

Molecular Dynamics Simulation of the Three-Phase Equilibrium Line of CO₂ Hydrate with OPC Water Model

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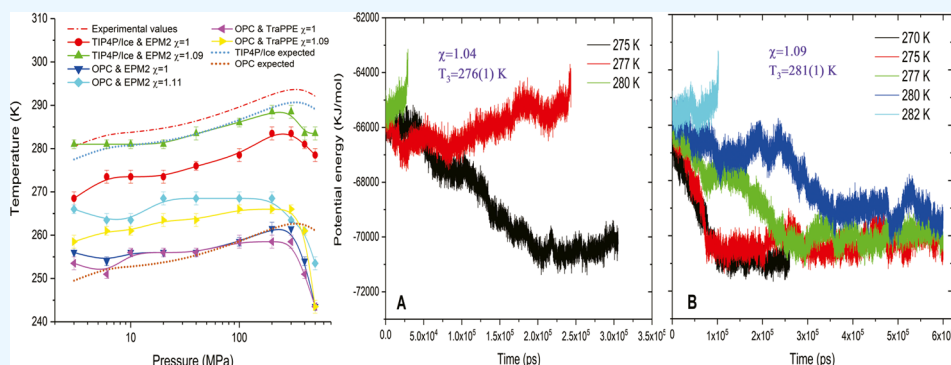
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ABSTRACT: The three-phase coexistence line of the CO₂ hydrate was determined using molecular dynamics (MD) simulations. By using the classical and modified Lorentz–Berthelot (LB) parameters, the simulations were carried out at 10 different pressures from 3 to 500 MPa. For the OPC water model, simulations with the classic and the modified LB parameters both showed negative deviations from the experimental values. For the TIP4P/Ice water model, good agreement with experimental equilibrium data can be achieved when the LB parameter is adjusted based on the solubility of CO₂ in water. Our results also show that the influence of the water model on the equilibrium prediction is much larger than the CO₂ model. Current simulations indicated that the influence of the H₂O–H₂O and H₂O–CO₂ cross-interactions’ parameters might contribute equally to the accurate prediction of T₃. According to our simulations, the prediction of T₃ values showed relatively higher accuracy while using the combination of TIP4P/Ice water and EPM2 CO₂ with modified LB parameter. Furthermore, varied χ values are recommended for accurate T₃ estimation over a wide pressure range. The knowledge obtained in this study will be helpful for further accurate MD simulation of the process of CO₂/CH₄ replacement.

1. INTRODUCTION

Clathrate hydrates are crystalline, nonstoichiometric ice-like compounds formed at conditions of relatively lower temperature and higher pressure.^{1,2} The structure of the gas hydrate is composed of a three-dimensional network of hydrogen-bonded water molecules, giving rise to well-defined “cages” where guest molecules can reside. According to the size and number of cages within a single unit cell of the hydrate crystal, hydrate structures can be classified as sI, sII, and sH. The structure type of the hydrate is mainly determined by the size of the guest molecule.

CO₂ hydrates received increasing attention in recent years because of the global concern about greenhouse gas emission. Since the CO₂ emission caused by global warming is recognized as one of the most severe environmental threats in the future, it is urgent to find some effective methods for carbon capture and sequestration. Under such circumstances, hydrate technology became an alternative solution to both the CO₂ capture and storage processes. Specifically, carbon dioxide

can be captured from flue gases in the form of hydrates,^{3,4} owing to its relatively higher thermodynamic stability.⁵ Moreover, methane recovery from natural methane hydrate by CO₂ replacement could provide a promising approach for simultaneous hydrocarbon retrieving and carbon sequestration.^{6–9} In order to explore the potential usage of carbon dioxide hydrates, it is necessary to develop reliable methodologies for better understanding and accurate prediction of both their kinetic and equilibrium properties.

Molecular dynamics (MD) simulation has long been proved to be a useful tool for the study of clathrate hydrates.¹⁰ It can

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not only provide insight into molecular phenomena but also offer a useful tool to predict the macroscopic hydrate properties when the experimental data are unavailable. In order to perform accurate simulations, the methodologies and force fields used in MD studies need validation first by comparing the MD results to reliable experimental data. The three-phase coexistence line is one of the most important parameters in hydrate studies. As methane hydrate is the most commonly found hydrate in nature, many previous MD simulations have been performed to study the phase equilibria of methane hydrate.^{11–16} In these works, the performance of different H₂O models was compared, finding that the consideration of positive deviation from the ideal Lorentz–Berthelot combination rules could yield a remarkable representation of experimental coexistence results. Except for these phase equilibrium-related studies, MD simulation has been applied in a variety of research areas of methane hydrate from a molecular point of view.^{10,17–28}

