

External Barium Affects the Gating of KCNQ1 Potassium Channels and Produces a Pore Block via Two Discrete Sites

GILAD GIBOR, DANIEL YAKUBOVICH, ASHER PERETZ, and BERNARD ATTALI

Department of Physiology and Pharmacology, Sackler Medical School, Tel Aviv University, Tel Aviv 69978, Israel

ABSTRACT The pore properties and the reciprocal interactions between permeant ions and the gating of KCNQ channels are poorly understood. Here we used external barium to investigate the permeation characteristics of homomeric KCNQ1 channels. We assessed the Ba²⁺ binding kinetics and the concentration and voltage dependence of Ba²⁺ steady-state block. Our results indicate that extracellular Ba²⁺ exerts a series of complex effects, including a voltage-dependent pore blockade as well as unique gating alterations. External barium interacts with the permeation pathway of KCNQ1 at two discrete and nonsequential sites. (a) A slow deep Ba²⁺ site that occludes the channel pore and could be simulated by a model of voltage-dependent block. (b) A fast superficial Ba²⁺ site that barely contributes to channel block and mostly affects channel gating by shifting rightward the voltage dependence of activation, slowing activation, speeding up deactivation kinetics, and inhibiting channel inactivation. A model of voltage-dependent block cannot predict the complex impact of Ba²⁺ on channel gating in low external K⁺ solutions. Ba²⁺ binding to this superficial site likely modifies the gating transitions states of KCNQ1. Both sites appear to reside in the permeation pathway as high external K⁺ attenuates Ba²⁺ inhibition of channel conductance and abolishes its impact on channel gating. Our data suggest that despite the high degree of homology of the pore region among the various K⁺ channels, KCNQ1 channels display significant structural and functional uniqueness.

KEY WORDS: permeation • ion channels • outer vestibule • inactivation • channel gating

INTRODUCTION

The KCNQ family of voltage-gated K⁺ channels comprises five members (KCNQ1-5) that are expressed in a wide variety of tissues and play a major role in brain and cardiac excitability (Jentsch, 2000; Robbins, 2001). The channel complex formed by the KCNQ2/Q3 α subunits produces the M-current, a slowly activating, noninactivating K⁺ current that has profound effects on brain excitability and acts as a brake for repetitive neuronal firing (Wang et al., 1998; Brown and Yu, 2000; Jentsch, 2000; Rogawski, 2000). The KCNQ1 α subunits coassemble with various KCNE proteins, an emerging family of auxiliary β subunits, to produce functionally distinct K⁺ currents (Abbott and Goldstein, 1998; Melman et al., 2002). The coassembly of KCNQ1 and KCNE1 produces the I_{KS} potassium current that is critical for the late repolarization of the cardiac action potential (Barhanin et al., 1996; Sanguinetti et al., 1996). Consistent with their physiological importance, mutations of human KCNQ genes lead to major cardiovascular and neurological disorders such as the cardiac long

QT syndrome or neonatal epilepsy (Lehmann-horn and Jurkat-Rott, 1999; Ashcroft, 2000; Brown and Yu, 2000; Jentsch, 2000; Lerche et al., 2001; Robbins, 2001).

Like all voltage-gated K⁺ channel α subunits, KCNQ proteins share a common core structure of six transmembrane segments and a P-loop forming the pore region of the channel (Yellen, 2002). They are built from two parts: (a) a voltage-sensing domain, each comprised of S1-S4 that is responsible for controlling gating behavior and for energy transduction; and (b) a pore domain (S5-P-S6) whose structure is likely to be similar to the crystal structures of the bacterial KcsA and MthK K⁺ channels, with the inner helix S6 lining much of the pore (Doyle et al., 1998; Jiang et al., 2002). Coexpression of KCNQ1 with KCNE1 produces voltage- and time-dependent K⁺ currents with very slow activation kinetics and with a positive shift in the voltage dependence of activation, in comparison to homomeric KCNQ1 channels (Barhanin et al., 1996; Sanguinetti et al., 1996). There is also an increase in unitary channel conductance, leading to increased macroscopic current amplitude (Pusch, 1998; Sesti and Goldstein, 1998; Yang and Sigworth, 1998). When associated with KCNE3, KCNQ1 channels yield K⁺ currents that are nearly instantaneous and voltage independent (Schroeder et al., 2000).

Homomeric KCNQ1 channels elicit a delayed-rectifier K⁺ current that undergoes inactivation (Pusch et

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Address correspondence to Bernard Attali, Department of Physiology and Pharmacology, Sackler Medical School, Tel Aviv University, Tel Aviv 69978, Israel. Fax: (972) 3640 9113; email: battali@post.tau.ac.il

al., 1998; Tristani-Firouzi and Sanguinetti, 1998). Inactivation of homomeric KCNQ1 is invisible from macroscopic inspection and could be revealed by a hook of the tail currents. The onset of KCNQ1 inactivation occurs with an apparent delay and voltage dependence and is markedly prevented by the presence of KCNE1 (Pusch et al., 1998; Tristani-Firouzi and Sanguinetti, 1998). Though the mechanism of KCNQ1 inactivation is unknown, its biophysical characteristics are different from classical N and C-type inactivation of *Shaker*-like potassium channels, where K⁺ ions modulate C-type inactivation at an external pore site and N-type inactivation involves a direct pore occlusion at an internal site (Yellen, 2002).

Barium ions have been widely used to probe the permeation mechanisms of a variety of K⁺ channels. The similar crystal radius of Ba²⁺ compared with K⁺ and its stronger charge allows it to fit tightly into the deep selectivity filter (Jiang and MacKinnon, 2000). When applied from the inside, Ba²⁺ blocks delayed-rectifiers at a site that is accessible only when the channel is open (Armstrong and Taylor, 1980; Eaton and Brodwick, 1980). In contrast, external Ba²⁺ block can occur both in the open and closed channel states (Hurst et al., 1995; Harris et al., 1998).

The pore properties and the reciprocal interactions between permeant ions and the gating of the KCNQ channel family are poorly understood. Here we used external barium as a probe to address this issue in homomeric KCNQ1 channels. We determined the Ba²⁺ binding kinetics and the concentration and voltage dependence of Ba²⁺ steady-state block. Our results indicate that external barium interacts with the permeation apparatus of KCNQ1 at two main sites: (a) A deep Ba²⁺ site that produces an occlusion of channel conductance and could be simulated by a model of voltage-dependent block; and (b) a superficial Ba²⁺ site that has a minor impact on channel block and mainly affects KCNQ1 channel gating by shifting rightward the voltage dependence of activation, slowing activation, speeding up deactivation kinetics, and inhibiting channel inactivation. While high external K⁺ (50 mM) attenuates the inhibitory effect of Ba²⁺ on channel conductance, it abolishes its impact on channel gating. The strong impact of Ba²⁺ on KCNQ1 channel gating could not be described by a simple model of voltage-dependent block.

MATERIALS AND METHODS

Channel Expression into *Xenopus* Oocytes

Female *Xenopus laevis* frogs were purchased from *Xenopus* 1 (Dexter). The procedures followed for surgery and maintenance of frogs were approved by the animal research ethics committee of Tel Aviv University and in accordance with the Guide for the Care and Use of Laboratory Animals (National Academy of Sci-

ences, Washington DC). Frogs were anesthetized with 0.15% tricaine (Sigma-Aldrich). Pieces of the ovary were surgically removed and digested with 1 mg/ml collagenase (type IA, Sigma-Aldrich) in Ca²⁺-free ND96 for about 1 h, to remove follicular cells. Stage V and VI oocytes were selected for cRNA or DNA injection and maintained at 18°C in ND96 (in mM: 96 NaCl, 2 KCl, 1.8 mM CaCl₂, 1 MgCl₂, and 5 HEPES titrated to pH = 7.5 with NaOH), supplemented with 1mM pyruvate and 50 µg/ml gentamycin. The human KCNQ1 cDNA (in pGEM vector) was linearized by NotI. Capped complementary RNA was transcribed by the T7 RNA polymerase, using the mMessage mMachine transcription kit (Ambion Corp). The cRNA size and integrity was confirmed by formaldehyde-agarose gel electrophoresis. Homomeric expression of human KCNQ1 was performed by injecting 40 nl per oocyte (5 ng cRNA) using a Nanoject injector (Drummond). Several expression experiments were also performed by microinjecting a recombinant DNA vector (pcDNA3) encoding the human KCNQ1 cDNA directly into *Xenopus* oocyte nuclei (1 ng into 10 nl). Very similar data were obtained for either cRNA or DNA injections.

Electrophysiology

Standard two-electrode voltage-clamp measurements were performed at room temperature (22–24°C) 2–5 d after cRNA or DNA microinjection. Oocytes were placed into a 100-µl recording chamber and superfused with a modified ND96 solution (containing 0.1 mM CaCl₂) under constant perfusion using a fast perfusion system at a rate of 0.48 ml/s (ALA VM8, ALA Scientific Instruments). Whole-cell currents were recorded using a GeneClamp 500 amplifier (Axon Instruments, Inc.). Stimulation of the preparation, and data acquisition were performed using the pCLAMP 6.02 software (Axon Instruments, Inc.) and a 586 personal computer interfaced with a Digidata 1200 interface (Axon Instruments, Inc.). Glass microelectrodes (A-M systems, Inc.) were filled with 3M KCl and had tip resistances of 0.2–1 MΩ. Current signals were digitized at 1 kHz and low pass filtered at 0.2 kHz. The holding potential was –80 mV. Leak subtraction was performed off-line, using steps from –120 to –90 mV, assuming that the channels are closed at –80 mV and below. Errors introduced by the series resistance of the oocytes were not corrected and were minimized by keeping expression of the currents below 10 µA. All BaCl₂ solutions were prepared in modified ND96 (containing 0.1mM CaCl₂) and were iso-osmotically changed for NaCl. Modified ND96 solutions containing high K⁺ concentrations (50 mM) were also iso-osmotically changed for NaCl.

Data Analyses

Data analysis was performed using the Clampfit program (pCLAMP 8; Axon Instruments, Inc.), Microsoft Excel 2002 (Microsoft Corporation), SigmaPlot 8.0 (SPSS, Inc.), and Prism (GraphPad software). To analyze the voltage dependence of channel activation, a double exponential fit was applied to the tail currents at –60 or –120 mV and the slow exponential component was extrapolated to the beginning of the repolarizing step. Chord conductance (G) was calculated by using the following equation:

$$G = I / (V - V_{rev}),$$

where I corresponds to the extrapolated tail current and V_{rev} to the reversal potential measured in each experiment and which was –98 ± 2 mV (n = 10). G was estimated at the tail voltage V and then, normalized to a maximal conductance value, G_{max}. Activation curves were fitted to a Boltzmann distribution:

$$G/G_{\max} = 1/\{1 + \exp[(V_{50} - V)/s]\},$$

where V_{50} is the voltage at which the current is half-activated and s is the slope factor. For a measure of inactivation kinetics, we used a three-pulse protocol (see Fig. 6 C) where the membrane potential was stepped to a +30-mV conditioning prepulse for increasing durations (200-ms increments) to progressively activate and inactivate the channel; then, a brief (15 ms) hyperpolarizing interpulse to -130 mV was used to allow channel recovery from inactivation before a +30 mV test pulse was applied to reopen and reinactivate KCNQ1 channels. The decaying current of the third test pulse (reinduction of inactivation) could be fitted by a single exponential function. For a quantitative measure of fractional inactivation that takes into account the electrochemical driving force, G was deduced from tail currents at either -60 or -120 mV, where tail relaxations reveal a recovery from inactivation (hook) faster than deactivation. Currents recorded at the voltage V were first converted to conductance G . The time course of the tail conductance was then fitted to a triple exponential function of the form:

$$G(t) = A_{s1} \exp(-t/\tau_{s1}) + A_{s2} \exp(-t/\tau_{s2}) - A_f \exp(-t/\tau_f) + A_0,$$

where a fast time constant τ_f reflects the recovery from inactivation at the tail potential and two slow time constants τ_{s1} and τ_{s2} describe the deactivation process and A_0 is a steady-state current. At the tail potential, the amplitude A_f is directly related to the degree of inactivation and the ratio $A_f/(A_{s1} + A_{s2})$ is a good indicator of the probability of the channels being in the inactivated state.

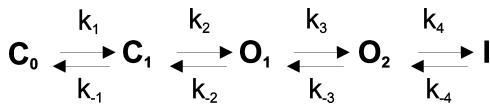
The voltage dependence of Ba^{2+} block was calculated according to the Woodhull's model (Woodhull, 1973) and the voltage dependence of the dissociation constant is given by:

$$K_D(V) = K_D(0) \exp(-z\delta VF/RT),$$

where $K_D(0)$ is the dissociation constant at zero voltage, z is the ion valency, F is the Faraday constant, δ is the fractional potential drop at the binding site (fractional electrical distance), V is the membrane potential (in volts), R is the gas constant, and T is the absolute temperature. All data were expressed as mean \pm SEM. Statistically significant differences between paired groups were assessed by Student's t test.

Model

Analysis in the absence of barium. The time course of channel activation could be described as a Markov process (Pusch et al., 1998) and was fitted to the following scheme:



SCHEME I

in which k_s are the rate constants in s^{-1} and C_0 to I are the corresponding state occupancies. Only O_1 and O_2 are assumed to conduct. Scheme I is sufficient for the description of macroscopic behavior of KCNQ1 channels (Pusch et al., 1998), assuming that the stochastic behavior of a single channel is reflected in the macroscopic behavior of a channel population (Blunck et al., 1998). We generated a set of differential equations based on Scheme I and fitted it to the observed time-course of channel activation in order to obtain the transition rate constants and rela-

tive occupancies of each state. For simplicity, we assumed that no openings occur when the cell is voltage-clamped to -80 mV (the holding potential) and consequently we fixed initial values of all states to zero, except C_0 .

