## Chromophores

# 2- and 2,7-Substituted para-N-Methylpyridinium Pyrenes: Syntheses, Molecular and Electronic Structures, Photophysical, Electrochemical, and Spectroelectrochemical Properties and Binding to Double-Stranded (ds) DNA 

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#### Abstract

Two $N$-methylpyridinium compounds and analogous $N$-protonated salts of 2- and 2,7-substituted 4-pyridylpyrene compounds were synthesised and their crystal structures, photophysical properties both in solution and in the solid state, electrochemical and spectroelectrochemical properties were studied. Upon methylation or protonation, the emission maxima are significantly bathochromically shifted compared to the neutral compounds, although the absorption maxima remain almost unchanged. As a result, the cationic compounds show very large apparent Stokes shifts


#### Abstract

of up to $7200 \mathrm{~cm}^{-1}$. The N -methylpyridinium compounds have a single reduction at ca. -1.5 V vs. $\mathrm{Fc} / \mathrm{Fc}^{+}$in MeCN . While the reduction process was reversible for the 2,7-disubstituted compound, it was irreversible for the mono-substituted one. Experimental findings are complemented by DFT and TD-DFT calculations. Furthermore, the $N$-methylpyridinium compounds show strong interactions with calf thymus (ct)-DNA, presumably by intercalation, which paves the way for further applications of these multi-functional compounds as potential DNA-bioactive agents.


## Introduction

Pyrene, a polycyclic aromatic hydrocarbon (PAH), is one of the most widely investigated fluorophores, as it possesses some unique properties, such as intense blue emission, a high fluorescence quantum yield, a long-lived singlet state, and propensity for excimer and exciplex formation. ${ }^{[1]}$ In addition to its rich photophysical properties, pyrene is chemically stable and can enhance intermolecular charge mobility. ${ }^{[1]}$ Many pyrene deriva-
tives have, therefore, been developed for a wide range of applications in modern scientific fields which include organic optoelectronic devices, such as organic light emitting diodes (OLEDs), organic field-effect transistors (OFETs), and organic photovoltaic cells (OPVs). ${ }^{[2]}$ The fluorescence of pyrene derivatives has been employed to investigate interactions in macromolecules and lipids, ${ }^{[3,4]}$ and also in probes for signalling of biomolecules, such as oligonucleotides and nucleic acids (e.g., DNA). ${ }^{[5]}$ In addition, their long fluorescence lifetimes and the
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possibility of excimer formation enable the exploration of advanced applications, such as the detection of DNA/RNA interactions both as a single label, ${ }^{[6]}$ as well as excimer-forming pairs or as multi-pyrene probes. ${ }^{[7]}$ The flat aromatic structure of pyrene derivatives facilitates its stacking with nucleobases and consequently intercalation into DNA/RNA. ${ }^{[8]}$ However, intercalative binding to DNA/RNA has limited selectivity and, therefore, some interesting applications of those pyrene derivatives rely on their interactions with DNA or RNA grooves ${ }^{[9]}$ or in combination with their turn-on and -off excimers. ${ }^{[10]}$ Pyrene derivatives have further been employed for applications including the determination of cellular oxygen concentrations or reactive oxygen species (ROS) in biological systems, ${ }^{[11]}$ and the detection of various metal ions in solutions. ${ }^{[12]}$ In addition, as the fluorescence of pyrene-based compounds can be influenced by external stimuli, this has been exploited for various applications such as sensing temperature, ${ }^{[13]}$ pressure, ${ }^{[14]} \mathrm{pH}_{,}^{[15]}$ and detection of small molecules. ${ }^{[16]}$ The ability of pyrene to exhibit strong $\pi \cdots \pi$ interactions makes it a suitable probe to examine arene-perfluoroarene interactions and associated electrostatic quadrupole-quadrupole interactions, even in dilute solutions, e.g., with hexafluorobenzene and octafluoronaphthalene. ${ }^{[17]}$ Furthermore, metal organic frameworks (MOFs) ${ }^{[18]}$ and covalent organic frameworks (COFs) ${ }^{[19]}$ have been reported with linkers containing a pyrene moiety.

Although pyrene is a symmetric molecule, its 10 peripheral reactive sites ( $\mathrm{C}-\mathrm{H}$ bonds) can be divided into three chemically inequivalent classes, depending on different electronic and steric features that correspond to different chemical reactivities. ${ }^{[20,21]}$ The 1-, 3-, 6-, and 8-positions of pyrene, having maximum contributions of the highest occupied molecular orbital (HOMO), are favourable for electrophilic aromatic substitution. ${ }^{[22]}$ Thus, the $\pi$-orbitals of substituents at these positions participate very efficiently with pyrene's HOMO and LUMO. The 2- and 7-positions of pyrene lie on a nodal plane in both the HOMO and LUMO, but are the least hindered sites. Therefore, sterically controlled reactions such as Ir-catalysed direct borylation, are favourable at these sites. ${ }^{[23]}$ Moreover, the substituents at the 2- or 2,7-positions do not interact with the HOMO/LUMO, but can interact strongly with the HOMO-1 and LUMO +1 of pyrene that have non-zero contributions at these positions. ${ }^{[21,24-26]}$ Consequently, the photophysical properties of pyrene derivatives with substituents at the 2 - or 2,7 - positions differ significantly from those with substituents at the 1-, 3-, 6-, and 8-positions. This is termed the "site-effect". ${ }^{[20,24]}$ It has also been demonstrated by us and others that judicious incorporation of a combination of strong $\pi$-donors/acceptors at the 2 or 2,7-positions of pyrene switches the order of its $\mathrm{HOMO} /$ HOMO-1 and LUMO/LUMO + 1, respectively, which in turn greatly influences the photophysical and electrochemical properties. ${ }^{[21,25,26]}$ The localised double bond $(\mathrm{C}=\mathrm{C})$ character at the $4-5-, 9-$, and 10 -positions can be exploited to derivatise pyrene at its K-region. ${ }^{[27,28]}$ Thus, various compounds with substituents at the 4-, 5-, 9-, and 10-positions with interesting photophysical and electrochemical properties have recently been reported. ${ }^{[20,29,30]}$

Methylviologens, or viologens in general, are di- $N, N$ '-functionalised 4,4'-bipyridyl salts, which have been extensively investigated in the past, and have attracted more attention recently. ${ }^{[31,32]}$ The possibility of having three stable redox states, efficient electron accepting capabilities (oxidizing agent), and tunability of the substituents at the nitrogen make them an important class of compounds. The most common viologen compound, $N, N^{\prime}$-dimethyl-4,4'-dipyridyl dichloride ( $\mathrm{MV}^{2+}$ ), commonly known as „paraquat", can undergo two reversible reductions, forming a radical-cation ( $\mathrm{MV}^{\mathrm{C}+}$ ) and a neutral species (MV). Consequently, the colour changes from colourless $\left(\mathrm{MV}^{2+}\right)$ to blue-violet ( $\mathrm{MV}^{\mathrm{C}}$ ), and finally to yellow-brown (MV) during its redox processes. ${ }^{[31]}$ Such compounds have wide applications in electrochromic ${ }^{[32-33]}$ and photochromic materials. ${ }^{[34]}$ Moreover, viologen-based materials have been explored over the years for their application in molecular shuttles and switches or machines. ${ }^{[35]}$ Recently, viologen-based compounds have gained much importance for their use in organic „green batteries". ${ }^{[36]}$ It should be noted that paraquat was one of the most widely used herbicides for decades; however, it has been banned in several countries due to its acute toxicity. ${ }^{[37,38]}$

Depending on the substituents attached at the nitrogen atoms, and the counteranions, the colour of viologen compounds can change dramatically. Various counteranions have been utilised, halides being the most common and, recently, non-coordinating anions such as triflate and hexafluorophosphate have been employed. The counteranions also play an important role in the stability and solubility of viologens. ${ }^{[39]}$ Significant efforts have been devoted to fine tune the electronic and photophysical properties of viologens by introducing further functionalisation such as additional bridging or incorporating a hetero-atom in the molecule. A few examples are sulfide-bridged viologens, ${ }^{[40]}$ phosphole-bridged "phosphaviologen" ${ }^{[41]}$ germanium-bridged viologen, ${ }^{[42]}$ thiophene-based, ${ }^{[43]}$ and $\pi$-extended viologen compounds. ${ }^{[44]}$

Recently, we have reported the syntheses and photophysical properties of several mono- and bis-pyridyl-pyrene compounds substituted at various positions of pyrene. ${ }^{[20]}$ It has been demonstrated that the protonation of pyridyl-pyrene compounds shifts the emission bathochromically. ${ }^{[45]}$ The protonation of other fluorophores containing a conjugated pyridyl moiety has also been demonstrated to exhibit bathochromically shifted emissions. ${ }^{[46]}$ Thus, polyaromatic hydrocarbon (PAH)-based viologen compounds or compounds containing pyridinium moieties have attracted attention recently due to their interesting photophysical properties which have various potential applications. ${ }^{[47]}$ Konishi and co-workers reported a pyrene-based viologen compound where the extended $\pi$-conjugation was elongated by incorporating an ethenyl group between the pyrene and pyridinium moieties. This acceptor- $\pi$-acceptor type dye exhibited an emission in the 'tissue transparent window' (6501100 nm ) and a two-photon absorption band around 1000 nm with a large cross-section, and thus performed better than analogous PAH dyes (e.g., with naphthalene and anthracene) in three-dimensional (3D) imaging of mitochondria in living cells. ${ }^{[48]}$ The photophysical properties of analogous 1,6-bis( $N$ -methyl-3'-pyridinium)pyrene, 1,6-bis(N-methyl-4'-pyridinium)-
pyrene and 1,3,6,8-tetrakis( $N$-methyl-4'-pyridinium)pyrene have been reported along with their energy transfer processes on a clay surface in aqueous media. ${ }^{[49]}$ However, detailed investigation of the photophysical properties along with electrochemical and spectroelectrochemical studies, and the interactions with biorelevant targets (e.g., DNA) of pyridinium-pyrene compounds have not been documented so far. Herein, we report the structures, photophysical and electrochemical properties of mono- and bis-( $N$-methylpyridinium)-pyrene and analogous protonated compounds (Scheme 1), as well as preliminary studies on DNA binding of cationic compounds $\mathbf{1 M}$ and $\mathbf{2 M}$.

