


 Cite this: *RSC Adv.*, 2021, 11, 16962

A two-step chemical vapor deposition process for the growth of continuous vertical heterostructure WSe₂/h-BN and its optical properties†

 M. Alahmadi,^{ab} F. Mahvash,^{ac} T. Szkopek^c and M. Sijaj^{id}*^a

The expansion of two-dimensional (2D) van der Waals heterostructure materials growth and synthesis leads to impressive results in the development and improvement of electronic and optoelectronic applications. Herein, a vertical WSe₂/hBN heterostructure was obtained *via* a dual CVD system, in which prior to the WSe₂ growth a continuous monolayer hBN was obtained on a SiO₂/Si substrate. Comparing growth on SiO₂/Si and quartz substrates, we found that the underlayer of hBN leads to a desorption/diffusion process of tungsten (W) and selenium (Se) producing high-quality and large-area WSe₂ growth. In contrast with WSe₂/SiO₂ and WSe₂/quartz heterostructures, the photoluminescence properties of WSe₂/hBN exhibit a sharp intense WSe₂ peak at 790 nm with a narrow full width at half-maximum (80 meV) due to no dangling bonds and dielectric effect of the hBN interface. The photoluminescence results suggest that the WSe₂/hBN heterostructure has high crystallinity with a defect-free interface.

Received 31st March 2021

Accepted 29th April 2021

DOI: 10.1039/d1ra02523f

rsc.li/rsc-advances

Introduction

Among all 2D materials, semiconducting transition metal dichalcogenides (TMDCs) such as WSe₂, WS₂, MoS₂, and MoSe₂ have offered useful physical and optical properties that can potentially be used for future optoelectronic and electronic applications.^{1–4} Among these materials, WSe₂ in its monolayer form is composed of three atomic layers; the outer top and bottom layers are dichalcogenide atoms (selenium Se), encompassing a transition metal atom (tungsten W) in the middle.^{5,6} Recent studies of monolayer WSe₂ materials have shown desirable properties including a small direct bandgap (1.65 eV),⁷ novel spin valley coupling,⁸ tunable charge carrier behaviour,⁹ and strong photoemission.^{10,11} In the last few years, significant studies have demonstrated that the substrate effect on the characteristics of as-grown 2D materials is the most important pillar in reducing surface impurities and structural defects. In addition, the substrate has a direct effect on the optical and electronic properties.^{12–15} For example, the mobility of charge carriers in graphene on the SiO₂ substrate is around 10 000 cm² V⁻¹ s⁻¹, while free-standing graphene at 5 K has a charge carrier mobility of 200 000 cm² V⁻¹ s⁻¹.^{16,17} Further, the

electron mobility of MoS₂ on SiO₂ is 30 to 60 cm² V⁻¹ s⁻¹ at room temperature, and the majority of devices display N-type behaviour. On the other hand, the mobility of MoS₂ on the polymethyl methacrylate (PMMA) dielectric substrate increases up to 470 cm² V⁻¹ s⁻¹.¹⁸

Recently, insulating hexagonal boron nitride has been considered an exemplary substrate because of its unique properties. The insulating two-dimensional layer, single atomic thick hBN film has sp² hybridized atoms (boron and nitrogen) in a hexagonal lattice, a wide bandgap (5.5 eV) and high optical phonon energy (twice than SiO₂).^{19–21} hBN has a robust hexagonal in-plane bonding structure resulting in an extremely flat surface free of charge impurities and free of charge-trap from the dangling bonds.¹³ Consequently, it becomes an excellent choice as a dielectric substrate for growth 2D materials and electric and photoelectric devices like a field-effect transistor.^{12,22,23} So far, many studies have reported the CVD growth of TMDC heterostructures on mechanically exfoliated hBN substrates. Mechanical exfoliation (scotch tape) could exhibit an excellent hBN crystalline structure.²⁴ However, it is relatively complicated process, the yield is very low, the interface between layers can be easily contaminated, and this method lacks producing large-area of 2D material.^{25,26} Therefore, the direct growth of out-of-plane heterostructures *via* bottom-down technique gas-phase chemical vapor deposition which has been elucidated to be compliant, scalable, easily operational, and precisely controllable,²⁷ is required to synthesize highly crystalline TMDC heterostructures on hBN substrates, rather than using mechanical exfoliation. Nevertheless, the growth of good quality, large-scale monolayer MX₂ heterostructures have been a great challenge. The growth temperature and introduction of

^aNanoQAM, Quebec Center for Functional Materials, Department of Chemistry, University of Quebec in Montreal, Succ CentreVille, CP8888, Montreal, Quebec H3C 3P8, Canada. E-mail: sijaj.mohamed@uqam.ca

^bDepartment of Chemistry, College of Science, Taibah University, Al-Madinah Al-Munawarah 41321, Saudi Arabia

^cDepartment of Electrical and Computer Engineering, McGill University, Montréal, H3A 2A7, Quebec, Canada

† Electronic supplementary information (ESI) available. See DOI: 10.1039/d1ra02523f



H₂ gas to the growth chamber can induce the decomposition of the hBN film before growing the TMDC layer.²⁸ For That, the growth time of semiconducting TMDC should be short to avoid the hBN decomposition.

In the present study, we realized a controllable synthesis of WSe₂/hBN vertical heterostructure, where a direct CVD WSe₂ atomic layers are grown on ultra-smooth CVD-hBN layer in ambient pressure and under different hydrogen flux concentrations and at high temperatures (950 °C). We found that the CVD growth of WSe₂ is largely dependent on the synthesis temperature and hydrogen concentration. The optimal condition for up to 300 μm² WSe₂ complete monolayer growth on insulating hBN substrates was found to be, 950 °C and 15 sccm H₂ concentration. Increasing the H₂ flux to 20 sccm leads to WSe₂ multilayer formation. The high quality and uniformly monolayer of WSe₂ growth on hBN substrates were characterized using Raman spectroscopy, photoluminescence and X-ray photoelectron spectroscopy. Not only were we able to effectively grow large-area WSe₂ on top of hBN at 950 °C under high H₂ gas concentration, but we also improve the optical properties. After the WSe₂ growth, the optical properties of the heterostructure WSe₂/hBN are compared to WSe₂/SiO₂ and WSe₂/quartz heterostructures. For WSe₂/hBN structure a sharp and intense photoluminescence peak of the WSe₂ at 790 nm with narrow FWHM (80 meV) was obtained. The sharp photoluminescence peak suggests that the WSe₂ materials have substantial crystallinity and high quality which is related directly to the dielectric constant and the flat surface of hBN substrate free from dangling bonds. Recently, Lui *et al.* showed substrates effect on the photoluminescence of chemical vapor deposition transferred WSe₂ monolayer on polyethylene terephthalate (PET), quartz, SiO₂/Si and sapphire substrates.²⁹ The PL emission intensities can drastically change depending on the substrate effect. From the analysis of the excitonic behavior, the results reveal that the spectral weight between the neutral and charged excitons in the PL spectra is significantly modified by the substrate types, and the charged exciton binding energy is inversely related to the substrate dielectric constant. These results suggest that the choice of the substrate plays a significant role in the modulation of the PL properties and exciton states of atomically thin WSe₂, which is in a good agreement with the present work. This work will provide a route to synthesize various TMDC/hBN heterojunctions at high temperatures and will encourage optical and electrical applications based TMDCs.

