FEG field emission scanning electron microscope. The TEM and HRTEM images analyses were carried out on a JEM-2010 microscope. XRD was performed on a D/MaxIIIA X-ray diffractometer (Rigaku Co., Japan), using Cu K α (λ K α 1 = 1.5418 Å) as the radiation source. The nitrogen adsorption and desorption isotherms at 77 K were measured with a Micromeritics ASAP 2020 M analyzer. The Brunauer, Emmett, and Teller (BET) equation was used to obtain the specific surface areas. The Pb(II) concentrations were determined in the liquid phase using inductively coupled plasma atomic emission spectrometry (ICP-AES, Jarrell-Ash model ICAP 9000). Electrochemical experiments were recorded using a CHI 660D computer-controlled

Figure 5 Schematics of how adsorptive nanoscale materials exposed with different crystal facets could be designed to enhance the performance of electrochemical sensing. (a), Bare GCE. (b), $Co₃O₄$ nanocubes modified GCE. (c), $Co₃O₄$ nanoplates modified GCE. Facet-dependent electrochemical properties of Co₃O₄ nanocrystals toward heavy metal ions (e.g. Pb²⁺) are investigated. The (111)-bound Co₃O₄ nanoplates are superior to (001)-bound $Co₃O₄$ nanocubes. Adsorption measurements reveal that $Co₃O₄$ (111) facet can adsorb more metal ions than the (001) facet. DFT calculations demonstrate that Pb on Co₃O₄(111) facet exhibits larger adsorption energies, more adsorption sites, and faster diffusion than on (001) facet.

potentiostat (ChenHua Instruments Co., Shanghai, China). A conventional threeelectrode system consisted of a glassy carbon working electrode (GCE, 3 mm diameter), an Ag/AgCl as the reference electrode and a platinum wire as the counter electrode.

Fabrication of modified electrode. 2 mg of the as-prepared $Co₃O₄$ nanocubes and nanoplates were dispersed in 4 mL water with ultrasonic agitation to give a homogeneous solution. Prior to each modification, the bare GCE was sequentially polished with 1.0 µm and 0.05 µm alumina power slurries to a mirror shiny surface, and then sonicated with $1:1$ HNO₃ solution, absolute ethanol and deionized water. The construction of $Co₃O₄$ on the surface of GCE was performed as follows: 5.0 µL of $Co₃O₄$ solution dripped onto the surface of a freshly polished GCE and then evaporating it at room temperature in the air.

Electrochemical measurements. The electrochemical measurement was carried out in SWASV mode for Pb(II) detection in 0.1 M NaAc-HAc solution. A deposition potential of -1.2 V was applied for 180 s to the working electrode under stirring. The SWASV responses were recorded between -1.0 to -0.2 V with step potential of 5 mV, amplitude of 20 mV, and frequency of 25 Hz. A desorption potential of 0 V for 210 s was performed to remove the residual metals under stirring condition.

Adsorption experiments. Experiments were carried out at 298 K in 10 mL polyethylene centrifuge tubes containing 1 $g \cdot L^{-1}$ adsorbent and various concentrations of Pb(II) for 24 h. The pH values of all these Pb(II) solutions 5.0 \pm 0.2. After adsorption equilibrium, the adsorbent was separated by centrifugation. The Pb(II) concentrations remaining in the solution were analyzed. The amount of metal adsorbed (q_e) was calculated according to the following equation:

$$
q_e = \frac{(C_0 - C_e)V}{m} \tag{3}
$$

where C_0 and C_e represent the initial and equilibrium Pb(II) concentrations (mg L^{-1}), respectively, V is the volume of the solutions (mL), and m is the amount (mg) of adsorbent. A Langmuir isotherms model was used to analyze the experimental data. The mathematical expressions of the Langmuir isotherm is

$$
\frac{C_{\rm e}}{q_{\rm e}} = \frac{1}{q_{\rm m}K_{\rm L}} + \frac{C_{\rm e}}{q_{\rm m}}\tag{4}
$$

where q_m and K_L are Langmuir constants, representing the maximum adsorption capacity of adsorbents (mg/g) and the energy of adsorption, respectively. The values of $q_{\rm m}$ and $K_{\rm L}$ can be calculated from the slope and intercept of plots of $\frac{C_{\rm e}}{q_{\rm e}}$ versus $C_{\rm e}$.

Computational methods. The DFT calculations were performed by using the Vienna
ab-initio simulation package (VASP)^{42,43} with the Perdew-Burke-Ernzerh parameterization of the generalized gradient approximation (GGA) adopted for the exchange correlation potential⁴⁴. An energy cut of 400 eV were consistently used in our calculations. The atomic positions are fully relaxed with the conjugate gradient procedure until the residual forces vanished within 0.02 eV/Å. A $2 \times 2 \times 1$ Monkhorst-Pack kpoint was used to sample the surface Brillouin zone. Because $Co₃O₄$ is a strongly correlated antiferromagnet and Pb ion is one of HMIs, we also performed test $\overline{\rm DFT}$ + $U^{45,46}$ (U = 3.3 eV and J = 0 eV according to previous reports⁴⁷) calculations as well as PBE calculations with relativistic corrections⁴⁸ to Pb. As shown in Table S1, we can confirm that our PBE results are qualitatively correct. The transition states were obtained by using the climbing image nudged elastic band (CI-NEB) method³⁶.

Computational models. The optimized lattice constant (a_0) of the bulk Co_3O_4 is 8.12 Å, which is close to the experimental value $(8.08 \text{ Å})^{47}$. The bulk Co_3O_4 shows antiferromagnetic properties. The $1 \times 1 \times 10$ and $1 \times 1 \times 11$ supercells were used to model the $Co₃O₄$ (001) and (111) surfaces, respectively, in which a 12 Å of vacuum is adopted. During the structural optimizations, we allowed the atoms in the top four layers in the supercells to relax fully, and fixed all other atoms in the bulk configuration. The side and top views of the optimized $Co₃O₄$ (001) and (111) surfaces are shown in Figures S9 and S10, respectively. These optimized geometric structures agree well with the previous reports on the surface atoms arrangement of the Co₃O₄ nanocrystals with different crystal planes^{10,47,49}.

Adsorption energy. To compare the Pb adsorption ability on two different crystal planes, we define the Pb adsorption energy as $E_{ads} = E_{Pb/surface} - (E_{surface} + E_{Pb})$. Here, E_{surface} and $E_{\text{Pb/surface}}$ are the total energies of the surface and the Pb atom adsorbing on surface, respectively, and $E_{\rm Pb}$ is the atomic energy of single isolated Pb atom. Under this definition, the more negative value stands for the more energetically stable adsorption.

Bader charge analysis and diffusion paths. To address the Pb-substrate interaction, we conduct the Bader charge analysis for Pb atom adsorbing on $Co₃O₄$ (111) and (001) surfaces. The corresponding results are presented in Figure S11. For clarity, the diffusion directions of Pb on $Co₃O₄$ (111) and (001) surfaces are illustrated in Figure S12.

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Author contributions

X.J.H. conceived the electrochemical studies of nanocrystals with different facets. Q.X.L. designed the calculations. X.Y.Y. carried out all the experiments and Q.Q.M. performed calculation and contributed equally to this work. X.Y.Y. and Q.Q.M. co-wrote the paper. T.L., Y.J., B.S., and J.H.L. contributed data analysis and interpretation equally.

Additional information

Supplementary information accompanies this paper at [http://www.nature.com/](http://www.nature.com/scientificreports) [scientificreports](http://www.nature.com/scientificreports)

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