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## The Role of CHPD and AIMI processing on enhancing $J_C$ and transverse connectivity of *in-situ* MgB<sub>2</sub> strand

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### Abstract

Research into *in-situ* MgB<sub>2</sub> strand has been focused on improvements in  $J_C$  through reduction of porosity. Both of cold-high-pressure-densification (CHPD) and advanced-internal-magnesium-infiltration (AIMI) techniques can effectively remove the voids in *in-situ* MgB<sub>2</sub> strands. This study shows the nature of the reduced porosity for *in-situ* MgB<sub>2</sub> strands lies on increases in transverse grain connectivity as well as longitudinal connectivity. The CHPD method bi-axially applying 1.0 GPa and 1.5 GPa yielded 4.2 K  $J_{C_{mid}}$ s of  $9.6 \times 10^4$  A/cm<sup>2</sup> and  $8.5 \times 10^4$  A/cm<sup>2</sup> at 5 T, respectively, with compared with  $6.0 \times 10^4$  A/cm<sup>2</sup> for typical powder-in-tube (PIT) *in-situ* strand. Moreover, AIMI-processed monofilamentary MgB<sub>2</sub> strand obtained even higher  $J_C$ s and transverse grain connectivity than the CHPD strands.

### 1. Introduction

MgB<sub>2</sub> superconducting strands are promising to the practical magnetic application due to its high transition temperature  $T_C$  (39 K) [1], high coherence length [2, 3], and low anisotropy of upper critical fields ( $B_{c2}$ ) [3–6]. The powder-in-tube (PIT) *in-situ* MgB<sub>2</sub> strands were fabricated by filling a mixture of Mg and B powder into a non-reactive metallic tube and then being cold-worked into wires or tapes. The PIT strands have large amount of voids elongated along longitudinal strand axis which were left behind by molten Mg powders after heat treatment [7]. The present of voids tends to limit the number of the continuous current path in the *in-situ* MgB<sub>2</sub> strand and therefore suppress the current-carrying-capacity of the strand.

The CHPD technique can effectively increase the transport properties of PIT *in-situ* MgB<sub>2</sub> strands by eliminating the pores [8–11]. The pre-reacted powder-in-tube composite was bi-axially cold-densified at room temperature to increase the mass density of the Mg + B mixture. In this case, the cold-densified MgB<sub>2</sub> strand can obtain higher grain connectivity after heat treatment. Additionally, the Mg reactive – liquid – infiltration (RLI) process,

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initiated by Giunchi *et al* [12], also has the ability to eliminate the pores and produce a dense  $\text{MgB}_2$  layers in *in-situ*  $\text{MgB}_2$  strands. For RLI process, a Mg rod is inserted axially into a boron-filled metallic tube. After wire drawing the heat treatment (H. T.),  $\text{MgB}_2$  layer is formed through the reactive diffusion of Mg into B layer. Since Mg is totally separated with precursor B layer before H. T., the RLI process can totally eliminate the “Mg-site porosity” from  $\text{MgB}_2$  layer but induce the formation of a big hole at the central region of the strand [13]. Furthermore, since the molar volume of  $\text{MgB}_2$  ( $17.46 \text{ cm}^3/\text{mol}$ ) is twice as that of two B atoms ( $9.18 \text{ cm}^3/\text{mol}$ ), the volume expansion associated with the reactive transformation from  $2\text{B}$  to  $\text{MgB}_2$  enables even better connections between  $\text{MgB}_2$  grains during heat treatment [13]. Our group named our RLI-processed strands as advanced – internal – magnesium – infiltration (AIMI) strands due to their optimized strand architecture and high  $J_{CS}$  [13].