According to Scheme I, the open channel probability can be calculated from Eq.1:

$$P_o = \frac{O_1 + O_2}{C_0 + C_1 + O_1 + O_2 + I}, \quad (1)$$

and the relative state occupancy (ϕ) is defined as:

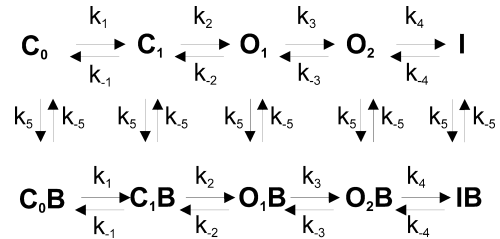
$$\phi = \frac{S_i}{\sum_i S_i}. \quad (2)$$

We fitted the activation time course from the holding potential (-80 mV) to the +30-mV voltage step. We fixed the k_4 and k_{-4} values (which are assumed to be voltage-independent) and also the initial C_0 to those obtained from the previous procedure, and then fitted the activation from the holding potential (-80 mV) to the -60-mV voltage step. This analysis was performed for KCNQ1 expressing oocytes that were superfused with low (2 mM) or high (50 mM) extracellular K^+ containing ND96 solutions in the absence of barium.

To validate the model, we used the fit data and simulated the tail envelope protocol. For the simulation, we used the averaged values of the rate constants and the relative state occupancies obtained from the fit of the activation time course during the +30-mV voltage step and allowed the simulation to run for pre-defined time intervals. Thereafter, we used the simulated values and changed the rate constants to those corresponding to the -60-mV voltage step. Based on the simulation of the different prepulse intervals, we generated a curve describing the dependence of the "hook" amplitude on the prepulse time and compared it to experimental data.

Analysis of Ba^{2+} Block in High External K^+ Solutions

For simulating Ba^{2+} block in high external K^+ solutions, we used a simplistic scheme of voltage-dependent block as illustrated in Scheme II.



SCHEME II

We assumed that the conductance block originates from the existence of one Ba^{2+} -occluding site. Thus, the channel can populate two different kinds of states, assumed to be in equilibrium: a Ba^{2+} -blocked state (B) and a Ba^{2+} -unblocked state. For simplicity, we assumed that all states (C_0 to I) bind Ba^{2+} with equal affinity. Indeed, the affinity of the closed states cannot be estimated from our records. However, the on-kinetics of Ba^{2+} are clearly dominated by a slow component that allows us to estimate Ba^{2+} affinity for the open states. To preserve microscopic reversibility, we also assumed that transitions through the blocked states occur with similar rates to their unblocked counterparts. Further analyzing the equilibrium, we defined f , the fractional inhibition as:

$$f = \frac{I_{hk} - I_{hk,Ba}}{I_{hk}}, \quad (3)$$

in which I_{hk} and $I_{hk,Ba}$ are the steady-state amplitudes of the recorded currents before and after Ba^{2+} application (in 50 mM external K^+). Consequently, according to the mass action law:

$$\frac{S_u \times Ba^{2+}}{S_b} = K_D = \frac{k_{off}}{k_{on}}, \quad (4)$$

in which S_u corresponds to unblocked state and S_b to the blocked state, while K_D is the dissociation constant of Ba^{2+} . Combining Eqs. 3 and 4 and rearranging gives:

$$K_D = Ba^{2+} \times (1 - f)/f, \quad (5)$$

thus enabling us to estimate the K_D value.

On the other hand, according to the properties of bimolecular reactions:

$$1/(k_{on} \times Ba^{2+} + k_{off}) = \tau_{on}, \quad (6)$$

in which τ_{on} is the blocking time constant obtained from experiments. Combining Eqs. 5 and 6 allows us to estimate the rates of block and unblock.

Time-course fitting and simulation were done with Berkeley Madonna software (Version 8.0.2 for Windows; Kagi Shareware).

Online Supplemental Material

A figure illustrating the lack of effects of 10 mM $MgCl_2$ on KCNQ1 currents, with representative current traces, current-voltage relations, normalized conductance curves, quantitative measures of inactivation, and deactivation.

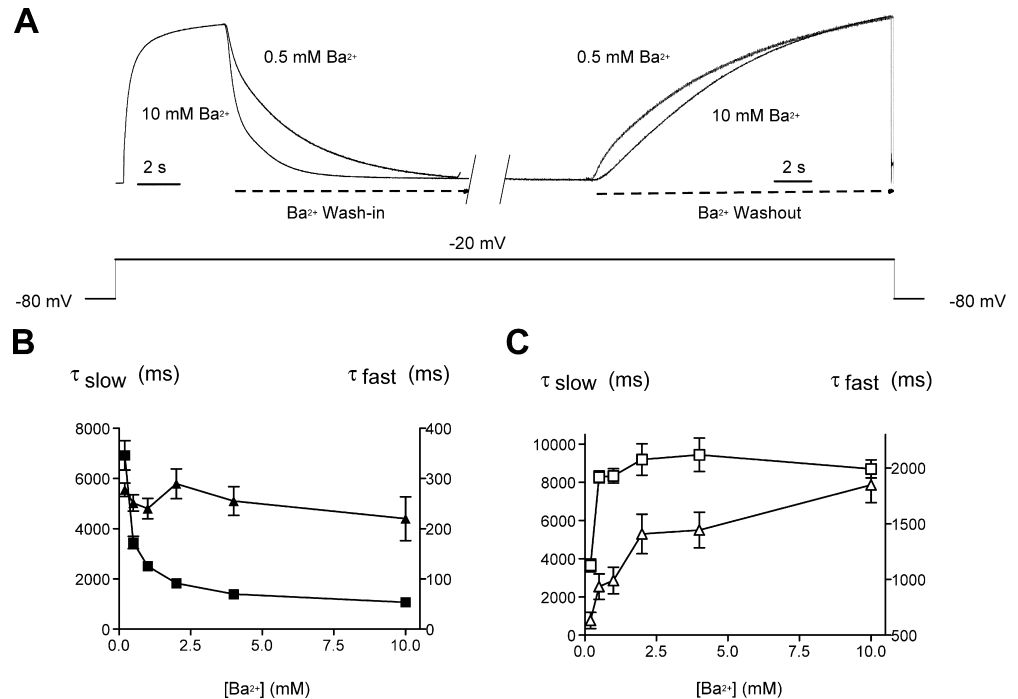
A table summarizing the rate constants of KCNQ1 activation in low and high external K^+ (without Ba^{2+}) obtained from the fit to Scheme I and the rate constants of KCNQ1 activation in high external K^+ in the presence of Ba^{2+} obtained from the fit to Scheme II. Online supplemental material is available at <http://www.jgp.org/cgi/content/full/jgp.200409068/DC1>.

RESULTS

Kinetics of External Barium Block

In low external K^+ (2 mM), extracellular Ba^{2+} produced a time-dependent inhibition of homomeric KCNQ1 K^+ currents when expressed in *Xenopus* oocytes. As shown in Fig. 1 A, Ba^{2+} potently inhibited the K^+ currents when applied to opened channels. From a holding potential of -80 mV, the KCNQ1 channels were opened by step depolarization and the oocyte was held constantly at this depolarizing potential until the end of the kinetic experiment. After 4 s of channel opening, Ba^{2+} was applied for 20 s and thereafter washed out for 16 s at the same depolarizing potential (e.g., at -20 mV in Fig. 1 A). It is clear from the wash-in experi-

FIGURE 1. Concentration dependence of Ba^{2+} wash-in and washout kinetics. The kinetic experiments were performed with various extracellular Ba^{2+} concentrations at -20 mV in low external K^+ solutions (2 mM). (A) Representative normalized traces are shown. From a holding potential of -80 mV, the KCNQ1 channels were opened by a -20-mV depolarizing step. After 4 s of channel opening, 0.5 mM or 10 mM Ba^{2+} was applied for 20 s and thereafter washed out for 16 s at the same depolarizing potential (-20 mV). (B) Ba^{2+} wash-in kinetics were fitted by a double exponential function with a fast (solid triangles) and slow (solid squares) time constants plotted as a function of Ba^{2+} concentrations ($n = 8$). (C) Ba^{2+} wash-out kinetics were fitted by a double exponential function with a fast (empty triangles) and slow (empty squares) time constants plotted as a function of Ba^{2+} concentrations ($n = 8$).



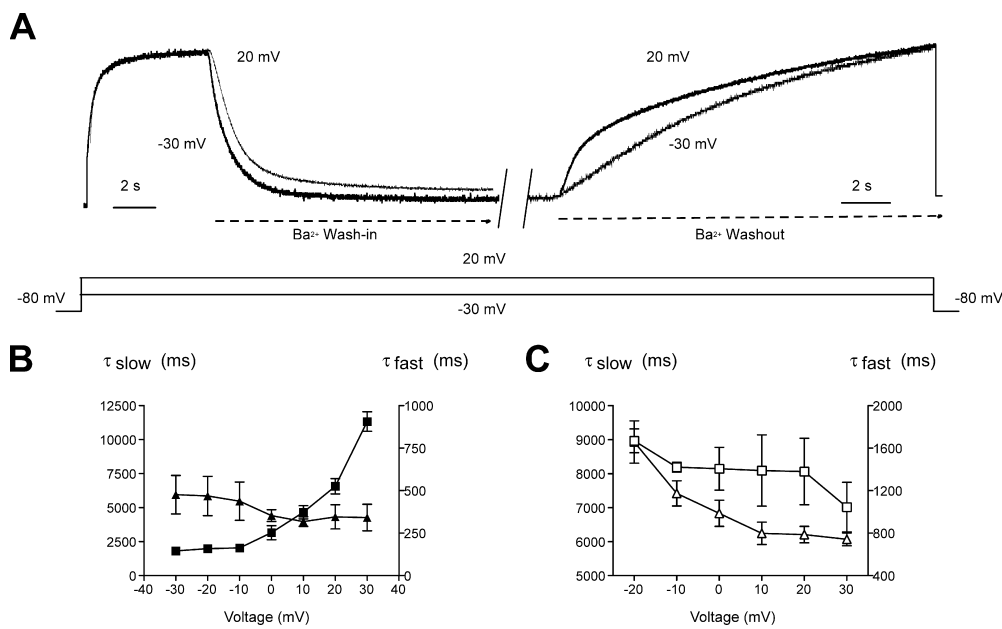


FIGURE 2. Voltage dependence of Ba²⁺ wash-in and washout kinetics. The kinetic experiments were performed with 10 mM Ba²⁺ at various depolarizing voltages in low external K⁺ solutions (2 mM). (A) Representative normalized traces are shown. From a holding potential of -80 mV, the KCNQ1 channels were opened at -30 or +20 mV. After 4 s of channel opening, 10 mM Ba²⁺ was applied for 20 s and thereafter washed out for 16 s at the same depolarizing potential. (B) Ba²⁺ wash-in kinetics were fitted by a double exponential function with a fast (solid triangles) and slow (solid squares) time constants plotted as a function of depolarizing voltages ($n = 5$). (C) Ba²⁺ washout kinetics were fitted by a double exponential function with a fast (empty triangles) and slow (empty squares) time constants plotted as a function of depolarizing voltages ($n = 5$).

ments that Ba²⁺ on-binding displayed a fast and slow kinetic component (Fig. 1, A–C). For example, 10 mM Ba²⁺ blocked at -20 mV the KCNQ1 K⁺ currents with fast and slow time constants $\tau_f = 220 \pm 43$ ms and $\tau_s = 1,073 \pm 128$ ms, respectively ($n = 8$) (Fig. 1 B). This suggested that Ba²⁺ binds to two sites. However, the fast kinetic component was too close to the rate of exchange of the bath to determine accurately its time course. This probably explains why the fast time constant did not vary significantly as a function of Ba²⁺ concentration (Fig. 1 B). The slow time constant increased markedly with depolarized potentials, illustrating the voltage dependence of Ba²⁺ binding to its slow site (Fig. 2 B). According to the Woodhull model, it was possible to calculate the electrical distance of the slow Ba²⁺ site and to show that it lies deep in the membrane electric field with $\delta = 0.32$ (Woodhull, 1973). Consistent with a voltage-dependent block at the slow Ba²⁺ site, the relative amplitude of the slow binding component decreased significantly with depolarized potentials from 0.45 at -30 mV to 0.21 at +20 mV (see also Fig. 3 B). The slow time constant τ_s (consequently, the slow blocking rate $1/\tau_s$) reached saturation at concentrations above 2 mM Ba²⁺ (Fig. 1 B). This was also accompanied by a saturation of the relative amplitude of the slow blocking component that drops from 0.81 at 0.5 mM Ba²⁺ to 0.66 at 2 mM Ba²⁺ and levels off at 0.55

at higher Ba²⁺ concentrations (unpublished data). This saturation suggests that barium binding does not follow a simple bimolecular reaction with just one Ba²⁺ binding site.

