

SCIENTIFIC REPORTS



OPEN

Zn Vacancy Formation Energy and Diffusion Coefficient of CVT ZnO Crystals in the Sub-Surface Micron Region

Narendra S. Parmar¹, Lynn A. Boatner², Kelvin G. Lynn^{3,4} & Ji-Won Choi^{1,5}

By using positron annihilation spectroscopy methods, we have experimentally demonstrated the creation of isolated zinc vacancy concentrations $>10^{20} \text{ cm}^{-3}$ in chemical vapor transport (CVT)-grown ZnO bulk single crystals. X-ray diffraction ω -rocking curve (XRC) shows the good quality of ZnO single crystal with (110) orientation. The depth analysis of Auger electron spectroscopy indicates the atomic concentrations of Zn and O are almost stoichiometric and constant throughout the measurement. Boltzmann statistics are applied to calculate the zinc vacancy formation energies (E_f) of $\sim 1.3\text{--}1.52 \text{ eV}$ in the sub-surface micron region. We have also applied Fick's 2nd law to calculate the zinc diffusion coefficient to be $\sim 1.07 \times 10^{-14} \text{ cm}^2/\text{s}$ at 1100°C . The zinc vacancies began annealing out at 300°C and, by heating in the air, were completely annealed out at 700°C .

Zinc oxide has been under intensive and long-term investigation in both the academic and industrial communities because of its varied actual and potential applications in electronic and optoelectronic-based devices. Point-defect engineering in this material has become of significant importance since theoretical and experimental results suggest that a zinc vacancy acts as a shallow acceptor and would, therefore, be of importance in forming a p - n junction for blue/UV LEDs. Accordingly, achieving an understanding of the thermodynamics and kinetics of intrinsic point defects in ZnO is of both fundamental and technological significance. Zinc migration has previously been considered, in particular, with respect to the degradation of varistor devices¹ whose function is believed to proceed through the migration of intrinsic defects - most likely zinc interstitials. Gaining an understanding of zinc defect diffusivities is also of importance in order to control the formation of unwanted and compensating defects that are likely to contribute to the well-known difficulties in synthesizing p -type zinc oxide^{2,3}. Until recently, the most prevalent way of creating zinc vacancies has been either by electron or laser irradiation^{4,5}, i.e., by non-equilibrium thermodynamic processes that lead to the simultaneous creation of compensating defects such as zinc interstitials, oxygen interstitials, and oxygen vacancies. The creation of such defects is relatively hard to avoid and control. Recently, Parmar *et al.*⁶, reported the formation of high concentrations of isolated zinc vacancies ($>10^{20} \text{ cm}^{-3}$) in thermodynamic equilibrium - findings that can prove to be of importance in obtaining p -type ZnO by controlling the formation of zinc vacancies⁶, as zinc vacancies are shallow acceptors in ZnO crystal.

Zinc vacancies created by electron irradiation⁷⁻⁹ and laser radiation¹⁰, have been the subject of a number of previously conducted studies designed to measure the self-diffusion of zinc in ZnO¹¹. The prior experimental data, however, exhibit a considerable spread that renders their interpretation difficult. In view of this situation, a theoretical approach can provide valuable insights into the various atomistic migration processes and, thereby, can help to quantify their respective contributions.

In this letter, we present results on the formation energy and diffusivity of zinc vacancies in ZnO, where a large number of zinc vacancies ($10^{17} < V_{\text{Zn}} \sim 10^{20}$) are created in thermodynamic equilibrium by oxygen annealing. Having such a concentration of zinc vacancies in thermodynamic equilibrium provides the basis for carrying out

¹Center for Electronic Materials, Korea Institute of Science and Technology, Seoul, 02792, Republic of Korea.

²Materials Science and Technology Division, Oak Ridge National Laboratory, Oak Ridge, TN, 37831, USA. ³Department of Physics and Astronomy, Washington State University, Pullman, WA, 99164, USA. ⁴Center for Materials Research, Washington State University, Pullman, WA, 99164, USA. ⁵Department of Nanomaterials Science and Engineering, Korea University of Science and Technology, Daejeon, 34113, Republic of Korea. Correspondence and requests for materials should be addressed to N.S.P. (email: nparmar@kist.re.kr) or J.-W.C. (email: jwchoi@kist.re.kr)