Experimental

CVD synthesis of hexagonal boron nitride (hBN)

We have conducted the hBN layer growth directly on 300 nm Si/SiO₂ and quartz substrates with different carrier gas flow rates and with fixed growth time by utilizing the ammonia borane (AB) precursor (Sigma-Aldrich, 97%). At a high temperature, 1000 °C, for 30 minutes, we tested different hydrogen/argon flux growth conditions, H₂ : Ar sccms (5 : 5; 10 : 10 and 15 : 15). The precursor (AB) was heated in a separated growth chamber to reach approximately 100 °C using a heating belt, and the

temperature of the heating belt was monitored using a thermocouple thermometer. At the temperature of 70 °C, the precursor slowly was introduced into the system growth chamber. Upon completion of the synthesis process, the systems' furnace was rapidly cooled to room temperature maintaining the growth gas flow. The as-synthesized hBN film was used as the substrate to grow WSe₂ films in the next step.³⁰

Direct CVD synthesis of the WSe₂/hBN/SiO₂ heterostructure

The two-zone CVD furnaces were used for large-scale area of WSe₂/hBN heterostructure film synthesis. Selenium (Se) and tungsten trioxide (WO₃) were used as precursors. Se powder (440 mg) was positioned in a ceramic crucible occupying the first CVD furnace section at a low temperature of 540 °C.³⁰ The as-grown hBN/SiO₂ film³⁰ was placed facing top-down above the ceramic crucible containing WO₃ powders (260 mg). In the middle of the furnace-heating zone, this ceramic crucible was placed at a temperature of 950 °C. The distance between these precursors was set to 45 cm. Thereafter, the system was heated until it reached a temperature of 950 °C at a rate of 40 °C min⁻¹. Then, the system remained under atmospheric pressure for 15 minutes. During this growth time, the hydrogen and argon (H₂/Ar) gases were utilized as carrier gases with flow rates of (320/20 sccm) and (320/15 sccm) for the synthesis of the multilayer and monolayer of WSe₂ films, respectively. Eventually, the samples were removed after the system was cooled down to room temperature.

Characterizations

Several characterization methods have been used to examine the resulting products. The optical spectroscopy images were taken by using an Olympus microscope (BX51). Raman and photoluminescence measurements were done in a confocal Raman system (Alpha300R) and (WITec UHTS 300) equipped with laser excitation of 532 nm wavelength. The chemical configurations were studied and examined using X-ray photoelectron spectroscopy using a PHI 5600-ci system (physical Electron, Eden Prairie, MN). AFM was implemented in Bruker's Scanasyst to determine the surface morphology and thickness of the resulting products.

Results and discussion

The illustration in Fig. 1a shows the CVD growth system for WSe₂/hBN van der Waals heterostructure film synthesis. hBN layers were grown on the top of Si/SiO₂ and quartz substrates under atmospheric pressure and at a high furnace temperature (1000 °C). The Si/SiO₂ substrates were ~1 cm² in size and the hBN layers fully covered substrates, as made evident by the slightly changed green color of the Si/SiO₂ substrate. XPS was carried out on hBN coated Si/SiO₂ substrates for chemical bonding and stoichiometric analysis. The survey spectrum is illustrated in Fig. 1b, which indicates the presence of B and N elements. A B/N ratio of 1.03 was determined, verifying that the stoichiometry of the film is in excellent agreement with that of hBN.³¹ The carbon and oxygen binding energies are attributes to

