



Extreme chemical sensitivity of nonlinear conductivity in charge-ordered LuFe_2O_4

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Nonlinear transport behaviors are crucial for applications in electronic technology. At the nonlinear critical turning point, the nonequilibrium states cause rich physics responses to environment. The corresponding study in this field is crucial for physics and industry application. Here nonlinear conductivity in charge-ordered (CO) LuFe_2O_4 has been demonstrated. Remarkable resistivity switching behavior was observed and the gas-sensing property can be reversibly tuned by a small alternation of partial pressure and/or chemical components of the environment. These facts allow us to use LuFe_2O_4 materials as a sensitive chemical gas sensor in technological applications. Careful analysis of the gas sensing process in LuFe_2O_4 suggests a novel sensing mechanism in sharp contrast with that discussed for the conventional gas sensors which depend fundamentally on surface chemical reactions.

In past decades, the precise detections of environmental gas components have gained great importance in a variety of fields, including those of combustion control, food preservation, instrumentation of experiments, biology, medical engineering, and industry¹. Conductometric sensor have been investigated intensively recently due to their simplicity, low cost, and small size^{2,3}. On the other hand, the nonlinear current-voltage (I-V) character and switching of resistivity states in charge-ordered (CO) systems have been found under a variety of environmental conditions, such as magnetic field⁴, stress field⁵, electric field⁶ and with irradiation by X-rays⁷, respectively. We herein report our discovery of the remarkable gas sensitivity of resistivity in the charge/spin frustrated LuFe_2O_4 which shows a remarkable charge-stripe order at room temperature⁸. Our experimental measurements of LuFe_2O_4 in a gaseous environment of O_2 , He, Ar and vacuum often give rise to visibly distinctive transport features. In particular, remarkable switching behaviors of the conductivity are found above the threshold values in association with the field-driven CO melting or collective motions of charge stripes. This gas-sensing property can be reversibly tuned by a small alternation of partial pressure and/or chemical components of the environment. These facts demonstrate that the CO LuFe_2O_4 materials could be used as chemical gas sensors with a fast response and a substantially higher sensitivity at room temperature. In sharp contrast with the conventional conductometric semiconducting gas sensors which depend on reversible interactions of the environmental gas with the surface of the materials^{9,10}, the gas sensitivity of LuFe_2O_4 is found to be in correlation with the nonlinear conductivity influenced by the environmental gases.

Results

In order to characterize the essential transport nonlinearity in the layered LuFe_2O_4 crystal, we first performed our resistivity measurements on a number of single crystal samples along the relevant directions. Figures 1a and 1b show the current density-voltage field (J-E) curves obtained from a single crystalline LuFe_2O_4 sample for the current flowing respectively parallel (E_c) and perpendicular (E_∥ ab plane) to the c-axis direction, illustrating the presence of visible nonlinear J-E characteristics along the c-axis and within the ab plane in a LuFe_2O_4 crystal. The nonlinear behavior in each measurement can be clearly recognizable as a clear anomaly cross the threshold electric field (E_t), the E_t in general appear at around 60 V/cm and 10 V/cm for E parallel and perpendicular to the c-axis direction. According to the empirical power-law relation of $I = kV^\alpha$, the nonlinear coefficients of the single crystal can be estimated to be about 200 and 30 for the I_∥c and I_⊥c directions, these coefficients depend moderately on the current density used in transport measurements.

It has been commonly noted in previous investigations that visible alternation of resistance state could be triggered in a number of the CO systems⁴⁻⁷. Our recent experimental investigations reveal that the nonlinear conductivity in association with CO melting in LuFe_2O_4 is extremely sensitive to the environmental conditions. Figure 2a shows nonlinear J-E behaviors taken from a well-characterized polycrystalline samples upon exposure

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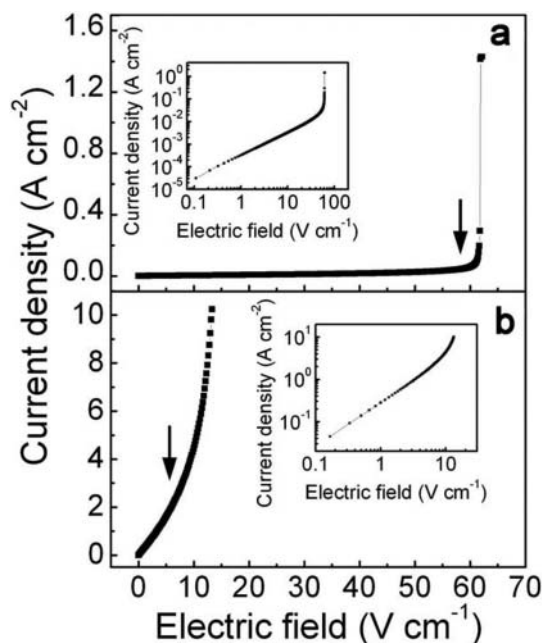


