


Can C_p Be Less Than C_V ?

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ABSTRACT: Can C_p be less than C_V ? This is a fundamental question in physics, chemistry, chemical engineering, and mechanical engineering. This question hangs in the minds of many students, instructors, and researchers. The first instinct is to answer “Yes, for water between 0 and 4 °C” if one knows that water expands as temperature decreases in this temperature range. The same question is asked in several Physical Chemistry and Physics textbooks. Students are supposed to answer that water contracts when heated at below 4 °C in an isobaric process. Because work is done to the contracting water, less heat is required to increase the water temperature in an isobaric process than in an isochoric process. Therefore, C_p is less than C_V . However, this answer is fundamentally flawed because it assumes, implicitly and incorrectly, that the internal energy change of water depends solely on its temperature change. Neglecting the variation of the internal energy with volume (internal pressure) will invalidate the Clausius inequality and violate the second law of thermodynamics. Once the internal pressure is properly taken into account, it becomes clear that C_p cannot be less than C_V for any substance at any temperature regardless of the sign of the thermal expansion coefficient of the substance.

- I. Thou shalt not violate the second law of thermodynamics.
- II. Substances shalt not expand as pressure increases.
- III. C_p shalt not be less than C_V .

INTRODUCTION

“The heat capacity $C_{p,m}$ is less than $C_{V,m}$ for $H_2O(l)$ near 4 °C. Explain this result”. This is one of the concept problems in Chapter 3 of Physical Chemistry (3rd ed.) by Engel and Reid.¹ Students are supposed to explain that liquid water contracts when heated at below 4 °C in an isobaric process. Because work is done to the contracting water, less heat is required to increase its temperature in an isobaric process than in an isochoric process. The above explanation is flawed because it assumes, implicitly and incorrectly, that the internal energy change of water in an isobaric process is the same as in an isochoric process. As a result of the flawed assumption, $C_{p,m}$ is less than $C_{V,m}$ for $H_2O(l)$ near 4 °C. The same flawed assumption also appears on page 252 of Physical Chemistry by Cooky.² “Why should C_p be different from C_V ? Briefly, because if the volume can change, some of the heat added to the system is used to do work by expansion, leaving less to raise the temperature of the sample”. This seemingly small mistake is beyond just a trivial omission of the variation of the internal energy with volume; it results in the violation of the second law of thermodynamics and thereby destroys the foundation of thermodynamics. Unfortunately, this severe error might have been infused into thousands of chemistry students and perhaps physics students as well. In University Physics by Ling, Sanny, and Moebs,³ Conceptual Question 3.13 asks: “Most materials expand when heated. One notable exception is water between 0 and 4 °C, which actually decreases in volume with the increase in temperature. Which is greater for water in this temperature region, C_p or C_V ?” The answer key is provided in this textbook: “Typically C_p is greater

than C_V because when expansion occurs under constant pressure, it does work on the surroundings. Therefore, heat can go into internal energy and work. Under constant volume, all heat goes into internal energy. In this example, water contracts upon heating, so if we add heat at constant pressure, work is done on the water by surroundings and therefore, C_p is less than C_V ”. In University Physics by Young and Freedman,⁴ Q19.10 asks “There are a few materials that contract when their temperature is increased, such as water between 0 and 4 °C. Would you expect C_p for such materials to be greater or less than C_V ? Explain?” The key is not provided but this question alone may be misleading. And this error has affected not only classrooms: a research article published in the Journal of Chemical Physics claims that C_p can be less than C_V for silicon at low temperatures.⁵ There seems to be a widespread misunderstanding about the difference between C_p and C_V in academia, which originates from the improper omission of the variation of the internal energy with volume and results in the violation of the second law of thermodynamics.

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RESULTS AND DISCUSSION

C_p cannot be less than C_V because this violates the second law of thermodynamics. Our proof starts with the following equation that is rigorous for all matters^{1,6–8}

$$C_p - C_V = T \left(\frac{\partial P}{\partial T} \right)_V \left(\frac{\partial V}{\partial T} \right)_P = TV \frac{\alpha^2}{\kappa}$$

where α is the isobaric thermal expansion coefficient and κ is the isothermal compressibility