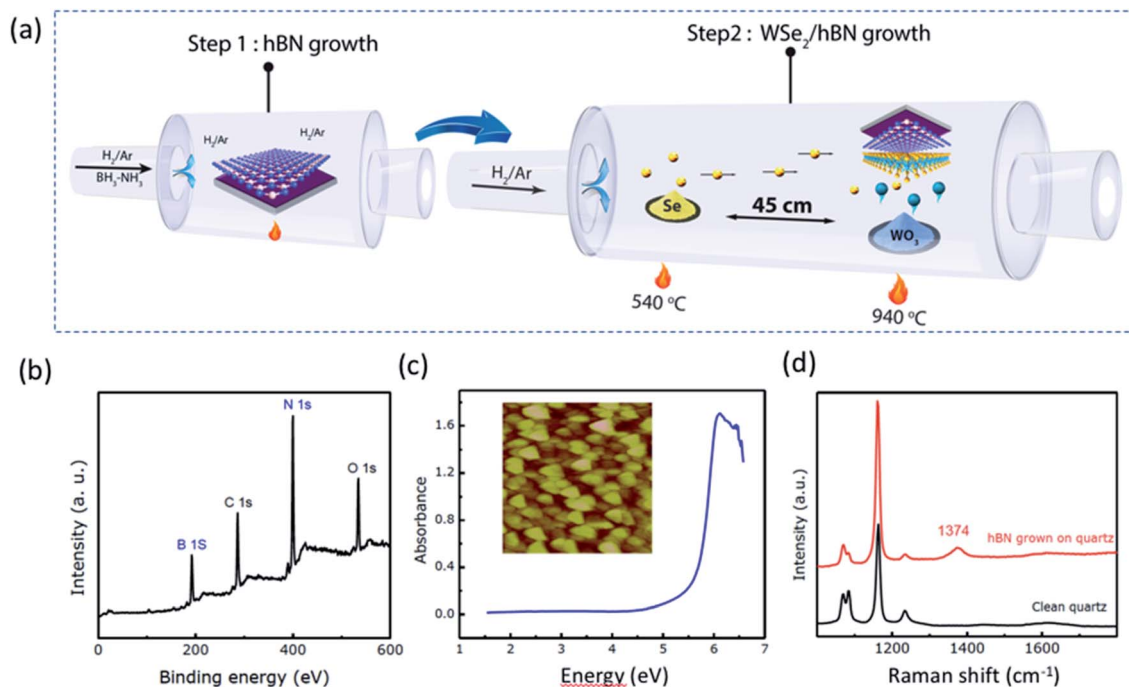


Fig. 1 (a) A schematic view shows the CVD process of the WSe_2/hBN van der Waals heterostructures using a two-step CVD method. CVD step I: hBN synthesis upon the SiO_2 substrates at a high growth temperature ($950\text{ }^\circ\text{C}$), the end-product of this step is illustrated; CVD step II: the growth of the WSe_2 film on the hBN film, the growth conditions and the end-product are schemed. Characterization of hBN film grown on Si/SiO_2 with the growth time of 1 h. (b) XPS survey. Characterization of hBN film grown on a quartz substrate with the growth time of 1 h. (c) The UV-visible absorption spectrum of hBN grown on the quartz substrate, measured at room temperature. Inset shows the AFM image of the sample. The image size is $1\ \mu\text{m}$. The root mean squared roughness R_q was calculated to be 8 nm. (d) Raman spectra of a clean quartz substrate and hBN film grown on a quartz substrate.

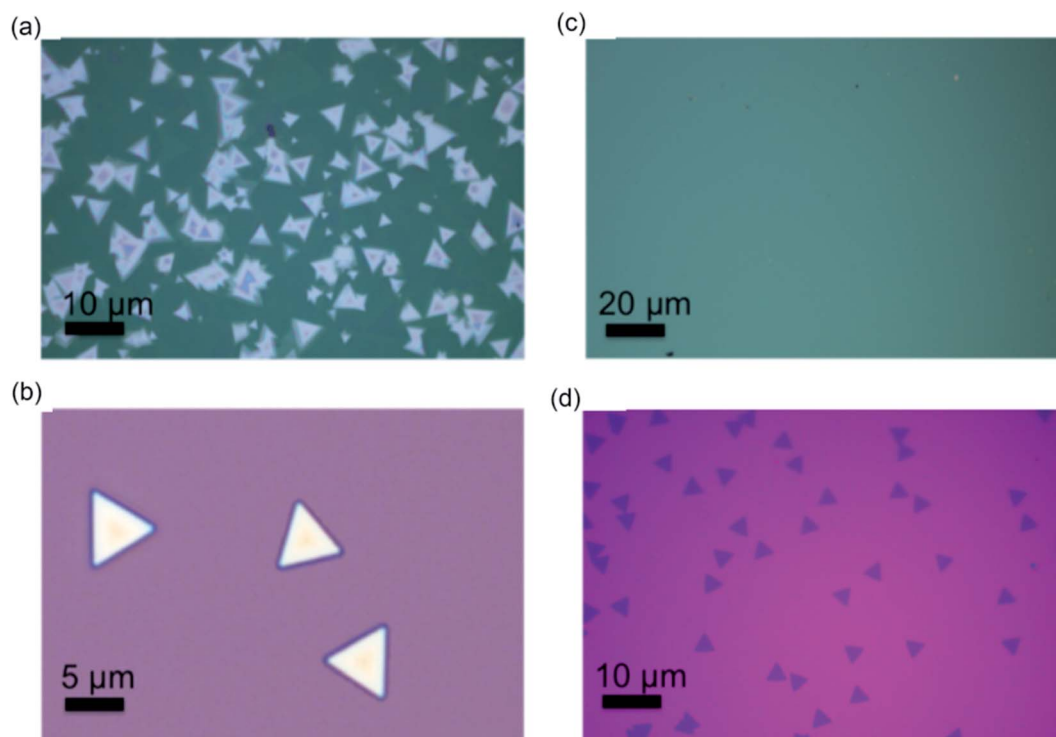


Fig. 2 (a and b) Show the optical microscopy (OM) images of the multilayered WSe_2 on the hBN film and the SiO_2 substrate, respectively, at the same growth condition. (c and d) Shows the OM images of the monolayer WSe_2/hBN and $\text{WSe}_2/\text{SiO}_2$, respectively.

the exposure of hBN film to air prior to XPS measurement. The hBN film grown on quartz during 1 h was characterized by Raman spectroscopy and UV-visible absorption spectroscopy. Fig. 1c illustrates the absorbance spectrum along with an AFM image. The quartz background was subtracted using a quartz blank. A sharp absorption peak at 6.1 eV has been observed, which is in excellent agreement with the literature value for hBN.³² A rough film was obtained on quartz, as was observed on SiO₂/Si, with a roughness of 8 nm. We could not estimate the thickness of as grown film on quartz using ellipsometry due to inadequate index contrast between hBN and quartz. The Raman spectrum is shown in Fig. 1d affirms the presence of hBN film *via* the presence of the E_{2g} Stokes peak at 1374 cm⁻¹.

For the WSe₂/hBN synthesis, we found that 15 min is the optimum growth time to prohibit the CVD hBN decomposition. Indeed, when we attempt to synthesize WSe₂ for a long period of time (40 min) on top of the CVD-hBN film, we found that the hBN film had completely decomposed. To show clearly the hBN effect as the substrate on the WSe₂ growth, we proceed to WSe₂ growth on SiO₂ and quartz substrates, under the same growth condition. Fig. 2a exhibits the optical images of WSe₂ grown on the hBN at 950 °C with a flowing rate of H₂ (20 sccm). The growth under this condition yielded to multilayered WSe₂ films, which is confirmed by the optical microscopy images as WSe₂ domains (yellow triangles) on top of WSe₂ film (dark-green areas) and Raman spectroscopy. Based on our observation, the triangular domains on top of the WSe₂ layer, which have domain sizes in the range of 10 to 20 μm, could merge to form

