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# Computational modelling elucidates the mechanism of ciliary regulation in health and disease

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## Abstract

**Background:** Ciliary dysfunction leads to a number of human pathologies, including primary ciliary dyskinesia, nephronophthisis, situs inversus pathology or infertility. The mechanism of cilia beating regulation is complex and despite extensive experimental characterization remains poorly understood. We develop a detailed systems model for calcium, membrane potential and cyclic nucleotide-dependent ciliary motility regulation.

**Results:** The model describes the intimate relationship between calcium and potassium ionic concentrations inside and outside of cilia with membrane voltage and, for the first time, describes a novel type of ciliary excitability which plays the major role in ciliary movement regulation. Our model describes a mechanism that allows ciliary excitation to be robust over a wide physiological range of extracellular ionic concentrations. The model predicts the existence of several dynamic modes of ciliary regulation, such as the generation of intraciliary  $Ca^{2+}$  spike with amplitude proportional to the degree of membrane depolarization, the ability to maintain stable oscillations, monostable multivibrator regimes, all of which are initiated by variability in ionic concentrations that translate into altered membrane voltage.

**Conclusions:** Computational investigation of the model offers several new insights into the underlying molecular mechanisms of ciliary pathologies. According to our analysis, the reported dynamic regulatory modes can be a physiological reaction to alterations in the extracellular environment. However, modification of the dynamic modes, as a result of genetic mutations or environmental conditions, can cause a life threatening pathology.

## Background

Cilia are cellular protrusions which have been conserved in a wide range of organisms ranging from protozoa to the digestive, reproductive and respiratory systems of vertebrates [1]. Mobile or immotile cilia exist on every cell of the human body [2] and the insufficiently recognised importance of the cilium compartment in human physiology has been recently highlighted [1,3]. Cilia are present on most eukaryotic cell surfaces with the exception of the cells of higher plants and fungi [4]. Ciliary motility is important for moving fluids and particles over epithelial surfaces, and for the cell motility of

vertebrate sperm and unicellular organisms. The cilium contains a microtubule-based axoneme that extends from the cell surface into the extracellular space. The axoneme consists of nine peripheral microtubule doublets arranged around a central core that may or may not contain two central microtubules (9+2 or 9+0 axoneme, respectively). Cilia can be broadly classified as 9+2 motile cilia or 9+0 immotile sensory cilia, although there are examples of 9+2 sensory cilia and 9+0 motile cilia. In mammals, motile 9+2 cilia normally concentrate in large numbers on the cell surface, beat in an orchestrated wavelike fashion, and are involved in fluid and cell movement. In contrast to motile cilia, primary cilia project as single immotile organelles from the cell surface. Primary cilia are found on nearly all cell types in mammals [5] and many are highly adapted to serve

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specialized sensory functions. The 9+2 cilia usually have dynein arms that link the microtubule doublets and are motile, while most 9+0 cilia lack dynein arms and are non-motile. In total, eight different types of cilia has been identified to date [6]. In this study, we investigate the mechanism of movement regulation for the motile type of cilia.

Although each individual cilium represents a tiny hair-like protrusion of only 0.25  $\mu\text{m}$  in diameter and approximately 5-7  $\mu\text{m}$  in length, cilia covering human airways can propel mucus with trapped particles of length up to 1 mm at a speed of 0.5 mm/second [7]. Such efficiency can be achieved due to the coordination between cilia and stimulus-dependent regulation of the rate of cilia beat. Dysfunction of ciliary regulation gives rise to pathologic phenotypes that range from being organ specific to broadly pleiotropic [3]. A link between ciliary function and human disease was discovered when individuals suffering from syndromes with symptoms including respiratory infections, anosmia, male infertility and situs inversus, were shown to have defects in ciliary structure and function [6].

Microscopic organisms that possess motile cilia which are used exclusively for either locomotion or to simply move liquid over their surface include *Paramecia*, *Karyorelictea*, *Tetrahymena*, *Vorticella* and others. The human mucociliary machinery operates in at least two different modes, corresponding to a low and high rate of beating. It has been shown that the high rate mode is mediated by second messengers [8], including purinergic, adrenergic and cholinergic receptors [9-19]. This mode enables a rapid response, which can last a significant period of time, to various stimuli by drastically increasing the ciliary beat frequency (CBF). At the same time, several ciliary movement modes have been reported in a ciliate *Paramecium caudatum* [20]. The remarkable conservation of ciliary mechanisms [21-25] creates grounds for the speculation that there can more than two ciliary beating modes in human tissues. It is, therefore, reasonable to suggest, that some human diseases, associated with aberrant ciliary motility, can arise due to modifications in the beating mode. Clearly, the development of therapeutic strategies against ciliary-associated pathologies will require advanced understanding of ciliary beating regulation mechanisms.