$$\alpha = \frac{1}{V} \left(\frac{\partial V}{\partial T} \right)_P$$

$$\kappa = \frac{-1}{V} \left(\frac{\partial V}{\partial P} \right)_T$$

Because T and V are always positive and α^2 is non-negative, κ will have to be negative if C_p is less than C_V . Experimentally, the value of κ of liquid water was found to be always positive.^{9–12} Specifically, under 1 atm pressure, the κ value for liquid water is $5.1 \times 10^{-10} \text{ Pa}^{-1}$ at 0°C or $4.5 \times 10^{-10} \text{ Pa}^{-1}$ at 25°C .¹² The value of κ must also be positive between 0°C and 25°C because the κ vs T curve must be continuous for liquid water within this temperature range. A single negative value of κ would result in κ having to be equal to zero at two different temperatures between 0 and 25°C . The difference between C_p and C_V would then have to be infinity at one of these two temperatures using the $C_p - C_V = TV \frac{\alpha^2}{\kappa}$ equation because α is zero only at one temperature near 4°C . Hereafter, we will prove in a more rigorous yet still concise manner that a negative κ for any substance would invalidate the Clausius inequality and consequently violate the second law of thermodynamics.

First, we prove that a negative κ would invalidate the Clausius inequality ($TdS > dq_{\text{irreversible}}$). The infinitesimal change to the state of a system, from (P, V, T, S) to $(P + dP, V + dV, T + dT, S + dS)$, can be made via a reversible process or an irreversible process. In the reversible process, $dU = TdS - PdV$. In the irreversible process, $dU = dq_{\text{irreversible}} + dw_{\text{irreversible}}$. Because U is a state function, we have $TdS - PdV = dq_{\text{irreversible}} - P_{\text{external}} dV$. Therefore, $TdS - dq_{\text{irreversible}} = (P - P_{\text{external}})dV$. If κ is negative, $\left(\frac{\partial V}{\partial P} \right)_T = -\kappa V$ must be positive. In this case, when P_{external} is greater than P , the system expands as its system pressure increases to match with P_{external} . Because $(P - P_{\text{external}})$ is negative and dV is positive, the Clausius inequality is invalidated: $TdS - dq_{\text{irreversible}} < 0$. This further violates the second law of thermodynamics:

$$TdS < dq_{\text{irreversible}} = -dq_{\text{surroundings}} = -TdS_{\text{surroundings}}$$

$$dS_{\text{universe}} = dS + dS_{\text{surroundings}} = TdS + TdS_{\text{surroundings}} < 0$$

We have proved that a negative κ would violate the second law of thermodynamics and thus $C_p - C_V = TV \frac{\alpha^2}{\kappa}$ cannot be negative.

It is well known that α can be zero, e.g., for water at $\sim 4^\circ\text{C}$. In this case, C_p equals C_V . Can κ be zero too? *No*. If any substance had a zero κ , its α value would have to be zero to avoid the difference between C_p and C_V approaching infinity. The volume of such a substance under such conditions would then be independent of pressure and temperature as $dV =$

$V\alpha dT - V\kappa dP = 0$. This is apparently unphysical. Therefore, we conclude that κ must be positive for all substances.

The above rigorous proof is supported by numerical calculations for water at 0 – 4°C . The α value of liquid water under 1 atm pressure is $-6.8 \times 10^{-5} \text{ K}^{-1}$ at 0°C or $2.57 \times 10^{-4} \text{ K}^{-1}$ at 25°C .¹² Its κ value under 1 atm pressure is $5.1 \times 10^{-10} \text{ Pa}^{-1}$ at 0°C or $4.5 \times 10^{-10} \text{ Pa}^{-1}$ at 25°C .¹² Using the