a continuous WSe₂ film. A high growth temperature tends to favor the multilayer growth of WSe₂. However, the WSe₂ crystal synthesis on top of SiO₂ exhibited smaller domains with an average domain size around 15 μm, which can be referred to the rough surface or the charge impurities of the SiO₂ substrate. Based on the above results, under an optimized experimental condition (950 °C and 15 sccm of H₂) a high continuous WSe₂ layer could be obtained on top of hBN/SiO₂ (Fig. 2c and d). Due to the elimination of the substrate defects such as charged impurities and substrate roughness, WSe₂ crystals growth on the h-BN layer leads to a large film size above 400 micrometer, which is larger than the WSe₂ crystals grown on SiO₂ and quartz substrates (Fig. 2b). Homogeneity in color contrast on the optical image obviously indicates that the large-area monolayer WSe₂ appeared with a flat and uniform surface. The growth of WSe₂ on the hBN film is also larger than other TMDC crystals such as WSe₂, WS₂, and MoSe₂ grown on top of SiO₂ substrates.^{33–36} Eventually, by relatively decreasing both growth temperature and H₂ gas flow rate (15 sccm), the layer number of WSe₂ is decreased, and WSe₂ triangular domains are not observed suggesting a homogeneous WSe₂ layer (Fig. 2c). Fig. S1† shows optical images of hBN on SiO₂/Si and quartz substrates, clearly showing continuous hBN layers (cm scale). The SEM and energy dispersive X-ray (EDX) elemental mapping are shown in Fig. S1c and d.† In addition, TEM and SAED, was performed on the edge of suspended pristine hBN layers by using a partial etch of the substrate.³⁷ A bright-field TEM image illustrated in Fig. S2a† reveals a continuous and transparent

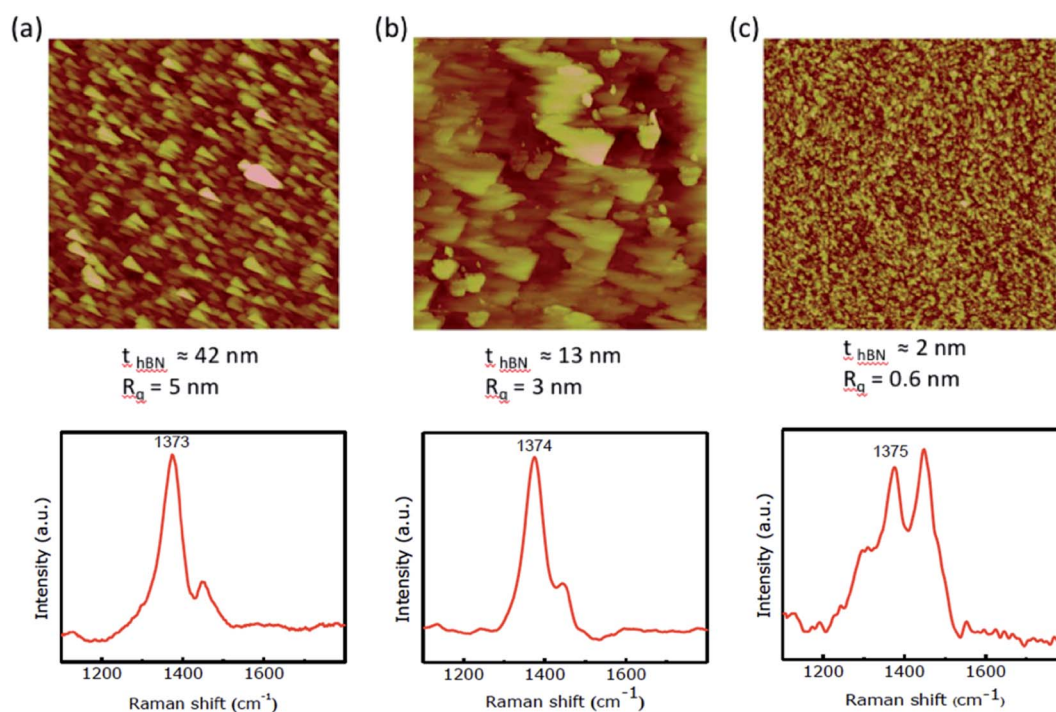


Fig. 3 Characterization of hBN film grown on Si/SiO₂ with the growth time of 1 h and variable carriers gas flow rates. AFM image and Raman spectrum of hBN film grown on Si/SiO₂ substrate in 5 sccm of H₂ and 5 sccm of Ar (a), 10 sccm of H₂ and 10 sccm of Ar (b), 15 sccm of H₂ and 15 sccm of Ar (c). The image size is 1 μm. The root mean squared roughness R_q was calculated using a Nanoscope analyzer. The thickness of the hBN film was estimated using spectroscopic ellipsometry.

layer and the SAED measurement shown in the Fig. S2b† gives the expected hexagonal lattice structure of hBN monolayer. The data shown in Fig. S1 and S2† confirm continuous hBN layer growth on SiO₂/Si substrates.

The growth mechanism could be explained by a simple dissociation, adsorption, and diffusion process in which the W and S adsorb on top of hBN and assemble into WSe₂ film. Indeed, the hBN growth on top of SiO₂ leads to a smoother surface. Fig. 3 illustrates the AFM images and Raman spectra of as-grown hBN films on Si/SiO₂ substrates with different carrier gas flow rates. The thickness of each film was estimated using spectroscopic ellipsometry. A single Stokes peak located at 1374 cm⁻¹ corresponding to the E_{2g} vibration mode of hBN is observed (Fig. 3a and b). The smaller peak at 1450 cm⁻¹ is assigned as the third order Si transverse optical (TO) phonon which is due to the underlying Si/SiO₂ substrate.³⁸ We have found that the hBN film is smoother and thinner as the carrier gas flow rate increases. We have observed different Raman shift for our hBN film in the range of 1370–1375 cm⁻¹. The shifting of hBN Raman peak has been observed in literature. It was found that the Raman peak frequency would shift to a higher frequency under a compressive stress.³⁹ It was reported that monolayer hBN exhibits a blue shifted Raman peak in comparison to its position in bulk hBN.⁴⁰ Interestingly, Pakdel

et al. found that the high-frequency E_{2g} phonon mode in hBN shifted to a higher frequency and broadened as the crystallite size decreased.⁴¹ They established relationships between the hBN frequency shift, the broadening of the mode, and the crystallite size. The hBN peak is slightly shifted to a higher frequency as the carrier gas flow rate increased (Fig. 3c). This may suggest that the crystalline size of hBN shrank. Using 15 sccm carrier gases, we have estimated a crystallite size of ~4 nm based on the reported dependence of Raman shift and crystallite size, however, the AFM measurements give the crystallite size average of ~18 nm. We have observed a decrease in hBN film thickness as the carrier gas flow rate increased. Moreover, as the film grew thicker, it roughened due to surface irregularities and multiple grain boundaries. Thus, as the first step, we succeeded in growing hBN on Si/SiO₂ and quartz substrates using CVD without the aid of any metal catalyst and the grown hBN film is thinner and smoother. Moreover, the carrier gas flow rate affects the CVD grown hBN film roughness and thickness. Large area hBN films with the thickness of 2 nm and rms roughness of 0.6 nm, as made evident by the slightly changed green color of the Si/SiO₂ substrate (Fig. 2c).