Barium also blocked KCNQ1 channels in the closed state. As already shown for *Shaker* channels (Hurst et al., 1995; Harris et al., 1998), external Ba²⁺ entered the pore of KCNQ1 at negative potentials where the channel is closed and remained bound to the deep slow site after washout, as recorded by a subsequent depolarizing step to open the gate and measure the KCNQ1 current block. While holding the membrane potential at -100 mV, 10 mM Ba²⁺ was extracellularly applied to the oocyte for 2 min (~ 70 times the slow time constant of Ba²⁺ wash-in at -30 mV), after which barium was washed out for 2 min (at -100 mV); then a single depolarizing step to +30 mV was applied to assay the level of current block (Fig. 4, A and B). Under these conditions, the KCNQ1 current was blocked by $51 \pm 7\%$ ($n = 6$) at the beginning of the depolarizing step (at 200 ms) and by $31 \pm 3\%$ ($n = 6$) at the end of the pulse (at 2,000 ms), thus providing strong evidence that barium could block KCNQ1 channels in the closed state. The slower rise in KCNQ1 current after Ba²⁺ exposure, as well as the progressive decrease in block during the depolarizing step, reflects the slow Ba²⁺ unbinding at +30 mV (Fig. 4, A and B). Though this protocol of Ba²⁺

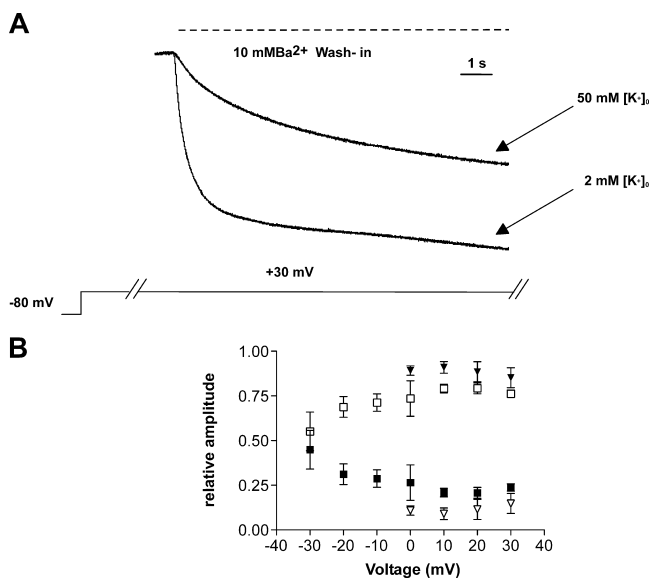


FIGURE 3. Comparison of Ba^{2+} wash-in kinetics in low and high external K^+ solutions. (A) Representative normalized traces of Ba^{2+} wash-in kinetics performed on the same oocyte in 2 mM $[\text{K}^+]_0$ and in 50 mM $[\text{K}^+]_0$. From a holding potential of -80 mV, the KCNQ1 channels were opened by a $+30$ -mV depolarizing step. After 4 s of channel opening, 10 mM Ba^{2+} was applied for 20 s and thereafter washed out for 16 s at the same depolarizing potential ($+30$ mV). (B) Relative amplitudes of Ba^{2+} wash-in fast (empty symbols) and slow (solid symbols) kinetic components in 2 mM $[\text{K}^+]_0$ (squares) and in 50 mM $[\text{K}^+]_0$ (triangles) plotted as a function of membrane voltage.

binding and washout at -100 mV kept Ba^{2+} bound to its deep slow site, it removed Ba^{2+} from its fast external site and thus prevented its impact on channel gating as we shall see later (see Fig. 4 and below). Holding constantly the membrane potential at $+30$ mV, we let Ba^{2+} dissociate from the open state by washing out the external solution. Results reveal that Ba^{2+} off-kinetics also comprised a fast and slow component (Figs. 1 and 2). Various concentrations of Ba^{2+} were applied to open channels at -20 mV until it reached steady-state inhibition after which Ba^{2+} dissociation was measured. Ba^{2+} wash-out exhibited significant dependence on blocker concentration with the fast and slow time constants varying, respectively, from $\tau_f = 0.63 \pm 0.07$ s and $\tau_s = 3.65 \pm 0.33$ s at 0.2 mM Ba^{2+} to $\tau_f = 1.85 \pm 0.16$ s and $\tau_s = 8.70 \pm 0.47$ s at 10 mM ($n = 8$) (Fig. 1, A and C). The wash-out rate does not only reflect the unbinding rate. Rather, it is the sum of binding and unbinding rate. Therefore, it should be expected to depend on Ba^{2+} concentration. Ba^{2+} dissociated faster with increasing depolarization. The fast time constant of Ba^{2+} unbinding displayed marked dependence on voltage, with $\tau_f = 1.66 \pm 0.20$ s at -20 mV and $\tau_f = 0.74 \pm 0.06$ s at $+30$ mV ($n = 5$) (Fig. 2, A and C).

Steady-state Barium Block

Various concentrations of extracellular Ba^{2+} were applied and current-voltage relations were constructed at steady-state, allowing normalized activation curves to be deduced from the tail currents. In 2 mM $[\text{K}^+]_0$, extracellular Ba^{2+} exerted on homomeric KCNQ1 channels a series of complex effects, including a voltage-dependent pore blockade as well as gating alterations (Fig. 5 A). Like with many K^+ channels, Ba^{2+} block of KCNQ1 currents was voltage-dependent, with weaker inhibition at depolarized potentials (Figs. 5 B and 11 B). For example, 2 mM Ba^{2+} produced a 90% and 51% block at -50 mV and $+30$ mV, respectively (Fig. 5 B). The slopes of Ba^{2+} dose-response curves were shallow at depolarized potentials, suggesting that Ba^{2+} bound to more than one site. Between -70 and -30 mV, the Ba^{2+} inhibitory curves could be fitted to one site with apparent K_D values of ~ 0.30 – 0.35 mM; however, the Hill coefficients were ranged from 2.8 to 2.1 (at -70 to -40 mV), suggesting the existence of at least two Ba^{2+} sites of very similar affinity at these negative potentials (Fig. 5 C). At more positive voltages from -20 to $+30$ mV, the curves became very shallow and needed to be fitted to two binding sites (Fig. 5 D). While the high-affinity site was not sensitive to membrane voltage changes, the low-affinity site was voltage dependent ($\delta = 0.45$) with increasing K_D values (from 0.67 to 6.4 mM) at increasing depolarized potentials (from -20 to $+30$ mV). Though being a rough approximation, application of the Woodhull's model for only one binding site did not fit properly the apparent K_D values over all potentials tested (Fig. 5 E). We had to split and fit the Ba^{2+} K_D values into two voltage regions. Between -70 and -20 mV, we got a shallow voltage dependence with a low fractional electrical distance $\delta_1 = 0.06$, whereas between -20 and $+30$ mV there was a steeper sensing of the membrane electric field with a higher electrical distance $\delta_2 = 0.33$ (Fig. 5 E). This latter value is in reasonable agreement with the δ value calculated above from the fit of two Ba^{2+} binding sites and with that of the slow Ba^{2+} site (0.32), derived from the on-kinetic experiments. As we shall also see below, these results suggest the existence of two Ba^{2+} sites, a fast superficial binding site that senses weakly the electric field and a slow Ba^{2+} binding site that lies deep in the membrane electric field.

In addition to its voltage-dependent block of the pore, Ba^{2+} also altered the gating of homomeric KCNQ1 channels, including slowing of the activation kinetics, rightward shift in the voltage dependence of activation, inhibition of channel inactivation and acceleration of deactivation kinetics. Ba^{2+} produced a $+34$ -mV gating shift in the voltage dependence of activation with $V_{50} = -26.0 \pm 0.6$ mV ($n = 7$) and $V_{50} = +8.0 \pm 1.3$ mV ($n = 7$) in the absence and presence of 10 mM

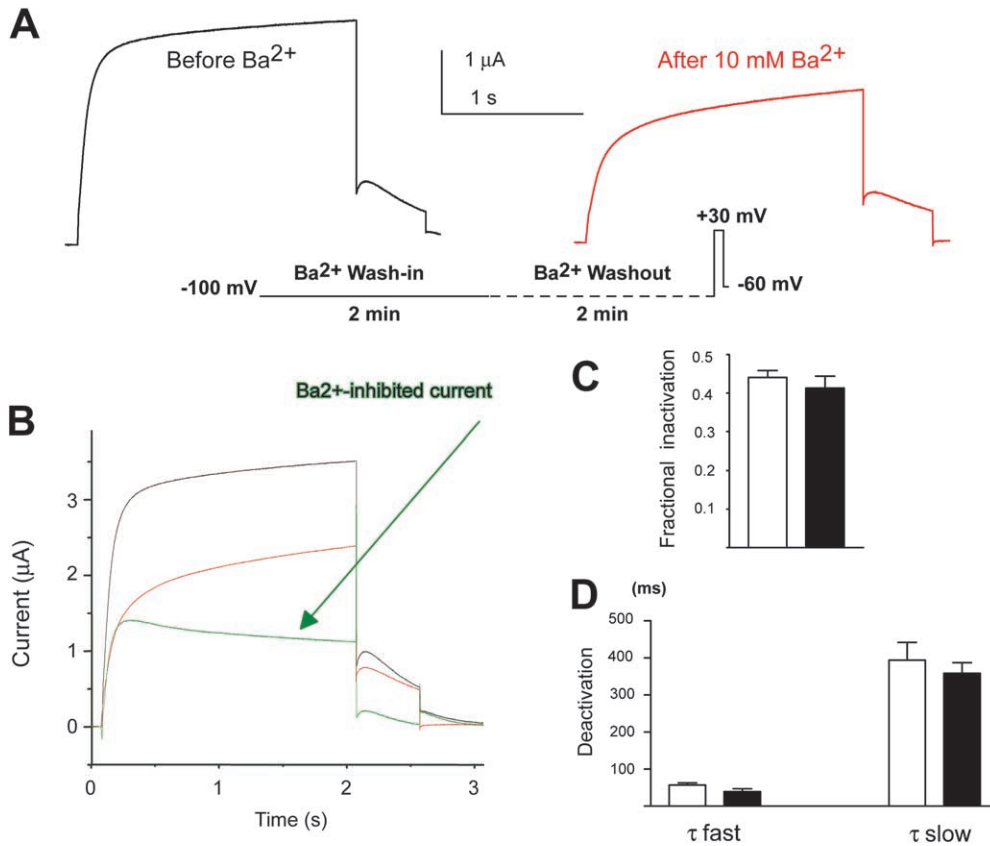


FIGURE 4. Effect of barium exposure to closed KCNQ1 channels. (A) Representative traces where from a -100 mV holding potential, KCNQ1 channels were opened for 2 s at $+30$ mV and repolarized for 0.5 s to a -60 -mV tail potential (left trace, before Ba²⁺); then, 10 mM Ba²⁺ was bound for 2 min to closed channels at -100 mV and subsequently washed out for 2 min at -100 mV, before channels were opened for 2 s at $+30$ mV as above (right trace, after 10 mM Ba²⁺). (B) Superimposed traces showing KCNQ1 current before Ba²⁺ exposure (black trace), after 10 mM Ba²⁺ (red trace), and the subtracted trace (green) corresponding to the Ba²⁺-inhibited current. (C) Fractional inactivation measured from the tail currents at -60 mV, before (empty bars) and after 10-mM Ba²⁺ application (solid bars) as described in A ($n = 8$). (D) Fast and slow deactivation time constants measured from the tail currents at -60 mV, before (empty bars) and after 10-mM Ba²⁺ application (solid bars) as described in A ($n = 8$).

Ba²⁺, respectively (Table I, Fig. 6 D). This rightward shift was dose-dependent with an EC₅₀ of 0.85 mM Ba²⁺ as deduced from the plot of the V₅₀ shift as a function of Ba²⁺ concentration (Table I). Barium produced a marked slowing of KCNQ1 activation kinetics that is well seen at all potentials (Figs. 5 A and 6). Activation kinetics were best fitted by a double exponential function, with a fast and slow component. Ba²⁺ increased by ~3.5-fold the fast time constant, from $\tau_f = 88 \pm 6$ ms to $\tau_f = 312 \pm 49$ ms at $+30$ mV in the absence and presence of 4 mM Ba²⁺, respectively ($n = 9$, $P < 0.01$; Fig. 6 A). The increase of the fast time constant displayed an EC₅₀ of 0.44 mM Ba²⁺ at $+30$ mV (Fig. 6 A) and the slowing of KCNQ1 activation was so pronounced that at 10 mM Ba²⁺, the fast kinetic component disappeared, leaving only one slow kinetic component (Fig. 6, B and C). The slowing of activation kinetics seems to arise mainly from a gating effect rather than from a slow Ba²⁺ dissociation at depolarized potentials. First, as already described in the protocol where 10 mM Ba²⁺ was bound and washed at -100 mV, a channel block was obtained (Fig. 4, A and B). However, only a mild slow-

ing of the activation kinetics was found, which contrasts to the marked slowing found for the same Ba²⁺ concentration in steady-state experiments (see Figs. 5 A, 6 C, and 4). Second, high external K⁺ (50 mM) prevented the slowing of KCNQ1 activation produced by 10 mM Ba²⁺, though Ba²⁺ washout kinetics were even slower in 50 mM [K⁺]₀ than in 2 mM [K⁺]₀ (Fig. 10 A, Table II and see below). Figs. 7 and 8 depict the effects of external Ba²⁺ on KCNQ1 deactivation kinetics that could be fitted by a double exponential function. Ba²⁺ dose-dependently speeded up deactivation, especially at tail potentials comprised from -30 to -70 mV, with an EC₅₀ of 1.4 mM Ba²⁺ as measured for τ_{fast} at -60 mV (Figs. 7 B and 8 A). The acceleration of deactivation was so marked at 10 mM Ba²⁺ (~10-fold) that a slow kinetic component could not be anymore measured at -60 mV (Fig. 7, A and B). As shown in Figs. 7 A and 8 A, the speeding effect of Ba²⁺ on deactivation was weaker at very negative tail potentials (compare tails at -60 and -140 mV). The Ba²⁺-induced acceleration of tail currents appears to involve a gating effect rather than a rapid block at repolarizing potentials (voltage-depen-