## Results and Discussion

## Synthesis

The neutral compounds 1 and 2 were synthesised via the Suzuki-Miyaura coupling reaction of 2-(Bpin)pyrene and 2,7bis(Bpin)pyrene, respectively, with 4-bromopyridine, as we have reported previously; ${ }^{[20]}$ however, $\left[\mathrm{Pd}(\mathrm{dppf}) \mathrm{Cl}_{2}\right]$ was used as the catalyst precursor instead of $\left[\mathrm{Pd}\left(\mathrm{PPh}_{3}\right)_{4}\right]$. The methylation reactions yielding compounds $\mathbf{1 M}$ and $\mathbf{2 M}$ were carried out with MeOTf in dry toluene under an argon atmosphere, and single crystals of $\mathbf{1 M}$ and $\mathbf{2 M}$ were obtained by slow evaporation of concentrated solutions of the compounds in MeCN. The salts $\mathbf{1 H}$ and $\mathbf{2 H}$ were obtained by reacting 1 and 2 , respectively, with $\mathrm{HNO}_{3}$ in MeCN , and their single crystals were obtained by slow evaporation from MeCN and DMF solutions, respectively. All compounds were characterised by various spectroscopic methods. In the ${ }^{1} \mathrm{H}$ NMR spectrum, the signals for the protons of both the pyridyl and pyrene moieties were found to be down-field shifted, as expected, upon protonation or methylation. For example, the signals for the pyridyl protons appear at 8.75 and 8.57 ppm for $\mathbf{2 M}$ (Figure S7, Supporting Information), and at 9.04 and 8.57 ppm for $\mathbf{2 H}$ (Figure S11, Supporting Information), and the signals for the pyrene protons appear at 8.85 and 8.40 ppm for $\mathbf{2 M}$, and at 9.0 and 8.43 for $\mathbf{2 H}$, which are significantly down-field shifted compared to the signals observed for 2 at 8.79 and 7.81 ppm (pyridyl protons), and 8.45 and 8.19 ppm (pyrene protons) (Figure S3, Supporting Information). Similar down-field shifts were also observed for
the signals of 1 M and 1 H compared to that of 1 (Figures S 1 , S5, S9, Supporting Information).

## Description of the crystal structures

The crystal structures of compounds $\mathbf{1 M}, \mathbf{2 M}, \mathbf{1 H}$ and $\mathbf{2 H}$ were analysed by single-crystal X-ray diffraction. Among these, only the structure of $\mathbf{1 M}$ is non-centrosymmetric, and it was solved in the monoclinic space group $P 2_{1}$. The unit cell of 1 M contains two symmetry-independent cations and two symmetryindependent anions. The structures of the two cations differ mainly in the torsion angles between the planes of the pyridinium cations and the pyrene moieties, which are $18.2^{\circ}$ and $22.7^{\circ}$, respectively (Figure 1 (a)). The cations stack in a parallel head-to-tail (anti) fashion and show $\pi$-stacking (cation- $\pi$ ) interactions between the pyridinium cation and pyrene moieties (interplanar separations (d) are in the range 3.34-3.41 Å). In addition, $\mathrm{C}-\mathrm{H} \cdots \pi$ interactions between these moieties with the methyl-H atom are observed. The triflate anions exert various weak interactions including $\mathrm{C}-\mathrm{H} \cdots \mathrm{O}$ and $\mathrm{C}-\mathrm{H} \cdots \mathrm{F}$ interactions with the cations (Figure S13, Supporting Information).

The salt 1 H crystallizes in the monoclinic space group $P 2_{1} / c$ and, in contrast to $\mathbf{1 M}$, only one symmetry-independent cation $\left(1-\mathrm{H}^{+}\right)$and nitrate ion are present in the unit cell. However, the molecular packing in 1 H is similar to that in 1 M , with the cations displaying parallel stacking in a head-to-tail fashion. The cation- $\pi$ interactions between the pyridinium cation and pyrene moieties ( $d=3.30-3.43 \AA$ ) are in a similar range to those in 1M. However, in contrast to 1 M , the cation remains almost planar in $\mathbf{1 H}$, as the torsion angle between pyridinium and pyrene moieties is only $2.5^{\circ}$. $\mathrm{N}-\mathrm{H} \cdots \mathrm{O}$ and weak $\mathrm{C}-\mathrm{H} \cdots \mathrm{O}$ hydrogen bonding interactions are observed between the nitrate anion and the pyridinium cation (Figure 1 (c)).

The crystal structure of $\mathbf{2 M}$ was refined in the monoclinic, centrosymmetric space group $C 2 / m$. The cation has $2 / m\left(C_{2 h}\right)$ symmetry, while the triflate anion is disordered and has mirrorsymmetry $\left(C_{s}\right)$. The dications arrange in infinite, offset parallel stacks with strong $\pi$-stacking (cation $-\pi$ ) interactions between the pyridinium cation and pyrene moieties ( $d=3.35-3.37 \AA$ ) (Figure 1 (b)). Similar to the monocation 1H, the dication of $\mathbf{2 M}$ is almost planar with a torsion angle between the pyridinium


Scheme 1. Compounds discussed in this study. Atom numbering of pyrene is also shown (left).


Figure 1. Molecular packing of (a) $\mathbf{1 M}$, (b) $\mathbf{2 M}$, (c) $\mathbf{1 H}$, and (d) $\mathbf{2 H}$ in the solid state at 100 K . Only selected hydrogen atoms are shown for clarity. The numbers indicate the torsion angles between pyridyl and pyrene planes.
and pyrene planes of $0.8^{\circ}$. The triflate anions interact with the dication via various weak $\mathrm{C}-\mathrm{F} \cdots \pi$ and $\mathrm{C}-\mathrm{H} \cdots \mathrm{O}$ interactions (Figure S15, Supporting Information).
The salt 2 H crystallizes in the monoclinic space group $P 2_{1} / \mathrm{c}$ and the dication has inversion symmetry. The dications form offset parallel stacks with $\pi$-stacking interactions between the pyridinium cation and pyrene moieties ( $d=3.37$ and $3.39 \AA$ ) (Figure 1 (d)), similar to the arrangement in $\mathbf{2 M}$. As in $\mathbf{1 H}$ and $\mathbf{2 M}$, the molecule is almost planar, as the torsion angle between the pyridinium and pyrene planes is $1.9^{\circ}$. The nitrate ions are involved in $\mathrm{N}-\mathrm{H} \cdots \mathrm{O}$ and weak $\mathrm{C}-\mathrm{H} \cdots \mathrm{O}$ hydrogen bonding interactions with the dication (Figure S23, Supporting Information).
In summary, the mono- or dications are almost planar in all compounds in the solid state (torsion angles: $0.8-2.5^{\circ}$ ), except for $\mathbf{1 M}$ (torsion angles: 18.2 and $22.7^{\circ}$ ). It is important to note that the observed torsion angles in the corresponding neutral compounds $1\left(38.3^{\circ}\right)$ and $2\left(12.1^{\circ}\right.$ for $\alpha$-form and $38.6^{\circ}$ for $\beta$ form $)^{[20]}$ are larger than those in the cationic compounds, which are strongly diminished by protonation or methylation (except for 1M). The smaller torsion angles may also be a
result of the large degree of $\pi$-stacking interactions in the near-planar cationic compounds. These show strong $\pi$-stacking interactions between pyridinium and pyrene moieties as a result of protonation, which determine the parallel head-to-tail stacking in $\mathbf{1 M}$ and $\mathbf{1 H}$ and the offset parallel stacking in $\mathbf{2 M}$ and 2 H . A similar $\pi$-stacking interaction between pyridyl and pyrene moieties is not observed in the neutral compounds 1 and 2 , in which $\pi$-stacking is only present between the pyrene moieties (1) or both the pyrene and the pyridyl moieties (2). ${ }^{[20]}$ In 1M, 1H, 2M and $\mathbf{2 H}$, cation stacks weakly interact with the anions, which are located in between the stacks. A Hirshfeld surface analysis reveals the area and relative strengths of the intermolecular interactions and is presented in the Supporting Information (Figures S14 (1M), S18 (2M), S21 (1H), and S24 $(\mathbf{2 H})$ ). The powder diffraction patterns of all samples were measured in order to examine the phase purity of the bulk samples. The observed powder diffraction patterns of all compounds agree very well with the respective simulated ones (Figures S25 and S26 in the Supporting Information). Small systematic shifts of the reflection positions can only be observed in $\mathbf{1 M}$ and $\mathbf{2 M}$, which may be due to the sample preparation.

## Photophysical properties

Detailed photophysical properties of all compounds are depicted in Figure 2. Absorption and emission were measured for 1M in THF, MeOH, MeCN, DMSO, and water, for 2M in MeOH, MeCN, DMSO, and water, as well as in the solid state for both compounds. For $\mathbf{1 H}$ and $\mathbf{2 H}$, the absorption and emission spectra in MeCN are presented. The key photophysical data are summarised in Table 1. The neutral compounds 1 and 2 were measured in MeCN and in the solid state for comparison with the cations.