Raman and PL were performed in order to investigate the crystal structure and the band gap of the resulting WSe₂/hBN film heterostructure.^{42,43} Fig. 4a shows Raman peak frequency

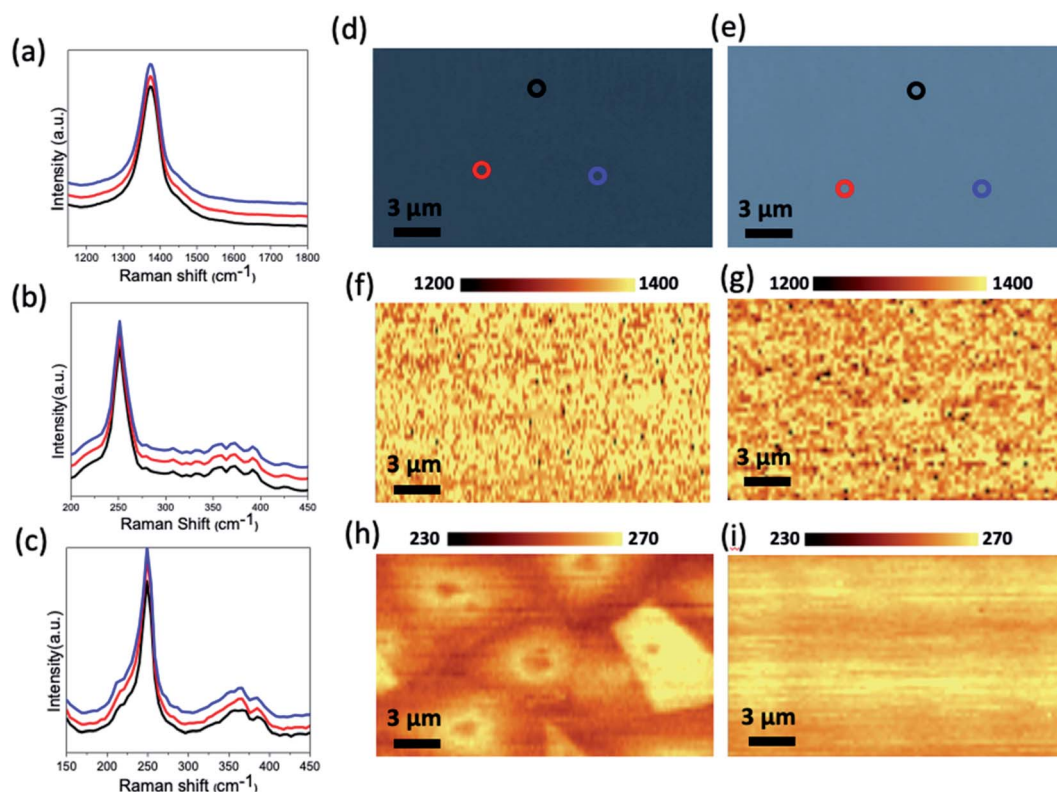


Fig. 4 Raman investigation of the WSe₂/hBN heterostructure. (a) Displays three Raman spectra of the underlying hBN film collected from different positions of the sample. (b) Shows the Raman spectrum of the multilayer WSe₂ grown at (950 °C/20 sccm H₂). (c) Displays the Raman spectrum of the monolayer WSe₂ synthesized at (940 °C/15 sccm H₂). (d and e) Are optical images of the WSe₂/hBN synthesized (950 °C/20 sccm H₂) and (950 °C/15 sccm H₂), respectively. (f and g) Are Raman mapping images of the underlying hBN materials taken from the regions of panels (d) and (e) following the Stokes E_{2g} (~1369 cm⁻¹) peak intensity of hBN integrated between 1300 and 1400 cm⁻¹. (h and i) Show the Raman map of the E_{2g} (~249 cm⁻¹) peak intensity of the WSe₂ taken from the regions of panels (d) and (e), respectively.

E_{2g} mode of the underneath hBN layer, which located at 1373 cm^{-1} . This is consistent with the CVD grown hBN binding energy previously reported.^{21,28} This peak evidently confirmed the presence of the underlying hBN after completing the second CVD growth process step of the WSe_2 materials. The Raman spectra of the WSe_2 film showed two modes E_{2g}^1 and A_{1g} at 249 cm^{-1} and 251 cm^{-1} , respectively. The layer number of WSe_2 can be measured by the difference between these intensity peaks. Fig. 4b demonstrates the typical Raman spectra taken from a multilayer CVD grown WSe_2 on the hBN layer at $950\text{ }^\circ\text{C}/20\text{ sccm H}_2$. The Raman spectra of the monolayer WSe_2 is shown in Fig. 4c grown at ($950\text{ }^\circ\text{C}/15\text{ sccm H}_2$). The sharp peak refers to the mode E_{2g}^1 at 249 cm^{-1} and the absence of the peak at around 307 cm^{-1} (A_{1g}^2 vibration mode) confirms that the

synthesized heterostructure in these growth parameters was a monolayered WSe_2 . The as-synthesized WSe_2/hBN heterostructures at different growth conditions were characterized using Raman mapping. Raman intensity mapping recorded over a $20\text{ }\mu\text{m}^2$ area for the underlying hBN film (Fig. 4f and g). These also reveal that the uniform color contrast over the whole area clearly shows a continuous and homogeneous hBN film on top of the SiO_2 substrate. A Raman map of hBN E_{2g}^1 ($\sim 1369\text{ cm}^{-1}$) peak intensity is shown in Fig. S3 (see ESI†). It is clear from this image that the intensity is very uniform indicating a uniform thickness. The Raman mapping of the as-synthesized multilayer WSe_2 films is also shown in Fig. 4h. The red regions on the Raman mapping correspond to the multilayer WSe_2 while the yellow regions correspond to the multilayer