The periodic beating of cilia is governed by the internal apparatus of the organelle [26]. Its core part, the axoneme, contains nine microtubule pairs encircling the central pair. The transition at the junction of the cellular body and the ciliary axoneme is demarcated by Y-shaped fibres, which extend from the microtubule outer doublets to the ciliary membrane. The transition area, in combination with the internal structure of the basal body, is thought to function as a filter for the cilium,

regulating the molecules that can pass into or out of the cilium. Ciliary motility is accomplished by dynein motor activity in a phosphorylation-dependent manner, which allows the microtubule doublets to slide relative to one another [12]. The dynein phosphorylation that controls ciliary activity is regulated by the interplay of calcium ( $\text{Ca}^{2+}$ ) and cyclic nucleotide pathways. The beating pattern of cilia consists of a fast effective stroke and a slower recovery stroke. During the effective stroke cilia are in an almost upright position, generating force for mucus movement. During the recovery stroke, the cilia are recovering from the power stroke to the original position by moving in the vicinity of the cell surface.

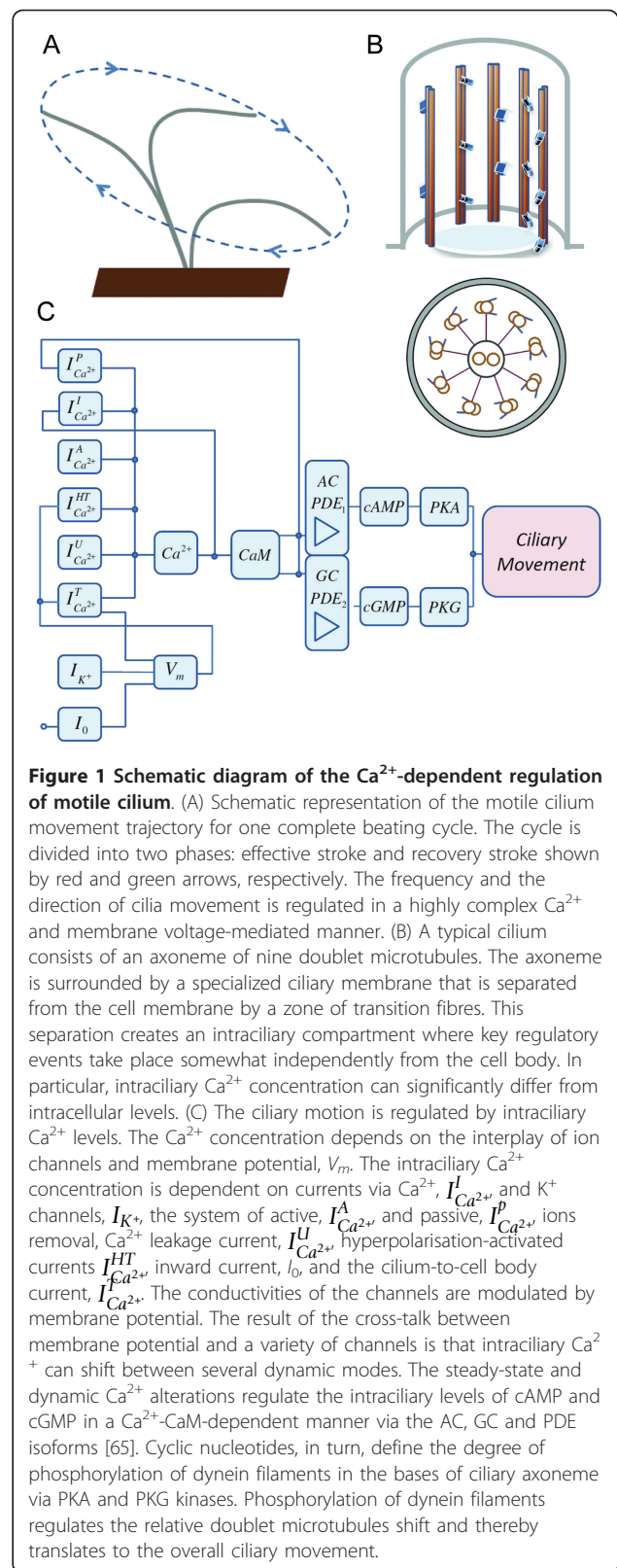
Current theories which attempt to explain the workings of the  $\text{Ca}^{2+}$ -dependent CBF regulation mechanism are incomplete and highly controversial. Elevation of intraciliary  $\text{Ca}^{2+}$  is one of the major regulators of ciliary movement. Calcium influx regulates ciliary activity by increasing intraciliary  $\text{Ca}^{2+}$  only, while the cytosolic bulk remains at a low level. Separate ciliary compartmentalisation for  $\text{Ca}^{2+}$  allows prolonged activation of ciliary beating without damaging the cell through high  $\text{Ca}^{2+}$  concentrations. It is well known that calcium fluxes via calcium channels lead to changes in organisms' swimming behaviour [27-29]. In mucus-transporting cilia,  $\text{Ca}^{2+}$  mediates CBF increase [19,30-32]. It has also been shown that there are some differences in the  $\text{Ca}^{2+}$ -dependent CBF regulation in single cell organisms and in humans [12]. Sustained CBF increase requires prolonged elevation of  $\text{Ca}^{2+}$  levels which can be lethal to the cell [33,34]. It has been suggested that  $\text{Ca}^{2+}$ -dependent