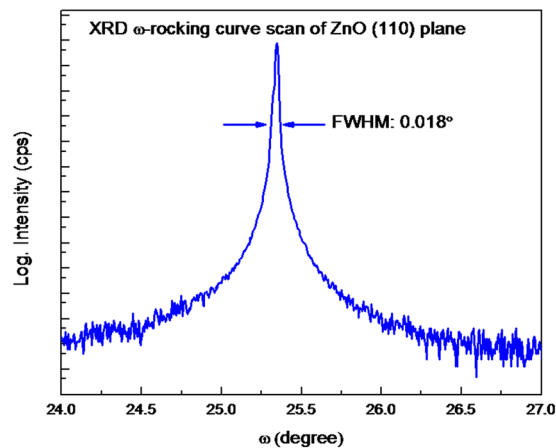


Figure 1. X-ray rocking curve (XRC) of the ZnO single crystal.

a study of the formation energy and diffusion coefficient that is more reliable relative to an approach where zinc vacancies are created by a non-thermodynamic-equilibrium process such as, an electron irradiation or bombardment using a laser exposure.

Experimental Methods

Chemical vapor-transport-grown ZnO crystals were placed in Heraeus high-purity quartz ampoules that were evacuated using roughing and turbo molecular pumps and that were then baked at $\sim 150^\circ\text{C}$ overnight to remove residual water vapor. The vacuum pressure was $\sim 10^{-7}$ Torr prior to back filling with ultra-high-purity oxygen to ~ 300 Torr. The ampoules were then sealed using an oxy-hydrogen torch and placed in a tube furnace. The annealing process was carried out at 1100°C or 1200°C for 24 hours, and the ampoule remained in the furnace during cooling.

Quenching experiments were also carried out by dipping the hot ampoule (1200°C) in water at room temperature (RT). This quenching process reduced the ZnO temperature from 1200°C to RT in a few seconds. Annealing (i.e., annihilation) out of the zinc vacancies was performed by heating the crystals in air for one hour at either 300°C , 500°C or 700°C .

Positron annihilation spectroscopy (PAS)¹² is a well-known tool to characterize negatively charged defects, such as zinc vacancies (V_{Zn}) in ZnO crystals¹³. Positrons are positively charged and become trapped in negatively charged native defects, which reduces their Doppler momentum. The trapped positrons eventually annihilate the surrounding electrons, emitting two photons of 511 keV energy¹⁴. Emitted photons with Doppler broadening are the signature of the annihilation site. Though, PAS is not very useful for investigating positively charged defects, it has been quite helpful in investigating neutral defects¹⁵. The signal-to-noise ratio can be increased by a significant amount by performing two-detector coincidence measurements¹⁶. However, the coincidence-mode process takes a much longer time to collect data. In this letter, depth-resolved positron annihilation spectroscopy (PAS) Doppler broadening measurements were performed at the Washington State University (WSU). The 511 keV annihilation peak was recorded using a liquid-nitrogen-cooled HPGe detector. The *S* parameter is sensitive to the annihilation fraction with low-momentum valence electrons and is proportional to the concentration of trapping centers¹⁷. The *W* parameter comprises the wings of the peak where higher momentum Doppler shifts dominate, and it relates chemical species to the annihilation site. Together, the *S* and *W* parameters were used to characterize positron-trapping centers in the ZnO crystals. Further positron experimental details and analysis are discussed elsewhere^{6,18}.

X-ray diffraction (XRD) measurement was carried out using a high-resolution ATX-G, Rigaku, triple-axis diffractometer system, using Cu K α radiation, with a scintillation counter (0-D detector).

The elemental composition along the depth direction was measured by Auger electron spectroscopy (AES). The crystal surfaces were measured with an AES, PHI 700 (ULVAC-PHI, INC) system, and the accelerating voltage of the first exciting electron beam was 5 kV.

Data Analysis

X-ray diffraction (XRD). Figure 1 shows the X-ray diffraction ω -rocking curve (XRC) of the as-grown CVT ZnO single crystal. An omega scan was performed for the reflection from the (110) crystal surface. The full-width at half-maximum (FWHM) was measured to be 0.018° , indicating the good crystallinity of the CVT-grown ZnO single crystal. Electron back scattering diffraction (EBSD) measurement, also suggested a good quality of ZnO crystal, with a (110) plane orientation (not shown).