Compared with methane hydrate, the MD prediction of phase equilibria of CO₂ hydrate has attracted limited attention so that only a few works have been done in this research field. Sarupria and Debenedetti²⁹ carried out MD simulations to study the relationship between the dissociation rate and cage occupancy of carbon dioxide hydrates. Using TIP4P/2005³⁰ for water and TraPPE³¹ for carbon dioxide, they found that the dissociation temperature is modestly related to the hydrate occupancy. As the configuration used in their study did not contain the gas phase of carbon dioxide, the dissociation temperature of the CO₂ hydrate was treated as the three-phase equilibrium temperature. They reported that the deviation between the melting point of CO₂ hydrate and the melting point of ice as predicted by TIP4P/2005³² is about 4–8 K at 3.05 MPa, which is close to the temperature difference of the experimental hydrate T₃ and ice melting values (~7 K). Tung et al.³³ used a three-phase configuration to simulate the growth of carbon dioxide hydrate and calculated the T₃ values at different pressure conditions. The TIP4P/Ew³⁴ water model and EPM2³⁵ CO₂ models were applied in the MD simulation, with Lennard–Jones cross-interaction parameters suggested by Sun and Duan.³⁶ At relatively lower pressure conditions (less than 30 MPa), their predictions agree well with the experimental values, while at higher pressure conditions (100 MPa), the deviations of T₃ become significant. Míguez et al.³⁷ studied the three-phase equilibria of carbon dioxide hydrate using the direct-phase coexistence method up to 500 MPa. The TIP4P/Ice³⁸ and TIP4P/2005 force fields were used for water and several different force fields were used for carbon dioxide. Their simulations showed that the evolution of the system was the same when all large cages were initially occupied (the cage occupancies range from 75 to 100%), and the influence of the water model on the T₃ prediction is obviously larger than the carbon dioxide model. At higher simulation pressure (300–500 MPa), both water models showed the retrograde behavior of the equilibrium line. For the TIP4P/Ice model, the introduction of a modified cross-interaction parameter $\chi = 1.13$ was effective in correcting the predicted T₃ up to 200 MPa. However, T₃ was still underestimated with the modified χ value while in the pressure range of 200–400 MPa. Costandy et al.³⁹ also used the direct-phase coexistence methodology to predict the T₃ value of carbon dioxide hydrate in the pressure range of 20–500 MPa. The authors used the TIP4P/Ice and TIP4P/2005 force fields for water and the TraPPE force field for carbon dioxide. Their simulations indicated that T₃ values

cannot be predicted accurately using the classic LB combination rules. The authors then modified the water–guest interaction parameters based on the solubility of carbon dioxide in water. In this case, they found that the obtained T₃ values are closely related to the melting point of ice, as predicted by the water model.

Although a variety of water models have been applied in previous MD studies on CO₂ hydrate, the water model referred to as the “optimal” point charge (OPC) water model developed by Izadi et al.⁴⁰ has not been included yet. Compared with the commonly used TIP4P/Ice water models in ice and gas hydrate simulations, the bulk properties of water (including density, self-diffusion, dielectric constant, heat capacity, etc.) reproduced by the OPC model are significantly more accurate over a wide range of temperatures. However, the melting point predicted by the OPC model exhibits a large deviation⁴¹ when compared with the TIP4P/Ice model. Our previous study¹⁶ has already shown the relatively good performance of the OPC model in the prediction of T₃ of methane hydrate especially at lower pressure, which encourages us to perform MD simulation to study the three-phase coexistence line of CO₂ hydrate with this water model. The present work focuses on the prediction of the phase equilibria of carbon dioxide hydrate with an emphasis on the role of water–guest interactions in the determination of the three-phase coexistence temperature. To this purpose, we follow the direct-phase coexistence methodology that was successfully used in our previous study of methane hydrates. In line with the work of Míguez et al.,³⁷ we test two different CO₂ force fields, namely, TraPPE and EPM2, in combination with the OPC force field for water. Also, the combination of the TIP4P/Ice water model and the EMP2 CO₂ model is included for comparison.

2. METHODOLOGY

The direct-phase coexistence method is used in the current study to determine the coexistence temperature in the CO₂–water–hydrate system. The system configuration of this study consists of a CO₂ hydrate slab, a CO₂ slab, and two liquid water slabs. The CO₂ gas slab is inserted into two water slabs to generate a hydrate–water–CO₂–water arrangement. A buffer distance of 0.1 nm is assigned between each slab to avoid bad contacts while processing energy minimization.