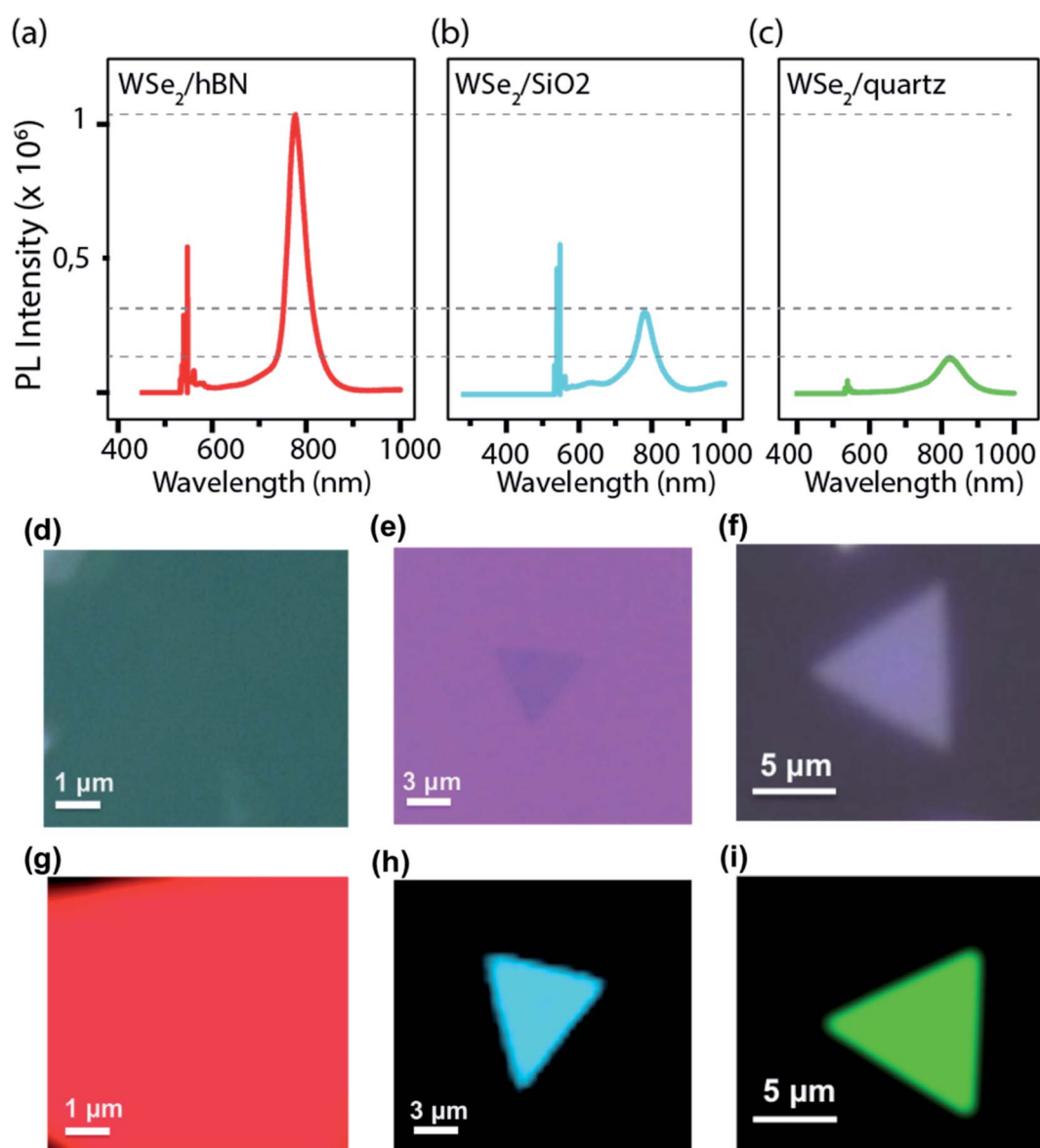


Fig. 5 PL investigation of the growth of WSe_2 on various substrates. (a–c) Display the monolayer PL spectra of the WSe_2/hBN , $\text{WSe}_2/\text{SiO}_2$ and $\text{WSe}_2/\text{quartz}$, respectively. (d–f) Demonstrate optical images of the WSe_2/hBN , $\text{WSe}_2/\text{SiO}_2$ and $\text{WSe}_2/\text{quartz}$, respectively. (g–i) Show PL intensity mapping of the monolayer WSe_2/hBN , $\text{WSe}_2/\text{SiO}_2$ and $\text{WSe}_2/\text{quartz}$, respectively.

continuous film WSe₂ domains, which grew overlapping. The Raman intensity map of the peaks (249 and 307 cm⁻¹) further demonstrates the synthesis of the uniform WSe₂ monolayer on top of the hBN film at (940 °C/15 sccm H₂) as is indicated in Fig. 4i. The Raman map demonstrates the same color throughout the entire monolayer, indicating a highly homogeneous crystal quality. In addition to Raman spectroscopy, the as-synthesized WSe₂/hBN heterostructure was further examined using X-ray photoelectron spectroscopy as shown in Fig. S1.† Fig. S1a† represents the B 1s peaks at binding energy ~190.2 eV for the B–N bond. Fig. S1b† shows the N 1s peak at binding energy ~397.6 eV for the N–B bonding. These peaks are consistent with those previously reported for monolayer hBN.^{44–46} Both B 1s and N 1s confirm that the hBN film still exists after the WSe₂ growth is completed in the second CVD step. The XPS spectra of the top layer CVD grown WSe₂ is present in Fig. S4c and d,† respectively. The binding energies at 32.8, 36.8, 39.4 eV are assigned to the W 4f_{7/2}, W 4f_{5/2} and W 5p_{3/2}, respectively. The presence of the two binding peaks at 55.8 and 56.7 eV are assigned to the Se₃ d_{3/2}, and Se₃ d_{5/2}, respectively. These resulting values are consistent with previously published values.^{47–49}