Figure 1 | Current density-electric field (J-E) curves of LuFe_2O_4 measured at room temperature: (a) nonlinear characteristics of the single-crystalline LuFe_2O_4 with the applied electric field parallel to the c-axis; (b) J-E curve of the single crystal with electric field perpendicular to the c-axis. The arrows denote the threshold electric fields for nonlinear conduction. In the insets of (a, b) the log-log plots obtained numerically from the corresponding J-E curves are also shown.

to different gaseous environments. It is recognizable that a jump-like transition of electric current yields a visible nonlinear conductivity at a threshold electric field, or corresponding threshold currents for each curve. The most striking phenomenon revealed in our measurements is the notable divergent behaviors above the threshold, demonstrating the existence of atmosphere-dependent J-E characteristics in the CO LuFe_2O_4 material. It is also noted that the sample in a He atmosphere shows a relatively high threshold current and a larger resistivity in comparison with the data obtained in O_2 and Ar atmosphere. According to the results shown in Fig. 2a, we can estimate the threshold current values for LuFe_2O_4 under different environmental conditions as: $J = 0.5 \text{ A/cm}^2$ for vacuum ($\sim 10^{-1} \text{ Pa}$), 0.7 A/cm^2 for O_2 , and 0.9 A/cm^2 for He, respectively. Moreover, it is also noted that the gas sensitivity depends strongly on applied current in experimental measurements, i.e. the larger electric current is used, and the bigger difference appears in data of the resistivity. This feature is also recognizable in Fig. 2b which shows a series of resistivity curves under different measuring currents above threshold value from 0.3 A/cm^2 to 3.0 A/cm^2 . This measurement was first performed on a LuFe_2O_4 sample under ambient conditions, and then the O_2 was pumped out to lower the chamber pressure progressively to 10^{-1} Pa . The remarkable effects of electrically driven non-linear resistivity and visible resistivity changes with alterations of the O_2 partial pressure are plainly illustrated.

Though there is no uniform definition for the gas-sensor sensitivity now, we can use the ratio of $S_R = \frac{R_{\text{oxy}} - R_{\text{vac}}}{R_{\text{oxy}}} \times 100\%$ to qualitatively illustrate the sensitivity in our measurements. According our experimental data shown in Fig. 2b, S_R depends strongly on the applied electric current, for instance, it is estimated to be $\sim 37.8\%$ and 72.7% for $J = 0.3 \text{ A/cm}^2$ and 3.0 A/cm^2 , respectively.

In order to understand the gas sensitivity in correlation with the charge order in LuFe_2O_4 , we have checked the temperature dependence of the resistivity switching in the temperature range from 295 K