The hydrate slab is composed of 2 × 2 × 2 cubic unit cells (368 water molecules and 64 CO₂ molecules). Each hydrate unit cell is built as follows: positions of the water oxygen atoms are obtained from X-ray crystallography,⁴² and the water hydrogen atoms are inserted by adjusting the orientations to obey the Bernal–Fowler rule and minimize the potential energy and net unit cell dipole moment.⁴³ CO₂ molecules are assumed to fully occupy the water polyhedral cages constituted by the water hydrogen bond network and situate at the cage center. Each water slab contains 368 water molecules, and the CO₂ gas slab contains 256 CO₂ molecules to avoid the generation of a gas bubble during hydrate formation. Totally 1104 water molecules and 320 CO₂ molecules are present in the initial simulation box with initial dimensions of 2.5 × 2.5 × 9.6 nm³. A snapshot of the initial configuration of the simulation box before pre-equilibrium is shown in Figure 1. The TIP4P/Ice and the OPC water models are adopted for water molecules, while the TraPPE and EMP2 models are used for CO₂. A total of six sets of simulations were conducted (combination of OPC with TraPPE or EPM2 and TIP4P/Ice

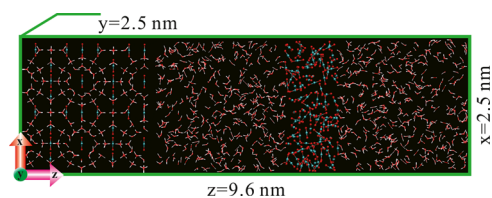


Figure 1. Initial configuration of the simulation system for the direct coexistence method before equilibrium. The red and white lines represent water molecules, and the cyan and red spheres represent CO₂ molecules.

with EPM2), of which three used the classical LB combination rules to determine the water–gas interaction, and the other three used a modified factor (χ) to correct the LB cross-interaction energy parameter, as shown in the following equation:

$$\epsilon_{\text{CO}_2-\text{H}_2\text{O}} = \chi \sqrt{\epsilon_{\text{CO}_2-\text{CO}_2} \epsilon_{\text{H}_2\text{O}-\text{H}_2\text{O}}} \quad (1)$$

The χ value is determined based on the solubility of CO₂ in water. To calculate the solubility dependence on χ , the simulation box is built with 2000 water molecules and 500 CO₂ molecules. The pressure and temperature conditions to calculate the solubility of CO₂ are 40 MPa and 286 K. After a 500 ps NVT temperature equilibration, NPT simulations with the duration of 200 ns were performed while the χ value varied from 1 to 1.18 with an increment of 0.02. Then, the mole fraction of carbon dioxide in the water slab away from the interface was calculated after the system reached its equilibrium state (after 50 ns). In order to reduce the statistical uncertainty, totally 100 measurements of carbon dioxide solubility in water for each χ value were used for solubility estimation. The atomic charges, Lennard–Jones parameters, and interatomic distances of the models used in this study are summarized in Table 1.

Molecular simulations are performed using the Gromacs package version 5.1.5.^{44,45} Newton's equation of motion is integrated using the leapfrog algorithm with a 2 fs time step. Periodic boundary conditions are introduced in all three dimensions. The long-range electrostatic interactions⁴⁶ are handled with the particle-mesh Ewald summation method and the cutoff distance is adjusted to 1.0 nm. The van der Waals

interactions are calculated using the Lennard–Jones potential with a cutoff radius of 1.0 nm. Velocity-rescale thermostat algorithm is adopted for temperature control, and the time constant is 0.5 ps. The pressure of the simulation systems is controlled by semi-isotropic Parrinello–Rahman barostat (for hydrate formation) and isotropic Parrinello–Rahman barostat (for solubility calculation), respectively, with a 1 ps relaxation time.

For hydrate formation simulations, energy minimization is first performed with the steepest descent algorithm to relax the liquid water at the hydrate surface. Then, a 500 ps simulation run under the NVT ensemble is carried out for temperature equilibrium and followed by the NPT production simulation. Totally 10 equilibrate pressure conditions are selected, namely, 3, 6, 10, 20, 40, 100, 200, 300, 400, and 500 MPa. The equilibrium temperature varies from 240 to 295 K for each independent simulation run. Each simulation will continue until complete formation or dissociation of CO₂ hydrate is achieved.