The previous researches mostly focused on investigating the effect of CHPD and AIMI techniques on the transport properties along longitudinal strand axis of *in-situ*  $\text{MgB}_2$  strands, such as transport  $J_{CS}$  and longitudinal grain connectivity. However, Shi and Susner pointed out that high anisotropic grain connectivity exists in PIT *in-situ* strands, which was resulted from the elongated voids between elongated  $\text{MgB}_2$  stringers [7]. Therefore longitudinal transport properties are different with the transport properties along transverse strand axis for  $\text{MgB}_2$  superconducting wires. In continuing along these lines, the transport properties along transverse strand axis were investigated for the 2.0 mol% C-doped *in-situ*  $\text{MgB}_2$  strands in this study. The anisotropic connectivity of the PIT strands results in the differences between perpendicular magnetic  $J_C$  ( $J_{CM\perp}$ ) and parallel  $J_C$  ( $J_{CM\parallel}$ ). The influence of aspect ratio ( $S = \text{length}/\text{diameter}$ ) on transverse and longitudinal  $J_C$  were investigated for the PIT strand (P00). The CHPD-processed strands were P10 (1.0 GPa cold-pressing) and P15 (1.5 GPa cold-pressing). According to the previous results of our group, the CHPD technique increased the transport  $J_C$  of the monofilamentary PIT *in-situ*  $\text{MgB}_2$  strand from  $3.0 \times 10^4 \text{ A/cm}^2$  to  $3.6 \times 10^4 \text{ A/cm}^2$  at 4.2 K and 10 T due to decreased porosity [11] and AIMI-processed  $\text{MgB}_2$  strands attained the 4.2 K, 10 T transport layer  $J_{CS}$  of  $1.0 \sim 1.5 \times 10^5 \text{ A/cm}^2$  [13–16]. In this study, we compared the  $J_{CM\parallel}$ s and transverse connectivity of the CHPD- and AIMI-processed strands with those of the P00 strand at 4.2 K and 20 K. The relationship between porosity and transverse flux pinning force density  $F_{p\parallel}$ , which is  $J_{CM\parallel} \times B$ , for the *in-situ*  $\text{MgB}_2$  strands was also discussed.

## 2. Experimental

### 2.1. Sample preparation

A series of pre-reacted powder-in-tube (PIT) *in-situ* strands, typically 0.834 mm diameter, with a Nb barrier and a Cu outer sheath were fabricated by Hyper Tech Research, Inc. (HTR). Two PIT strands (designated P10 and P15) were bi-axially densified with 1.0 GPa and 1.5 GPa at room temperature, respectively. The other strand (designed A00) manufactured through AIMI technique were also provided by HTR. The AIMI-processed strand, with 0.55 mm diameter, has a Nb barrier and a Monel outer sheath. The powders used for the present strands were 2 mol% C-doped amorphous B (10 – 100 nm) from Specialty Materials Inc. (SMI). The cold-densified PIT strands were heat-treated at  $675 \text{ }^\circ\text{C}$

for 1 h and the AIMI strand was heat – treated at 625 °C for 16 h. The specification and heat treatment (H. T.) conditions of the strands are presented in table 1.

## 2.2. Transport and Magnetic Measurements

The transport  $I_C$  ( $I_{CT}$ ) test was conducted in perpendicular magnetic field up to 13 T in a pool of liquid Helium at 4.2 K on the MgB<sub>2</sub> strands with a total length of 50 mm and a gauge length of 5 mm. The electric criterion used for determining  $I_{CT}$ s is 1.0 μV/cm. The magnetizations versus perpendicular and parallel magnetic fields ( $M-H$ ) loops were measured by a Quantum Design Model 6000 Physical Property Measuring System (PPMS) for all strands with a sample length of 3 – 5 mm.

## 3. Results

The  $J_{CT}$ s of the PIT strands were the transport critical current normalized by MgB<sub>2</sub> core area. As shown in Figure 1(a) the MgB<sub>2</sub> core of the typical PIT *in-situ* strand is a solid cylinder. Figure 1(b)–(d) shows the shape of MgB<sub>2</sub> cores for the CHPD- and AIMI-processed strands are cuboid and hollow cylinder, respectively. The  $J_{CT}$  of the AIMI-processed strands, which is also named as transport layer  $J_C$ , were calculated by dividing the  $I_{CT}$  by the area of annulus MgB<sub>2</sub> layer. Values of magnetic  $J_C$  ( $J_{CM\perp}$  and  $J_{CM\parallel}$ ) for the MgB<sub>2</sub> strands were extracted from the full  $M-H$  loops heights  $M$ , using the standard Bean model equations [7, 17]:

For the PIT *in-situ* wire P00:

$$\text{Perpendicular Magnetic } J_C: J_{CM\perp} = \frac{3\pi\Delta M}{8R_0} \quad (1)$$

$$\text{Parallel Magnetic } J_C: J_{CM\parallel} = \frac{3\Delta M}{2R_0} \quad (2)$$

Here  $R_0$  is the radius of the cylinder MgB<sub>2</sub> core in the PIT wire.

For the densified wires P10 and P15:

$$\text{Parallel Magnetic } J_C: J_{CM\parallel} = \frac{2\Delta M}{b\left(1 - \frac{b}{3a}\right)} \quad (3)$$

Here a, b are both lengths of the transverse cross – sectional area of cuboid MgB<sub>2</sub> core,  $a > b$ .