ciliary regulation takes place locally in the vicinity or within the ciliary compartment, almost independently from intracellular  $\text{Ca}^{2+}$  concentration [35]. Given that the gradient of free  $\text{Ca}^{2+}$  in the cytosol dissipates within 1-2 seconds [36], it appears more likely that cilia form their own compartment where  $\text{Ca}^{2+}$  is regulated by active  $\text{Ca}^{2+}$  transport in a similar fashion to the intracellular  $\text{Ca}^{2+}$  regulatory system. This hypothesis resolves the problem of maintaining physiological levels of intracellular  $\text{Ca}^{2+}$  concentration. A number of experimental studies have reported several controversial results relating to the  $\text{Ca}^{2+}$ -dependent mechanism of cilia regulation. For example, it has been reported that spontaneous cilia beat does not require alterations in  $\text{Ca}^{2+}$  [31,35], while nucleotide-dependent CBF increase requires  $\text{Ca}^{2+}$  [8]. It has also been shown that uncoupling between  $\text{Ca}^{2+}$  and CBF can be achieved by inhibition of  $\text{Ca}^{2+}$ -dependent protein calmodulin (CaM) or the cyclic nucleotide pathway [19,32,37].

These findings suggest that intraciliary  $\text{Ca}^{2+}$  does not regulate cilia beat in isolation, but instead does so as part of more complex signalling network. Although it was originally believed that  $\text{Ca}^{2+}$ , cyclic adenosine

monophosphate (cAMP) and guanosine monophosphate (cGMP) regulate ciliary beat independently, numerous reports now strongly indicate that all three pathways are tightly interconnected [12,19,32,38-42]. The cAMP-dependent protein kinase (PKA) phosphorylates dynein in the bases of cilia and thereby increases the forward swimming speed in *Paramecium* [43-45]. Similar effects have been reported for PKA-dependent phosphorylation of axonemal targets in mammalian respiratory cilia [46]. Several lines of evidence indicate that PKA and cGMP-dependent kinase (PKG) both phosphorylate specific axonemal targets in a cAMP and cGMP-dependent manner. A schematic diagram for the underlying biochemical machinery for cilia movement regulation is shown in Figure 1C. It is striking that, despite significant experimental characterisation of this system, there is still rather limited mechanistic understanding of how intraciliary  $Ca^{2+}$  and nucleotide interplay relates to CBF.