Auger electron spectroscopy (AES). AES depth analysis was done by etching a ZnO crystal with an Ar-ion gun, and the variation of O(KLL) and Zn (LMM) intensity along the depth direction of the crystals was measured. In all the depth profiles, the main elements present in ZnO, i.e. zinc (Zn) and oxygen (O) were detected. As shown in Fig. 2 the atomic concentrations of Zn and O were found to be almost stoichiometric

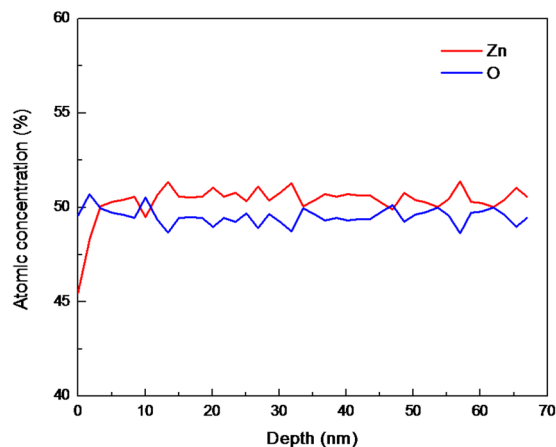


Figure 2. Auger electron spectroscopy (AES) depth profile spectra of an as-grown CVT ZnO crystal.

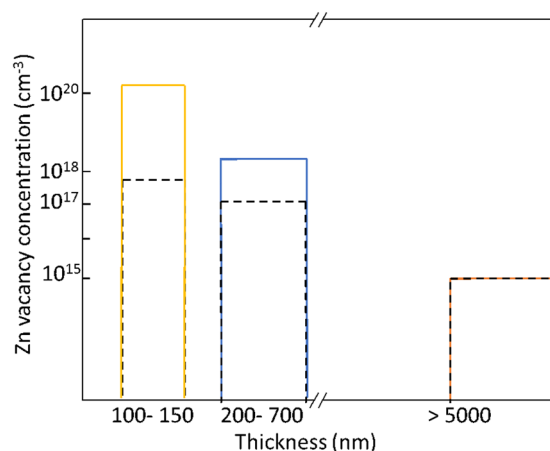


Figure 3. Schematic for the zinc vacancy concentration profile after oxygen annealing the CVT crystal at 1100 °C (broken lines) and 1200 °C (solid lines) (not to scale)⁶.

($\text{Zn}_{0.505}\text{O}_{0.495}$) and constant throughout the measurement, which clearly indicates a chemically uniform and stoichiometric ZnO crystal.

Defect concentration. The concentration of a point defect depends on its formation energy. In thermodynamic equilibrium and in the dilute regime (i.e., neglecting defect-defect interactions), the concentration of a point defect is given by

$$c = N_{\text{sites}} \exp\left(-\frac{E^f}{k_B T}\right) \quad (1)$$

where, E^f is the formation energy, N_{sites} is the number of sites the defect can be incorporated on, k_B the Boltzmann constant, and T is the temperature. Equation (1) shows that defects with high formation energies will occur in low concentrations.

Annealing CVT-grown ZnO crystals at 1200 °C created $\sim 5 \times 10^{20} \text{ cm}^{-3}$ defects in the top (100–150 nm) and $\sim 1 \times 10^{18} \text{ cm}^{-3}$ in the mid (200–700 nm) crystal region. The bulk crystal region ($> 3 \mu\text{m}$) remains with the minimal zinc vacancy concentration ($\sim 10^{15} \text{ cm}^{-3}$), thereby matching the pristine state (i.e., as-grown) of the ZnO crystals⁶ [Fig. 3]. Details of the calculation of the zinc vacancy (V_{Zn}) concentration are described elsewhere⁶.

Zinc vacancy formation energy. The zinc vacancy formation energies were calculated using the Boltzmann Equation – Eq. 1. A CVT sample annealed at 1200 °C yields $E^f \sim 0.64 \text{ eV}$ in the top region and ~ 1.39 in the mid region. The discrepancy in these formation energies can be understood since Eq. 1 is only valid in the dilute limit. Positron measurements showed that, in the top region, the SW data deviated from a straight line - suggesting a saturation of positrons with excess ($V_{\text{Zn}} > 10^{20} \text{ cm}^{-3}$) for the 1200 °C oxygen-annealed sample. The zinc vacancy formation energy (E^f) is 1.38 eV (top) and 1.52 eV (mid) region for the 1100 °C annealed CVT sample. The deviation of $\sim 0.14 \text{ eV}$ for this sample in two regions is quite reasonable, and Eq. 1 can be assumed to be

T (K)	C_a (cm^{-3}) (100–150 nm, top region)	C_b (cm^{-3}) (200–700 nm, mid region)	E_a^f (eV)	E_b^f (eV)
1473	5.3×10^{20}	1.5×10^{18}	0.64	1.39
1373	7.2×10^{17}	2.1×10^{17}	1.38	1.52

Table 1. Formation energy, where subscripts *a* and *b* denote the top and the mid regions, respectively.