For the AIMI wire A00:

$$\text{Parallel Magnetic } J_C: J_{CM\parallel} = \frac{3\Delta M}{2} \frac{R_0^2 - R_i^2}{R_0^3 - R_i^3} \quad (4)$$

Here  $R_i$  is the inner diameter of the annulus MgB<sub>2</sub> layer and  $R_0$  is the outer diameter of the annulus MgB<sub>2</sub> layer.

### 3.1. Transport and magnetic critical current densities, $J_{CT}$ and $J_{CM}$

Figure 2(a) shows the  $J_{CT}$  and  $J_{CM_{\perp}}$  versus  $B$  at 4.2 and 20 K for the strand P00. It can be seen that  $J_{CM_{\perp}}$ s agree with  $J_{CT}$  at low fields, whereas the bifurcation of  $J_{CT}$  and  $J_{CM_{\perp}}$  happened at high fields. Moreover,  $J_{CM_{\perp}}$ s were greatly affected by the aspect ratio  $S$ , especially at high magnetic fields. The relationships among  $J_{CT}$ ,  $J_{CM_{\perp}}$ , and aspect ratio were fully discussed in ref [7] and [18]. Therefore, the values of  $J_{CM_{\perp}}$  are not only determined by the intrinsic properties but also the extrinsic properties of the *in-situ* MgB<sub>2</sub> strands. As shown in Figure 2(b)  $J_{CM_{\parallel}}$ s were independent on the aspect ratio  $S$ . Based on this result, we can know the values of  $J_{CM_{\parallel}}$ s as well as  $J_{CT}$ s are merely determined by the intrinsic properties of the *in-situ* MgB<sub>2</sub> strands. The MgB<sub>2</sub> macrostructure of the reacted PIT *in-situ* strand is characterized by elongated polycrystalline MgB<sub>2</sub> fibers and elongated pores partially separating the fibers [7, 11, and 18]. The  $J_{CT}$ s and  $J_{CM_{\parallel}}$ s can represent the longitudinal and transverse connectivity of the MgB<sub>2</sub> fibers, respectively. In summary, we can investigate the effect of CHPD and AIMI technique on the current-carrying-capacity of the *in-situ* MgB<sub>2</sub> strands through  $J_{CM_{\parallel}}$ s as well as  $J_{CT}$ .

Figure 3 shows the  $J_{CT}$  and  $J_{CM_{\parallel}}$  versus  $B$  at 4.2 and 20 K for all strands. The CHPD technique significantly enhanced  $J_{CM_{\parallel}}$ s of PIT *in-situ* strands at 4.2 K. Strand P10 and P15 attained  $9.5 \times 10^4$  A/cm<sup>2</sup> and  $8.5 \times 10^4$  A/cm<sup>2</sup> at 4.2 K and 5 T, with respect to  $6.0 \times 10^4$  A/cm<sup>2</sup> for strand P00. 4.2 K  $J_{CT}$ s were slightly enhanced by the CHPD technique. Therefore, the increases in  $J_{CT}$  and  $J_{CM_{\parallel}}$  indicated that the both longitudinal and transverse connections between MgB<sub>2</sub> fibers were enhanced. With the formation of high density MgB<sub>2</sub> layer, strand A00 attained 4.2 K, 10 T  $J_{CT}$  of  $9.4 \times 10^4$  A/cm<sup>2</sup>, which is 180% higher than those of CHPD-processed strands. On the other hand, the AIMI strand obtained 5 T  $J_{CM_{\parallel}}$ s of  $3.1 \times 10^5$  A/cm<sup>2</sup> at 4.2 K and  $1.7 \times 10^4$  A/cm<sup>2</sup> at 20 K, which are about 240% higher than those of CHPD strands.

### 3.2. Transverse Flux Pinning, Porosity and Transverse Grain Connectivity

Figure 4 shows the  $F_{p\parallel}/F_{p,max\parallel}$  versus  $B/B_{irr}$  for all strands, where  $F_{p\parallel} = J_{CM_{\parallel}} \times B$ . It can be seen that the peak pinning occurred at  $b = B/B_{irr}$  close to 0.2 at 4.2 and 20 K, which is in agreement with the Dew-Hughes/Kramer model [19, 20]. In other words, the dominant pinning centers for the all strands are also grain boundaries for the direction along transverse strand axis. According to the previous work [11], the densified wires have decreased porosities, the values of porosity are  $\sim 50\%$  ( $p = 0$ ) and  $\sim 30\%$  ( $p = 1.0$  or  $1.5$  GPa). Since the Mg-site porosity was totally eliminated for AIMI-processed strands, the porosity of the AIMI strand is close to 0. Therefore, both CHPD and AIMI processes enable the resulting MgB<sub>2</sub> phases to be denser and more connected. For the discussion, it can be concluded that the enhanced  $J_{CT}$  and  $J_{CM_{\parallel}}$  in densified wires and AIMI wire is correlated with lower porosity (higher grain connectivity).