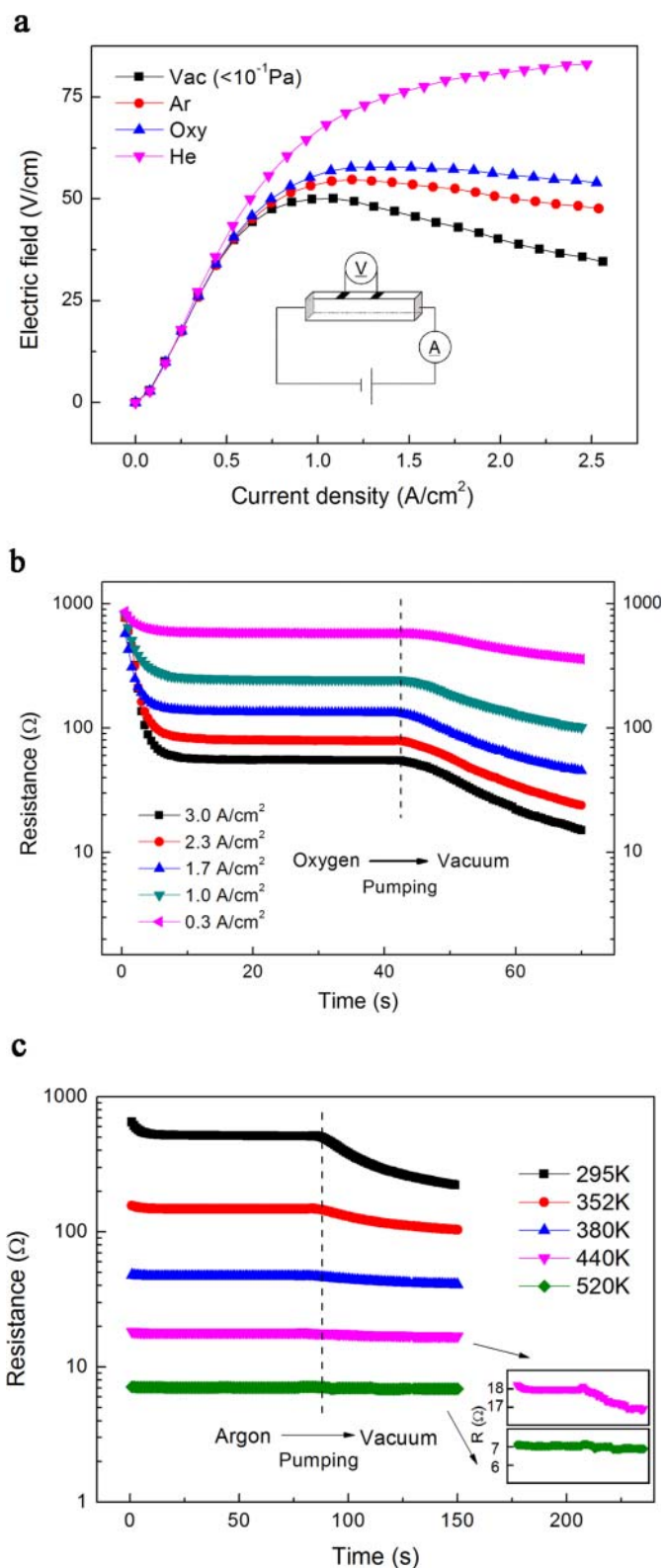


Figure 2 | Resistance switching character of LuFe_2O_4 at room-temperature. (a) Nonlinear current density-voltage field (J-E) behaviors at different atmosphere. Pulsed voltages with the width 200 ms and interval 600 ms were used. Vacuum of about 0.1 Pa was obtained by a mechanical pump. (b) A series of resistance curves under different testing current from 0.3 A/cm^2 to 3.0 A/cm^2 . Current is continuously applied on sample by four probe methods for resistance testing. (c) Resistance curves at different temperatures from 295 K to 520 K between vacuum and argon under testing current 0.3 A/cm^2 . Inset shows details of $R(T)$ curve at 440 K and 520 K respectively.



to 520 K and in the gaseous environment of Ar/vacuum in which the LuFe_2O_4 sample can be safely heated to high temperatures above the charge ordering transition temperature ($T_{\text{co}} \approx 500$ K). As shown in Fig. 2c, the switching of resistance state occurs when the sample chamber is pumped from argon to vacuum below 440 K, suggesting the environmental condition plays important role for the resistivity switching. This switching feature gradually disappears as the temperature rises to the CO transition value ($T_{\text{co}} \approx 500$ K) above which LuFe_2O_4 has a charge-disordered state with a relatively high conductivity. These facts demonstrate the presence of an essential correlation between charge ordering and gas sensitivity in present system.

In Fig. 3 the reversibility of the resistance switching in LuFe_2O_4 is discussed. We performed experimental measurements on a sample in the O_2 chamber. Decrease of oxygen pressure by pumping out the O_2 from the chamber to vacuum leads to a rapid resistance switch into low resistance states; then when the O_2 is reintroduced into the chamber, the experimental data switches reversibly to higher resistance state. Figure 3 shows a series of the experimental results as the surrounding chemical components was cycled between vacuum (≤ 0.1 Pa) and O_2 , illustrating the remarkable resistance switching and high gas sensitivity for LuFe_2O_4 at room temperature. It is clearly recognizable that a rapid and reversible change in the resistance oscillates in response to the alternating environmental pressure. Similar resistivity switching behavior has been also observed by using the LuFe_2O_4 single crystal samples, demonstrating an essential bulk property of gas sensitivity instead of effects from surface or grain boundaries as commonly discussed for the conventional gas sensors. Besides, measurements in oxygen environment often show some additional complex changes, which are believed to be connected with the interstitial and/or deficient oxygen atoms in the layered LuFe_2O_4 .