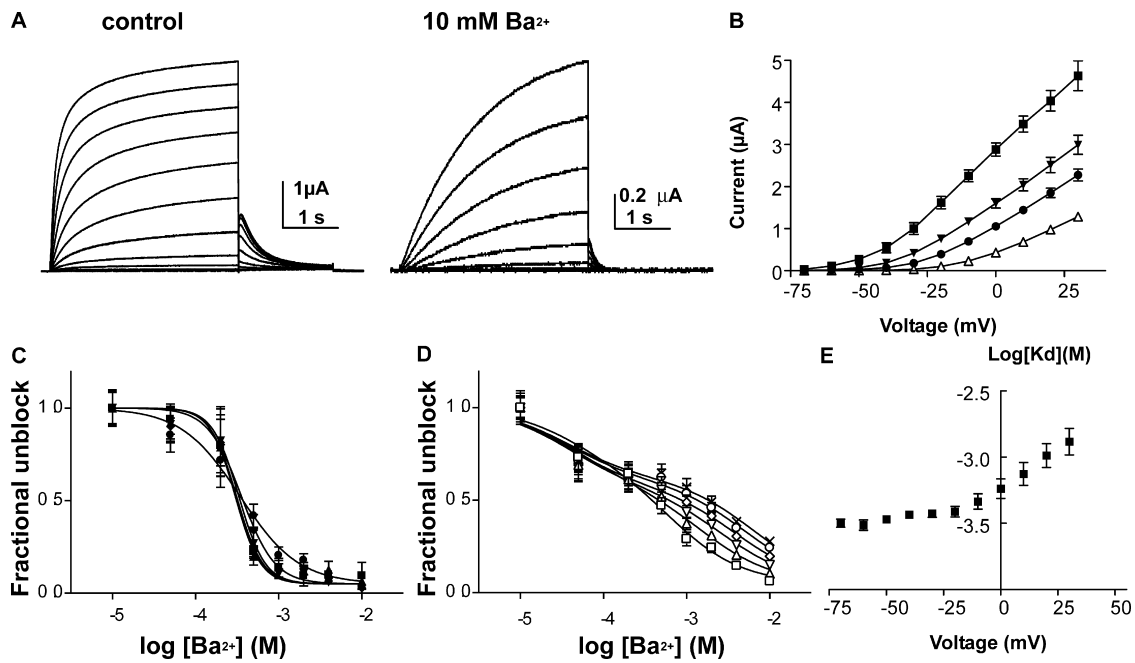


FIGURE 5. Barium produces a concentration- and voltage-dependent channel block. (A) Representative KCNQ1 current traces recorded in the absence (control) or presence of 10 mM Ba²⁺. From a holding potential of -80 mV, oocytes were stepped for 3 s from -70 to +30 mV in 10-mV increments and then repolarized for 1.5 s at -60-mV tail potential. (B) Current-voltage relations ($n = 7$) of control (solid squares), 0.5 mM Ba²⁺ (solid triangles), 2 mM Ba²⁺ (solid circles), and 10 mM Ba²⁺ (empty triangles). (C) Ba²⁺ inhibitory curves expressed as fractional channel unblock versus log₁₀ of Ba²⁺ concentrations and measured at -70 mV (closed squares), -60 mV (upward closed triangles), -50 mV (downward closed triangles), -40 mV (closed diamonds), and -30 mV (closed circles). The K_D values were obtained from the fit of a sigmoidal dose-response function ($n = 7$). (D) Ba²⁺ inhibitory curves expressed as in C and measured at -20 mV (open squares), -10 mV (upward open triangles), 0 mV (downward open triangles), +10 mV (open diamonds), +20 mV (open circles), and +30 mV (X labels). The K_D values were obtained from the fit of a two site competition function. (E) Semilogarithmic plot of the K_D values obtained from one site fit of the Ba²⁺ inhibitory curves versus membrane voltage ($n = 7$).

dent block) as speed-up of deactivation kinetics is less prominent at more negative potentials (-110–140 mV), where Ba²⁺ produces a very potent block (~70% and 90% inhibition at -60 and -140 mV, respectively). In addition, the impact of Ba²⁺ on deactivation could be prevented by high external K⁺ (50 mM) (see Fig. 9 A), a feature similar to the effect of Ba²⁺ on activation kinetics.

Homomeric KCNQ1 channels were previously shown to undergo an inactivation process that occurs with an apparent delay and voltage dependence (Pusch et al., 1998; Tristani-Firouzi and Sanguinetti, 1998). This inactivation is invisible macroscopically but could be revealed by hooked tail currents whose ascending phase reflects recovery from channel inactivation. Kinetic analysis suggested that KCNQ1 could populate two distinct open states and a voltage-independent inactivation state (Pusch et al., 1998; Tristani-Firouzi and Sanguinetti, 1998). Extracellular Ba²⁺ dose-dependently depressed KCNQ1 inactivation as revealed by the progressive suppression of the hook in tail currents (Fig. 7, A and C). The fractional inactivation was measured and quantified from the hooked tail currents (see MATERIALS

AND METHODS and Fig. 7 C). When measured upon a 3-s test potential at +30 mV and from a -60-mV tail potential, 10 mM Ba²⁺ reduced by ~61% the fractional inactivation of KCNQ1 from 0.54 ± 0.07 to 0.21 ± 0.03 ($n = 7$) with an $EC_{50} = 0.94 \pm 0.07$ mM (Fig. 7 C). Barium virtually suppressed inactivation when measured by a

TABLE I
Voltage Activation Parameters of KCNQ1 Channels in the Presence of Barium

| | Slope | V_{50} |
|-------------------------|----------------|----------------------|
| | mV | mV |
| Control | 16.6 ± 0.6 | -26.0 ± 0.6 (18) |
| 0.2 mM Ba ²⁺ | 22.1 ± 2.5 | -22.6 ± 1.8 (6) |
| 0.5 mM Ba ²⁺ | 18.5 ± 1.5 | -12.8 ± 1.3 (8) |
| 1 mM Ba ²⁺ | 15.4 ± 1.3 | -5.0 ± 1.4 (8) |
| 2 mM Ba ²⁺ | 16.3 ± 1.2 | -3.2 ± 1.3 (7) |
| 4 mM Ba ²⁺ | 15.2 ± 1.3 | 1.5 ± 1.7 (6) |
| 10 mM Ba ²⁺ | 13.9 ± 0.8 | 8.0 ± 1.3 (7) |

KCNQ1 channels were exposed to the indicated external Ba²⁺ concentrations until reaching steady-state and then the G - V curves were constructed from the tail currents, normalized, and fitted by a Boltzmann function. The numbers in parentheses represent the number of cells.

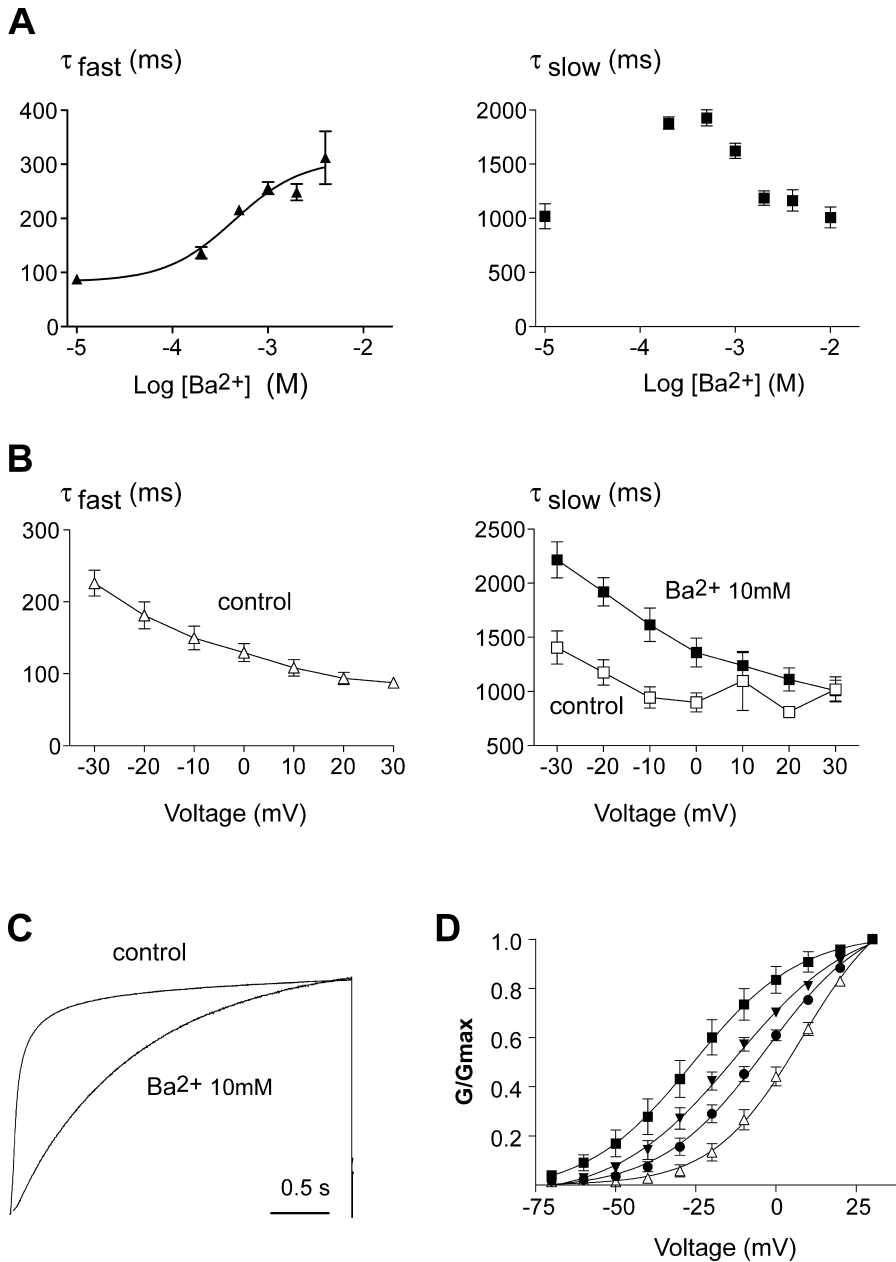


FIGURE 6. Effect of Ba^{2+} on activation gating. (A) Activation kinetics at +30 mV were fitted by a double exponential function, with a fast (left) and slow time constant (right), plotted as a function of external Ba^{2+} concentrations ($n = 10$). A sigmoidal dose-response function yielded an $\text{EC}_{50} = 0.44$ mM for the fast time constant. (B) The fast (left) and slow (right) activation time constant ($n = 11$) were plotted against voltage steps and were measured in the absence (empty symbols) or the presence of 10 mM Ba^{2+} (solid squares). (C) Representative normalized traces of current recorded at +30 mV before (control) and after application of 10 mM Ba^{2+} . (D) The normalized conductance, deduced from the tail currents was plotted as a function of step voltage for control (closed squares), 0.5 mM Ba^{2+} (solid triangles), 2 mM Ba^{2+} (closed circles), and 10 mM Ba^{2+} (empty triangles). The normalized conductance curves were fitted using a Boltzmann function.

triple pulse protocol (Fig. 8 C). In fact, KCNQ1 inactivation is an intrinsically voltage-independent and fast process with a $\tau_{\text{inact}} = 20 \pm 3$ ms ($n = 7$) (Fig. 8 C). The decaying current of the third test pulse which reflects the reinduction of inactivation is strongly depressed in the presence of 10 mM Ba^{2+} (Fig. 8 C).

External Magnesium Does Not Affect KCNQ1 Conductance and Gating

The strong impact of Ba^{2+} on KCNQ1 gating could in principle result from a surface screening charge effect. Ba^{2+} may screen diffuse negative charge and thus shift channel gating. Our experiments indicate that this is not the case, for two reasons. First, in case of a surface

screening charge effect, one would expect other divalent cations to produce similar effects. External Mg^{2+} (10 mM) neither blocked KCNQ1 conductance nor affected channel gating (online supplemental material, available at <http://www.jgp.org/cgi/content/full/jgp.200409068/DC1>). Steady-state experiments indicate that current amplitude was not inhibited by 10 mM Mg^{2+} . When measured at +30 mV, the steady-state current amplitude was 5.0 ± 0.8 μA ($n = 6$) and 6.1 ± 0.9 μA ($n = 6$) in the absence and presence of external Mg^{2+} (10 mM), respectively. In addition, external Mg^{2+} (10 mM) did not shift KCNQ1 activation curve ($V_{50} = -18.3 \pm 0.6$ mV and $V_{50} = -22.9 \pm 0.6$ mV with and without 10 mM Mg^{2+} , respectively, $n = 6$). External

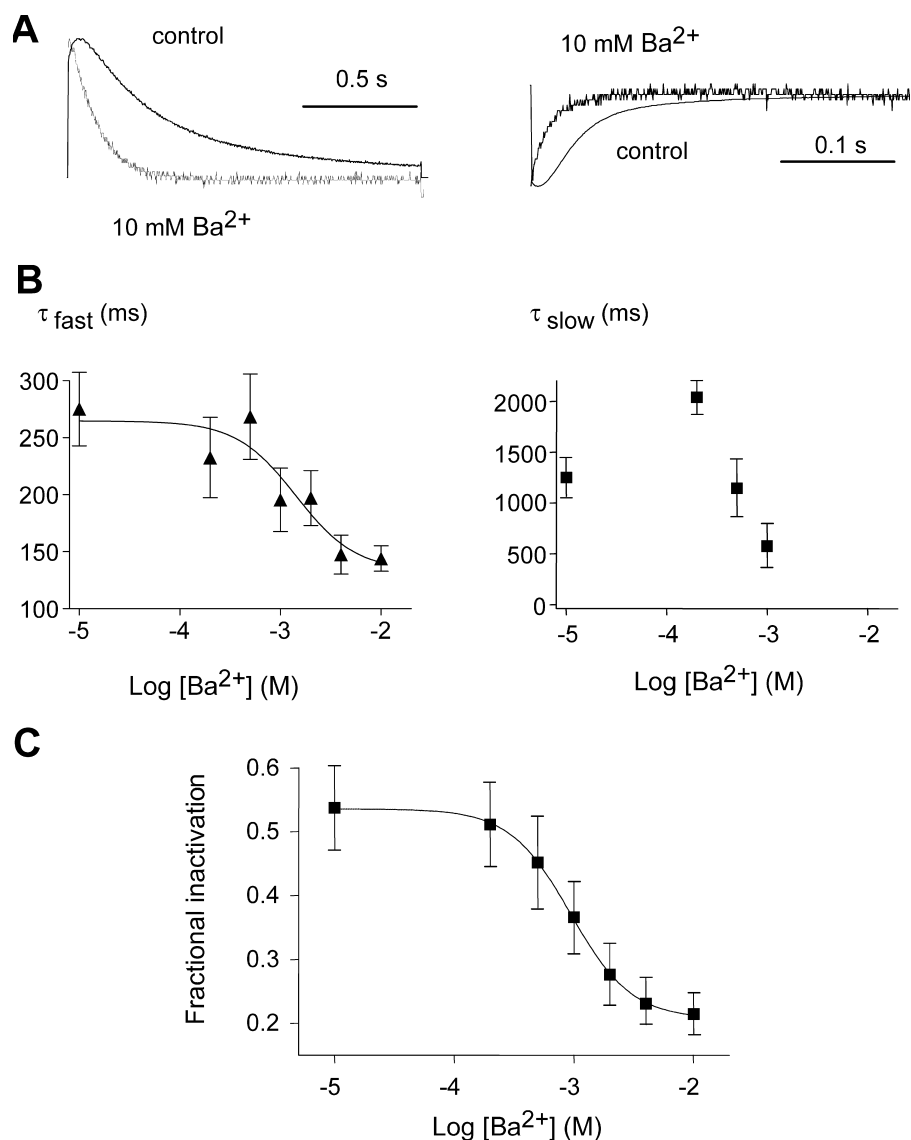


FIGURE 7. Concentration-dependence of Ba²⁺ action on deactivation and inactivation gating. (A) Representative normalized traces of tail current recorded at -60 mV (left) or at -140 mV (right) after a $+30$ -mV prepulse, before (control) and after application of 10 mM Ba²⁺. (B) The deactivation kinetics were recorded at -60 mV after a $+30$ -mV prepulse. Tail currents were fitted by a double exponential function, with a fast (left) and slow time constant (right), plotted as a function of external Ba²⁺ concentrations ($n = 7$). A sigmoidal dose-response function yielded an $EC_{50} = 1.4 \pm 0.2$ mM for the fast time constant ($n = 7$). (C) Dose-dependent inhibition of KCNQ1 inactivation by increasing concentrations of Ba²⁺ as measured by the hook of the tail currents at -60 mV ($n = 7$). A sigmoidal dose-response function yielded an $EC_{50} = 0.94 \pm 0.07$ mM ($n = 7$).