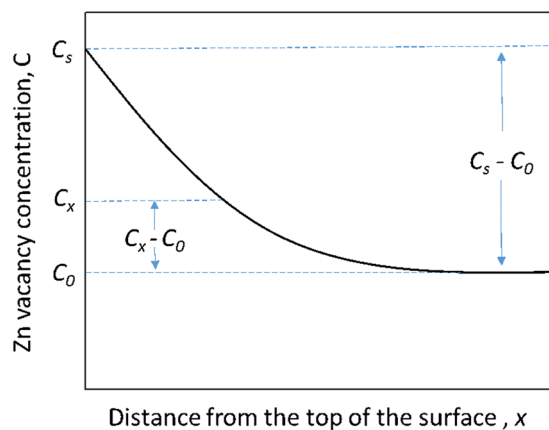


Figure 4. Demonstration of the Gaussian error function for the estimation of the diffusion coefficient.

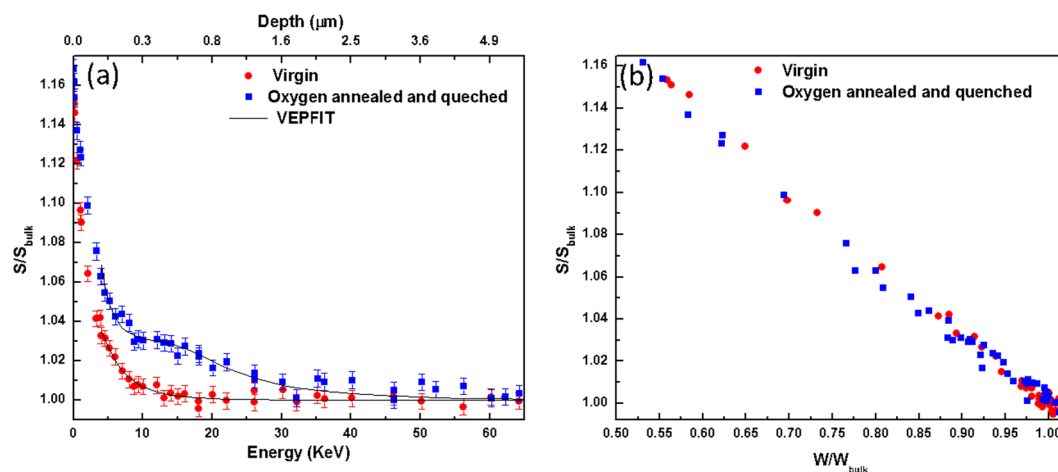


Figure 5. (a) Normalized S-vs-energy/depth (b) normalized S-W plot, in an as-grown (virgin) and oxygen annealed then quenched CVT ZnO crystal.

valid for the calculation of zinc vacancy formation energies, since (V_{zn}) is $<10^{18} \text{ cm}^{-3}$. The zinc vacancy formation energies are summarized in the Table 1.

Diffusion coefficients. The zinc diffusion coefficient at 1100 °C was calculated by using Fick's 2nd law:

$$\frac{C_x - C_0}{C_s - C_0} = 1 - \text{erf}(z), \quad \text{where } z = \frac{x}{2\sqrt{Dt}} \quad (2)$$

where, *C* is the zinc vacancy concentration [Fig. 4], erf(*z*) is the Gaussian error function, *x* is the diffusion distance, *z* is the approximated value of the Gaussian error function, *D* is the diffusion coefficient or diffusivity, and *t* is the diffusion time. The diffusion coefficients are summarized in the Table 2.

Quenching. Positron depth-resolved data were fit (from 4 keV onwards) using the VEPFIT¹⁹ computer code. The S-parameter is ~3% higher with respect to the bulk in the top layer [Fig. 5(a)]. The positron diffusion length is 20 nm in the top layer. The S-W data [Fig. 5(b)] lie on a straight line suggesting the presence of one type of defect related to (V_{zn}) .

T (K)	C_x (200–700 nm, mid region)	C_s (100–150 nm, top region)	C_o (bulk)	$\frac{C_x - C_o}{C_s - C_o}$	erf(z)	z	x (cm)	D (cm ² /s)
1373	2.5×10^{-4}	8.5×10^{-4}	1.2×10^{-6}	2.93×10^{-1}	0.70	0.74	4.5×10^{-5}	1.07×10^{-14}

Table 2. Diffusion coefficient calculation for CVT samples that were oxygen annealed at 1100 °C for 24 hours. Diffusion distance, $x = 4.5 \times 10^{-5}$ cm. Where, C_x , C_s and C_o values were calculated using the percentage ratio of the zinc vacancy concentration in their respective region to the atomic density (8.3×10^{22} cm⁻³). From Fick's 2nd law, the zinc diffusion coefficient at 1100 °C was calculated to be $\sim 1.07 \times 10^{-14}$ cm²/s.