3. RESULTS AND DISCUSSION

3.1. Equilibrium Conditions with Classical LB Combination Rule. The combination of the OPC water with EPM2 or TraPPE carbon dioxide was examined first using classical LB combining rules. Meanwhile, the use of TIP4P/Ice and EPM2 was also included for comparison. The T₃ at 10 different pressure conditions were calculated. As shown in Figure 2, two

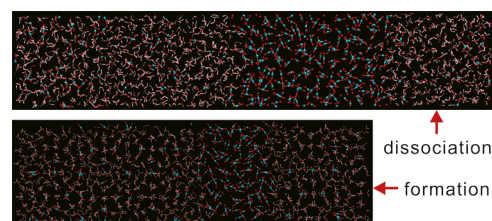


Figure 2. Snapshots of the simulation box at 40 MPa, using the TIP4P/Ice model for water, the EPM2 model for CO₂, and the modified LB combining rules (T₃ = 283.5 K), and the two final states at 282 and 285 K correspond to complete hydrate dissociation and formation, respectively. The water molecules are represented by the red and white lines, while the carbon dioxide molecules are represented by the cyan (carbon) and red (oxygen) spheres.

Table 1. Potential Parameters of TIP4P/Ice (Water), OPC (Water), EPM2 (CO₂), and TraPPE (CO₂) Models^a

force field	atom	σ_{LJ} (Å)	ϵ_{LJ} (kJ/mol)	q (e)	geometry
TIP4P/Ice	O	3.1668	0.88217	0	$d_{\text{OH}} = 0.9572 \text{ \AA}$
	H	0	0	0.5897	$d_{\text{OM}} = 0.1577 \text{ \AA}$
	M	0	0	-1.1794	H–O–H = 104.52°
OPC	O	3.16655	0.89036	0	$d_{\text{OH}} = 0.8724 \text{ \AA}$
	H	0	0	0.6791	$d_{\text{OM}} = 0.1594 \text{ \AA}$
	M	0	0	-1.3582	H–O–H = 103.6°
EMP2	C	2.757	0.23388	0.6512	$d_{\text{OC}} = 1.149 \text{ \AA}$
	O	3.033	0.66937	-0.3256	O–C–O = 180°
TraPPE	C	2.80	0.22459	0.70	$d_{\text{OC}} = 1.16 \text{ \AA}$
	O	3.05	0.65712	-0.35	O–C–O = 180°

^aThe geometry, partial charges (q), and the L–J parameters σ and ϵ are given for all of the atoms and the negative charge site of the water model (M).

indicative snapshots of the simulations performed at 40 MPa represented the final states of the system: complete CO₂ hydrate dissociation and formation occurred at 285 and 282 K, respectively. The representative evolution of the potential energy of the simulation system at different temperatures is presented in Figure 3. The decrease in the total potential

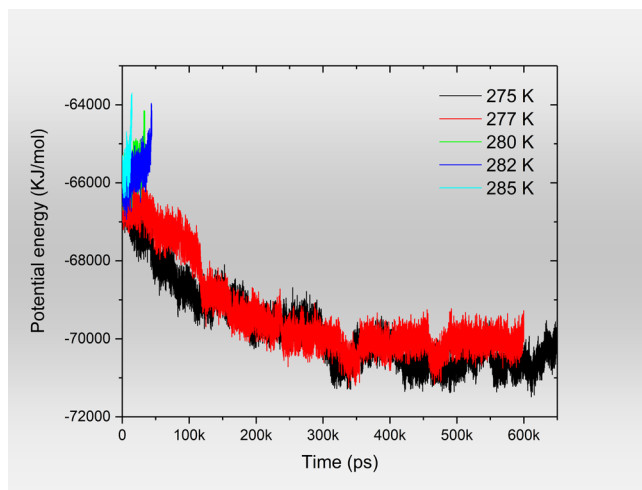


Figure 3. Evolution of the potential energy as a function of time for the NPT runs for the TIP4P/Ice and EPM2 models at 100 MPa.

energy indicates the growth of CO₂ hydrate, while the increase indicates hydrate dissociation. Since we only conducted one single simulation at each temperature–pressure condition, the value of T_3 is assigned as the arithmetic average of the lowest hydrate dissociation temperature and the highest hydrate formation temperature. Take the temperature scan shown in Figure 3 as an example. The T_3 at 100 MPa would be estimated as 278.5 K.

The calculated T_3 values for the three sets of simulations are given in Table 2 and are presented in Figure 4. For the pressures examined, the use of the OPC water model with the classical LB rules generates large deviations of T_3 , which are approximately 25–50 K for both EPM2 and TraPPE CO₂ models. Our results indicate that the usage of different CO₂ models makes only a small difference in T_3 prediction, which agrees well with the conclusion made by Míguez et al.³⁷ Comparably, simulations using TIP4P/Ice water with LB rules