Considering the notable effects of interstitial oxygen ions on both microstructure and transport properties of LuFe_2O_4 , we have also examined the changes of microstructure in association with non-linear transition by using *in situ* TEM observation, which could be also important for understanding the sensing mechanism in present system. Figure 4 clearly shows the superstructure modulations in correlation with charge order and oxygen order as discussed in our previous investigations¹¹. Following the increase of applied electric field, the superstructure spots in the diffraction pattern gradually become invisible and disappear as illustrated in Fig. 4b. Based on our TEM experimental data, we estimate that the threshold electric

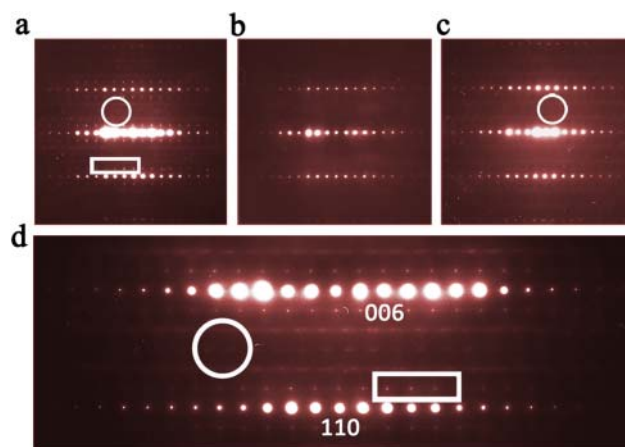


Figure 4 | The *in situ* TEM electron diffraction patterns taken along the [1–10] zone axis direction. (a) The diffraction spots in circle and rectangular is charge ordering and oxygen modulations respectively. (b) The charge-ordered (CO) state and oxygen modulations disappear when applied electric field is larger than E_{th} (~ 45 V/cm). (c) The CO state has recovered after the applied electric field is removed while the oxygen modulations fail to reappear. (d) The CO and oxygen modulations reemerge when the sample is placed in air for more than 24 hrs.

field is approximately 45 V/cm. After the external electric field is removed, the CO satellite spots in the diffraction pattern re-emerge in the same crystal but the superstructure from oxygen order shows clear irreversibility in a high-vacuum column of TEM. When the sample was placed in air for more than 24 hrs, the superstructure spots of oxygen order appear again as showed in Fig. 4d. These facts suggest that certain microstructure features in present system also vary under the applied current, and their reversibility could be also important for technological applications of gas sensing materials.

Discussion

The gas sensing process and fundamental mechanism for the gas response in the CO LuFe_2O_4 as essential issues were also concerned in our investigations. Based on localized electron variable-range hopping (VRH) picture¹², the electric conductivity in the CO materials could be enormously sensitive to the modification of the

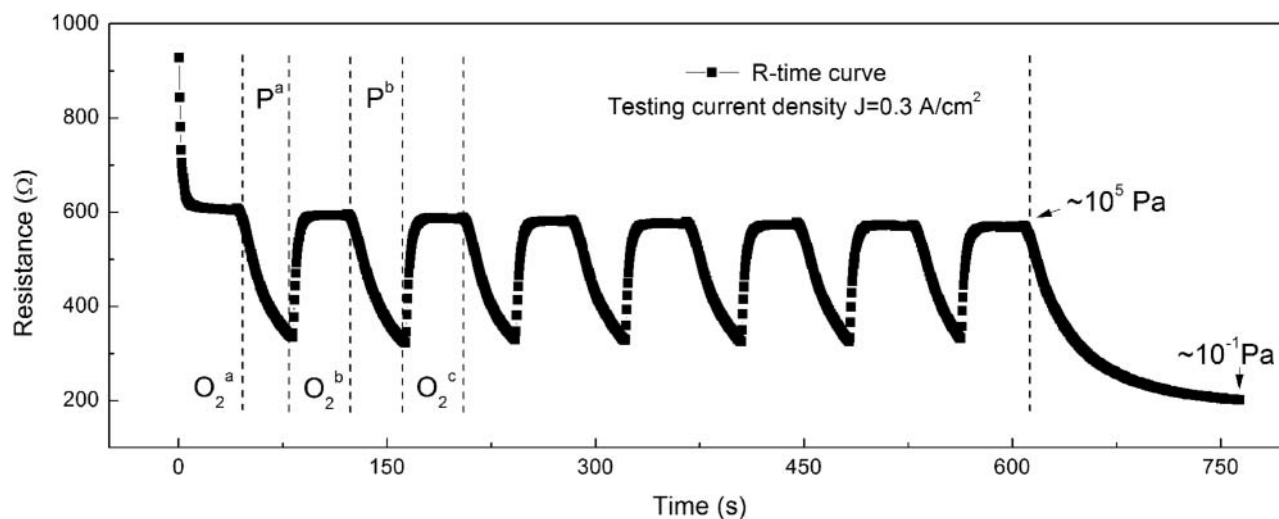


Figure 3 | Period of resistance change when switching atmosphere between oxygen and vacuum regularly. The testing current is 0.3 A/cm². In the different regions divided by dashed lines, the O_2^{a} is the oxygen region where sample was placed initially; P^{a} is the region of pumping to vacuum; O_2^{b} is the region of introducing oxygen again; P^{b} is the region of re-pumping vacuum.