λ_b^{-1} (ps)	μ (s ⁻¹)	S_b	S_{defect}	S_{ave}	c (atom ⁻¹)	C (cm ⁻³)
170	3×10^{15}	1	1.07	1.03	1.39×10^{15}	1.15×10^{17}

Table 3. Zinc vacancy concentration in the quenched sample.

T (K)	C (cm ⁻³) (500 nm–1 μm)	E_f' (eV)
1473	1.15×10^{17}	1.71

Table 4. Zinc vacancy formation energy under quenched conditions.

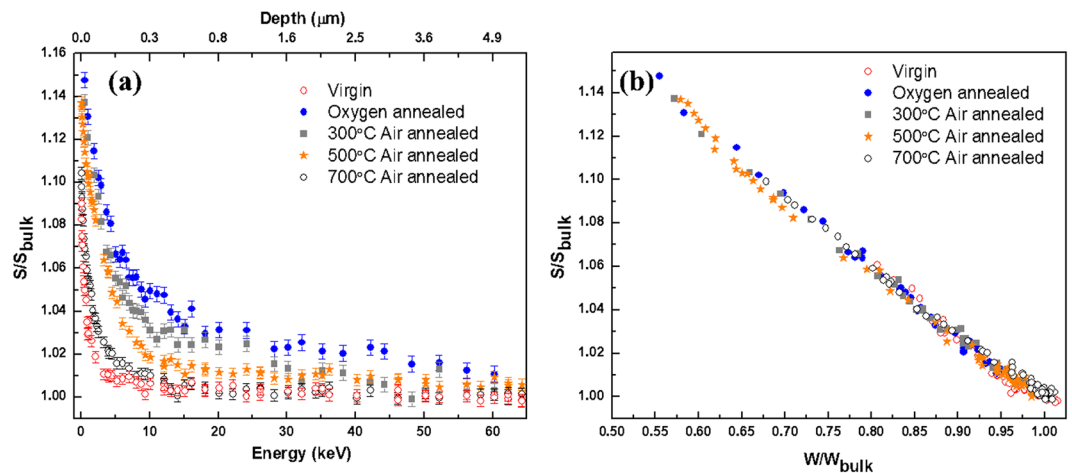


Figure 6. Positron data (a) (S vs Energy/depth) (b) S vs W, of an air-annealed (1 hour) CVT crystal that was then oxygen annealed at 1100 °C.

The zinc vacancy concentration (V_{zn}) (atom⁻¹) was calculated by using the equation:

$$c = \frac{\lambda_b}{\mu} \cdot \frac{S_{ave} - S_b}{S_{defect} - S_{ave}}$$

where, μ is the specific trapping rate, and $[\lambda_b^{-1} = \tau_b]$ is the positron lifetime in the bulk semiconductor. The subscripts refer to the measured τ or S in the case of average (*ave*), the bulk value for (*b*) and the defect specific value in the case of *defect*. The ZnO atomic density was assumed ($\rho \sim 8.3 \times 10^{22}$ cm⁻³).

In a quenched sample, S_{ave} increased by only $\sim 3\%$, i.e., a much lower value than for the slowly cooled samples - leading to a value of only (V_{zn}) $\sim 1.15 \times 10^{17}$ cm⁻³ [Table 3]. Generally, it is believed that the actual concentration of vacancies will be higher than the equilibrium value if the crystal is annealed at an elevated temperature and then cooled suddenly, thereby freezing in the vacancies. Our observation, however, was exactly the opposite since the zinc vacancy formation energy in the quenched condition [Table 4] is somewhat higher than that compared to the slow cooling case. A precise comparison is not justified, however, since the quenching process does not follow thermodynamic equilibrium conditions.