The connectivity  $K$  defined by Rowell [21] can represent the grain connectivity of MgB<sub>2</sub> strands. The connectivity  $K$  is calculated by the equation:

$$K = \frac{\Delta\rho_{SC}}{\Delta\rho} \quad (5)$$

Here  $\rho$  is the difference between the sample's resistivity at 300 K and the sample's resistivity at 40 K and  $\rho_{SC}$  is the resistivity difference for an ideal single crystal. However, transverse connectivity  $K_{\perp}$  is difficult to be determined for the MgB<sub>2</sub> wires. It has been reported that the maximum flux pinning force densities of fully-connected MgB<sub>2</sub> superconductor, where grain connectivity  $K = 100\%$ , are estimated to be 90 GN/m<sup>3</sup> at 4.2 K and 22 GN/m<sup>3</sup> at 20 K [22]. Therefore, we can roughly estimate the transverse grain connectivities  $K_{\perp}$ s of the *in-situ* MgB<sub>2</sub> strands by normalizing the  $F_{p,max\parallel}$  with the  $F_{p,max}$  of fully-connected MgB<sub>2</sub>. Table 2 shows the transverse maximum flux pinning force densities  $F_{p,max\parallel}$  at 4.2 K and 20 K and estimated transverse grain connectivities for all the strands. The connectivity of 5% is achieved by the 2 mol% C-doped PIT strand, P00. The cold-densification increases the connectivity of PIT strand by 20%. The AIMI strand A00 obtained the highest transverse grain connectivities of 20%.

#### 4. Conclusion

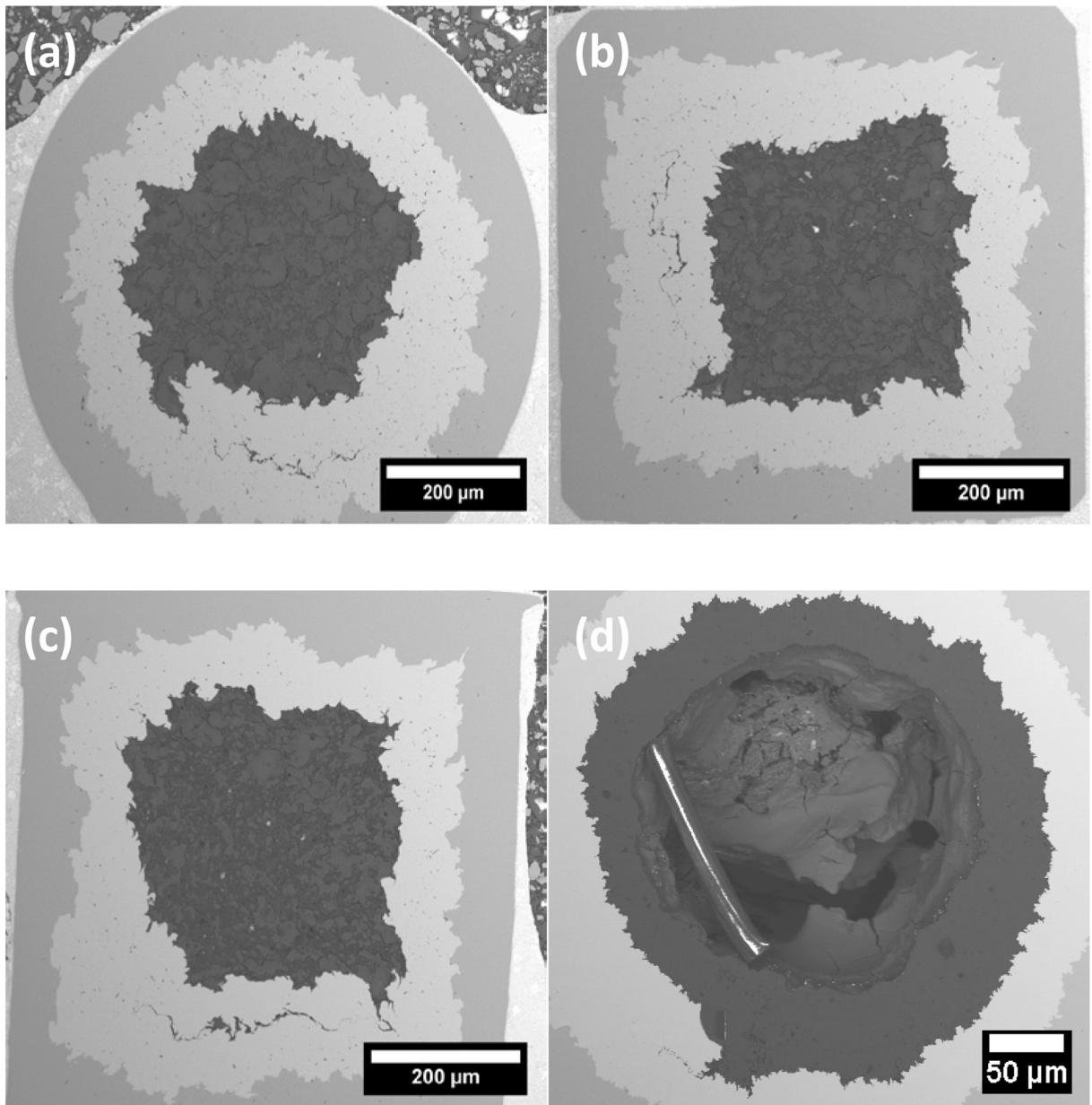
$J_{CM\parallel}$ s is merely determined by the intrinsic properties of *in-situ* MgB<sub>2</sub> strands, so the current-carrying capacity of *in-situ* MgB<sub>2</sub> strands can be represented by  $J_{CM\parallel}$  as well as  $J_{CT}$ . The CHPD of 1.0 GPa and 1.5 GPa enhanced the 4.2 K, 5 T  $J_{CM\parallel}$  from  $6.0 \times 10^4$  A/cm<sup>2</sup> to  $9.6 \times 10^4$  A/cm<sup>2</sup> and  $8.5 \times 10^4$  A/cm<sup>2</sup>, respectively. AIMI strand attained the highest  $J_{CT}$  and  $J_{CM\parallel}$  at 4.2 and 20 K due to the formation of a high dense MgB<sub>2</sub> layer. By eliminating the voids in *in-situ* MgB<sub>2</sub> strands through CHPD and AIMI technique, better connections between MgB<sub>2</sub> grains along transverse strand axis can be obtained in *in-situ* MgB<sub>2</sub> strands