Mg²⁺ (10 mM) did not affect channel inactivation (fractional inactivation $f = 0.58 \pm 0.09$ and $f = 0.55 \pm 0.08$ with and without 10 mM Mg²⁺, respectively, $n = 6$). The activation and deactivation kinetics were not altered by 10 mM Mg²⁺ (online supplemental material). Similarly, external Ca²⁺ (10 mM) neither blocked channel conductance nor affected KCNQ1 gating (unpublished data). Second, replacing partially external Na⁺ by K⁺ (50 mM) eliminated the impact of 10 mM Ba²⁺ on KCNQ1 gating, including activation and deactivation kinetics, inactivation and rightward shift of the activation curve (see below). Thus, these results do not stand with a surface screening charge hypothesis.

High External K⁺ Attenuates Ba²⁺ Block and Eliminates its Impact on Channel Gating

Previous studies have shown that elevating extracellular K⁺ can relieve the Ba²⁺ block of voltage-gated K⁺ cur-

rents, including those of *Shaker* channels (Armstrong and Taylor, 1980; Wollmuth, 1994; Hurst et al., 1995). To check whether the effects of extracellular Ba²⁺, including those related to gating, are mediated through its interaction with the permeation pathway and could be competed out by K⁺, we examined the Ba²⁺ action at high external K⁺ (50 mM [K⁺]_o) (Figs. 9 and 10). The voltage-dependent block of KCNQ1 current produced by external Ba²⁺ was significantly attenuated in 50 mM [K⁺]_o, as measured in steady-state experiments. For example, at $+30$ mV 10 mM Ba²⁺ led to a $54 \pm 2\%$ block of KCNQ1 outward currents in 50 mM [K⁺]_o, compared with a $72 \pm 3\%$ inhibition in 2 mM [K⁺]_o ($P < 0.01$, $n = 10$; Fig. 9, A and B and Fig. 11, A and B). Similarly, high external K⁺ attenuated the Ba²⁺-induced inhibition of inward tail currents when measured at $+30$ -mV test pulse and -120 -mV tail potential with a $58 \pm 4\%$ block in 50 mM [K⁺]_o, compared with a $89 \pm$

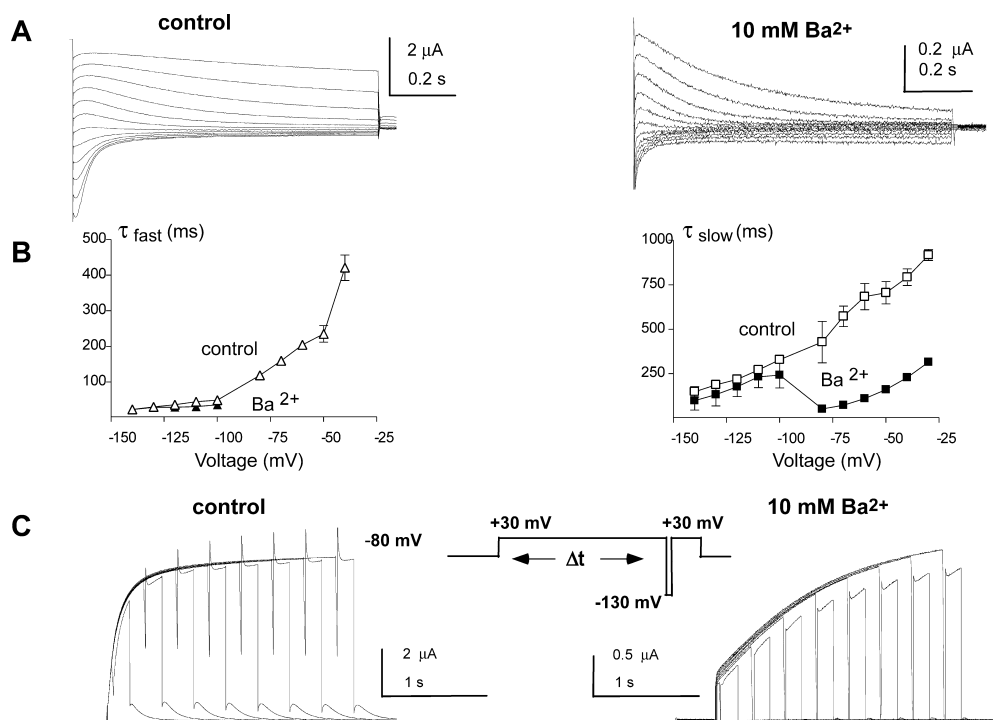


FIGURE 8. Effect of barium on time-dependent inactivation and voltage-dependent deactivation kinetics. (A) Representative tail currents obtained after a 1.5-s prepulse at +30 mV from -140 to -30 mV in 10-mV increments before (left) and after (right) application of 10 mM Ba²⁺. (B) The deactivation kinetics were determined in the absence (control) or presence of 10 mM Ba²⁺ at various tail potentials following a 1.5-s prepulse at +30 mV and were fitted by a double exponential function, with a fast (left) and slow (right) time constant ($n = 5$). (C) Representative current traces ($n = 7$) where the membrane potential was stepped to +30 mV for increasing durations by increments of 200 ms to progressively activate and inactivate the channel; then, a brief (15

ms) hyperpolarizing interpulse to -130 mV was used to allow channel recovery from inactivation before a +30-mV test pulse was applied to reopen and reinactivate KCNQ1 channels. The decaying current of the third test pulse (reinduction of inactivation) could be fitted by a single exponential function.

6% block in 2 mM [K⁺]₀ ($P < 0.01$, $n = 6$; Fig. 9, C and D). Like in low external K⁺, the Ba²⁺ block in 50 mM [K⁺]₀ was voltage dependent with weaker inhibition at increasing depolarization from -10 mV to +30 mV (Fig. 11, A and B). Extracting the apparent Ba²⁺ K_D values from the fractional current inhibition (see MATERIALS AND METHODS and Fig. 11 C) allowed us to derive the electrical distance of the Ba²⁺ blocking site in high external K⁺ (50 mM) with $\delta = 0.28$, a value similar to that found above for the Ba²⁺ slow site ($\delta = 0.32$). In contrast to low external K⁺, the voltage dependence of the Ba²⁺ block in 50 mM [K⁺]₀ was different at negative potentials from -40 to -60 mV, with a decreasing inhibition (Fig. 11 A). The increase in apparent K_D values at these negative potentials probably reflects the propensity of Ba²⁺ to dissociate in the inward direction.

Increasing external K⁺ did not only attenuate the voltage-dependent Ba²⁺ block but also eliminated all Ba²⁺ effects on channel gating. First, while at low external K⁺ (1 μ M and 2 mM) Ba²⁺ virtually suppressed KCNQ1 inactivation, Ba²⁺ left KCNQ1 inactivation intact at 50 mM [K⁺]₀, as revealed by the presence of the hook in tail currents (Fig. 9 A) with a fractional inactivation of ~ 0.6 without and with Ba²⁺ (Fig. 10 C). It is also important to notice that in contrast to C-type inactivation, KCNQ1 inactivation is not affected by high external K⁺. Second, in 50 mM [K⁺]₀ Ba²⁺ did not pro-

duce a significant gating shift in the voltage dependence of activation that is observed at 1 μ M or 2 mM [K⁺]₀ (Fig. 9 E). Ba²⁺ (10 mM) shifted by only +2.4 mV the KCNQ1 activation curve at 50 mM [K⁺]₀, while it produced a gating shift of +28.2 and +34 mV at 1 μ M [K⁺]₀ and 2 mM [K⁺]₀, respectively. Third, high external K⁺ prevented the slowing of KCNQ1 activation kinetics produced by Ba²⁺ (Fig. 10 A). This relief did not result from an acceleration of Ba²⁺ dissociation in high external K⁺. On the contrary, Ba²⁺ washout kinetics were even slower in 50 mM [K⁺]₀ compared with 2 mM [K⁺]₀ (Table II). Fourth, elevation of external K⁺ from 2 to 50 mM prevented the accelerating effect of Ba²⁺ on deactivation kinetics (Fig. 10 B).

The electrical distance of the Ba²⁺ blocking site in high external K⁺ ($\delta = 0.28$) is similar to that of the slow Ba²⁺ site ($\delta = 0.32$) as deduced from kinetic experiments in low external K⁺. Thus, we hypothesized that the fast Ba²⁺ site observed in 2 mM [K⁺]₀ could be involved in the gating effects mediated by Ba²⁺. Accordingly, we measured the wash-in kinetics of Ba²⁺ in high external K⁺ (50 mM) (Fig. 3). Under these conditions, Ba²⁺ on-binding displayed a fast and slow kinetic component. However, the wash-in time course was dominated by the slow kinetic component (Fig. 3, A and B; Table II). While in low external K⁺, the relative amplitude of the slow component decreased from 0.45 to

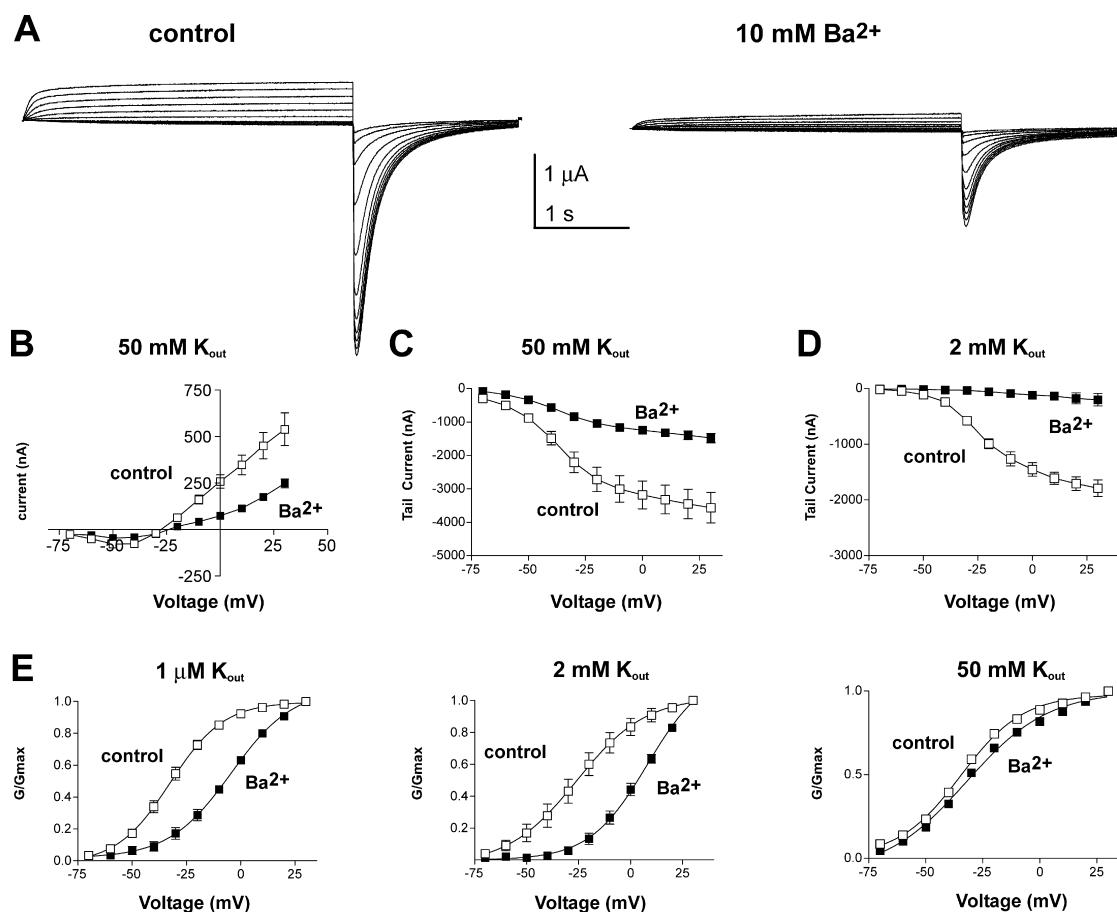


FIGURE 9. Barium effects in high external K^+ solutions. (A) Representative KCNQ1 current traces recorded in high external K^+ (50 mM) solution, in the absence (left) or presence (right) of 10 mM Ba^{2+} . From a holding potential of -80 mV, oocytes were stepped for 3 s from -70 to $+30$ mV in 10-mV increments and then repolarized for 1.5 s at -120 -mV tail potential. (B) Current-voltage relations ($n = 8$) before (control, empty squares) and after application of 10 mM Ba^{2+} (solid squares) measured at steady-state in high external K^+ (50 mM). (C) Current-voltage relations of the same experiments as in B ($n = 8$), before (control, empty squares) and after application of 10 mM Ba^{2+} (solid squares) measured at -120 mV tail potential in high external K^+ (50 mM). (D) Current-voltage relations ($n = 7$), before (control, empty squares) and after application of 10 mM Ba^{2+} (solid squares) measured at -120 -mV tail potential in low external K^+ (2 mM). (E) The normalized conductance, deduced from the tail currents at -120 mV, was measured before (control, open squares) and after application of 10mM Ba^{2+} (solid squares) in $1 \mu\text{M}$ (left), 2 mM (middle), and 50 mM (right) external K^+ . The normalized conductance curves were fitted using a Boltzmann function.