Annealing out zinc vacancies. A CVT crystal, oxygen annealed at 1100 °C, [$(V_{zn}) \sim 1 \times 10^{18}$ cm⁻³ (top layer)], was air annealed at various temperature steps (300 °C, 500 °C and 700 °C) for 1 hour each and this procedure was followed by performing the positron measurements [Fig. 6] after every air-annealing step. The positron data suggest that the (V_{zn}) completely anneals out (i.e., to the as-grown (virgin) level) after a 700 °C air anneal for

Defect	Charge state	Formation energy (eV)				
		ref. ²¹	ref. ²⁴	ref. ²⁵	ref. ²⁶	ref. ²⁷
V_{Zn}	-2	6.6	7.06	5.8	5.1	5.9
	-1	5.8	5.96	5.7	5.0	5.8
	0	6.0	5.6	5.8	≥5.1	6.0

Table 5. Calculated formation enthalpies for V_{Zn} defects in bulk zinc oxide for zinc-rich and p -type conducting conditions [$E_F = 0$ eV, VBM]. ref.²¹; DFT, LDA, ultra-soft pseudopotentials; ref.²⁴: GGA + U; ref.²⁵: DFT, LDA, norm-conserving pseudopotentials; ref.²⁶: DFT, GGA, ultra-soft pseudopotentials; and ref.²⁷: DFT, LDA, norm-conserving pseudopotentials.

Defect	Charge state	Formation energy (eV)				
		ref. ²¹	ref. ²⁴	ref. ²⁵	ref. ²⁶	ref. ²⁷
V_O	0	0	1.71	1.5	—	0.9
	+1	0.2	0.71	0.8	—	—
	+2	-0.3	-0.73	-0.5	-0.9	-0.5
$Zn_{i,oct}$	0	1.7	4.25	3.4	1.2	—
	+1	1.3	1.69	1.5	≥0.4	—
	+2	0.9	0.02	-0.2	-0.6	—

Table 6. Calculated formation enthalpies for V_O and $Zn_{i,oct}$ defects in bulk zinc oxide for zinc-rich and p -type conducting conditions [$E_F = 0$ eV, VBM].

1 hour. The S-W data followed a straight-line trend suggesting that no zinc vacancy-related clusters were formed during the annihilation process.

A comparative study. Janoti *et al.*²⁰, have reported a zinc vacancy formation energy of ~ 3.7 eV using first-principles calculations under p -type conditions (E_F at the VBM), in an oxygen-rich atmosphere. Kohan *et al.*²¹, have suggested that the zinc vacancy formation energy is $\sim (Zn(0/-1) = 5.8$ eV & $Zn(-1/-2) = 6.6$ eV) under a high zinc partial pressure and $\sim (Zn(0/-1) = 2.42$ eV and $Zn(-1/-2) = -0.2$ eV) under a low zinc pressure condition - assuming E_F at the VBM. Various other first-principles calculations are summarized in the Tables 5 and 6.

Also, Lany *et al.*²², have performed first principle calculations and showed that the intrinsic defects (V_O and Zn_i) can't lead to shallow donors due to a large formation energy, when E_F is close to the CBM. Tomlins *et al.*¹¹, performed zinc self-diffusion experiment in single crystal ZnO using nonradioactive ⁷⁰Zn as the tracer isotope and reported zinc vacancies formation energy > 3.51 eV. Azarov *et al.*²³, did zinc-diffusion measurements using isotopically modulated ZnO crystals and showed zinc vacancies formation energy varies as, $E_{V_{Zn}}^f = 1.1 - 2 \times (E_F - E_C)$, (where, $(E_F - E_C)$ is the difference in the Fermi level and conduction band energy) and depends on the position of the Fermi level. It can be seen that there exist a large discrepancies in zinc vacancies formation energy values calculated theoretically or experimentally.

Based on the positron analysis for the zinc vacancy concentration, the ZnO crystal region was divided into 3 parts: (1) top region (100–150 nm), (2) mid-region (200–700 nm) and, (3) bulk region ($> 3 \mu m$). Our calculated values are based on the near-surface (~ 100 –700 nm) experimental data (V_{zn}), that requires less formation energy as compared to the bulk region. The calculated (V_{zn}) formation energy is much lower than any previously reported values. This was observed experimentally, since there was no increase in the V_{zn} concentration in the bulk region, which is consistent with the high formation energy V_{zn} in the bulk region.

Conclusions

The good quality of ZnO crystal with (110) orientation was used in this study as suggested by X-ray diffraction ω -rocking curve (XRC). The depth analysis of AES indicates the atomic concentrations of Zn and O are almost stoichiometric and uniform throughout the measurement. In general, the diffusivity is greater through the less restrictive structural regions, such as the near surface (micron region), as compared to the bulk. Based on the zinc vacancy concentration and using Fick's 2nd law, we have calculated zinc diffusion coefficient of $\sim 1.07 \times 10^{-14}$ cm²/s at 1100 °C in the sub-micron region. The zinc vacancy formation energy (E^f) is calculated to be 1.38 eV (100–150 nm) and 1.52 eV (200–700 nm) region for 1100 °C oxygen anneal samples. These values are significantly lower than the reported values as obtained by first-principles calculations based on the bulk region of the ZnO crystal and few other reported experimental results.