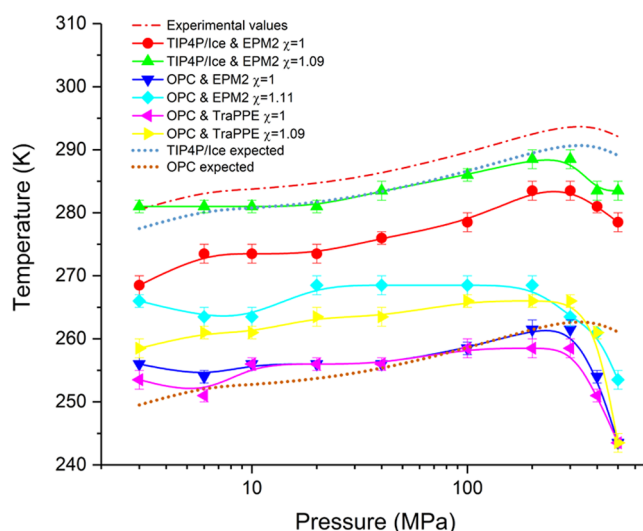


Figure 4. Experimental (red dash-dot line) and predicted values of T_3 for carbon dioxide hydrates from this work. The blue and brown dashed lines represent the expected values of T_3 for the TIP4P/Ice (experimental values – 3 K) and the OPC (experimental values – 31 K) force fields, respectively. Experimental values were obtained from ref 2.

result in a much smaller T_3 deviation of only 10–14 K. Both TIP4P/Ice and the OPC models cannot reproduce experimental T_3 of CO₂ hydrate combined with a CO₂ model using ideal Lorentz–Berthelot cross-combining parameters. However, the equilibrium curve of each simulation set exhibits re-entrant behavior at high-pressure conditions, and the shape of equilibrium lines are similar with that of the experimental line, especially for TIP4P/Ice model.

3.2. Solubility-Based χ -value Calculation. The above results indicated that both the OPC and the TIP4P/Ice water models cannot reproduce the experimental T_3 of the CO₂ hydrate using ideal LB combination parameters. Previous studies on the nucleation of CO₂ hydrate have revealed the importance of the modification of LB parameters in MD simulation to obtain proper CO₂ concentration and to trigger nucleation.⁴⁷ Thus, a proper modification is needed for T_3 prediction of the CO₂ hydrate using MD simulation. Míguez et al.³⁷ and Costandy et al.³⁹ used different approaches to modify the interaction between molecules to correct the deviation between the predicted and experimental T_3 values of gas

Table 2. Experimental and Calculated T_3 of the Six Sets of Simulations Using Ideal ($\chi = 1.00$) and Modified LB Combining Rules

pressure (MPa)	T_3 (K)						
	experimental	TIP4P/Ice and EPM2		OPC and EPM2		OPC and TraPPE	
		$\chi = 1$	$\chi = 1.09$	$\chi = 1$	$\chi = 1.11$	$\chi = 1$	$\chi = 1.09$
3	280.5	268.5(1.5)	281(1)	256(1)	266(1)	253.5(1.5)	258.5(1.5)
6	283.4	273.5(1.5)	281(1)	254(1)	263.5(1.5)	251(1)	261(1)
10	283.7	273.5(1.5)	281(1)	256(1)	263.5(1.5)	256(1)	261(1)
20	284.6	273.5(1.5)	281(1)	256(1)	268.5(1.5)	256(1)	263.5(1.5)
40	286.2	276(1)	283.5(1.5)	256(1)	268.5(1.5)	256(1)	263.5(1.5)
100	289.5	278.5(1.5)	286(1)	258.5(1)	268.5(1.5)	258.5(1.5)	266(1)
200	292.6	283.5(1.5)	288.5(1.5)	261.5(1.5)	268.5(1.5)	258.5(1.5)	266(1)
300	293.8	283.5(1.5)	288.5(1.5)	261.5(1.5)	263.5(1.5)	258.5(1.5)	266(1)
400	293.6	281(1)	283.5(1.5)	254(1)	261(1)	251(1)	261(1)
500	292.1	278.5(1.5)	283.5(1.5)	243.5(1.5)	253.5(1.5)	243.5(1.5)	243.5(1.5)

hydrate. Míguez et al.³⁷ performed simulations at the same temperature and pressure conditions and tested different cross-interaction energy parameter values, until a good consistency of predicted T_3 and experimental T_3 was achieved (the corresponding χ value is 1.13). Differently, Costandy et al.³⁹ used the solubility of CO_2 in water to modify the cross-interaction parameter. They tested the deviation between the predicted solubility and the values reported by Duan et al.⁴⁸ at the same P–T conditions and found that the deviation could be minimized when the modification parameter χ value was 1.08. Using this modified LB parameter, Costandy et al.³⁹ confirmed that the deviation between predicted and experimental T_3 of CO_2 hydrate was equal to the bias of the water model in predicting the melting temperature of ice.