#### 5. References

- [1]. Nagamatsu J, Nakagawa N, Muranaka T, Zenitani Y, and Akimitsu J2001Nature41063–64 [PubMed: 11242039]
- [2]. Finnemore DK, Ostenson JE, Bud'ko SL, Lapertot G and Canfield PC 2001 Rhys. Rev. Lett 86 2420–2422
- [3]. Vinod K, Abhilash Kumar RG and Syamaprasad U 2007 Supercond. Sci. Technol 20 R1–R3
- [4]. Eltsev Y, Lee S, Nakao K, Chikumoto N, Tajima S, Koshizuka N and Murakami M 2002 Phys. Rev. B 65 140501
- [5]. Eisterer M, Zehetmayer M and Weber HW 2003 Phys. Rev. Lett 90 247002 [PubMed: 12857216]
- [6]. De Liam OF, Ribeiro RA, Avila MA, Cardoso CA and Coelho AA 2001 Phys. Rev. Lett 86 5974 [PubMed: 11415407]
- [7]. Shi ZX, Susner MA, Majoros M, Sumption MD, Peng X, Rindfleisch M, Tomsic MJ and Collings EW 2010 Supercond. Sci. Technol 23 045018
- [8]. Flukiger R, Hossain MSA and Senatore C 2009 Supercond. Sci. Technol 22 085002
- [9]. Hossain MSA, Senatore C, Flukiger R, Rindfleisch MA, Tomsic MJ, Kim JH and Dou SX 2009 Supercond. Sci. Technol 22 095004
- [10]. Hossain MSA. et al. 2014; Supercond. Sci. Technol. 27 :095016.
- [11]. Wan F, Sumption MD, Rindfleisch MA and Collings EW 2017 IOP Conf. Series: Materials Science and Engineering 279 012024
- [12]. Giunchi G, Ripamonti G, Perini E, Cavallin T and Bassani E 2007 IEEE TRANSACTIONS ON APPLIED SUPERCONDUCTIVITY 17 2761
- [13]. Li GZ, Sumption MD, Susner MA, Yang Y, Reddy KM, Rindfleisch MA, Tomsic MJ, Thong CJ and Collings EW 2012 Supercond. Sci. Technol 25 115023
- [14]. Ye SJ, Matsumoto A, Zhang YC and Kumakura H 2014 Supercond. Sci. Technol 27 085012