The absorption spectra of 1 and 2 in MeCN are very similar and demonstrate "pyrene-like" absorptions in that the $\mathrm{S}_{1} \leftarrow \mathrm{~S}_{0}$ absorption ( $370-420 \mathrm{~nm}$ ) is very weak with extinction coefficients $\varepsilon$ below $1400 \mathrm{~m}^{-1} \mathrm{~cm}^{-1}$ (Figure 2(a)). In general, the $\mathrm{S}_{1} \leftarrow \mathrm{~S}_{0}$ absorption of pyrene and 2 - and 2,7 -substituted pyrene derivatives is attributed to a short-axis-polarised and transition-dipole-forbidden transition. ${ }^{[24]}$ The $\mathrm{S}_{2} \leftarrow \mathrm{~S}_{0}$ transition was observed at 337 nm with vibrational progressions of 1382 and $1411 \mathrm{~cm}^{-1}$ for 1 , and at 339 nm with vibrational progressions of 1365 and $1187 \mathrm{~cm}^{-1}$ for 2 , which are much more intense ( $\varepsilon=37000$ and $35000 \mathrm{~m}^{-1} \mathrm{~cm}^{-1}$ ) than the $\mathrm{S}_{1} \leftarrow \mathrm{~S}_{0}$ absorption, as also observed for unsubstituted pyrene ( $\lambda_{\text {abs }}=334 \mathrm{~nm}$ for $\mathrm{S}_{2} \leftarrow \mathrm{~S}_{0}$ absorption, $\varepsilon=55000 \mathrm{~m}^{-1} \mathrm{~cm}^{-1}$ ). ${ }^{[50]}$ The $\mathrm{S}_{1} \leftarrow \mathrm{~S}_{0}$ absorption of 2 is slightly bathochromically shifted compared to $\mathbf{1}$, by $175 \mathrm{~cm}^{-1}$. The absorption spectra of the cationic compounds $\mathbf{1 M}$ and $\mathbf{2 M}$ in MeCN are also very similar to that of pyrene, but show a stronger absorbance compared to 1 and 2 (Figure 2 (a)). Thus, their $\mathrm{S}_{1} \leftarrow \mathrm{~S}_{0}$ absorption is still a weak band with $\varepsilon=3000-4300 \mathrm{~m}^{-1} \mathrm{~cm}^{-1}$; however, this band is bathochromically shifted and broad, covering ca. 100 nm (375-475 nm). The corresponding $\mathrm{S}_{2} \leftarrow \mathrm{~S}_{0}$ transitions of both $\mathbf{1 M}$ and $\mathbf{2 M}$ are less influenced as they are also observed at 337 nm in MeCN with vibrational progressions of 1872 and $1576 \mathrm{~cm}^{-1}$, respectively. The absorption of 1 H appears to be the combination of the


Figure 2. (a) Absorption, and (b) emission spectra of 1, 2, 1M, 2M, 1H, and 2H in MeCN ( $\lambda_{\mathrm{ex}}=337 \mathrm{~nm}$ ); (c) absorption (solid lines), and emission (dashed lines) spectra of $\mathbf{1 M}$ in THF, MeOH, MeCN, DMSO, water and in the solid state; and (d) absorption (solid lines) and emission (dashed lines) spectra of $\mathbf{2 M}$ in $\mathbf{M e O H}$, MeCN, DMSO, water and in the solid state. Very weak emission for $\mathbf{1 M}$ at $405-430 \mathrm{~nm}$ and for $\mathbf{2 M}$ at $420-440 \mathrm{~nm}$ might be due to the presence of traces of neutral compounds 1 and $\mathbf{2}$, respectively. All measurements were carried out in deoxygenated solvents.
absorbance of 1 and $\mathbf{1 M}$. All of the absorption bands in 1 and $\mathbf{1 M}$ are present in 1 H with proportionate intensity. It seems that 1 H undergoes dissociation in MeCN solution which is even more evident in the emission spectrum (Figure 2 (b)). The absorption of 2 H resembles that of $\mathbf{2 M}$ in that its $\mathrm{S}_{1} \leftarrow \mathrm{~S}_{0}$ absorption ( $\lambda_{\text {abs }}=428 \mathrm{~nm}, \varepsilon=1530 \mathrm{~m}^{-1} \mathrm{~cm}^{-1}$ ) is also transitiondipole forbidden and broad. However, this band in $\mathbf{2 M}$ is further bathochromically shifted and more allowed compared to that in 2 H . The $\mathrm{S}_{2} \leftarrow \mathrm{~S}_{0}$ transition of 2 H occurs at 341 nm ( $\varepsilon=$ $30500 \mathrm{~m}^{-1} \mathrm{~cm}^{-1}$ ) with a vibrational progression of $1827 \mathrm{~cm}^{-1}$.
Furthermore, the absorption spectra of 1M and 2M were measured in various solvents to investigate possible solvent effects (see Table 1 and Figure 2 ( $c$ and d)). A very small shift (from 337 nm to 340 nm ) was observed for 2M in DMSO, which shows a bathochromic shift of only $430 \mathrm{~cm}^{-1}$ compared to its absorption in MeOH or MeCN. Thus, no significant solvatochromism was observed in the absorption spectra, which was further confirmed by theoretical calculations (vide infra).

Although these compounds exhibited similar absorption characteristics, their emission characteristics are quite different and display interesting features. The emission of the neutral compounds 1 and 2 are pyrene-like, which agrees with their
absorption properties. Thus, the emission of 1 at $\lambda_{\mathrm{em}}=405 \mathrm{~nm}$ in MeCN shows a vibrational progression of $1250 \mathrm{~cm}^{-1}$. The emission of 2 at $\lambda_{\mathrm{em}}=418 \mathrm{~nm}$ is slightly bathochromically shifted compared to that of 1 and also possesses a vibrational progression of $1300 \mathrm{~cm}^{-1}$. On the other hand, the emissions of the cationic compounds $\mathbf{1 M}$ and $\mathbf{2 M}$ are different than that of pyrene. Their emission consists of a single broad band, which is strongly bathochromically shifted compared to 1 and 2 , respectively. The emission of 1 M occurs at $\lambda_{\mathrm{em}}=599 \mathrm{~nm}$ in MeCN with an apparent Stokes shift of $7200 \mathrm{~cm}^{-1}$, which is the most bathochromically shifted in this series. It can be noted that this apparent Stokes shift is very large for 2- or 2,7-substituted pyrenes, ${ }^{[20,21,24]}$ which indicates a large change of the dipole moment in the excited state. Considering the energy of maximum emission intensity of $1\left(\lambda_{\mathrm{em}}=405 \mathrm{~nm}\right)$ and $1 \mathrm{M}\left(\lambda_{\mathrm{em}}=\right.$ 599 nm ), the emission of 1 M is bathochromically shifted by $8000 \mathrm{~cm}^{-1}$, which is a large effect that is introduced by one methyl group. The emission of 2 M occurs at $\lambda_{\mathrm{em}}=538 \mathrm{~nm}$ with an apparent Stokes shift of $4300 \mathrm{~cm}^{-1}$. Considering the energy of maximum emission intensity of $2\left(\lambda_{\text {em }}=418 \mathrm{~nm}\right)$ and 2 M , ( $\lambda_{\text {em }}=538 \mathrm{~nm}$ ), the emission of $\mathbf{2 M}$ is bathochromically shifted by $5300 \mathrm{~cm}^{-1}$. Interestingly, the emission of $\mathbf{1 M}$ is even more

Table 1. Photophysical data for $\mathbf{1}, \mathbf{2}, \mathbf{1 M}, \mathbf{2 M}, \mathbf{1 H}$, and $\mathbf{2 H}$ both in solution and in the solid state.

| Compound | Solvent or Solid | $\begin{aligned} & \lambda_{\text {abs }} \\ & {[\mathrm{nm}]} \end{aligned}$ | $\begin{aligned} & \varepsilon \\ & {\left[\mathrm{m}^{-1} \mathrm{~cm}^{-1}\right]} \end{aligned}$ | $\begin{aligned} & \lambda_{\mathrm{em}} \\ & {[\mathrm{~nm}]} \end{aligned}$ | Apparent Stokes shift ${ }^{[a]}$ $\left[\mathrm{cm}^{-1}\right]$ | $\Phi$ | $\begin{aligned} & \tau \\ & {[\mathrm{ns}]} \end{aligned}$ | $\tau_{\text {avg }}{ }^{[b]}[\mathrm{ns}]$ | $\begin{aligned} & \tau_{0}^{[c]} \\ & {[\mathrm{ns}]} \end{aligned}$ | $k_{r}$ $\left[10^{6} \mathrm{~s}^{-1}\right]$ | $k_{n r}$ $\left[10^{6} s^{-1}\right]$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | MeCN | 362 | 1100 | 451 | 2900 | 0.22 | 36.1 | - | 164 | 6.1 | 22 |
|  |  | 337 | 37000 | 427 |  |  |  |  |  |  |  |
|  |  | 322 | 21000 | 405 |  |  |  |  |  |  |  |
|  |  | 308 |  |  |  |  |  |  |  |  |  |
|  | solid | - | - | 452 |  | 0.05 | 3.4 (32.6) | 6.9 |  |  |  |
|  |  |  |  | 430 |  |  | 8.9 (37.6) |  |  |  |  |
|  |  |  |  |  |  |  | 23.5 (29.8) |  |  |  |  |
| 1M | THF | 418 |  | 586 | 6800 | 0.08 | 16 | - | 200 | 5 | 58 |
|  |  | 337 |  |  |  |  |  |  |  |  |  |
|  |  | 319 |  |  |  |  |  |  |  |  |  |
|  | MeOH | 418 |  | 593 | 7000 | 0.05 | 12.6 | - | 252 | 4 | 75 |
|  |  | 337 |  |  |  |  |  |  |  |  |  |
|  |  | 318 |  |  |  |  |  |  |  |  |  |
|  | MeCN | 418 | 2900 | 599 | 7200 | 0.10 | 25 | - | 250 | 4 | 36 |
|  |  | 337 | 45000 |  |  |  |  |  |  |  |  |
|  |  | 317 | 68500 |  |  |  |  |  |  |  |  |
|  | DMSO | 416 |  | 592 | 7100 | 0.11 | 22.6 | - | 205 | 5 | 39 |
|  |  | 342 |  |  |  |  |  |  |  |  |  |
|  |  | 321 |  |  |  |  |  |  |  |  |  |
|  | water | 425 |  | 608 | 7080 | 0.01 | 3 | - | 300 | 3.3 | 330 |
|  |  | 340 |  |  |  |  |  |  |  |  |  |
|  |  | 318 |  |  |  |  |  |  |  |  |  |
|  |  | 309 |  |  |  |  |  |  |  |  |  |
|  | solid | - | - | 535 | - | 0.08 | 18.6 (60.7) | 25.5 |  |  |  |
|  |  |  |  |  |  |  | 59.6 (39.3) |  |  |  |  |
| 1H | MeCN | 421 | [d] | 590 | 6800 | - | - | - | - | - | - |
|  |  | 337 |  | 426 |  |  |  |  |  |  |  |
|  |  | 319 |  | 404 |  |  |  |  |  |  |  |
|  | solid | - | - | 545 |  | 0.05 | 8.4 (69.6) | 10.3 |  |  |  |
|  |  |  |  |  |  |  | 23.2 (30.4) |  |  |  |  |
| 2 | MeCN | 363 | 1400 | 469 | 3600 | 0.41 | 56.6 | - | 138 | 7.2 | 10 |
|  |  | 339 | 35000 | 442 |  |  |  |  |  |  |  |
|  |  | 324 |  | 418 |  |  |  |  |  |  |  |
|  |  | 312 |  |  |  |  |  |  |  |  |  |
|  | solid | - | - | 482 | - | 0.02 | 1.3 (24.8) | 3.5 |  |  |  |
|  |  |  |  | 453 |  |  | 5.6 (34.3) |  |  |  |  |
|  |  |  |  |  |  |  | 13.8 (40.9) |  |  |  |  |
| 2M | MeOH | 438 |  | 536 | 4100 | 0.09 | 15 | - | 167 | 6 | 61 |
|  |  | 337 |  |  |  |  |  |  |  |  |  |
|  |  | 320 |  |  |  |  |  |  |  |  |  |
|  | MeCN | 437 | 4300 | 538 | 4300 | 0.19 | 27 | - | 142 | 7 | 30 |
|  |  | 337 | 80000 |  |  |  |  |  |  |  |  |
|  |  | 320 | 120000 |  |  |  |  |  |  |  |  |
|  | DMSO | 440 |  | 537 | 4100 | 0.22 | 22.2 | - | 101 | 10 | 35 |
|  |  | 342 |  |  |  |  |  |  |  |  |  |
|  |  | 325 |  |  |  |  |  |  |  |  |  |
|  | water | 443 |  | 540 | 4050 | 0.12 | 15.5 | - | 129 | 7.7 | 57 |
|  |  | 339 |  |  |  |  |  |  |  |  |  |
|  |  | 321 |  |  |  |  |  |  |  |  |  |
|  | solid | - | - | 612 | - | 0.03 | 3.3 (79.5) | 3.9 |  |  |  |
|  |  |  |  | 560 |  |  | $13.4(20.5)$ |  |  |  |  |
|  |  |  |  | 525 |  |  |  |  |  |  |  |
| 2 H | MeCN | 428 | 1530 | 555 | 5300 | 0.07 | 15.6 | - | 223 | 6.4 | 60 |
|  |  | 341 | 30500 |  |  |  |  |  |  |  |  |
|  |  | 321 | 50000 |  |  |  |  |  |  |  |  |
|  | solid | - | - | 642 | - | 0.0 | 1.2 (69) | 4 | - | - | - |
|  |  |  |  |  |  |  | 4.8 (19.9) |  |  |  |  |
|  |  |  |  |  |  |  | 19.9 (11.1) |  |  |  |  |