A similar methodology as described in Costandy et al.³⁹ was applied in our work to modify the cross-parameter between water and carbon dioxide. The results of the modification simulation are presented in Figure 5. It is shown that the

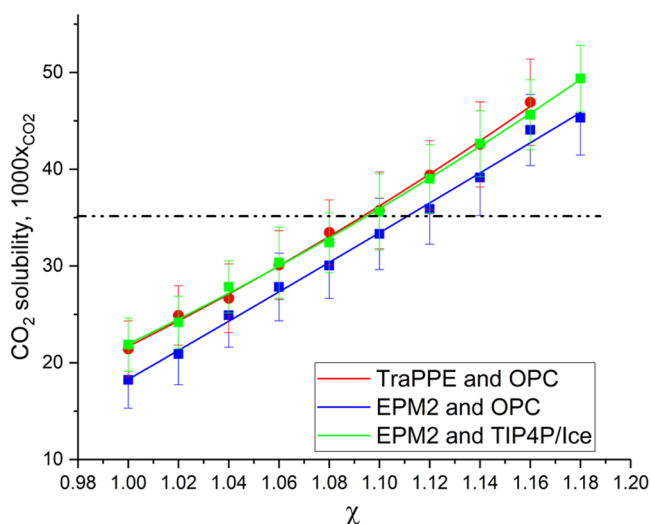


Figure 5. Dependence of the solubility of CO_2 in water on the modification factor of the LB cross-interaction energy parameter for the three sets of simulations with different molecule models. The dashed line shows the solubility of CO_2 obtained by Duan's equation⁴⁸ at 40 MPa and 286 K.

solubility of CO_2 in water can be predicted accurately with χ value of 1.09 for the combination of TIP4P/Ice and EPM2 models; while for the OPC model, the χ values are 1.11 and 1.09 for EPM2 and TraPPE CO_2 models, respectively. As described in Section 2, these χ values were all obtained at 40 MPa and 286 K. In order to validate the suitability of the χ values in the wide range of P–T conditions performed in our study, the CO_2 solubility around hydrate equilibrium points from 3 to 200 MPa was re-examined, and the results are shown in Figure 6. At a pressure range from 6 to 100 MPa, the calculated solubilities by three sets of simulation are in good agreement with the values obtained using Duan's equation⁴⁶ along the three-phase equilibrium line. However, the CO_2 solubility is overestimated at lower pressure (3 MPa) and underestimated at higher pressure (200 MPa) when compared with the experimental values. Note that the predicted CO_2 solubility data above 200 MPa were not provided. One reason is that Duan's equation does not support solubility calculations above 200 MPa for comparison. Another reason is that we observed the generation and collapse of a CO_2 gas bubble in

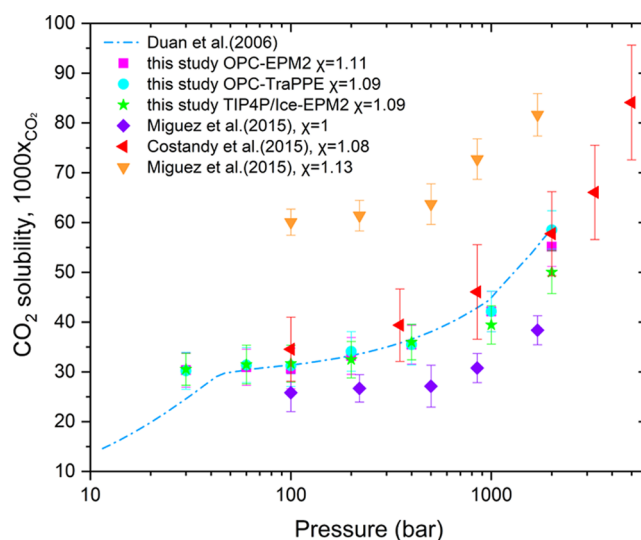


Figure 6. CO_2 Solubility in water as a function of pressure under hydrate equilibrium conditions. The dashed cyan line is the data calculated based on the work of Duan et al.⁴⁸ Purple diamonds and brown down-triangles are results from the work of Miguez et al.³⁷ with $\chi = 1$ and $\chi = 1.13$, respectively. Red triangles are results from the work of Costandy et al.³⁹ with $\chi = 1.08$ using TIP4P/Ice–TraPPE combination.

our simulation box when the pressure exceeds 200 MPa, so the estimation of the CO_2 solubility becomes inconvincible.