- [15]. Li GZ, Sumption MD, Zwayer JB, Susner MA, Rindfleisch MA, Thong CJ, Tomsic MJ and Collings EW 2013 *Supercond. Sci. Technol* 26 095007
- [16]. Kumakura H, Hur J, Togano K, Matsumoto A, Wada H and Kimura K 2011 *IEEE Trans. Appl. Supercond* 21 2643
- [17]. Sumption MD, Peng X, Lee X, Wu X and Collings EW 2004 *Cryogenics* 44 711–725
- [18]. Susner MA, Daniels TW, Sumption MD, Rindfleisch MA, Thong CJ and Collings EW 2012 *Supercond. Sci. Technol* 25 065002
- [19]. Dew-Hughes D 1974 *Philos. Mag* 30 293–305
- [20]. Kramer EJ 1973 *J. Appl. Phys* 44 1360–70
- [21]. Rowell J 2003 *Supercond. Sci. Technol* 16 R17–27
- [22]. Matsushita T, Kiuchi M, Yamamoto A, Shimoyama J and Kishio K 2008 *Physica C* 468 1833





**Figure 1.**

(a) Back scattered SEM images of (a) strand P00, (b) strand P10 (1.0 GPa cold pressure), (c) strand P15 (1.5 GPa cold pressure), and (d) AIMI strand A00.

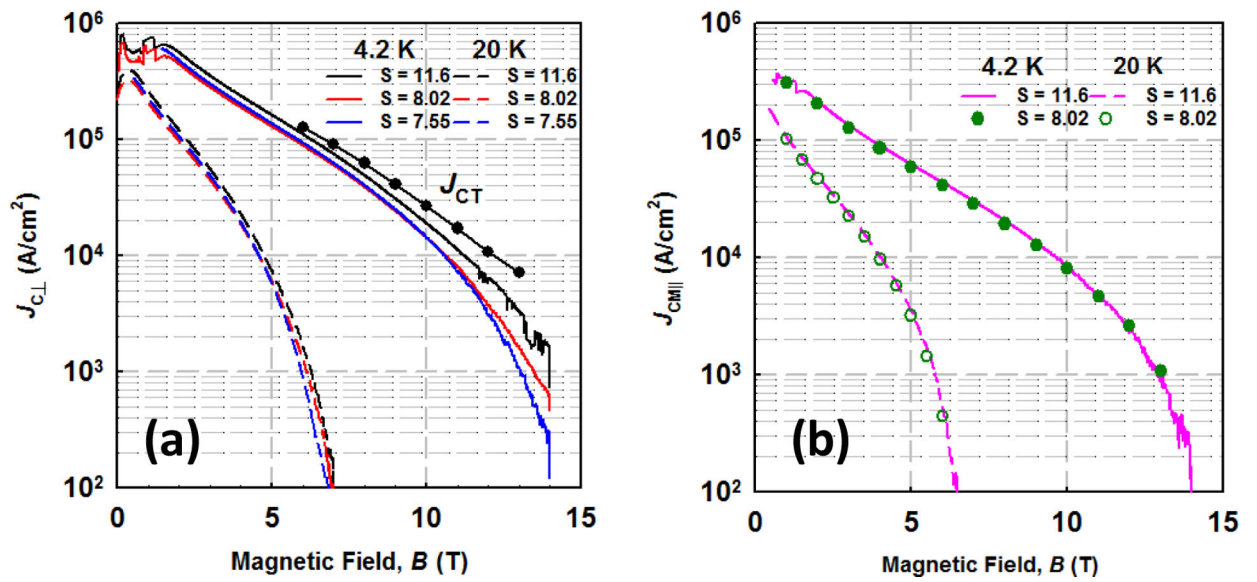


Figure 2.

(a)  $J_{CT}$  and  $J_{c\perp}$  versus  $B$  at 4.2 and 20 K (b)  $J_{c\parallel}$  versus  $B$  at 4.2 and 20 K for strand P00