[a] The Stokes shift is defined as the energy difference between the 0-0 transitions of the absorption and the emission. However, as the $S_{1} \leftarrow S_{0}$ transitions in substituted pyrene derivatives are usually broad with an undefined vibrational pattern, the band maxima of the absorption and emission have been used to determine the Stokes shift; consequently, we use the term „apparent Stokes shift". [b] For multi-exponential decays, the experimental average lifetime is given by $\tau_{\mathrm{avg}}=\Sigma \tau_{\mathrm{n}} B_{\mathrm{n}} / \Sigma B_{\mathrm{n}}$, where $B_{\mathrm{n}}$ are pre-exponential factors of the respective lifetime $\tau_{\mathrm{n}}$. [c] The pure radiative lifetime, $\tau_{0}=\tau / \Phi ; k_{\mathrm{r}}=1 / \tau_{0} ; k_{\mathrm{nr}}=$ $(1-\Phi) / \tau$. [d] Partially dissociates in solution, see text.
bathochromically shifted than $2 \mathbf{M}$, which is an effect of a larger stabilization of the excited state in $\mathbf{1 M}$ due to the larger dipole moment compared to $\mathbf{2 M}$, which has a symmetrical charge distribution. It can be noted that the emission of $\mathbf{2 M}$ is hypsochromically shifted compared to analogous 1,6-bis( $N$ -methyl-3'-pyridinium) pyrene dichloride $\left(\lambda_{\mathrm{em}}=591 \mathrm{~nm}\right)^{[49 \mathrm{a}]}$ and bathochromically shifted compared to [1,3,6,8-tetrakis( $N$ -methyl-4-pyridinium)pyrene][PF $\left.{ }_{6}\right]_{4}\left(\lambda_{\mathrm{em}}=485 \mathrm{~nm}\right) .{ }^{[49 \mathrm{~b}]}$

To investigate the possible charge-transfer (CT) nature in derivatives $\mathbf{1 M}$ and $\mathbf{2 M}$, we conducted emission measurements in various solvents (see Table 1 and Figures 2 (c) and (d)). Indeed, the solvatochromic effect is very small. Very little change of $\lambda_{\mathrm{em}}$ with the polarity of the solvent was observed for $\mathbf{1 M}$. The difference of $\lambda_{\text {em }}$ between THF ( 586 nm ) and water ( 608 nm ) is 22 nm , which is only $617 \mathrm{~cm}^{-1}$. However, it can be concluded that $\mathbf{1 M}$ has greater $C T$ nature than $\mathbf{2 M}$, which is also in line with the lower extinction coefficients of the first transition. Interestingly, the solid-state emission of $1 \mathbf{M}\left(\lambda_{\text {em(solid })}=535 \mathrm{~nm}\right)$ is considerably hypsochromically shifted compared to its solution emission (Figure 2(c)). On the other hand, the solid-state emission of $\mathbf{2 M}\left(\lambda_{\text {em(solid })}=612 \mathrm{~nm}, 560 \mathrm{~nm}, 525 \mathrm{~nm}\right)$ is comparable to its solution-phase features (Figure 2(d)).

The emission of 1 H possesses two bands, one band at $\lambda_{\mathrm{em}}=$ 405 nm that shows a fine structure, which is identical to compound 1, and also a broad band at 590 nm , which is similar to that of $1 \mathbf{M}$. Thus, it is evident that $\mathbf{1 H}$ undergoes dissociation and maintains an equilibrium between the neutral and protonated forms. The emission of 2 H occurs at $\lambda_{\text {em }}=555 \mathrm{~nm}$ with an apparent Stokes shift of $5300 \mathrm{~cm}^{-1}$. Considering the energy of maximum emission intensity of $2\left(\lambda_{\mathrm{em}}=418 \mathrm{~nm}\right)$ and $2 \mathrm{M}\left(\lambda_{\mathrm{em}}=\right.$ 538 nm ), the emission of $\mathbf{2 H}$ is bathochromically shifted by $5900 \mathrm{~cm}^{-1}$ compared to that of 2 , and by $570 \mathrm{~cm}^{-1}$ compared to that of $\mathbf{2 M}$. Furthermore, we again observe that the monosubstituted derivative $\mathbf{1 H}$ has stronger CT character than its disubstituted analogue $\mathbf{2 H}$, as its emission is bathochromically shifted by $1070 \mathrm{~cm}^{-1}$.

The lifetime and quantum yields were measured in deoxygenated solutions. Monoexponential decays were obtained for the neutral and cationic compounds, and it is found that the neutral compounds possess a longer-lived excited state compared to their cationic analogues (see Table 1). Furthermore, the non-radiative decay rates of the latter ones are significantly increased $\left(k_{\mathrm{nr}}=5.8-6.1 \times 10^{7} \mathrm{~s}^{-1}\right)$, which is a result of their larger apparent Stokes shifts, fully in line with the energy gap law. ${ }^{[51]}$ Similarly, the quantum yields of the neutral compounds are much higher compared to their cationic analogues. All of the compounds possess very low quantum yields in the solid state, which might be due to the absence of an effective suppression of the nonradiative decay, in the presence of strong $\pi$-stacking interactions. This would also explain their shorter lifetimes in the solid state.

We found in our previous work that the solid state emissions of 1 and 2 strongly depend on crystallinity and particle size. ${ }^{[20]}$ In addition, two polymorphic phases of 2 ( $\boldsymbol{\alpha}$ and $\boldsymbol{\beta}$ ) have also been reported. While $2 \boldsymbol{\beta}$ displays a large $\pi-\pi$ overlap area, these interactions are much weaker in $\mathbf{2 \alpha}$. It is feasible to assume that the solid-state emission is strongly influenced by
the local environment and by the proportion of this local environment to the overall sample (Figures S29 and S30 in the Supporting Information). The overall solid-state emission spectrum is therefore caused by an overlap of the emission spectra of the different polymorphs. However, these differences disappear when the samples are dissolved in solution.