0.24 at potentials between -30 and $+30$ mV, in high external K^+ it predominated at 0.89–0.85 at potentials between 0 and $+30$ mV (Fig. 3 B). In contrast, the relative amplitude of the fast component had a minor contribution in 50 mM $[K^+]_0$, ranging from 0.1 to 0.15 at increasing depolarization (Fig. 3 B). These results suggest that the slow Ba^{2+} site ($\delta = 0.32$) measured in wash-in kinetics in low external K^+ and which predominates in high external K^+ mainly accounts for the voltage-dependent block of Ba^{2+} . The fast Ba^{2+} site ($\delta = 0.06$) that has a minor contribution in high external K^+ likely accounts for the gating effects of Ba^{2+} observed in low external K^+ solutions. In support of this hypothesis, the experiment depicted in Fig. 4 indicates that external Ba^{2+} could bind to a deep occluding site in the closed state (-100 mV) and remained bound to it dur-

ing washout at -100 mV, thus leading to channel blockade. In contrast, the washout removed Ba^{2+} from its fast external site, hence preventing its impact on channel gating. Indeed, despite a substantial channel block no effects on inactivation, deactivation, and activation were induced by Ba^{2+} exposure (Fig. 4, C and D).

Tail Current Analysis and Modeling

Recent studies suggested the presence of multiple open states for homomeric KCNQ1 channels (Pusch et al., 1998; Tristani-Firouzi and Sanguinetti, 1998). They proposed a sequential gating scheme (Scheme I) assuming two closed states, two voltage-dependent open-states, and a fast, voltage-independent inactivated state (Pusch et al., 1998).

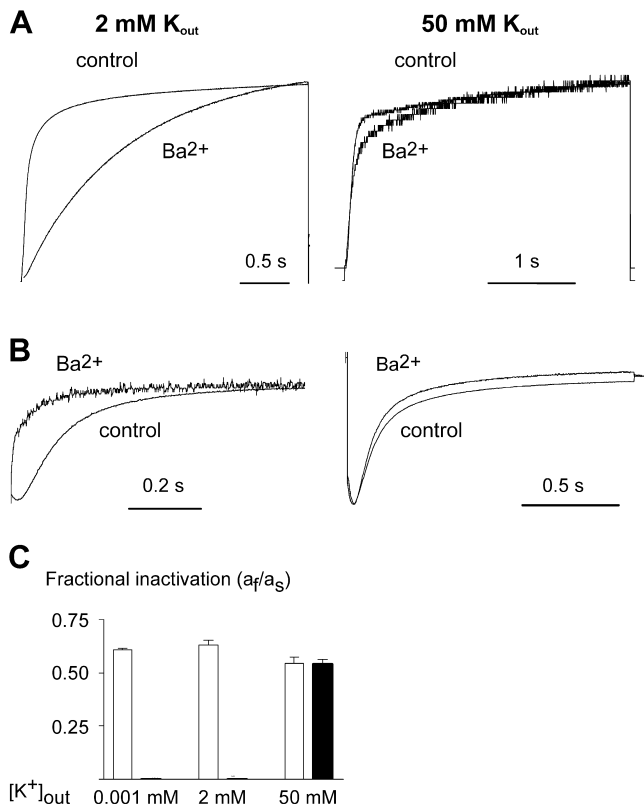


FIGURE 10. High external K^+ relieves the impact of Ba^{2+} on KCNQ1 activation, inactivation, and deactivation gating. (A) Representative normalized traces of KCNQ1 activation measured at $+30$ mV before (control) and after application of 10 mM Ba^{2+} in 2 mM $[K^+]_o$ (left), or 50 mM $[K^+]_o$ (right). (B) Representative normalized traces of KCNQ1 deactivation measured from a $+30$ -mV prepulse at -120 -mV tail potential before (control) and after 10 mM Ba^{2+} in 2 mM $[K^+]_o$ (left), or 50 mM $[K^+]_o$ (right). (C) Fractional KCNQ1 inactivation measured from the tail currents before (empty bars) and after application of 10 mM Ba^{2+} (filled bars) in 1 μ M, 2 mM, and 50 mM external K^+ . Tail currents were recorded at -60 mV for 1 μ M and 2 mM $[K^+]_o$ and at -120 mV for 50 mM $[K^+]_o$. The fractional inactivation was corrected for the change in driving force.

Using a tail current analysis, we tested whether KCNQ1 channels could populate multiple open states in the presence of 10 mM Ba^{2+} at steady-state in low external K^+ (2 mM). We measured the kinetics of current deactivation at -60 mV tail potential for various prepulse durations to $+30$ mV in the absence and presence of 10 mM Ba^{2+} (Fig. 12, A and B). KCNQ1 deactivation kinetics exhibited a fast and slow component and in order to take into account the recovery from inactivation, we fitted the whole tail current time course by a triple exponential function (see MATERIALS AND METHODS). The rate of channel deactivation progressively decreased with increasing prepulse durations (Fig. 12, C and D). For example, the fast and slow time constants of KCNQ1 deactivation increased, respec-

TABLE II
Kinetic Parameters of Barium Binding to KCNQ1 Channels in High External K^+ Solutions

| | Wash-in | | Washout | |
|-------|---------------|--------------------|-----------------|--------------------|
| | Fast τ | Slow τ | Fast τ | Slow τ |
| | ms | | ms | |
| 0 mV | 319 ± 37 | $6,472 \pm 806$ | 511 ± 49 | $11,810 \pm 1,879$ |
| 10 mV | 368 ± 58 | $7,615 \pm 1,269$ | 576 ± 88 | $13,774 \pm 2,038$ |
| 20 mV | 532 ± 149 | $8,173 \pm 1,567$ | 713 ± 145 | $11,958 \pm 1,715$ |
| 30 mV | 781 ± 276 | $12,047 \pm 1,703$ | $1,013 \pm 437$ | $13,134 \pm 1,769$ |

On- and off-binding kinetics of extracellular barium were measured in high external K^+ solutions (50 mM) by applying 10 mM Ba^{2+} to KCNQ1 channels previously opened at the indicated potentials. Wash-in and washout kinetics were fitted by a double exponential function and the fast and slow time constants are indicated (mean \pm SEM, $n = 5$).

tively, from $\tau_f = 91 \pm 16$ ms and $\tau_s = 156 \pm 22$ ms at 0.05 s prepulse ($n = 12$) to $\tau_f = 160 \pm 17$ ms and $\tau_s = 406 \pm 41$ ms at 2.25 -s prepulse ($n = 12$). It is interesting to note that the tail currents were hooked for prepulse durations >100 ms and the hook amplitude increased with the length of the prepulse (Fig. 12 E). The progressive decrease in deactivation rate as a function of prepulse duration, is indicative of the existence of multiple open states (Tristani-Firouzi and Sanguinetti, 1998). The sigmoidal time course of the tail currents at negative potentials further supports this idea (Fig. 8 A, left). In the presence of 10 mM Ba^{2+} , the rate of channel closure was faster but it did not decrease significantly with increasing prepulse duration (Fig. 12 C). In the presence of Ba^{2+} , the acceleration of deactivation was such that the slow kinetic component disappeared. As shown before, Ba^{2+} potently depressed inactivation by virtually eliminating the hook of the tail currents (Fig. 12 D). The sigmoidal shape of the tail currents at negative potentials disappeared in the presence of Ba^{2+} (Fig. 8 A, right). These results suggest that in the presence of Ba^{2+} and in low external K^+ , unblocked KCNQ1 channels transit to one open state from which inactivation is not permissible (see DISCUSSION).

For modeling KCNQ1 behavior in low external K^+ (2 mM) and in the absence of Ba^{2+} , the time course of channel activation was fitted according to Scheme I, following the previously described gating model (Pusch et al., 1998), where only O_1 and O_2 are assumed to conduct and for which we assumed for simplicity that no openings could occur when the cell is voltage-clamped to -80 mV. A fitting analysis was performed for KCNQ1 activation in low and high external K^+ in the absence of barium (see table of online supplemental material and MATERIALS AND METHODS). To validate the model, we used the fit data and simulated a $+30$ -mV activation time course and a tail envelope protocol.

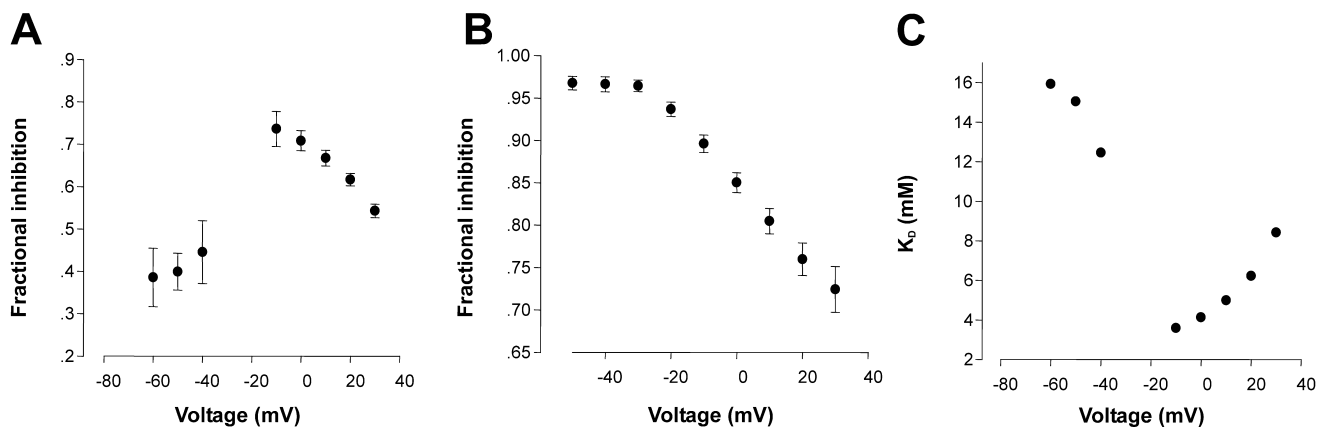


FIGURE 11. Effect of membrane voltage on the fractional current inhibition produced by barium. (A) Fractional current inhibition produced by 10 mM Ba^{2+} in high external K^+ (50 mM) and plotted as a function of step voltages. (B) Fractional current inhibition produced by 10 mM Ba^{2+} in low external K^+ (2 mM) and plotted as in A. (C) Extraction of the Ba^{2+} K_D values from the fractional current inhibition (see MATERIALS AND METHODS) and plotted a function of step voltages.

The simulated activation time course in low external K^+ closely matched the experimental trace (Fig. 13 C). Simulating the tail currents for the different prepulse intervals allowed us to generate a simulated trace that described faithfully the dependence of the “hook” amplitude on the prepulse duration as observed in experimental data (Fig. 13, A and B). Similarly, the simulated time course of KCNQ1 activation in high external K^+ and in the absence of Ba^{2+} closely matched the averaged experimental traces (Fig. 13 D).

Since high external K^+ prevented the impact of Ba^{2+} on channel gating, we used a simplistic model of voltage-dependent block as illustrated in Scheme II in order to simulate Ba^{2+} block in high external K^+ solutions.

We assumed that the conductance reduction originates from the existence of one Ba^{2+} -occluding site. Thus, the channel could populate two different kinds of states, assumed to be in equilibrium: a Ba^{2+} -blocked state (B) and a Ba^{2+} -unblocked state. For simplicity, we assumed that all states (C_0 to I) bind Ba^{2+} with equal affinity (see MATERIALS AND METHODS). As the kinetics of Ba^{2+} in 50 mM $[K^+]_0$ are clearly dominated by a slow component, we could estimate Ba^{2+} affinity for the open states. To preserve microscopic reversibility, we also assumed that transitions through the blocked states occur with similar rates to their unblocked counterparts. Simulating the time course of KCNQ1 activation in the presence of 10 mM Ba^{2+} and in high external K^+ solutions (50 mM) faithfully matched the kinetics of the averaged experimental traces, including the decrease of the current amplitude (Fig. 13 D). In contrast, the simple model of voltage-dependent block failed to simulate the complex Ba^{2+} actions on KCNQ1 channels in low external K^+ conditions (see DISCUSSION).

DISCUSSION

The main findings of the present work are that extracellular Ba^{2+} exerts on homomeric KCNQ1 channels a series of complex effects, including a voltage-dependent pore block as well as unique gating alterations. Our results indicate that external barium interacts with the permeation pathway of KCNQ1 at two distinct sites. (a) A slow deep Ba^{2+} site that occludes the channel pore and could be simulated by a model of voltage-dependent block. (b) A fast superficial Ba^{2+} site that barely influences channel block and mainly affects channel gating by shifting rightward the voltage dependence of activation, slowing activation, speeding up deactivation kinetics, and inhibiting channel inactivation. Both sites appear to reside in the permeation pathway as high external potassium attenuates the Ba^{2+} inhibition of channel conductance and abolishes its impact on channel gating.