The optical property of the monolayer WSe₂/hBN heterostructure film grown at (950 °C to 15 sccm H₂) was further characterized using photoluminescence (PL) as shown in Fig. 5. The WSe₂ films exhibited strong PL spectra localized at approximately 760 nm with a narrow FWHM (80 meV), in comparison with PL intensity spectra of the WSe₂ on the SiO₂ and the quartz substrate as seen in Fig. 5b and c, respectively. Ordinarily, the quality of the materials and the lifetime of electron–hole coupling are related to the FWHM of the PL emission peak.^{12,50} The narrow FWHM of WSe₂/hBN can be a result of the high crystallinity, clean interfaces and reduced charged impurities of the continuous films when grown on hBN. The PL mapping of the synthesized WSe₂ materials on different substrates, such as hBN, SiO₂, and quartz is shown in Fig. 5d–f, respectively. The difference in the PL mapping colors does not reflect the emission color. We can clearly conclude that the underlying insulating hBN is an ideal substrate for reducing structural defects, improving optical properties, and promising for growth scalability.

Conclusions

We have accomplished the synthesis of a large area, uniform, and high quality two-dimensional atomic layer WSe₂ directly on hBN substrates using a two-step CVD approach without the need for the transfer process. It is concluded that the surface morphology of the substrate influences the WSe₂ growth materials. By using an hBN film as an insulating substrate, a large area (400 μm²) WSe₂/hBN heterostructure was synthesized, which was much larger than those obtained using a SiO₂ substrate. To obtain an atomically smooth surface WSe₂/hBN van der Waals heterostructure, the growing time, temperature, and the amount of H₂ were optimized for WSe₂ material growth. The growth time was optimized to be 15 min at high temperature 950 °C and high concentration of H₂, 15 sccm. The growth

mechanism could be explained by a simple dissociation, adsorption, and diffusion process in which the W and S adsorb on top of hBN and assemble into WSe₂ film. The properties and the nanostructure of the van der Waals layers were confirmed using different techniques including optical microscopy, Raman, PL, and XPS. Raman and photoluminescence mapping showed that the as-synthesized heterostructure expresses optical modulation. The narrow FWHM of the emission peak (80 meV) of the WSe₂/hBN indicates the high crystallinity and the quality of this heterostructure. This work indicates that our two-step growth process can be used as a common approach to synthesize other wafer-scale and contamination-free out of plane heterostructure 2D materials.

Conflicts of interest

There are no conflicts to declare.

Acknowledgements

This work was supported through funding from the Natural Science and Engineering Research Council of Canada (NSERC), the Canada Research Chairs program (CRC) and Canada Foundation for Innovation. We acknowledge NanoQAM center for the characterization facilities. Also, a special thanks goes to the Taibah University, Kingdom of Saudi Arabia for Alahmadi's Fellowship.

References

- 1 R. F. Frindt and A. D. Yoffe, *Proc. R. Soc. London, Ser. A*, 1963, **273**, 69–83.
- 2 K. Wang, B. Huang, M. Tian, F. Ceballos, M.-W. Lin, M. Mahjouri-Samani, A. Boulesbaa, A. A. Puretzky, C. M. Rouleau and M. Yoon, *ACS Nano*, 2016, **10**, 6612–6622.
- 3 J. C. Shaw, H. Zhou, Y. Chen, N. O. Weiss, Y. Liu, Y. Huang and X. Duan, *Nano Res.*, 2014, **7**, 1–511.
- 4 Z. Li, S. Yang, R. Dhall, E. Kosmowska, H. Shi, I. Chatzakis and S. B. Cronin, *ACS Nano*, 2016, **10**, 6836–6842.
- 5 Z. Li, G. Ezhilarasu, I. Chatzakis, R. Dhall, C.-C. Chen and S. B. Cronin, *Nano Lett.*, 2015, **15**, 3977–3982.
- 6 K. Xu, Z. Wang, X. Du, M. Safdar, C. Jiang and J. He, *Nanotechnology*, 2013, **24**, 465705.
- 7 S. M. Eichfeld, L. Hossain, Y. Lin, A. F. Piasecki, B. Kupp, a G. Birdwell, R. a. Burke, N. Lu, X. Peng, J. Li, A. Azcatl, S. McDonnell, R. M. Wallace, M. J. Kim, T. S. Mayer, J. M. Redwing and J. a Robinson, *ACS Nano*, 2015, **9**, 2080–2087.
- 8 G. Clark, S. Wu, P. Rivera, J. Finney, P. Nguyen, D. H. Cobden and X. Xu, *APL Mater.*, 2014, **101101**, 1–6.
- 9 H.-J. Chuang, X. Tan, N. J. Ghimire, M. M. Perera, B. Chamlagain, M. M.-C. Cheng, J. Yan, D. Mandrus, D. Tománek and Z. Zhou, *Nano Lett.*, 2014, **14**, 3594–3601.
- 10 Q. H. Wang, K. Kalantar-Zadeh, A. Kis, J. N. Coleman and M. S. Strano, *Nat. Nanotechnol.*, 2012, **7**, 699.