## Electrochemistry

Cyclic voltammetry studies were performed on the neutral and methylated compounds to investigate their electrochemical properties. The voltammograms were measured at a scan rate of $250 \mathrm{mV} \mathrm{s}^{-1}$ with $0.1 \mathrm{~m}\left[n \mathrm{Bu}_{4} \mathrm{~N}\right]\left[\mathrm{PF}_{6}\right]$ relative to the $\mathrm{Fc} / \mathrm{Fc}^{+}$ couple. The first scan and then ten consecutive potential sweeps were measured for 1 in $\mathrm{CH}_{2} \mathrm{Cl}_{2}$ as presented in Figure 3 (a). For the very first scan, the formal potentials are $E_{\mathrm{pa}}=$ +0.87 and +1.02 V , in the cathodic direction; however, no reduction event was observed. The peak currents for the redox processes increased steadily with each successive scan. This was an indication that some chemical reaction (e.g., dimerisation or polymerisation) might take place and there was deposition on the electrode surface. It should be noted that electrochemical polymerisation of pyrene compounds has been reported previously. ${ }^{[52]}$ Electropolymerisation via anodic oxidation of unprotected bis(triarylamine)s leading to polymers with benzidine units is also known. ${ }^{[53]}$ In the case of 2, ten consecutive potential sweeps are shown in Figure 3 (b). The formal potential $E_{\mathrm{pa}}=+1.30 \mathrm{~V}$ was observed in the anodic direction; no reduction event was observed on the first scan. However, a reduction event can be observed only from the second scan. The reduction potential for the successive scans shifts to more negative values with a steady increase in the peak current ( $E_{\mathrm{pc}}=$ -1.60 V after 10 scans). This also indicates some chemical reaction or deposition on the electrode surface, as was observed for $1 .{ }^{[52,53]}$ On the other hand, $\mathbf{1 M}$ shows an irreversible reduction (Figure 3 (c)), which takes place at $E_{\mathrm{pc}}=-1.53 \mathrm{~V}$, and the corresponding oxidation process completes at -0.94 V . In the case of $\mathbf{2 M}$, a reversible reduction was observed with $E_{1 / 2}=$ -1.47 V , which can be attributed either to a two-electron reduction or two very closely spaced consecutive one-electron reductions leading to the formation of the neutral species. No further reduction processes for either $\mathbf{1 M}$ or $\mathbf{2 M}$ were observed within the accessible potential window for the solvent (even in MeCN ). The observed formal potentials for these two compounds are comparable; however, only $\mathbf{2 M}$ shows a reversible reduction. It can be noted that these observed reduction potentials for $\mathbf{1 M}$ and $\mathbf{2 M}$ occur at more negative values compared to the first one-electron reduction of paraquat $\left(\mathrm{MV}^{2+}\right)$, which occurs at -1.076 V vs. $\mathrm{Fc}^{+} / \mathrm{Fc}$ couple (the second oneelectron reduction of paraquat occurs at -1.51 V vs. $\mathrm{Fc}^{+} / \mathrm{Fc}$ couple in $\mathrm{MeCN} /\left[\mathrm{Et}_{4} \mathrm{~N}\right]\left[\mathrm{PF}_{6}\right]$ ]. ${ }^{[31 \mathrm{c}, \mathrm{d}, 41 \mathrm{a}, 54]}$ Moreover, it can be noted that the extended methyl viologen 1,4-bis(4'-N-methylpyridinium)benzene with $\left[\mathrm{BF}_{4}\right]^{-}$counteranions, in which the pyridinium groups are separated by a phenylene moiety, showed a single two-electron reduction signal at -0.91 V vs. SCE (i.e., -1.29 V vs. $\mathrm{Fc}^{+} / \mathrm{Fc}$ couple) ${ }^{[54]}$ with $0.1 \mathrm{~m} \mathrm{MeCN} /\left[\mathrm{Et}_{4} \mathrm{~N}\right]\left[\mathrm{ClO}_{4}\right]$. ${ }^{[31 \mathrm{e}]}$ Another $\pi$-extended viologen, i.e., 4,4'-bis(4-pyridinium-n-oc-


Figure 3. Cyclic voltammograms measured at $250 \mathrm{mV} \mathrm{s}^{-1}$ with $0.1 \mathrm{~m}\left[n B u_{4} N\right]\left[P F_{6}\right]$ relative to the $\mathrm{Fc} / \mathrm{Fc}{ }^{+}$couple: (a) the first scan (top) and ten consecutive potential sweeps (bottom) of $\mathbf{1}$ in $\mathrm{CH}_{2} \mathrm{Cl}_{2}$; (b) ten consecutive potential sweeps of $\mathbf{2}$ in $\mathrm{CH}_{2} \mathrm{Cl}_{2}$; (c) $\mathbf{1 M}$ in $\mathbf{M e C N}$; (d) $\mathbf{2 M}$ in MeCN .
tyl)biphenyl with $\left[\mathrm{PF}_{6}\right]^{-}$counteranions, showed one reversible two-electron reduction at -1.48 V vs . $\mathrm{Fc} / \mathrm{Fc}{ }^{+}$in THF , in a range comparable to that of $\mathbf{2 M}$. ${ }^{[44 a]}$ It can be further noted that the reduction of $\mathbf{2 M}$ occurs at a much lower potential compared to that of our previously reported 2,7-bis(BMes $)_{2}$ pyrene, an analogous acceptor-pyrene-acceptor system, which exhibits two reversible reduction events at -2.17 and -2.45 V vs. Fc/ $\mathrm{Fc}^{+}$in THF, indicating that the N -methylpyridinium group is a stronger acceptor than -BMes ${ }_{2}$. ${ }^{[25 b]}$

## Spectroelectrochemistry

UV/Vis/NIR spectroelectrochemical measurements were carried out in $\mathrm{CH}_{2} \mathrm{Cl}_{2}$ for the neutral compounds 1 and 2 , and in MeCN for the methylated compounds $\mathbf{1 M}$ and $\mathbf{2 M}$, with 0.1 m $\left[n B u_{4} N\right]\left[\mathrm{PF}_{6}\right]$ as the supporting electrolyte. Since reversible voltammograms were not observed for $\mathbf{1 , 2}$, and $\mathbf{1 M}$ in the cyclic voltammetry, we have presented their spectra in the Supporting Information (Figures S31-S33). Upon stepwise reduction of 2M, no absorption in the NIR region was observed; however, some changes in their UV/Vis absorption patterns were noticed (Figure 4). The absorption of reduced-2M is very broad in the UV region with an extended tail in the visible region. The weaker broad peak at $400-450 \mathrm{~nm}$ increased, while the peak at 320 nm decreased upon reduction, and an intense peak at 295 nm appeared. Such an observation is attributed to the fact


Figure 4. Spectroelectrochemical measurements of the stepwise reduction process of $\mathbf{2 M}$ with $0.1 \mathrm{~m}\left[n B u_{4} N\right]\left[\mathrm{PF}_{6}\right]$ in MeCN .
that $\mathbf{2 M}$ undergoes either a two-electron reduction or two very closely spaced consecutive one-electron reductions, leading to the formation of the neutral species, without detecting the radical monocation.

## Theoretical calculations

Density functional theory (DFT) and time-dependent density functional theory (TD-DFT) calculations were carried out to rationalise the observed spectroscopic and electrochemical properties. The structures of these compounds were optimised at the PBEO/LANL2DZP level in THF including GD3BJ empirical dispersion (see Theoretical Computations in the Experimental Section). Selected experimentally determined and theoretically optimised structural parameters are listed in Tables S2 and S3 (Supporting Information). Overall, the optimised structural parameters for $\mathbf{1}, \mathbf{2}, \mathbf{1 M}, \mathbf{2 M}, \mathbf{1 H}$, and $\mathbf{2 H}$ match well with the crystallographically determined parameters, except for the torsion angles for $\mathbf{1 H}$ and $\mathbf{2 H}$, which deviate (ca. $25-30^{\circ}$ ) probably due to crystal packing effects in the solid state. Despite rather strong $\pi$-stacking interactions observed between molecules in the cationic compounds (vide supra), periodic calculations on compound 2M, for example, (see Theoretical Computations in the Experimental Section) seem to indicate that this stacking hardly alters the electronic properties of each molecule. Indeed, the density of states (DOS) of 2M (Figure S34 in the Supporting Information) shows sharp peaks indicating weak electronic interactions between the dications in the solid.
HOMO-LUMO energy gaps are substantially larger for compounds 1 and 2 than for the cationic species ( 4.11 eV vs. ca. 3.15 eV , respectively), mainly due to some stabilization of the LUMO in the case of the latter ones. This suggests that the cationic compounds should absorb light at lower energies, as observed experimentally. This is supported by the maximum ab-
sorption and emission wavelengths which were also calculated at the TD-DFT PBE0-GD3BJ/LANL2DZP level in both MeCN and THF. No solvatochromic effect was observed in the calculations (see the comparison in Table S5 in the Supporting Information), which is in line with our experimental results. The main computed absorption wavelengths obtained in MeCN are given in Table 2 and compared with the experimentally observed values (see Table S5, in the Supporting Information, for values computed in THF). The corresponding simulated UV/Vis absorption spectra are shown in Figure S35 (Supporting Information). It can be noted that the calculated results agree better for the neutral compounds, with only a small deviation of a few nm , than for the methylated compounds, for which computed values are somewhat bathochromically shifted (ca. 850$4500 \mathrm{~cm}^{-1}$ ). Good agreement is also observed between computed and observed emission bands, although some blue shift is seen for the former ones with respect to the latter ones, especially for 1M in MeCN (Table S6, Supporting Information). The predominant frontier molecular orbitals which are involved in the main electronic transitions are reported in Table 2. Our calculations show that the pyrene-like nature of the transitions is generally maintained. In compounds 1 and 2, pyrene's $\mathrm{S}_{1} \leftarrow \mathrm{~S}_{0}$ transition mostly involves a mixture of $\mathrm{HOMO} \rightarrow$ LUMO +1 and HOMO- $1 \rightarrow$ LUMO electronic transitions and is transition dipole forbidden, and therefore has a very small oscillator strength, in agreement with the very small extinction coefficients observed experimentally (Table 2). In the cationic derivatives, the $S_{1} \leftarrow S_{0}$ transitions remain forbidden as well with weak oscillator strengths, which is in agreement with the

Table 2. Absorption wavelengths ( $\lambda_{\text {calc, }} \mathrm{nm}$ ), oscillator strengths ( $f$ ), and main electronic transitions calculated at the PBEO-GD3BJ/LANL2DZP level in MeCN for compounds $\mathbf{1}, \mathbf{2}, \mathbf{1 M}, \mathbf{2 M}, \mathbf{1 H}$, and $\mathbf{2 H}$. Corresponding excited state numbers $\left(S_{n}\right)$ are given in brackets. Experimental wavelengths $\left(\lambda_{\text {exp }}, n m\right)$ and their extinction coefficients $\left(\varepsilon, \mathrm{M}^{-1} \mathrm{~cm}^{-1}\right)$ measured in MeCN are given for comparison.