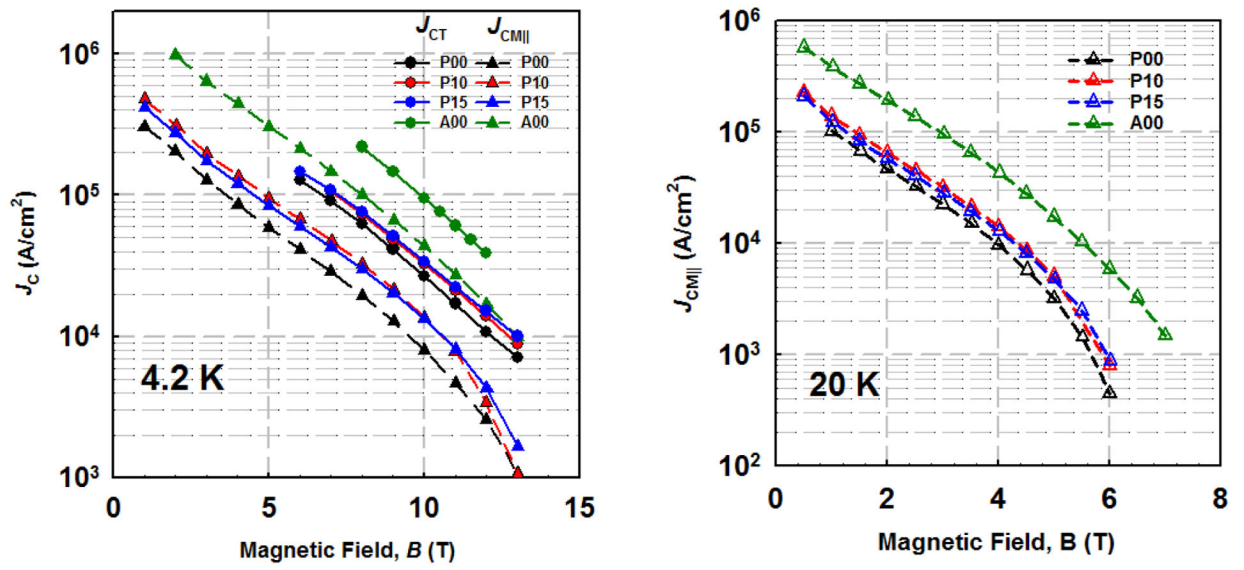
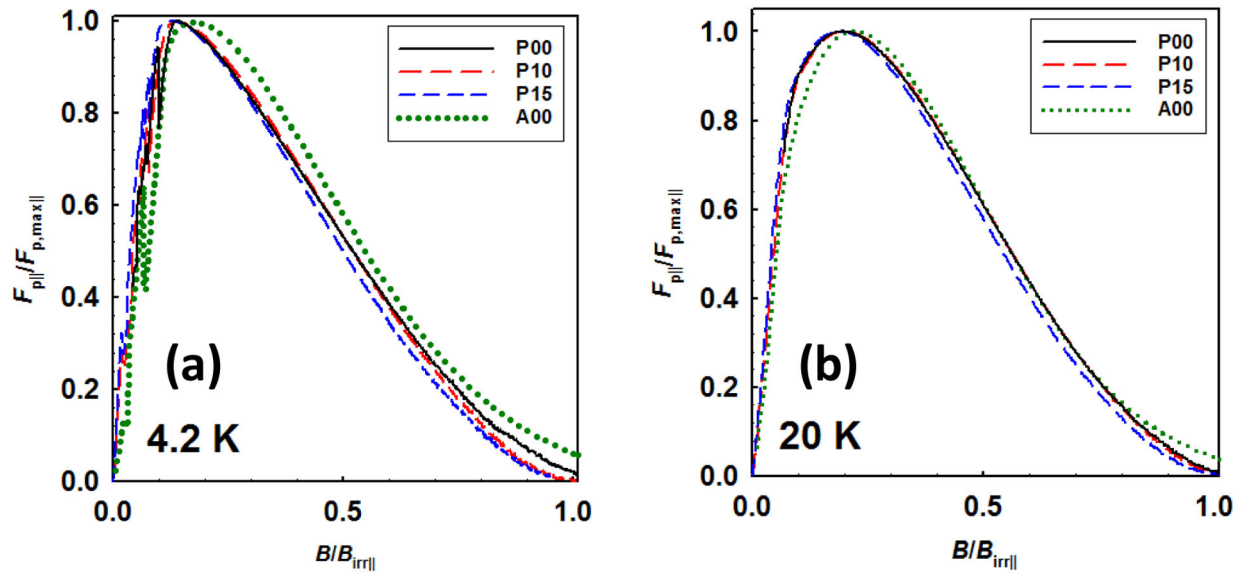
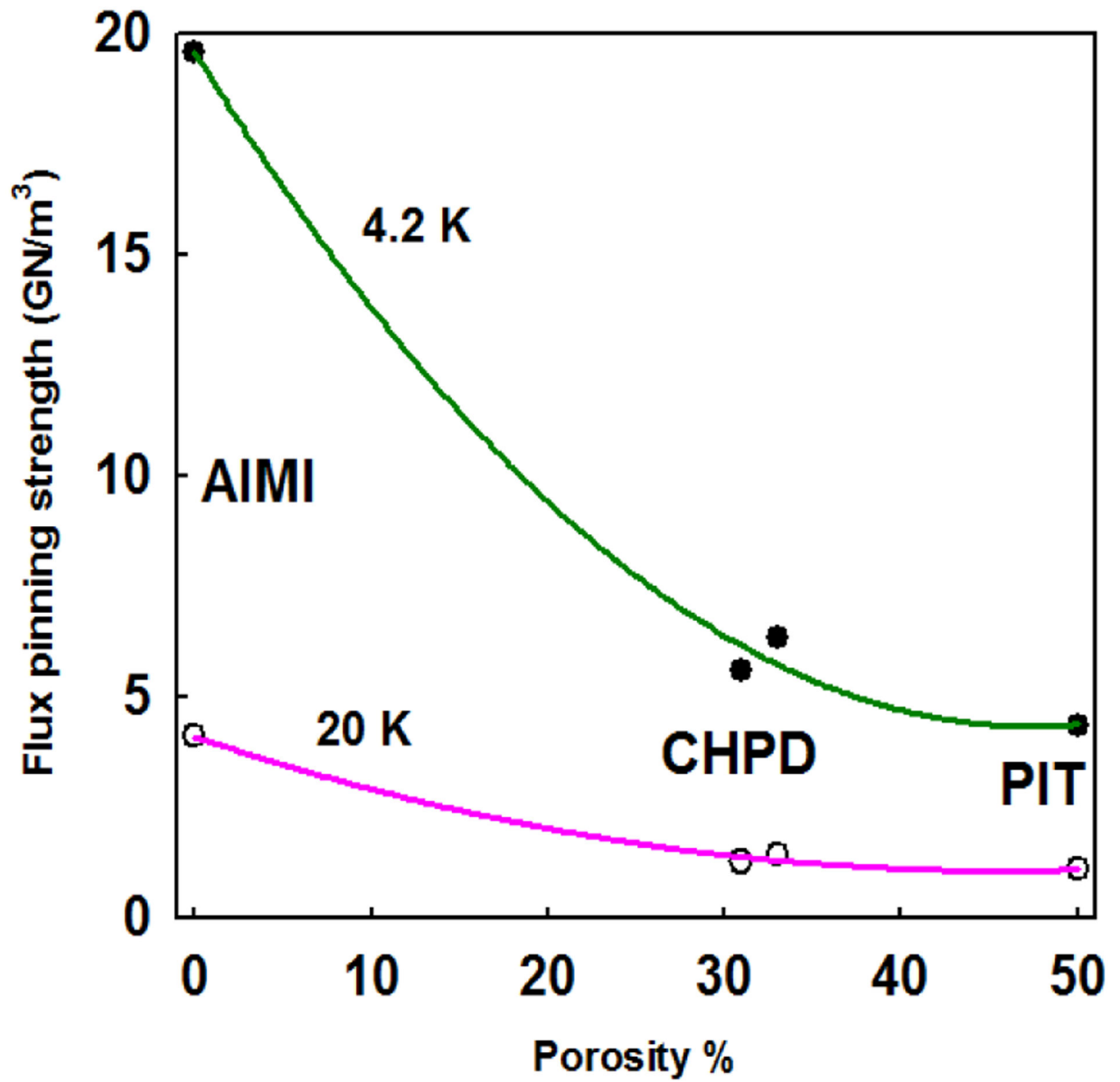