References

- Ramanachalam, M. S., Rohatgi, A., Schaffer, J. P. & Gupta, T. K. Characterization of ZnO varistor degradation using lifetime positron-annihilation spectroscopy. *Journal of Applied Physics* **69**, 8380–8386, <https://doi.org/10.1063/1.347402> (1991).
- Look, D. C., Clafin, B., Alivov, Y. I. & Park, S. J. The future of ZnO light emitters. *Physica Status Solidi a-Applications and Materials Science* **201**, 2203–2212, <https://doi.org/10.1002/pssa.200404803> (2004).
- McCluskey, M. D. & Jokela, S. J. Defects in ZnO. *Journal of Applied Physics* **106**, 13, <https://doi.org/10.1063/1.3216464> (2009).

4. Tuomisto, F., Saarinen, K. & Look, D. C. Irradiation-induced defects in ZnO studied by positron annihilation spectroscopy. *Physica Status Solidi a-Applied Research* **201**, 2219–2224, <https://doi.org/10.1002/pssa.200404809> (2004).
5. Tuomisto, F., Look, D. C. & Farlow, G. C. Defect studies in electron-irradiated ZnO and GaN. *Physica B-Condensed Matter* **401**, 604–608, <https://doi.org/10.1016/j.physb.2007.09.032> (2007).
6. Parmar, N. S., Choi, J. W., Boatner, L. A., McCluskey, M. D. & Lynn, K. G. Formation of high concentrations of isolated Zn vacancies and evidence for their acceptor levels in ZnO. *Journal of Alloys and Compounds* **729**, 1031–1037, <https://doi.org/10.1016/j.jallcom.2017.09.239> (2017).
7. Look, D. C., Reynolds, D. C., Hemsley, J. W., Jones, R. L. & Sizelove, J. R. Production and annealing of electron irradiation damage in ZnO. *Applied Physics Letters* **75**, 811–813, <https://doi.org/10.1063/1.124521> (1999).
8. Knutsen, K. E. *et al.* Zinc vacancy and oxygen interstitial in ZnO revealed by sequential annealing and electron irradiation. *Physical Review B* **86**, <https://doi.org/10.1103/PhysRevB.86.121203> (2012).
9. Tuomisto, F., Ranki, V., Saarinen, K. & Look, D. C. Evidence of the Zn vacancy acting as the dominant acceptor in n-type ZnO. *Physical Review Letters* **91**, 4, <https://doi.org/10.1103/PhysRevLett.91.205502> (2003).
10. Khan, E. H., Weber, M. H. & McCluskey, M. D. Formation of Isolated Zn Vacancies in ZnO Single Crystals by Absorption of Ultraviolet Radiation: A Combined Study Using Positron Annihilation, Photoluminescence, and Mass Spectroscopy. *Physical Review Letters* **111**, 5, <https://doi.org/10.1103/PhysRevLett.111.017401> (2013).
11. Tomlins, G. W., Routbort, J. L. & Mason, T. O. Zinc self-diffusion, electrical properties, and defect structure of undoped, single crystal zinc oxide. *Journal of Applied Physics* **87**, 117–123, <https://doi.org/10.1063/1.371832> (2000).
12. Selim, F. Gamma Induced Positron Annihilation: History, Current, and Future Developments. *Acta Physica Polonica A* **132**, 1450–1455, <https://doi.org/10.12693/APhysPolA.132.1450> (2017).
13. Selim, F. A., Weber, M. H., Solodovnikov, D. & Lynn, K. G. Nature of native defects in ZnO. *Physical Review Letters* **99**, <https://doi.org/10.1103/PhysRevLett.99.085502> (2007).
14. Krause-Rehberg, R., Leipner, H. S. & Abgarjan, T. & Polity, A. Review of defect investigations by means of positron annihilation in II-VI compound semiconductors. *Applied Physics a-Materials Science & Processing* **66**, 599–614, <https://doi.org/10.1007/s003390050721> (1998).
15. Kilanski, L. *et al.* Native vacancy defects in Zn_{1-x}(Mn,Co)(x)GeAs₂ studied with positron annihilation spectroscopy. *Journal of Applied Physics* **106**, <https://doi.org/10.1063/1.3168440> (2009).
16. Tuomisto, F. & Makkonen, I. Defect identification in semiconductors with positron annihilation: Experiment and theory. *Reviews of Modern Physics* **85**, 1583–1631, <https://doi.org/10.1103/RevModPhys.85.1583> (2013).
17. Schultz, P. J. & Lynn, K. G. Interaction of positron beams with surfaces, thin-films, and interfaces. *Reviews of Modern Physics* **60**, 701–779, <https://doi.org/10.1103/RevModPhys.60.701> (1988).
18. Parmar, N. S. & Lynn, K. G. Sodium doping in ZnO crystals. *Applied Physics Letters* **106**, <https://doi.org/10.1063/1.4905594> (2015).
19. Vanveen, A., Schut, H., Devries, J., Hakvoort, R. A. & Ijma, M. R. In *4th International Workshop on Slow-Positron Beam Techniques for Solids and Surfaces*. 171–196 (Aip Press, 1990).
20. Janotti, A. & Van de Walle, C. G. Native point defects in ZnO. *Physical Review B* **76**, <https://doi.org/10.1103/PhysRevB.76.165202> (2007).
21. Kohan, A. F., Ceder, G., Morgan, D. & Van de Walle, C. G. First-principles study of native point defects in ZnO. *Phys. Rev. B* **61**, 15019–15027 (2000).
22. Lany, S. & Zunger, A. Dopability, intrinsic conductivity, and nonstoichiometry of transparent conducting oxides. *Physical Review Letters* **98**, <https://doi.org/10.1103/PhysRevLett.98.045501> (2007).
23. Azarov, A. *et al.* Self-diffusion measurements in isotopic heterostructures of undoped and *in situ* doped ZnO: Zinc vacancy energetics. *Physical Review B* **94**, <https://doi.org/10.1103/PhysRevB.94.195208> (2016).
24. Erhart, P., Albe, K. & Klein, A. First-principles study of intrinsic point defects in ZnO: Role of band structure, volume relaxation, and finite-size effects. *Physical Review B* **73**, <https://doi.org/10.1103/PhysRevB.73.205203> (2006).
25. Zhang, S. B., Wei, S. H. & Zunger, A. Intrinsic n-type versus p-type doping asymmetry and the defect physics of ZnO. *Physical Review B* **63**, <https://doi.org/10.1103/PhysRevB.63.075205> (2001).
26. Oba, F., Nishitani, S. R., Isotani, S., Adachi, H. & Tanaka, I. Energetics of native defects in ZnO. *Journal of Applied Physics* **90**, 824–828, <https://doi.org/10.1063/1.1380994> (2001).
27. Erhart, P., Klein, A. & Albe, K. First-principles study of the structure and stability of oxygen defects in zinc oxide. *Physical Review B* **72**, <https://doi.org/10.1103/PhysRevB.72.085213> (2005).