| Compound | $\lambda_{\text {exp }}$ | $\begin{aligned} & \varepsilon \\ & {\left[\mathrm{M}^{-1} \mathrm{~cm}^{-1}\right]} \end{aligned}$ | $\lambda_{\text {calc }}$ | $f\left(S_{n}\right)$ | Main electronic transitions [weight \%] ${ }^{\text {[a] }}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 362 | 1100 | 355 | $0.013\left(S_{1}\right)$ | $\mathrm{H} \rightarrow \mathrm{L}+1$ (77\%); $\mathrm{H}-1 \rightarrow \mathrm{~L}$ ( $21 \%$ ) |
|  | 337 | 37000 | 338 | $0.331\left(S_{2}\right)$ | $\mathrm{H} \rightarrow \mathrm{L}$ (90\%) |
|  | 276 | 52000 | 276 | $1.572\left(\mathrm{~S}_{5}\right)$ | $\mathrm{H}-1 \rightarrow \mathrm{~L}+1$ (91\%) |
| 1M | 418 | 2900 | 462 | $0.017\left(S_{1}\right)$ | $\mathrm{H} \rightarrow \mathrm{L}$ (98\%) |
|  | 337 | 45000 | 353 | $0.031\left(S_{2}\right)$ | $\mathrm{H}-1 \rightarrow \mathrm{~L}(58 \%) ; \mathrm{H} \rightarrow \mathrm{L}+1$ (40\%) |
|  | 269 | 29000 | 329 | $1.444\left(S_{3}\right)$ | $\mathrm{H} \rightarrow \mathrm{L}+1$ (53\%); $\mathrm{H}-1 \rightarrow \mathrm{~L}$ (41 \%) |
|  |  | [b] | 270 | $0.142\left(\mathrm{~S}_{8}\right)$ | $\mathrm{H} \rightarrow \mathrm{L}+3$ (36\%); $\mathrm{H}-1 \rightarrow \mathrm{~L}+1$ (26\%); $\mathrm{H}-1 \rightarrow \mathrm{~L}+2$ (20\%); $\mathrm{H}-3 \rightarrow \mathrm{~L}$ (13\%) |
| 1H | 421 |  | 469 | $0.016\left(S_{1}\right)$ | $\mathrm{H} \rightarrow \mathrm{L}$ (98\%) |
|  | 337 |  | 354 | $0.034\left(S_{2}\right)$ | $\mathrm{H}-1 \rightarrow \mathrm{~L}$ (61\%); $\mathrm{H} \rightarrow \mathrm{L}+1$ (37\%) |
|  | 274 |  | 328 | $1.303\left(S_{3}\right)$ | $\mathrm{H} \rightarrow \mathrm{L}+1$ (55\%); $\mathrm{H}-1 \rightarrow \mathrm{~L}$ (37\%) |
|  |  |  | 269 | $0.146\left(S_{8}\right)$ | $\mathrm{H} \rightarrow \mathrm{~L}+3 \text { (34\%); } \mathrm{H}-1 \rightarrow \mathrm{~L}+1 \text { (26\%); }$ |
| 2 | 363 | 1400 | 371 | $0.017\left(\mathrm{~S}_{1}\right)$ | $\mathrm{H} \rightarrow \mathrm{L}+1$ (79\%); $\mathrm{H}-1 \rightarrow \mathrm{~L}$ (20\%) |
|  | 339 | 35000 | 342 | $0.190\left(S_{2}\right)$ | $\mathrm{H} \rightarrow \mathrm{L}$ (83\%); $\mathrm{H}-1 \rightarrow \mathrm{~L}+1$ (15\%) |
|  | 291 | 104000 | 296 | $2.553\left(\mathrm{~S}_{5}\right)$ | $\mathrm{H}-1 \rightarrow \mathrm{~L}+1$ (83\%); H $\rightarrow$ L (14\%) |
| 2M | 437 | 4300 | 468 | $0.027\left(S_{1}\right)$ | $\mathrm{H} \rightarrow \mathrm{L}$ (97\%) |
|  | 337 | 80000 | 399 | $0.000\left(S_{2}\right)$ | $\mathrm{H} \rightarrow \mathrm{L}+1$ (99\%) |
|  | 320 | 120000 | 362 | $0.436\left(S_{3}\right)$ | $\mathrm{H}-1 \rightarrow \mathrm{~L}(77 \%) ; \mathrm{H} \rightarrow \mathrm{L}+2$ (22\%) |
|  |  |  | 332 | $2.126\left(S_{4}\right)$ | $\mathrm{H} \rightarrow \mathrm{L}+2$ (73\%); H-1 $\rightarrow$ L (21 \%) |
| 2H | 428 | 1530 | 469 | $0.027\left(S_{1}\right)$ | $\mathrm{H} \rightarrow \mathrm{L}$ (97\%) |
|  | 341 | 30500 | 401 | $0.000\left(S_{2}\right)$ | $\mathrm{H} \rightarrow \mathrm{L}+1$ (99\%) |
|  | 321 | 50000 | 359 | $0.426\left(S_{3}\right)$ | $\mathrm{H}-1 \rightarrow \mathrm{~L}(80 \%) ; \mathrm{H} \rightarrow \mathrm{L}+2$ (19\%) |
|  |  |  | 328 | 1.887 ( $\mathrm{S}_{4}$ ) | $\mathrm{H} \rightarrow \mathrm{L}+2$ (75\%); $\mathrm{H}-1 \rightarrow \mathrm{~L}$ (19\%) |

[^0]

Figure 5. Representations of the frontier molecular orbitals and their energies (eV) calculated at the PBE0-GD3BJ level in MeCN (isocontour value: $\pm 0.03$ au), involved in various electronic transitions.
small extinction coefficients $\varepsilon$ that we observed experimentally ( $1500-4300 \mathrm{~m}^{-1} \mathrm{~cm}^{-1}$ ).
As noted above, substituents at the 2- and 2,7-positions do not contribute to pyrene's HOMO or LUMO. ${ }^{[21,24,25 a, b, 26 a]}$ However, they can have a profound effect on pyrene's HOMO-1 and LUMO +1 . Indeed, orbitals of the pyridine moiety in the neutral compounds 1 and 2 mix with pyrene's HOMO-1/LUMO + 1 , although not very efficiently. Consequently, the influence on the first transition is very small for $\mathbf{1}$ and $\mathbf{2}$ and, therefore, remains of HOMO- $1 \rightarrow$ LUMO and $\mathrm{HOMO} \rightarrow \mathrm{LUMO}+1$ character. However, protonation (compounds $\mathbf{1 H} / \mathbf{2 H}$ ) or methylation ( $\mathbf{1 M} / \mathbf{2 M}$ ) of the nitrogen, leads to a drastic stabilization of the LUMO + 1, below the pyrene-like LUMO (Figure 5). Hence, unlike pyrene, the LUMO of the cationic derivatives $\mathbf{1 H}, \mathbf{2 H}$, $\mathbf{1 M}$, and $\mathbf{2 M}$ does not have a nodal plane through the long axis of the molecule, as observed previously with - $\mathrm{BMes}_{2}$ as the acceptor substituent. ${ }^{[21,24,256]}$ This explains the stronger bathochromic shift of the first transition and strongly altered emission in these cationic compounds compared to 1 and 2. Our TD-DFT calculations also show that the $\mathrm{S}_{1} \leftarrow \mathrm{~S}_{0}$ transitions in $\mathbf{1 H}, \mathbf{2 H}, \mathbf{1 M}$, and $\mathbf{2 M}$ no longer involve a mixture of HOMO$1 \rightarrow$ LUMO and HOMO $\rightarrow$ LUMO +1 electronic transitions as in pyrene or the neutral compounds $\mathbf{1}$ and $\mathbf{2}$, but are nearly pure HOMO $\rightarrow$ LUMO electronic transitions. The CT character of the
excitations is nicely illustrated by the difference density plots between the $S_{1}$ and $S_{0}$ states shown in Table 3. The transferred charges ( $q$ ) for $\mathbf{1 H}\left(q^{c \top}=0.862\right.$ e) and $\mathbf{1 M}\left(q^{C \top}=0.878\right.$ e) are higher than for their disubstituted (symmetrical) analogues 2 H ( $q^{C \top}=0.743$ e) and 2 M ( $q^{C \top}=0.761 \mathrm{e}$ ), respectively. This is also fully in line with our observations which indicate that the $\mathrm{S}_{1} \leftarrow \mathrm{~S}_{0}$ transitions of 1 H and 1 M have smaller extinction coefficients, and their emissions are further bathochromically shifted compared to those of $\mathbf{2 H}$ and $\mathbf{2 M}$, respectively, and that $\mathbf{1 M}$ shows some degree of solvatochromism compared to non-solvatochromic 2M (Tables 1 and S5, Supporting Information).