Figure 3.  
 $J_{CT}$  and  $J_{c||}$  versus  $B$  at 4.2 and 20 K for all strands



**Figure 4.**

(a)  $F_{p||}/F_{p,max||}$  versus  $B/B_{irr||}$  for all strands at 4.2 K, (b)  $F_{p||}/F_{p,max||}$  versus  $B/B_{irr||}$  for all strands at 20K



**Figure 5.** Relationship between porosity and transverse flux pinning force densities for *in-situ* MgB<sub>2</sub> strands

**Table 1.**

Strand specification and H. T. conditions

Strand Name	Strand Type	CHPD (GPa)	H. T. (h/°C)
<b>P00</b>	PIT	0.0	1/675
<b>P10</b>	PIT	1.0	1/675
<b>P15</b>	PIT	1.5	1/675
<b>A00</b>	AIMI	0.0	16/625

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**Table 2.**

Transverse flux pinning force densities and estimated connectivity of the strands.

Strand Name	$F_{p,max1}$ , 4.2 K (GN/m <sup>3</sup> )	$F_{p,max1}$ , 20 K (GN/m <sup>3</sup> )	Transverse Connectivity, $K_f$ %
P00	4.3	1.1	5.0
P10	6.3	1.4	6.0
P15	5.6	1.3	6.0
A00	19.6	4.1	20

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