- 11 J. S. Ross, P. Klement, A. M. Jones, N. J. Ghimire, J. Yan, D. G. Mandrus, T. Taniguchi, K. Watanabe, K. Kitamura and W. Yao, *Nat. Nanotechnol.*, 2014, **9**, 268.
- 12 M. Okada, T. Sawazaki, K. Watanabe, T. Taniguchi, H. Hibino, H. Shinohara and R. Kitaura, *ACS Nano*, 2014, **8**, 8273–8277.
- 13 H. Wang, T. Taychatanapat, A. Hsu, K. Watanabe, T. Taniguchi, P. Jarillo-herrero and T. Palacios, *Nano*, 2011, **32**, 1209–1211.
- 14 S. Bertolazzi, D. Krasnozhan and A. Kis, *ACS Nano*, 2013, **7**, 3246–3252.
- 15 H. Lim, S. I. Yoon, G. Kim, A.-R. Jang and H. S. Shin, *Chem. Mater.*, 2014, **26**, 4891–4903.
- 16 K. I. Bolotin, K. J. Sikes, Z. Jiang, M. Klima, G. Fudenberg, J. Hone, P. Kim and H. L. Stormer, *Solid State Commun.*, 2008, **146**, 351–355.
- 17 J.-H. Chen, C. Jang, S. Xiao, M. Ishigami and M. S. Fuhrer, *Nat. Nanotechnol.*, 2008, **3**, 206.
- 18 W. Bao, X. Cai, D. Kim, K. Sridhara and M. S. Fuhrer, *Appl. Phys. Lett.*, 2013, **102**, 042104.
- 19 C. R. Dean, a. F. Young, I. Meric, C. Lee, L. Wang, S. Sorgenfrei, K. Watanabe, T. Taniguchi, P. Kim, K. L. Shepard and J. Hone, *Nat. Nanotechnol.*, 2010, **5**, 722–726.
- 20 N. Jain, T. Bansal, C. a. Durcan, Y. Xu and B. Yu, *Carbon*, 2013, **54**, 396–402.
- 21 S. Behura, P. Nguyen, S. Che, R. Debbarma and V. Berry, *J. Am. Chem. Soc.*, 2015, **137**, 13060–13065.
- 22 I. Meric, C. R. Dean, N. Petrone, L. Wang, J. Hone, P. Kim and K. L. Shepard, *Proc. IEEE*, 2013, **101**, 1609–1619.
- 23 G. Lu, T. Wu, Q. Yuan, H. Wang, H. Wang, F. Ding, X. Xie and M. Jiang, *Nat. Commun.*, 2015, **6**, 6160.
- 24 J. Bao, K. Jeppson, M. Edwards, Y. Fu, L. Ye, X. Lu and J. Liu, *Electron. Mater. Lett.*, 2016, **12**, 1–16.
- 25 S. Wang, X. Wang and J. H. Warner, *ACS Nano*, 2015, **9**, 5246–5254.
- 26 Y. Gong, J. Lin, X. Wang, G. Shi, S. Lei, Z. Lin, X. Zou, G. Ye, R. Vajtai, B. I. Yakobson, H. Terrones, M. Terrones, B. K. Tay, J. Lou, S. T. Pantelides, Z. Liu, W. Zhou and P. M. Ajayan, *Nat. Mater.*, 2014, **13**, 1135–1142.
- 27 Y. Feng, K. Zhang, F. Wang, Z. Liu, M. Fang, R. Cao, Y. Miao, Z. Yang, W. Mi, Y. Han, Z. Song and H. S. P. Wong, *ACS Appl. Mater. Interfaces*, 2015, **7**, 22587–22593.
- 28 S. Wang, X. Wang and J. H. Warner, *ACS Nano*, 2015, **9**, 5246–5254.
- 29 D. Liu, X.-Q. Yan, H.-W. Guo, Z.-B. Liu, W.-Y. Zhou and J.-G. Tian, *J. Appl. Phys.*, 2020, **128**, 43101.
- 30 B. Liu, M. Fathi, L. Chen, A. Abbas, Y. Ma and C. Zhou, *ACS Nano*, 2015, **9**, 6119–6127.
- 31 M. Merisalu, T. Kahro, J. Kozlova, A. Niilisk, A. Nikolajev, M. Marandi, A. Floren, H. Alles and V. Sammelseg, *Synth. Met.*, 2015, **200**, 16–23.
- 32 B. V. A. Rao, M. Y. Iqbal and B. Sreedhar, *Corros. Sci.*, 2009, **51**, 1441–1452.
- 33 C. Jung, S. M. Kim, H. Moon, G. Han, J. Kwon, Y. K. Hong, I. Omkaram, Y. Yoon, S. Kim and J. Park, *Sci. Rep.*, 2015, **5**, 15313.
- 34 Y. Rong, Y. Fan, A. L. Koh, A. W. Robertson, K. He, S. Wang, H. Tan, R. Sinclair and J. H. Warner, *Nanoscale*, 2014, **6**, 12096–12103.
- 35 X. Ma, R. Zhang, C. An, S. Wu, X. Hu and J. Liu, *Chin. Phys. B*, 2019, **28**, 37803.
- 36 J. Li and M. Östling, *Electronics*, 2015, **4**, 1033–1061.
- 37 W. D. McFall, D. R. McKenzie and R. P. Netterfield, *Surf. Coat. Technol.*, 1996, **81**, 72–78.
- 38 F. Mahvash, S. Eissa, T. Bordjiba, A. C. Tavares, T. Szkopek and M. Sijaj, *Sci. Rep.*, 2017, **7**, 42139.
- 39 S. Reich, A. C. Ferrari, R. Arenal, A. Loiseau, I. Bello and J. Robertson, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2005, **71**, 205201.
- 40 T. Taira, S. Obata and K. Saiki, *Appl. Phys. Express*, 2017, **10**, 55502.
- 41 A. Pakdel, C. Zhi, Y. Bando, T. Nakayama and D. Golberg, *ACS Nano*, 2011, **5**, 6507–6515.
- 42 S. Das, M. Kim, J. Lee and W. Choi, *Crit. Rev. Solid State Mater. Sci.*, 2014, **39**, 231–252.
- 43 S.-Y. Chen, C. Zheng, M. S. Fuhrer and J. Yan, *Nano Lett.*, 2015, **15**, 2526–2532.
- 44 D. Lee and S. H. Song, *RSC Adv.*, 2017, **7**, 7831–7835.
- 45 G. R. Bhimanapati, D. Kozuch and J. A. Robinson, *Nanoscale*, 2014, **6**, 11671–11675.
- 46 F. Mahvash, E. Paradis, D. Drouin, T. Szkopek and M. Sijaj, *Nano Lett.*, 2015, **15**, 2263–2268.
- 47 J. K. Huang, J. Pu, C. L. Hsu, M. H. Chiu, Z. Y. Juang, Y. H. Chang, W. H. Chang, Y. Iwasa, T. Takenobu and L. J. Li, *ACS Nano*, 2014, **8**, 923–930.
- 48 J. Huang, L. Yang, D. Liu, J. Chen, Q. Fu, Y. Xiong, F. Lin and B. Xiang, *Nanoscale*, 2015, **7**, 4193–4198.
- 49 B. Nie, T. Chen, M. Liang, Y. Wang, J. Zhong, Y. Zhu and E. Mollwo, *J. Appl. Polym. Sci.*, 2011, **30**, 9.
- 50 Z. Yao, J. Liu, K. Xu, E. K. C. Chow and W. Zhu, *Sci. Rep.*, 2018, **8**, 5221.