thermodynamic equilibrium. In a CO system, when nonlinear transition occurs above the threshold value, the melting of the CO state, together with the collective movement of charge stripes, can result in a strong electron-phonon interaction which contributes the Joule heating effect. Though the Joule heating could not change the essential feature of nonlinear conductivity, it indeed can induce visible alterations in the resistivity as discussed in our previous work on LuFe_2O_4 materials¹³. In the present study, LuFe_2O_4 was exposed to variable gaseous environments, and the thermal equilibrium states of samples depend mainly on thermal balance between Joule heating and thermal diffusion.

If the applied current density does not reach to the threshold value to trigger nonlinear conductivity, no clear gas sensitivity can be observed in our measurements. On the other hand, if the applied current is larger than the threshold value, nonlinear conductivity depends strongly on temperature and the thermal equilibrium. Because the thermal conductivity/diffusion coefficients of gases are different for O_2 , Ar and He, alternation of the composition and partial pressure of gases could lead to certain changes in the thermal equilibrium, which sequentially yields a visible change in resistivity as revealed in our experiments. Under a constant applied current, it is expected that the higher conductivity/diffusion coefficient of a gas should result in a bigger resistance state, which is in a good agreement with our experimental data. For example, the conductivity coefficients for He^{14} , O_2^{15} and Ar^{16} are 156.7, 26.3 and 17.9 mW/mK at 300 K (1atm), respectively. The corresponding resistances are 847.7, 641.8 and 587.7 ohms in these gas environments under a constant testing current density of 1.5 A/cm².

Careful analysis of the gas sensing process suggests a novel sensing mechanism based on the thermal conductivity/diffusion coefficients of environment gases in stark contrast with the conventional gas sensors that depend chiefly on surface reactions. Moreover, *in situ* TEM observation suggests that certain microstructure features in present system also are sensitive to the vacuum environment, and their reversibility could also be important for technological applications of gas sensing materials.

The experiments on other CO materials also reveal very similar sensitivity phenomena as shown in supplementary materials. These observations open up new prospects for finding new functional applications and understanding interesting physics in CO materials and other nonlinear materials.

Methods

The well-characterized LuFe_2O_4 polycrystal and single crystal samples used in present study were synthesized by a conventional solid-state reaction⁵. Polycrystalline samples of LuFe_2O_4 material were synthesized from stoichiometric mixtures of Lu_2O_3 (99.99%) and Fe_2O_3 (99.99%) under a controlled oxygen partial pressure atmosphere using a CO_2 - H_2 mixture at 1200°C for 48 hours. In order to obtain the single crystals, the polycrystalline LuFe_2O_4 powder was heated up to 1620°C in a platinum crucible and then the melted solid was cooled to 900°C at a rate of 1°C/min. The typical size of the polycrystalline sample used in our measurements is $2 \times 5 \times 1 \text{ mm}^3$. A thermocouple as commonly used in previous studies was attached on the sample holder to detect the temperature, and an additional thermometer was attached directly on the sample for certain measurements. Silver contact pads with a radius of ~1 mm were deposited on samples. Current density-electric field (J-E) curves and transport testing were obtained by the two-probe and/or four-probe methods with a Keithley 2400/2611A source meter. In order to protect the samples, a compliance current of 200 mA was used during the testing of nonlinear J-E curves. Transmission electron microscopy (TEM) investigations were performed at room temperature on a Tecnai F20

(200 kV) double-tilt TEM holder on which the *in situ* observations upon applied electric field can be performed.

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Author contributions

J.Q.L., S.C. planned and coordinated the experiments. S.C., J.L. carried out gas sensing resistivity experiments. Z.W., S.C. performed TEM measurement and analysis. L.J.Z. prepared the figure 1. Y.B.Q. grew LuFe_2O_4 polycrystal and single crystal samples. H.F.T., C. M. and H.X.Y. helped with data interpretation. J.Q.L., S.C. wrote the paper.

Additional information

Supplementary information accompanies this paper at <http://www.nature.com/scientificreports>

Competing financial interests: The authors declare no competing financial interests.

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