Acknowledgements

The authors thank Prof. Matthew McCluskey, and Dr. M. H. Weber at WSU for useful discussions. This research was supported by the Korea Research Fellowship Program through the National Research Foundation of Korea (NRF) funded by the Ministry of Science, ICT and Future Planning (2016H1D3A1909335). This research was also supported by the ATC program (No. 10048059) funded by the Ministry of Trade, Industry & Energy, and the KIST Future Research Program (2E28210). The authors acknowledge CMR at WSU support for positron beam facility. Research at the Oak Ridge National Laboratory for one author (LAB) was supported by the U.S. Department of Energy, Office of Science, Basic Energy Sciences, Materials Sciences and Engineering Division.

Author Contributions

The manuscript was written through contributions of all the authors. All the authors have given approval to the final version of the manuscript. N.S.P. conceived the idea and performed the experiments and major data analysis. J.-W.C. supervised the manuscript, L.A.B. provided CVT ZnO crystals, and K.G.L. provided the positron beam facility. All the authors discussed the results and contributed in the analysis of the data in the manuscript.

Additional Information

Competing Interests: The authors declare no competing interests.

Publisher's note: Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.



Open Access This article is licensed under a Creative Commons Attribution 4.0 International License, which permits use, sharing, adaptation, distribution and reproduction in any medium or format, as long as you give appropriate credit to the original author(s) and the source, provide a link to the Creative Commons license, and indicate if changes were made. The images or other third party material in this article are included in the article's Creative Commons license, unless indicated otherwise in a credit line to the material. If material is not included in the article's Creative Commons license and your intended use is not permitted by statutory regulation or exceeds the permitted use, you will need to obtain permission directly from the copyright holder. To view a copy of this license, visit <http://creativecommons.org/licenses/by/4.0/>.

© The Author(s) 2018