3.3. Equilibrium Conditions with Corrected Cross-Parameters. Using the above modified interaction parameters, all of the simulation runs were performed one more time to calculate the T_3 values of carbon dioxide hydrate. The results are shown in Table 2 and are also presented in Figure 4. It is quite clear that the correction of the CO_2 solubility significantly affected the predicted T_3 values. The deviations between the predicted T_3 and the experimental T_3 range from 0.5 to -10 , -20 to -32 , and -22 to -49 K for TIP4P/Ice and EPM2, OPC and EPM2, and the corresponding OPC and TraPPE force fields, respectively.

For the TIP4P/Ice model, the deviation of T_3 from the experimental values was approximately -3 K in the pressure range (6 to 100 MPa) that the CO_2 solubility in water can be accurately modified. This value is equal to the deviation of the predicted melting point of ice (-3 K) as determined by TIP4P/Ice. The result strongly supports that the accuracy of the prediction of T_3 by TIP4P/Ice can be improved on the grounds of an accurate description of the water–guest interactions based on the gas solubility in water. For the OPC model, the deviation of the predicted T_3 from the experimental values is significantly smaller than the difference between the estimated and experimental values of Ice Ih melting temperature (~ 31 K at 0.1 MPa) except for the simulations performed at high pressure (400 and 500 MPa). As depicted in Figure 4, the predicted T_3 values with the OPC and ideal LB parameters appear to be closer to the expected values (T_3 , experimental -31 K) within a larger pressure range. In other words, the OPC model does not agree with previous findings³⁹ that the deviation of T_3 is dictated by the deviation of the predicted melting temperature of ice, which is consistent with our previous findings on methane hydrate.¹⁶ We speculated that the predicted hydrate T_3 value is probably correlated with not only the ice point of the specific water model as reported by Costandy et al.³⁹ but also the accuracy of

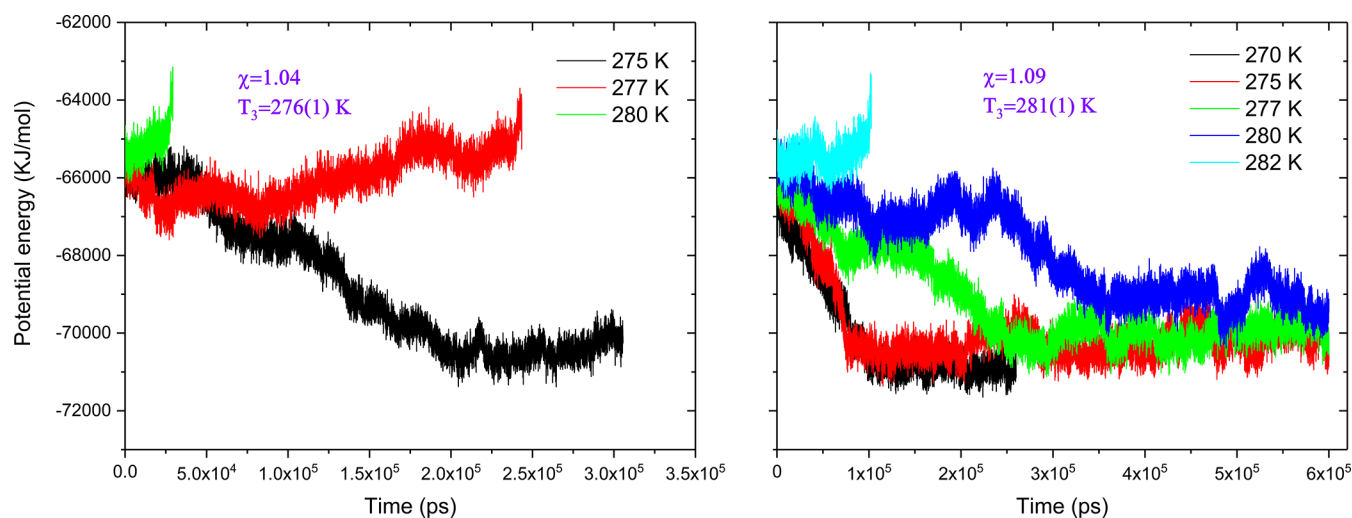


Figure 7. Evolution of the potential energy as a function of time for the NPT simulations at 3 MPa with two different χ values.

the reproduced bulk properties of water. The influence of the latter is comparable to the amendment of the H₂O–CO₂ interaction based on the CO₂ solubility.