## Interactions with ds-DNA

As many pyrene derivatives showed intriguing biorelevant interactions, acting as probes for various DNA/RNA sequences ${ }^{[6-10]}$ and intracellular mitochondrial recognition, ${ }^{[48]}$ we performed preliminary studies on interactions of permanently charged $N$-methylpyridinium derivatives $\mathbf{1 M}$ and $\mathbf{2 M}$ with the most commonly used ds-DNA representative-calf thymus (ct)DNA. This DNA is characterised by a typical B-helical secondary structure and almost equivalent number of GC-(48\%) and AT( $52 \%$ ) base-pairs, thus representing a good target for both intercalative and minor groove binding mode. For the purpose

Table 3. Density-difference plots between relaxed excited state $S_{1}$ and ground state $S_{0}$ (with $S_{1}$ geometry) (red $=$ increase, blue $=$ decrease of electron density) calculated at the PBEO-GD3BJ/LANL2DZP level. Isocontour value: $0.03 \mathrm{au} . q^{C T}$ (e) and $d^{C T}(\AA)$ ) correspond to the transferred charge and the distance between the positive (blue cross) and the negative (red cross) centroids of charge, respectively.
Compd
of the study under biorelevant conditions, we prepared stock solutions of $\mathbf{1 M}$ and $\mathbf{2 M}$ in DMSO ( 1 mm ), stored them at $+4^{\circ} \mathrm{C}$, and aliquots were used for dilution in buffer (sodium cacodylate buffer, $I=50 \mathrm{~mm}, \mathrm{pH} 7.0$ ) prior to each measurement. The stability of the stock solutions was confirmed by checking the UV/Vis absorption or emission spectra of the freshly prepared solutions in buffer. The UV/Vis absorption properties in the buffer (Figure S37, Supporting Information) revealed very similar shape, as noted for other organic solvents (Figure 2). However, the change in ratio of absorbances at wavelengths $316 \mathrm{~nm} / 306 \mathrm{~nm}$ (1M) and $336 \mathrm{~nm} / 319 \mathrm{~nm}$ (2M) at higher concentrations ( $>5 \times 10^{-6} \mathrm{~m}$ ) suggests intermolecular aggregation (Figure S37, Supporting Information). Therefore, all further experiments were performed at concentrations $<5 \times$ $10^{-6} \mathrm{~m}$.
As a preliminary screening of the interactions of our compounds with ds-DNA, we performed thermal denaturation experiments. It is well known that, upon heating, the ds-helices of polynucleotides dissociate into two single stranded polynucleotides at a well-defined temperature ( $T_{\mathrm{m}}$ value). Non-covalent binding of small molecules to ds-polynucleotides usually increases the thermal stability of the ds-helices, thus resulting in
an increased $T_{\mathrm{m}}$ value, and this increase ( $\Delta T_{\mathrm{m}}$ ) can (in corroboration with other methods) be related to the various binding modes. ${ }^{[55]}$ For example, if pyrene compounds intercalate into the ds-DNA, $\Delta T_{\mathrm{m}}>2-5^{\circ} \mathrm{C}$ can be expected, whereas the binding of pyrenes within the polynucleotide groove should have a negligible stabilising effect. ${ }^{[56]}$ We observed that the addition of 1 M monocation caused moderate stabilisation of the ctDNA against thermal denaturation $\left(\Delta T_{m}=2^{\circ} \mathrm{C}\right.$ for $\mathrm{r}_{[1 M) /[c t D N A]}=$ 0.2 ), whereas the dication $\mathbf{2 M}$ yielded strong stabilisation ( $\Delta T_{m}>16^{\circ} \mathrm{C}^{\text {for }} \mathrm{r}_{[2 M / / \text { /ctoNA }]}=0.2$ ), and such stabilisation is indicative of an intercalative binding mode.
The titration of $\mathbf{1 M}$ or $\mathbf{2 M}$ with ct-DNA resulted in pronounced changes of their UV/Vis spectra (Figure 6(a) and Figures S40 and S41, Supporting Information), characterised by strong hypochromic and bathochromic effects typical of aromatic stacking interactions. Different responses were observed for $\mathbf{1 M}$ and $\mathbf{2 M}$ in the fluorimetric titrations with the ct-DNA. For example, the emission of $\mathbf{1 M}$ was primarily quenched up to ratio $r_{[1 M / /[t D N A]}=0.3$ and, upon further additions of ct-DNA, its emission intensity increased (Figure S42, Supporting Information). The quenching of emission occurred at an excess of 1 M over the ct-DNA binding sites ( $r_{[1 M) /(C t D N A]}>0.3$ ) and, there-


Figure 6. Fluorimetric titration of $\mathbf{2 M}$ with ct-DNA at pH 7.0 , sodium cacodylate buffer, $I=0.05 \mathrm{M}$ : a) changes in UV/Vis absorption of $2 \mathbf{M}\left(c=5 \times 10^{-6} \mathbf{M}\right)$; and b) changes in relative fluorescence intensity of $\mathbf{2 M}\left(c=3 \times 10^{-8} \mathrm{M}\right)$.
fore, it was attributed to aggregation of $\mathbf{1 M}$ along the ct-DNA. The increased emission intensity at $\mathrm{r}_{[1 \mathrm{M} / \text { /CtDNA }]}<0.3$ indicates binding of single $1 \mathbf{M}$ molecules to separate ct-DNA binding sites. For the larger dicationic 2M, the addition of ct-DNA caused only quenching of emission (Figure 6(b) and Figure S43, Supporting Information), which can be attributed mainly to the binding of single 2M molecules to separate ctDNA binding sites. Analysis of the titration data by non-linear fitting to Scatchard equation ${ }^{[57]}$ gave binding constants of $\log K s=5.5$ (for 1 M , data fitted only for $\mathrm{r}_{[1 \mathrm{M} / /(\operatorname{cosNA}]}<0.3$ ) and $\log K s=8.8$ (for 2M), respectively, for Scatchard ratio $n_{[b b o u n d ~ d y e] ~}$ $[$ [CDNA] $=0.3$. The difference in binding constants matched well with the observed difference in thermal denaturation $\left(\Delta T_{m}\right)$ values. Thus, it can be inferred that the dicationic compound 2M interacts more strongly with the ct-DNA compared to the monocationic compound 1M.
In order to obtain additional insight into the structural properties of small molecule/polynucleotide complexes, we have chosen the circular dichroism (CD) spectroscopic method, which is highly sensitive to conformational changes in the sec-


Figure 7. CD titration of ct-DNA ( $c=2 \times 10^{-5} \mathrm{M}$ ) with: a) $\mathbf{1 M}$; and b) $\mathbf{2 M}$; at molar ratios $\mathrm{r}=$ [compound]/[polynucleotide] and at pH 7.0 , buffer sodium cacodylate, $I=0.05 \mathrm{M}$.
ondary structure of polynucleotides. In addition, achiral small molecules, such as $\mathbf{1 M}$ or $\mathbf{2 M}$, can acquire an induced CD spectrum (ICD) upon binding to polynucleotides, which can give useful information about the modes of interaction. ${ }^{[58]}$ Indeed, CD spectra (Figure 7) revealed weak negative ICD bands at $\lambda=$ $300-350 \mathrm{~nm}$ range for both, $\mathbf{1 M}$ and $\mathbf{2 M}$. As pyrene absorbs in this range, such ICD bands are characteristic of the intercalative binding mode, which agrees well with the bathochromic and hypochromic changes in UV/Vis titrations, high binding constants and thermal denaturation results.

## Conclusions

We have presented structural, photophysical, and electrochemical studies on 2- and 2,7-substituted pyridyl-pyrenes and their $N$-methylated compounds and analogous protonated salts, which are acceptor- $\pi$ and acceptor- $\pi$-acceptor systems, respectively. Their structure determination by X-ray crystallography revealed that the torsion angle between the planes of pyrene and pyridyl moieties tends to become very small upon protonation or methylation, leading to almost planar structures, which agree well with the calculated optimised structures obtained at the PBE0-GD3BJ/LANL2DZP level of theory. Although the absorption spectra show rather small changes upon protonation or methylation, their emission spectra were found to display large bathochromic shifts. The magnitude of such shifts is larger for the dipolar molecule $\mathbf{1 M}$ than for the symmetric molecule $2 \mathbf{M}$. Thus, upon methylation of $1\left(\lambda_{\mathrm{em}}=405 \mathrm{~nm}\right)$ and $2\left(\lambda_{\mathrm{em}}=418 \mathrm{~nm}\right)$ to $1 \mathrm{M}\left(\lambda_{\mathrm{em}}=599 \mathrm{~nm}\right)$ and $2 \mathrm{M}\left(\lambda_{\mathrm{em}}=538 \mathrm{~nm}\right)$, respectively, their corresponding emission maxima shifted bathochromically by $8000 \mathrm{~cm}^{-1}$ and $5300 \mathrm{~cm}^{-1}$, respectively, in MeCN. Moreover, dipolar 1M shows an apparent Stokes shift of $7200 \mathrm{~cm}^{-1}$, which is very large for 2- and 2,7-substituted pyrenes, indicating a rather large change of the dipole moment in the excited state. Symmetric, dicationic 2M shows an apparent Stokes shift of $4300 \mathrm{~cm}^{-1}$, which is smaller than that of 1 M , due to its lack of a dipole moment. Compound 1 M showed an irreversible reduction at -1.53 V , while $\mathbf{2 M}$ showed a reversible reduction at -1.47 V vs. $\mathrm{Fc}^{+} / \mathrm{Fc}$ couple in MeCN . These observed formal potentials are in accordance with other similar
$\pi$-conjugated extended viologens reported in the literature. However, the reduction of $\mathbf{2 M}$ occurs at a lower potential compared to 2,7-bis $\left(\mathrm{BMes}_{2}\right)$ pyrene, indicating that the $N$-methylpyridinium group is a stronger acceptor than $-\mathrm{BMes}_{2}$. The reduction of $\mathbf{2 M}$ can be attributed either to a two-electron reduction or to two closely spaced consecutive one-electron reductions, leading to the formation of the neutral species, without producing the intermediate radical-cation, which was further established by investigating the step-wise reduction processes by spectroelectrochemical measurements. Our experimental observations are further confirmed by theoretical calculations performed at both the PBEO-GD3BJ and CAM-B3LYP levels of theory. The HOMO of both the neutral and methylated/protonated compounds remain rather unaffected. However, upon methylation or protonation, the LUMOs of the neutral compounds are destabilised, while the LUMO + 1 orbitals are greatly stabilised and become the LUMOs. Thus, the new LUMOs of the methylated/protonated compounds do not have a nodal plane passing through the long axis of the molecule. Furthermore, preliminary studies reveal that $\mathbf{1 M}$ and $\mathbf{2 M}$ interact strongly with ct-DNA by intercalation into the double helix, which makes them very interesting compounds for further studies of their biological activity. Detailed investigation of these interactions and their significance in biological systems are ongoing. Our research findings are expected to have significant impact on further investigations on the photophysics and electrochemistry of other PAH-derived $\pi$-conjugated viologen compounds and their potential applications in biological systems.