$C_{p,m} - C_{V,m} = TV_m \frac{\alpha^2}{\kappa}$ equation, we find that $C_{p,m}$ is greater than $C_{V,m}$ by $0.04 \text{ J mol}^{-1} \text{ K}^{-1}$ at 0°C or $0.8 \text{ J mol}^{-1} \text{ K}^{-1}$ at 25°C . $C_{p,m} - C_{V,m}$ is positive at both 0 and 25°C (and in between). The difference, however, is much smaller than the $C_{p,m}$ ($75.3 \text{ J mol}^{-1} \text{ K}^{-1}$) of water at 25°C and 1 atm pressure. Therefore, when we refer to the heat capacity of liquid water or many other liquids or solids, there is often no need to specify the isobaric or isochoric condition. Is it always true that $C_{p,m}$ is close to $C_{V,m}$ for liquids and solids or at least for liquids that are similar to water such as methanol? *No*. Methanol has a similar molecular structure as water but the difference between $C_{p,m}$ and $C_{V,m}$ of methanol is expected to be much greater than that of water. This is because its isobaric thermal expansibility ($1.49 \times 10^{-3} \text{ K}^{-1}$) is ~ 6 times that of water under the standard condition, while its isothermal compressibility ($1.2 \times 10^{-9} \text{ Pa}^{-1}$) is only ~ 3 times that of water at 25°C .¹ Moreover, the molar volume of methanol (40.7 mL mol^{-1}) is more than twice that of water due to its greater molar mass (32 g mol^{-1}) and lower density (0.787 g mL^{-1})¹³ at 25°C . The result is that the calculated difference (ca. $22 \text{ J mol}^{-1} \text{ K}^{-1}$) between $C_{p,m}$ ($81 \text{ J mol}^{-1} \text{ K}^{-1}$)¹ and $C_{V,m}$ for methanol is much greater than that of water at 25°C and even greater than the ideal gas constant. The $C_{V,m}$ for methanol is thus estimated to be $81 - 22 = 59 \text{ J mol}^{-1} \text{ K}^{-1}$ at 25°C and 1 atm. At the room temperature, the experimental data of $C_{V,m}$ for methanol is unavailable at 1 atm; it is $72 \text{ J mol}^{-1} \text{ K}^{-1}$ under 137 atm.¹⁴ Another example of $C_{p,m}$ being much larger than $C_{V,m}$ is liquid benzene. The values of $C_{p,m}$ and $C_{V,m}$ of benzene are 136 and $95 \text{ J mol}^{-1} \text{ K}^{-1}$, respectively, at room temperature and under ~ 0.14 atm of saturated benzene vapor.¹⁵ For benzene, $C_{p,m}$ is $41 \text{ J mol}^{-1} \text{ K}^{-1}$ greater than $C_{V,m}$. This vast difference between $C_{p,m}$ and $C_{V,m}$ of benzene is due to its very large molar volume (90 mL mol^{-1}) and significant thermal expansibility ($1.14 \times 10^{-3} \text{ K}^{-1}$).¹

There are caveats associated with the above analysis of the difference between C_p and C_V . One caveat is that $\left(\frac{\partial P}{\partial T} \right)_V$ must be negative for water between 0 and 4°C because $\left(\frac{\partial V}{\partial T} \right)_P$ is negative in $C_p - C_V = T \left(\frac{\partial P}{\partial T} \right)_V \left(\frac{\partial V}{\partial T} \right)_P$. This seems counter-intuitive, but, when temperature decreases between 0 and 4°C , the pressure of liquid water does increase because water at lower temperatures resembles more like ice and thus needs to occupy more space when the pressure is held constant or exerts (and feels) higher pressure when the volume is held constant. Another explanation is that $\left(\frac{\partial S}{\partial V} \right)_T = \left(\frac{\partial P}{\partial T} \right)_V$ is negative. The entropy of liquid water decreases when its volume increases because, given more space, water can form a more ordered structure (to reduce energy) that resembles the crystalline structure of ice. One may ask, if $\left(\frac{\partial S}{\partial V} \right)_T$ is negative for water, wouldn't the entropy of the universe decrease as the water volume increases while the volume of its surroundings decreases? *No*. While water expands at a constant temperature

between 0 and 4 °C, its internal energy decreases because the hydrogen bonding between water molecules, given more space, becomes more significant. This internal energy decrease becomes heat that flows into the surroundings and consequently increases the entropy of the surroundings. Overall, the entropy increase of the surroundings would be large enough to compensate the entropy loss of the expanding water; the second law of thermodynamics is not in jeopardy. A second caveat is that some materials such as NaNH_2BH_3 may expand along one or two (out of three) dimensions under compression.^{16–21} In these cases, the overall three-dimensional volume of such a material must decrease, dictated by the second law of thermodynamics.

CONCLUSIONS

When the temperature is held constant, no substance can expand against higher pressures because this would invalidate the Clausius inequality and violate the second law of thermodynamics. Consequently, C_p can never be less than C_v , even for water at temperatures between 0 and 4 °C.

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Notes

The authors declare no competing financial interest.