It is worth noting that at lower (3 MPa) and higher (over 200 MPa) pressure conditions, the T_3 deviation from experimental values becomes larger. This trend is consistent with the predicted CO₂ solubility. To further verify the influence of CO₂ solubility, we recalculated the χ value for the combination of TIP4P/Ice and EMP2 models at 3 MPa. The result showed that the CO₂ solubility can be correctly predicted with an LB modification parameter of $\chi = 1.04$. Using this parameter, the equilibrium temperature of CO₂ hydrate at 3 MPa was calculated to be 276 ± 1 K (Figure 7), which is about 3 K lower than the experimental value. The results strongly confirmed the importance of CO₂ solubility in accurate T_3 prediction. Unfortunately, there are no convincing CO₂ solubility data at higher pressure (over 200 MPa). Therefore, the determination of χ values at higher pressure conditions becomes unavailable. However, we could also note the trend of underestimation of CO₂ solubility while the pressure goes higher, which might imply a larger χ (over 1.09) is needed for accurate T_3 prediction. Our simulations indicated that a fixed modification parameter χ is not enough for the accurate prediction of carbon dioxide solubility under all pressure conditions. Thus, a varied χ is needed to correctly describe the three-phase equilibrium line within the large pressure range. Our results are somewhat different from Costandy's findings which used only one fixed χ to reproduce the experimental T_3 in the pressure range from 20 to 500 MPa. One possible reason might be the accuracy of the TraPPE CO₂ model in solubility estimation at higher pressure. As shown in Figure 6, the CO₂ solubility predicted by the combination of TraPPE and TIP4P/Ice models (red triangles) seems to have smaller deviations compared with our calculations. Moreover, simulations conducted at a pressure lower than 20 MPa were not included in Costandy's work, under which conditions the CO₂ solubility might be overestimated. In light of this, we conducted the solubility and T_3 prediction at 3 MPa using exactly the same parameters as reported in Costandy et al.³⁹ The obtained CO₂ solubility was basically consistent with those values we previously obtained using different models under the same condition. Moreover, the predicted T_3 value at this specific pressure was ~ 279 K. The deviation between the

predicted T_3 and experimental value was slightly smaller than the reported bias (-2.2 to -3.9 K). The above simulations implied that the usage of an χ value smaller than 1.08 might generate a more precise prediction of both the CO₂ solubility and T_3 values while using TIP4P/Ice and TraPPE models.

4. CONCLUSIONS

In this work, the direct coexistence method was used to calculate the three-phase coexistence line of the CO₂ hydrate. Totally, six sets of simulations were conducted to calculate the coexistence temperatures T_3 for 10 pressure conditions ranging from 3 to 500 MPa. For simulations using ideal LB combination parameters, all three sets of simulations showed significantly underestimation of T_3 when compared with the experimental values. Meanwhile, the re-entrant behavior of the equilibrium curve can be observed. The three-phase equilibrium line obtained by the combination of TIP4P/Ice and EPM2 models almost strictly obeys the experimental trend with a temperature shift of about 10–14 K. This discrepancy is speculated to be aroused by the combined underestimation of CO₂ solubility in water and ice melting point. Based on the solubility of CO₂ in water at 40 MPa and 286 K, the cross-interaction parameters χ were calculated to be 1.09, 1.11, and 1.09 for the combination of TIP4P/Ice and EPM2, OPC and EMP2, and OPC and TraPPE models, respectively. Using the above modified LB parameters, the deviations between the corresponding T_3 values and experimental values all become smaller. For TIP4P/Ice and EPM2 models at a pressure range from 20 to 200 MPa, the deviation is almost equal to the difference in the predicted melting point of ice from the respective experimental value. On the contrary, the simulations using the OPC water model and ideal LB parameters agree well with the idea that the deviation is correlated with the difference in the prediction of the melting point of ice from the respective experimental value. Considering the relatively higher ability of the OPC model in reproducing bulk properties of water, we speculated that the H₂O–H₂O and H₂O–CO₂ cross-interactions' parameters might contribute equally to the prediction of T_3 value.

Additionally, the solubility of CO₂ in water was overestimated at a lower pressure and underestimated at a higher pressure, indicating that a fixed χ value may not be suitable for the T_3 prediction over a large pressure range. Based on our

simulations, we suggested a small χ value for the prediction of CO₂ hydrate T₃ at lower pressures (lower than 10 MPa), specifically 1.04 for TIP4P/Ice and EPM2 models at 3 MPa. At high-pressure conditions, for example, over 200 MPa, the estimation of both T₃ and CO₂ solubility implied that a relatively larger χ value (>1.09) is needed for accurate T₃ prediction. However, the lack of CO₂ solubility data prevents us from giving the exact χ value.

Finally, the large deviation in the T₃ prediction of the OPC model means that this model may not be suitable for the MD simulation of the CO₂ hydrate involving phase equilibrium conditions. The suggested models for CO₂ hydrate simulation are TIP4P/Ice and TraPPE with proper LB parameters, which can predict the CO₂ solubility in water accurately at corresponding P–T conditions.

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Notes

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