## Experimental Section

## General

The catalyst precursor $[\operatorname{lr}(\mathrm{OMe})(\mathrm{COD})]_{2}$ was prepared according to a literature procedure ${ }^{[59]}$ and $\mathrm{B}_{2} \mathrm{pin}_{2}$ was a gift from AllyChem Co. Ltd. The starting materials 2-(Bpin)pyrene and 2,7-bis(Bpin)pyrene were synthesised by following methods developed by Marder and co-workers. ${ }^{[23 c]}$ The syntheses of $\mathbf{1 , 2 , 1 M}$, and $\mathbf{2 M}$ were carried out using standard Schlenk or glovebox (Innovative Technology Inc.) techniques under an argon atmosphere. All of the compounds are, in general, stable, and can be handled in air. All solvents used for reactions were HPLC grade, dried using an Innovative Technology Inc. Solvent Purification System, and were further deoxygenated and stored under argon. All other chemicals were purchased from common commercial sources and were used without further purification.
${ }^{1} \mathrm{H}$ and ${ }^{13} \mathrm{C}\left\{{ }^{1} \mathrm{H}\right\}$ solution NMR spectroscopic data were obtained at ambient temperature using a Bruker Avance 300 NMR spectrometer, operating at 300 MHz for ${ }^{1} \mathrm{H}$ and 75 MHz for ${ }^{13} \mathrm{C}\left\{{ }^{1} \mathrm{H}\right\}$. Chemical shifts ( $\delta$ ) were referenced to solvent residual peaks. High resolution mass spectrometry (HRMS) was performed with a Thermo Fisher Scientific Exactive Plus Orbitrap MS System with either an Atmospheric Sample Analysis Probe (ASAP) or Atmospheric Pressure Chemical lonization (APCI) or an electrospray ionization (ESI) probe. Elemental analyses were performed on an Elementar vario MICRO cube elemental analyser.
tion effects by employing the Bruker software packages. The structures were solved using the intrinsic phasing method (SHELXT). ${ }^{[60]}$ All non-hydrogen atoms were refined in anisotropic approximation, with hydrogen atoms 'riding' in idealised positions, by full-matrix least squares against $F^{2}$ on all data, using SHELX[ ${ }^{[60]}$ software. Hirshfeld surfaces were calculated and analysed using the Crystal Explorer program (version 17.5). ${ }^{[61]}$ Various structural information was extracted and graphics were produced using Mercury software. ${ }^{[62]}$ Crystal data and experimental details are listed in Table S1 in the Supporting Information. Deposition numbers 2016793 (1H), 2016794 (1M), 2016795 (2H), and 2016796 (2M) contain the supplementary crystallographic data for this paper. These data are provided free of charge by the joint Cambridge Crystallographic Data Centre and Fachinformationszentrum Karlsruhe Access Structures service.

## General photophysical measurements

All measurements were performed in standard quartz cuvettes ( $1 \mathrm{~cm} \times 1 \mathrm{~cm}$ cross-section) fitted with screw caps. UV/Visible absorption spectra were recorded using an Agilent 8453 diode array UV/Vis spectrophotometer. Emission and excitation spectra were recorded using an Edinburgh Instruments FLSP920 spectrometer equipped with a cooled red PMT detector from Hamamatsu (R13456-P), and with double monochromators, operating in rightangle geometry mode. A 450 W continuous xenon arc lamp (Xe900) was used as the excitation source and was focused on the sample. All solutions used in photophysical measurements had concentrations lower than $5 \times 10^{-6} \mathrm{M}$ to eliminate the chances of excimer formation during fluorescence measurements.

## Fluorescence quantum yield measurements

The fluorescence quantum yields were measured using a calibrated integrating sphere ( 150 mm inner diameter) from Edinburgh Instruments combined with the above FLSP920 spectrometer. The Ion-gest-wavelength absorption maximum was chosen as the excitation wavelength, unless otherwise stated. The corresponding blank solvent was used as a reference for solution-phase measurements and $\mathrm{BaSO}_{4}$ was used as reference for solid state measurements.

## Fluorescence lifetime measurements

Lifetimes were measured using a 315.8 nm emitting picosecond pulsed diode laser as the light source in the time-correlated singlephoton counting (TCSPC) mode using the Edinburgh Instruments FLS920 spectrometer. Emission decay was detected with a 3 nm emission-slit band-width. A 5000 kHz (every 200 ns ) pulse was generated. The time range was set to 200 ns , and 1024 channels were used. Instrument response functions (IRFs) were measured from the scatter of an aqueous suspension of Ludox at the same excitation wavelength. A reconvolution-fit was performed using Software F900 version 7.2.1. The quality of the decay fits was assessed by the calculated values of the reduced $\chi^{2}$ and Durbin-Watson parameters, and visual inspection of the weighted and autocorrelated residuals.

## Electrochemistry

Cyclic voltammetry measurements were carried out using a Gamry Instruments Reference 600 potentiostat. A standard arrangement of a three-electrode cell configuration was set up using a platinum disk working electrode, a platinum wire counter electrode, and a silver wire, separated by a Vycor tip, as the reference electrode.

Formal redox potentials are referenced to the ferrocene/ferrocenium $\left(\left[\mathrm{Cp}_{2} \mathrm{Fe}\right]^{+/ 0}\right)$ redox couple by using ferrocene as an internal standard. Tetra-n-butylammonium hexafluorophosphate ([nBu $\left.{ }_{4} \mathrm{~N}\right]$ [ $\mathrm{PF}_{6}$ ]) was employed as the supporting electrolyte. Compensation for resistive losses (iR drop) was employed for all measurements.

## Spectroelectrochemical measurements

Spectroelectrochemical experiments were carried out using an Agilent Cary 5000 spectrometer in combination with a designed sample compartment consisting of a cylindrical PTFE cell with an Infrasil ${ }^{\circledR}$ wedge window with an angle of $0.5^{\circ}$. An adjustable three-in-one electrode ( 6 mm platinum disc as working electrode, 1 mm platinum as counter electrode, and a pseudo reference electrode) was used. The potentials were adjusted with a Gamry 600 potentiostat. All experiments were measured at room temperature under an argon atmosphere.

## Theoretical computations

All DFT calculations were carried out using the Gaussian09-D01 suite of programs. ${ }^{[63]}$ The geometries of all the compounds were optimised without symmetry constraints using the PBEO function$\mathrm{al}^{[64,65]}$ including Grimme's dispersion correction with Becke-Johnson damping (GD3BJ), ${ }^{[66,67]}$ and the LANL2DZ basis set augmented with polarisation functions for all atoms, except hydrogen (xyz coordinates of the optimised geometries are given in Table S4, Supporting Information). The nature of the stationary points after optimisation was checked by calculations of the harmonic vibrational frequencies to ensure that genuine minima were obtained. Timedependent density functional theory (TD-DFT) calculations were performed at the same level of theory using the previously optimised geometries. For comparison, computations were also carried out using the CAM-B3LYP functional ${ }^{[68]}$ (Tables S5 and S6, Supporting Information). A better agreement with the experimentally observed values was obtained with the PBEO functional. Solvent effects were taken into account by means of the polarisable continuum model (PCM). ${ }^{[69]}$ The molecular orbitals and theoretical absorption spectra were drawn using the GaussView program. ${ }^{[70]}$
Periodic DFT calculations were performed on compound 2M using the VASP code. ${ }^{[71]}$ The exchange-correlation interaction was described within the generalised gradient approximation in the parametrisation of the Perdew-Burke-Ernzerhof functional ${ }^{[64]}$ Projectoraugmented wave potentials were used for all atoms. ${ }^{[72]}$ Calculations were performed with a cut-off energy of 530 eV . The electronic wave functions were sampled on dense densities in the irreducible BZ using the Monkhorst-Pack method. ${ }^{[73]}$ Geometry optimisation including cell parameters and atomic positions were carried out without any symmetry constraints (see 2 M -sol-opt.mol file in the Supporting Information).

## Study of interactions with DNA

Stock solutions of $\mathbf{1 M}$ and $\mathbf{2 M}$ were prepared in DMSO ( 1 mm ) and were diluted prior to use, whereby the DMSO content was always maintained at $<0.1 \%$. The UV/Vis spectra were recorded on a Varian Cary 100 Bio spectrophotometer, fluorescence spectra were recorded on an Agilent Eclipse fluorimeter, and CD spectra were recorded on a JASCO J 815 spectropolarimeter at $25.0^{\circ} \mathrm{C}$ using appropriate quartz cuvettes of 1 cm path-length. Calf thymus (ct)DNA was purchased from Aldrich, dissolved in sodium cacodylate buffer, $I=0.05 \mathrm{~m}, \mathrm{pH} 7.0$, additionally sonicated and filtered through a $0.45 \mu \mathrm{~m}$ filter to obtain uniform rod-like fragments of the B-helical DNA. ${ }^{[74]}$ Polynucleotide concentration was determined
spectroscopically, by measuring the absorbance at 260 nm for molar extinction coefficient $6600 \mathrm{~m}^{-1} \mathrm{~cm}^{-1}$, whereby the concentration per nucleobase was obtained (traditionally referred as the concentration of phosphates).
For the study of interactions with DNA, the aqueous solutions of the compounds were buffered to pH 7.0 (sodium cacodylate buffer, $I=0.05 \mathrm{~m}$ ). Spectrophotometric titrations were performed by adding portions of DNA solution into the solution of the compound under investigation and CD experiments were performed by adding portions of stock solution of our compounds into the solution of DNA. In all titration experiments, dilution effects were taken into account.
Thermal melting curves for ds-DNA and its complexes with 1 M and 2M were determined as previously described, ${ }^{[55]}$ by following the absorption change at 260 nm as a function of temperature. The absorbance of the ligands was subtracted from every curve and the absorbance scale was normalised. The $T_{\mathrm{m}}$ values, the midpoints of the transition curves, were determined from the maximum of the first derivative and were checked graphically by the tangent method. ${ }^{[55,75]}$ The $\Delta T_{\mathrm{m}}$ values were calculated by subtracting $T_{\mathrm{m}}$ values of the free nucleic acid from that of the complex. Every $\Delta T_{\mathrm{m}}$ value reported was the average of at least two measurements. The error in $\Delta T_{\mathrm{m}}$ can be $\pm 0.5^{\circ} \mathrm{C}$.

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## Conflict of interest

The authors declare no conflict of interest.
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[^0]:    [a] $\mathrm{H}=\mathrm{HOMO}$ and $\mathrm{L}=$ LUMO. [b] Partially dissociates in solution, see text.