External Ba^{2+} Interacts with the Pore of KCNQ1 at Two Distinct Sites

The existence of at least two distinct Ba^{2+} binding sites is supported by the biphasic time course of Ba^{2+} wash-in and washout kinetics. Furthermore, the Ba^{2+} dose-response curves need to be fitted to two binding sites. Although at negative potentials (-70 to -30 mV) the curves could be fitted assuming one site, the values of the Hill coefficients of ~ 2.1 – 2.8 suggest the existence of at least two Ba^{2+} sites of very similar affinity (Fig. 5 C). However, with increasing depolarization, the progressive appearance of a lower affinity site reflects the voltage dependence of the slow deep site (Fig. 5 D). The Ba^{2+} affinities for the two sites probably overlap, especially at negative potentials; however, it is worth noting that the EC_{50} values obtained for the effects of

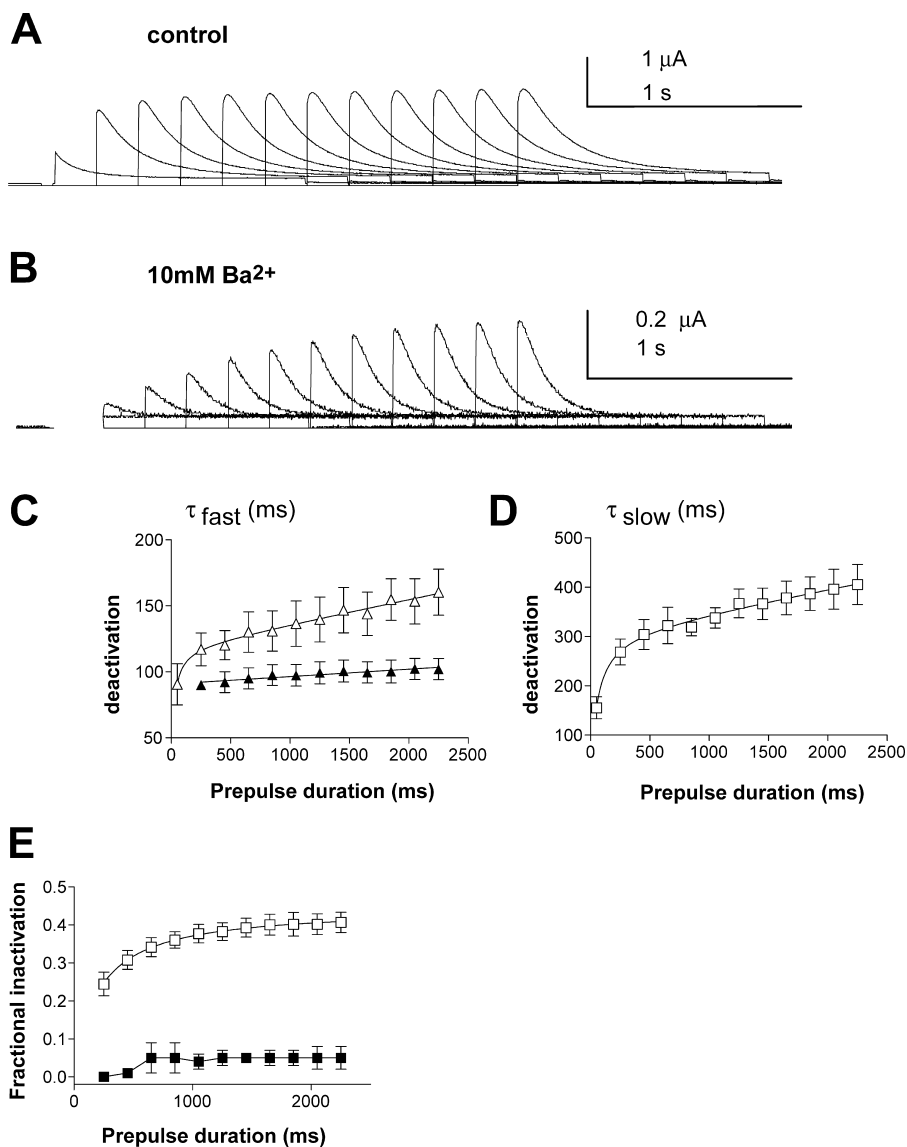
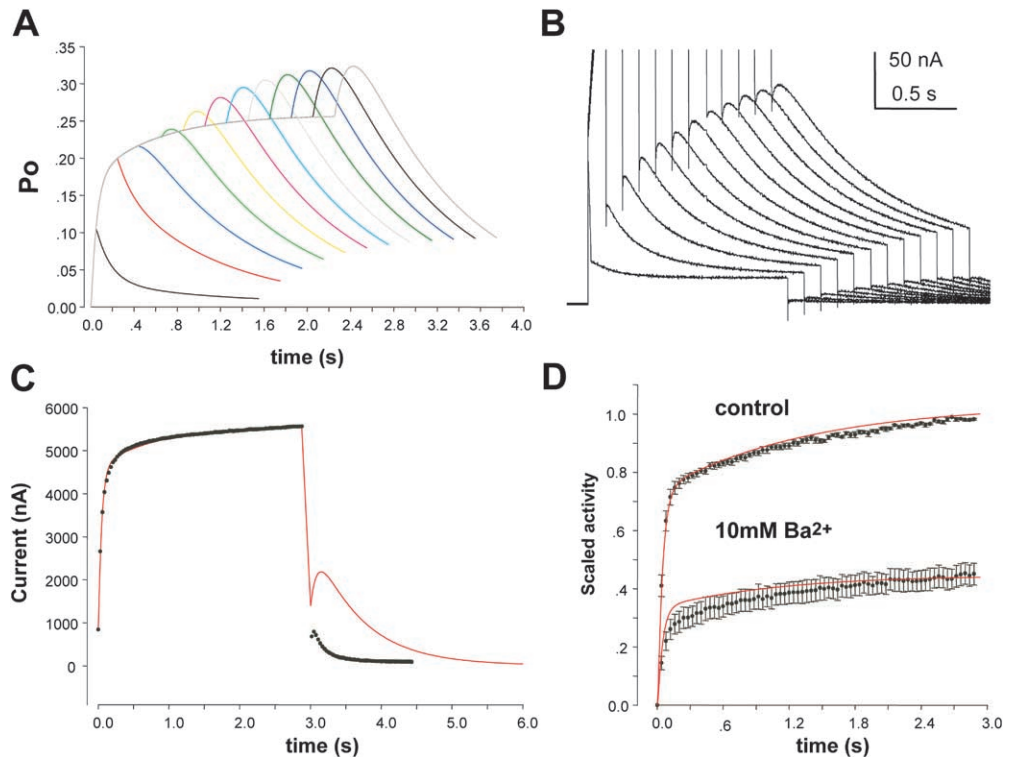


FIGURE 12. Effect of Ba²⁺ on tail current analysis. (A and B) Representative tail envelopes from the same oocyte showing the kinetics of current deactivation, measured at a -60 mV tail potential after the progressive activation of KCNQ1 channels by various prepulse durations to +30 mV from 50 ms to 2.25 s in 200-ms increments, in the absence (A) and presence of 10 mM Ba²⁺ (B). (C and D) Fast and slow time constant of channel deactivation before (empty symbols) and after 10-mM Ba²⁺ application (solid triangles) plotted as a function of prepulse duration ($n = 12$). (E) Fractional inactivation measured from the same tail current analysis before (empty squares) and after 10-mM Ba²⁺ application (solid squares) plotted as a function of prepulse duration ($n = 12$).

Ba²⁺ on channel gating are within the same range (0.44–1.4 mM). A good discrimination between two distinct Ba²⁺ sites relies on the data obtained in high external K⁺ solutions. In steady-state experiments, elevating external K⁺ from 2 to 50 mM substantially preserved the voltage-dependent block of channel conductance though weakening the amplitude of Ba²⁺ block, but abolished its impact on channel gating. Along the same line, Ba²⁺ wash-in kinetics performed in high external K⁺ were dominated by the slow kinetic component while the fast component had a minor contribution. Finally, binding of Ba²⁺ to the closed channel state adds further support to the existence of two discrete binding sites. Prolonged wash-in and washout of external Ba²⁺ from closed channels at -100 mV leaves Ba²⁺ bound to its deep site but removes Ba²⁺ from its superficial site, leading to a significant channel block and preventing Ba²⁺ impact on

KCNQ1 channel gating. Taken together, our results indicate that external barium interacts with the permeation pathway of KCNQ1 at two sites. (a) A superficial Ba²⁺ site that mainly affects channel gating and senses weakly the electric field ($\delta_1 = 0.06$). It corresponds to the fast kinetic component of Ba²⁺ wash-in that virtually disappears in high external K⁺. Knowing the strong rightward shift effect of Ba²⁺ on KCNQ1 activation curve (+34 mV at 10 mM Ba²⁺), it is not surprising that the fast kinetic component depresses substantially the current amplitude in wash-in experiments, a feature that in fact reflects the impact of Ba²⁺ on gating rather than a direct channel block. (b) A deep Ba²⁺ site that occludes the channel pore and lies profound in the membrane electric field ($\delta_2 = 0.33$). It corresponds to the slow kinetic component of Ba²⁺ wash-in in low external K⁺ that predominates in high external K⁺.

FIGURE 13. Simulation of KCNQ1 gating. (A) Simulation of the tail envelope protocol as described in Fig. 12. The simulation was done according to Scheme I. All rate constants represent averaged values obtained from the least square fit of currents at +30 mV and -60 mV in the absence of Ba²⁺ ($n = 7$). The ordinate represents the calculated open probability (fractional occupancy of the open states). (B) Representative experimental traces of the tail envelope protocol. (C) Simulation (red line) of KCNQ1 activation (+30 mV) and deactivation (-60 mV) according to Scheme I, superimposed on an experimental trace. (D) Scaled simulation (red lines) of KCNQ1 activation in high external K⁺ (50 mM) according to Scheme II, in the absence and presence of 10mM Ba²⁺. The simulation was superimposed on experimental average traces obtained from seven oocytes.



It was previously found that external Ba²⁺ can interact with two distinct and sequential binding sites in *Shaker* K⁺ channels (Hurst et al., 1995). In contrast to KCNQ1 channels where the slow deep site is mainly involved in the voltage-dependent Ba²⁺ block, both fast and slow sites were found to be responsible for the Ba²⁺ block of *Shaker* channels with apparent electrical distances of 0.14 and 0.25, respectively. Furthermore, it was suggested that Ba²⁺ had first to populate a fast external site before being able to populate a slow deep site to produce the voltage-dependent block (Hurst et al., 1995). Consistent with two sites in series, the washout time course of Ba²⁺ in *Shaker* channels exhibited only one kinetic component (slow). In KCNQ1 channels, the features of external Ba²⁺ binding are different and the washout kinetics display a biphasic time course suggesting that the two sites are not sequential and exhibit distinct functional properties. In further support of the latter conclusion, high external K⁺ preserves the voltage-dependent block at the deep slow site, while virtually abolishing the superficial fast site and the concomitant gating impact of Ba²⁺. This suggests that one does not need Ba²⁺ to bind to the external fast site in order to get Ba²⁺ binding to the slow deep site. The δ value obtained for the deep slow site of Ba²⁺ block of KCNQ1 is similar to that found for *Shaker* K⁺ channels ($\delta = 0.35$) (Harris et al., 1998), but is lower than that

obtained for the high affinity Ba²⁺ block of the inward-rectifier Kir2.1 channels ($\delta = 0.62$) (Alagem et al., 2001) or that found for the two-pore domain K⁺ channel KCNK0 ($\delta = 0.60$) (Ilan and Goldstein, 2001). Thus, despite the high degree of conservation of the pore region among various classes of K⁺ channels, subtle differences in the architecture of the permeation apparatus can confer different voltage-sensitivity for Ba²⁺ block. In other words, the same Ba²⁺ binding site may sense the voltage drop differently depending on the precise geometry of a given K⁺ channel pore.

Two external K⁺ sites located above the deep Ba²⁺ site were found to affect Ba²⁺ binding (Neyton and Miller, 1988a,b). One rapidly equilibrating site, of high K⁺ affinity (lock-in site), whose occupancy by K⁺ prevents outward Ba²⁺ dissociation, and one site of lower K⁺ affinity (enhancement site), whose occupancy by K⁺ repulses Ba²⁺ in the inward direction. All three sites were found to be accessible from the external solution to the closed channel state and in *Shaker* channels the deep Ba²⁺ site lies between the activation gate and the C-type inactivation gate (Harris et al., 1998). Given the high degree of topological similarity and the strong voltage dependence of Ba²⁺ block, we assume that the Ba²⁺ binding site of voltage-gated K⁺ channels in general, and that of KCNQ channels in particular, is analogous to that of KcsA (Jiang and MacKinnon, 2000). Mutagen-

esis studies performed on *Shaker* K⁺ channels showed that the residue T441, located at the inner boundary of the selectivity filter, may represent the deep high-affinity Ba²⁺ binding site (Harris et al., 1998). In the inward-rectifier Kir 2.1, an amino acid T141, located just one residue COOH-terminal to the equivalent T441 in *Shaker* was suggested to stabilize the high affinity Ba²⁺ binding (Alagem et al., 2001). These results fit well with the crystallographic study performed on KcsA channels (Jiang and MacKinnon, 2000). In KCNQ1, the homologous residue is also a threonine, T311. In *Shaker* channels, it was suggested that the residue V443 within the selectivity filter facilitates the entry and exit of Ba²⁺ to and from the deep site, while the amino acid T449 in the external vestibule may represent the external lock-in site (Harris et al., 1998). Preliminary studies performed in V319C KCNQ1 (coexpressed with KCNE1) indicate that the valine mutation V319C located in the external vestibule at the analogous *Shaker* position T449 potentially prevents the voltage-dependent Ba²⁺ block but does not affect the action of Ba²⁺ on channel gating (unpublished data). However, knowing the impact of KCNE1 on the biophysical properties of KCNQ1, these data may not be directly applicable to homomeric KCNQ1 V319C channels. Taken together, these results emphasize the structural and functional differences existing between KCNQ1 and other voltage-gated K⁺ channels in their permeation pathway, despite the high degree of homology of their pore region. They also suggest that in homomeric KCNQ1, the two Ba²⁺ sites exhibit distinct functional properties.