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REFERENCES

- (1) Engel, T.; Reid, P. *Physical Chemistry*, 3rd ed.; Pearson: Boston, 2012.
- (2) Cooksy, A. *Physical Chemistry: Thermodynamics, Statistical Mechanics, & Kinetics*; Pearson: Boston, 2014.
- (3) Ling, S. J.; Sanny, J.; Moebs, W. *University Physics*; Rice University, 2018; Vol. 2.
- (4) Young, H. D.; Freedman, R. A. *University Physics with Modern Physics*, 14th ed.; Pearson: Boston, 2016.
- (5) Fine, M. E. C_p – C_v in Silicon and Germanium. *J. Chem. Phys.* **1953**, *21*, 1427.
- (6) Smith, N. O. The Difference between C_p and C_v for Liquids and Solids. *J. Chem. Educ.* **1965**, *42*, 654–655.
- (7) McQuarrie, D. A.; Simon, J. D. *Physical Chemistry: A Molecular Approach*; University Science Books: Sausalito, CA, 1997.
- (8) Levine, I. N. *Physical Chemistry*, 3rd ed.; McGraw-Hill: New York, 1988.

(9) Millero, F. J.; Curry, R. W.; Drost-Hansen, W. Isothermal Compressibility of Water at Various Temperatures. *J. Chem. Eng. Data* **1969**, *14*, 422–425.

(10) Fine, R. A.; Millero, F. J. Compressibility of Water as a Function of Temperature and Pressure. *J. Chem. Phys.* **1973**, *59*, 5529–5536.

(11) Bradshaw, A.; Schleicher, K. Compressibility of Distilled Water and Seawater. *Deep-Sea Res. Oceanogr. Abstr.* **1976**, *23*, 583–593.

(12) Kell, G. S. Density Thermal Expansivity, and Compressibility of Liquid Water from 0. °C to 150. °C. Correlations and Tables for Atmospheric Pressure and Saturation Reviewed and Expressed on 1968 Temperature Scale. *J. Chem. Eng. Data* **1975**, *20*, 97–105.

(13) Wei, I. C.; Rowley, R. L. Binary Liquid Mixture Viscosities and Densities. *J. Chem. Eng. Data* **1984**, *29*, 332–335.

(14) Kuroki, T.; Kagawa, N.; Endo, H.; Tsuruno, S.; Magee, J. W. Specific Heat Capacity at Constant Volume for Water, Methanol, and Their Mixtures at Temperatures from 300 K to 400 K and Pressures to 20 MPa. *J. Chem. Eng. Data* **2001**, *46*, 1101–1106.

(15) *CRC Handbook of Chemistry and Physics*, 97th ed.; Haynes, W. M., Ed.; CRC Press, 2016.

(16) Cairns, A. B.; Goodwin, A. L. Negative Linear Compressibility. *Phys. Chem. Chem. Phys.* **2015**, *17*, 20449–20465.

(17) Magos-Palasyuk, E.; Fijalkowski, K. J.; Palasyuk, T. Chemically Driven Negative Linear Compressibility in Sodium Amidoborane, $\text{Na}(\text{NH}_2\text{BH}_3)$. *Sci. Rep.* **2016**, *6*, No. 28745.

(18) Serra-Crespo, P.; Dikhtiarenko, A.; Stavitski, E.; Juan-Alcañiz, J.; Kapteijn, F.; Coudert, F.-X.; Gascon, J. Experimental Evidence of Negative Linear Compressibility in the MIL-53 Metal–Organic Framework Family. *CrystEngComm* **2015**, *17*, 276–280.

(19) Cairns, A. B.; Catafesta, J.; Levelut, C.; Rouquette, J.; van der Lee, A.; Peters, L.; Thompson, A. L.; Dmitriev, V.; Haines, J.; Goodwin, A. L. Giant Negative Linear Compressibility in Zinc Dicyanoaurate. *Nat. Mater.* **2013**, *12*, 212–216.

(20) Li, W.; Probert, M. R.; Kosa, M.; Bennett, T. D.; Thirumurugan, A.; Burwood, R. P.; Parinello, M.; Howard, J. A. K.; Cheetham, A. K. Negative Linear Compressibility of a Metal–Organic Framework. *J. Am. Chem. Soc.* **2012**, *134*, 11940–11943.

(21) Cai, W.; Katrusiak, A. Giant Negative Linear Compression Positively Coupled to Massive Thermal Expansion in a Metal–Organic Framework. *Nat. Commun.* **2014**, *5*, No. 4337.