Barium Affects Channel Gating

Compared with *Shaker*-like K⁺ channels and other K⁺ channels, barium uniquely affects KCNQ1 gating. In addition to its voltage-dependent block of the pore, Ba²⁺ also alters the gating of homomeric KCNQ1 channels by favoring the closed state. Barium produces a marked rightward shift of the voltage-dependence of activation, a slowing of activation, and an acceleration of deactivation kinetics. In addition, Ba²⁺ potentially depresses inactivation gating of KCNQ1, an effect that opposes current inhibition. This latter impact of Ba²⁺ on KCNQ1 inactivation appears to be very different from its effects on *Shaker* channels, where it does not alter C-type inactivation (Harris et al., 1998). In HERG channels, Ba²⁺ also produced a voltage-dependent block. However, at positive potentials and concomitant with the development of C-type inactivation, Ba²⁺ block is completely relieved (Weerapura et al., 2000), a feature that is not observed with an inactivation-deficient HERG mutant. In HERG channels, Ba²⁺ does not prevent inactivation; in fact Ba²⁺ does not bind to the HERG-inactivated state (Weerapura et al., 2000). This feature contrasts with that obtained for KCNQ1 chan-

nels and further suggests that despite the high degree of similarity of the pore region, substantial structural and functional differences exist between the various voltage-gated K⁺ channel subtypes. Interestingly, all these gating alterations are mediated by the same superficial fast site of Ba²⁺ binding. The effects of Ba²⁺ are specific and do not result from a surface-charge-screening effect, as Mg²⁺ and Ca²⁺ neither affect the conductance nor alter the gating of KCNQ1. Furthermore, the relief of the impact of Ba²⁺ on channel gating by high external K⁺ cannot be accounted for by a charge-screening effect. We believe that the right-shift produced by Ba²⁺ is not due to a voltage-dependent block: (a) A model of voltage-dependent block does not explain the marked slowing of the activation kinetics that is observed with increasing Ba²⁺ concentrations (Fig. 6 A, left), as the unblocking rate is not expected to depend on the blocker concentration; (b) in high external K⁺, there is still a voltage-dependent Ba²⁺ block (Fig. 11A), precisely under conditions where the right-shift has been lost (Fig. 9 E, right); (c) in low external K⁺, a simple model of voltage-dependent block was unable to simulate the time course of KCNQ1 activation and deactivation, precisely under conditions where Ba²⁺ produces the right-shift; (d) the decrease in maximum conductance was much more significant than the shift in V₅₀ produced by all Ba²⁺ concentrations; and (e) the activation, deactivation, and V₅₀ parameters are not affected to the same extent. For example, 4 mM Ba²⁺ shifts the fast time constant of the activation kinetics by more than 60 mV (Fig. 6, A and B), while it shifts by only 27 mV the voltage dependence of activation (Table I).

The mechanisms by which Ba²⁺ affects the gating machinery of KCNQ1 are unknown, however, we suggest that they are coupled to the permeation process as we have shown that the gating effects of Ba²⁺ are virtually abolished by high external K⁺. The gating impact of Ba²⁺ probably involves a fast-equilibrating superficial site of the KCNQ1 external vestibule that remains to be identified. Interestingly, we found that external acidification also alters the gating of homomeric KCNQ1 channels by inhibiting inactivation, slowing activation kinetics, and causing a rightward shift of the voltage dependence of activation (Peretz et al., 2002). It remains to determine whether external protons and Ba²⁺ ions affect the gating of KCNQ1 by similar mechanisms. In principle, the basic operation of an ion channel relies on the combination of two distinct properties, gating and permeation. However, these processes are not independent. Permeating ions are known to interact with the pore that undergoes conformational changes upon channel opening, inactivation, and closure. Are there regulatory sites within the permeation pathway and remote from the narrow selectivity filter that affect per-

meation and impact on channel gating? Recent studies have shown that residues of the pore region do interact with the voltage-sensor module of *Shaker* channels (Li-Smerin et al., 2000). This protein interface may be a crucial determinant for the molecular coupling of the voltage-sensor movement and the opening of the pore. A direct interaction between charges of the voltage sensor and pore residues has already been suggested for the coupling between S4 and the slow inactivation gate of the *Shaker* channel (Gandhi et al., 2000; Larsson and Elinder, 2000; Loots and Isacoff, 2000; Ortega-Saenz et al., 2000; Elinder et al., 2001). In KCNQ1, it is possible that Ba²⁺ binds to an external vestibule site that is involved in sensor-gate coupling. Alternatively, the features of the fast superficial Ba²⁺ site may be also consistent with Ba²⁺ binding to a site remote from the permeation pathway and residing in a crevice in which K⁺ but not Mg²⁺ or Ca²⁺ can compete with. Along the same line, barium was previously found to directly affect the gating of *Shaker* channels by speeding up the rate of the off-gating current (Hurst et al., 1996), which is in agreement with the proposal that Ba²⁺ destabilizes the open state and favors the closed state.

KCNQ1 α subunits produce very different gating phenotypes when coexpressed with the KCNE β subunits (Abbott and Goldstein, 1998; Melman et al., 2002). In the heart, coassembly of KCNQ1 with KCNE1 elicits voltage- and time-dependent K⁺ currents (I_{KS}) with very slow activation kinetics, a positive shift in the voltage dependence of activation, and an increase in unitary channel conductance, in comparison to homomeric KCNQ1 channels (Barhanin et al., 1996; Sanguinetti et al., 1996). When associated with KCNE3, KCNQ1 channels yield K⁺ currents that are nearly instantaneous and voltage independent (Schroeder et al., 2000). Regarding the strong impact of KCNE β subunits on the biophysical properties of KCNQ1, it will be important in future experiments to investigate in depth the effects of external Ba²⁺ on these heteromeric channel complexes. Our preliminary data indicate that compared with homomeric KCNQ1, external Ba²⁺ has a similar action on KCNQ1/KCNE1 channels and differential effects on KCNQ1/KCNE3 complexes (unpublished data).

Modeling Barium Action

A sequential gating scheme for homomeric KCNQ1 channels was proposed (Scheme I), assuming two closed states, two voltage-dependent open-states, and a fast, voltage-independent inactivated state (Pusch et al., 1998). More recently, inactivation of KCNQ1 was suggested to be related to a fast flicker of the open channel (Seeböhm et al., 2003).

Simulating KCNQ1 behavior in low external K⁺ (2 mM) and in the absence of Ba²⁺ according to Scheme I generated an activation time course and a tail-current

envelope that described faithfully the dependence of the inactivation (hook amplitude) on prepulse duration that is observed in experimental data (Fig. 13, A and B). The progressive decrease in deactivation rate as a function of prepulse duration, further supports the existence of multiple open states (Pusch et al., 1998; Tristani-Firouzi and Sanguinetti, 1998). Using the same Scheme I, we could simulate the time course of KCNQ1 activation in high external K⁺ and absence of Ba²⁺ (Fig. 13 D). This suggests that elevating external K⁺ does not alter significantly the gating transitions of KCNQ1 channels. In contrast to C-type inactivation, high external K⁺ does not directly affect the inactivation process of KCNQ1. Paradoxically, high external K⁺ relieves Ba²⁺ inhibition of inactivation gating. As high external K⁺ prevented the impact of Ba²⁺ on channel gating, we could use a simplistic model of voltage-dependent block, as shown in Scheme II, in order to simulate Ba²⁺ block in high external K⁺.

Assuming that the conductance block originates from the existence of one Ba²⁺-occluding site (i.e., the deep slow Ba²⁺ site), we could simulate the time course of KCNQ1 activation, including the decrease of current amplitude, in the presence of 10 mM Ba²⁺ under high external K⁺ conditions (Fig. 13 D). In contrast, the simple model of voltage-dependent block failed to simulate the complex Ba²⁺ actions on KCNQ1 channels in low external K⁺ conditions. On the other hand, inhibition of inactivation gating by Ba²⁺ in low external K⁺ cannot be accounted for by a general slowing of the sequential gating Scheme I, since very long depolarizing pulses (+30 mV for 6 min) performed in the presence of 4 mM Ba²⁺ did not restore KCNQ1 inactivation, as measured by the hooked tail currents (unpublished data). The loss of inactivation does not reflect a lower occupancy of the preceding open state O₂ at a given activation potential since when the membrane was stepped to larger depolarizing potentials of up to +80 mV, inactivation did not reappear (unpublished data). The tail current analysis performed in the presence of Ba²⁺ and in low external K⁺ conditions reveals that the rate of channel closure was faster but did not decrease with increasing prepulse duration. These results suggest that in the presence of Ba²⁺ and in low external K⁺, unblocked KCNQ1 channels transit to an open state from which inactivation is not permissible. Thus, we have to assume that external Ba²⁺ ions interact with the permeation pathway of KCNQ1 channels at two distinct and nonsequential sites. (a) A deep binding site that occludes the pore and is responsible for the voltage-dependent block of the channel conductance. This Ba²⁺-bound site leads to nonconducting channels, is well described by Scheme II and is partially competed out by high external K⁺. (b) A superficial binding site that barely contributes to channel block but produces

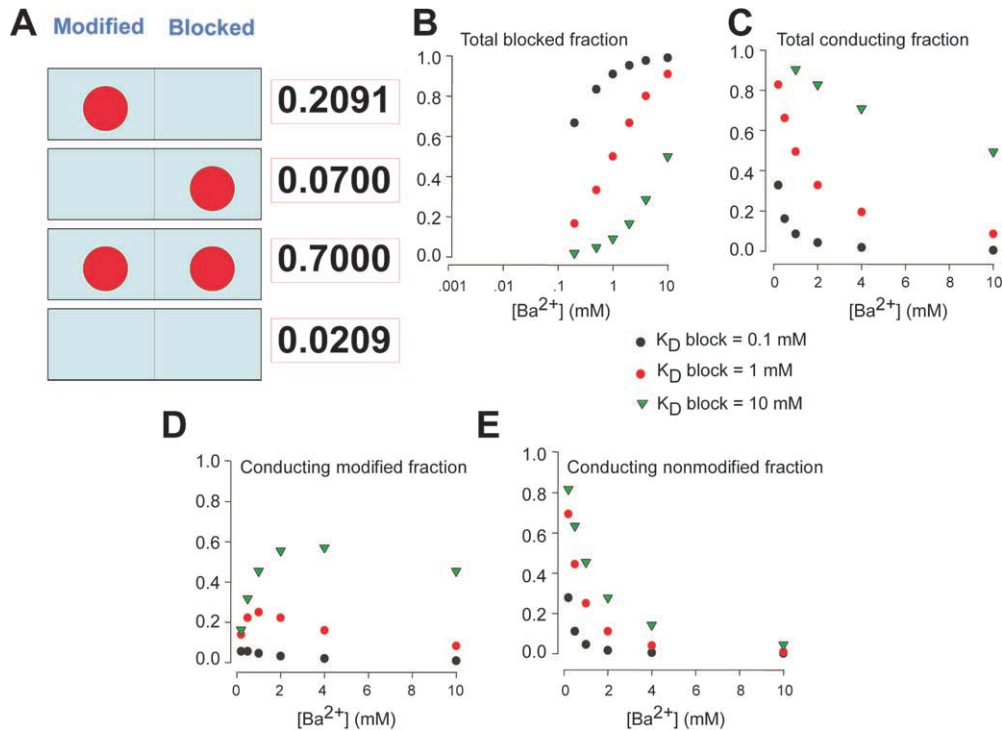


FIGURE 14. Schematic calculation of KCNQ1 channel distribution among different populations in the presence of Ba^{2+} . (A) Diagram of the different channel populations in the presence of 10 mM Ba^{2+} , assuming a $K_D = 3$ mM for channel block and a $K_D = 1$ mM for channel gating modification. It is assumed that external Ba^{2+} interact with the permeation pathway of KCNQ1 channels at two distinct and nonsequential sites. (1) A deep site that accounts for the voltage-dependent block of the pore and leads to blocked nonconducting channels. (2) A superficial site that modifies the gating transitions states of KCNQ1 and leads to modified-conducting channels. (B–E) Illustration of how the relative fractions of nonconducting (blocked), and conducting (modified and nonmodified) channels are predicted to evolve, depending on the K_D values for Ba^{2+} block.

the gating alterations on KCNQ1 and is totally competed out by high external K^+ . This latter Ba^{2+} -bound site modifies the gating transitions states of KCNQ1, leads to modified-conducting channels and cannot be accounted for by a model of voltage-dependent block. Assuming a $K_D = 3$ mM for channel block (via occlusion) and a $K_D = 1$ mM for channel modification (via gating), we can predict that in the presence of 10 mM Ba^{2+} and in low external K^+ (2 mM), 21% of KCNQ1 channels are conducting but modified, 7% are blocked, 70% are both modified and blocked and only 2% have no Ba^{2+} bound (Fig. 14 A). Overall, 77% of KCNQ1 channels are nonconducting (of which 70% are also modified) and 91% are modified (of which 70% are nonconducting). Fig. 14, B–E, illustrates how the relative fractions of nonconducting (blocked) and conducting (modified and nonmodified) channels are predicted to evolve, depending on the K_D values for Ba^{2+} block.

In all, our data suggest that despite the high degree of homology of the pore region of voltage-gated K^+ channel subtypes, significant structural and functional differences exist with KCNQ1 channels. Compared with *Shaker*-like K^+ channels, barium uniquely affects KCNQ1 channel gating, especially inactivation.

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