Another major regulator of ciliary beating is the membrane potential. A number of studies have reported the voltage-dependent effects of ciliary beating. The ciliate *Didinium Nasutum* has been shown to respond both to hyper- and de-polarization of the membrane [47]. The transmembrane potential alterations were shown to be mediated via the potential-dependent  $Ca^{2+}$  channels [48]. Electrophysiological studies in *Paramecium caudatum* have revealed complex relationships between ciliary  $Ca^{2+}$  currents, intraciliary  $Ca^{2+}$  concentration and transmembrane potential in the regulation of ciliary motility [49-55].

A number of previous computational studies have analysed various aspects of cilia movement regulation. One earlier model assessed the degree of synchronization between small ciliary areas [56]. The effects of viscosity have been investigated in mucus propelling cilia in [57]. The authors found that increasing the viscosity not only decreases CBF, but also changes the degree of correlation and synchronization between cilia. The mechanical properties of cilia motion were studied in an attempt to understand the ciliary dynamics in [58]. The authors concluded that bending and twisting properties of the cilium can determine self-organized beating patterns. While these reports offer valuable insights into the regulatory mechanisms of cilia, a number of essential questions remain unresolved. For example, there has not been a detailed analysis of how individual  $Ca^{2+}$  currents influence intraciliary  $Ca^{2+}$  levels. It also remains unclear how  $Ca^{2+}$  modulates nucleotide levels and membrane potential, and how such regulation affects ciliary movement. None of these reports have elucidated the underlying mechanisms governing the interplay between intraciliary  $Ca^{2+}$  and nucleotide alterations and CBF.



In this study, we integrate the available experimental information on the molecular pathways that regulate intraciliary  $\text{Ca}^{2+}$  concentration into a comprehensive mathematical model. By applying systems analysis, we elucidate the mechanisms of intraciliary  $\text{Ca}^{2+}$  spike generation, analyse the properties of such spikes and demonstrate the conditions under which the  $\text{Ca}^{2+}$  surges can become repetitive. We carry out detailed investigations of the individual current contributions to the regulation of the intraciliary  $\text{Ca}^{2+}$  concentrations and elucidate both steady-state and dynamic responses of  $\text{Ca}^{2+}$  currents and intraciliary  $\text{Ca}^{2+}$  concentration dynamics in response to the altered transmembrane potential shift. The model allows detailed elucidation of transmembrane potential and intraciliary  $\text{Ca}^{2+}$  coupling.

We employ the proposed model in order to understand the underlying molecular mechanisms of the crosstalk between  $\text{Ca}^{2+}$ , membrane potential and nucleotide pathways that regulate ciliary movement. The systems model allows detailed analysis of the individual current contributions to the intraciliary homeostatic  $\text{Ca}^{2+}$  levels. Furthermore, we establish specific regulatory mechanisms for  $\text{Ca}^{2+}$  and cyclic nucleotide-dependent cilia movement characteristics. Crucially, our model predicts the possibility of several ciliary beating modes and describes specific conditions that initiate them. Specifically, we describe intraciliary  $\text{Ca}^{2+}$  dynamic modes that regulate healthy and pathologic cilia beating. We use these findings in order to propose experimentally testable hypotheses for possible therapeutic interventions in human diseases associated with pathologic cilia motility.

## Results

### A new model for the interplay between $\text{Ca}^{2+}$ and $\text{K}^+$ currents and transmembrane potential alterations

A new model for the regulation of ciliary movement that combines multiple  $\text{Ca}^{2+}$  and  $\text{K}^+$  currents [59-62] and transmembrane potential has been developed. In this model, the intraciliary  $\text{Ca}^{2+}$  levels are modulated by  $\text{Ca}^{2+}$  currents through the channels of passive and active  $\text{Ca}^{2+}$  transport, the current from the cilium into the cell body, the  $\text{Ca}^{2+}$  leakage current, and depolarisation and hyperpolarisation-activated currents. Variable extracellular conditions have continuous impact on the transmembrane potential which is intertwined with transmembrane ion currents and intraciliary  $\text{Ca}^{2+}$  homeostasis.

The overall network that regulates ciliary movement is divided into several functional modules (Figure 1C). One module combines all  $\text{Ca}^{2+}$  and  $\text{K}^+$  currents that define intraciliary  $\text{Ca}^{2+}$  homeostasis and the transmembrane potential. One of the most essential intraciliary  $\text{Ca}^{2+}$  binding proteins, CaM [63,64], selectively regulates the activities of adenylate cyclase (AC), guanylate cyclase

(GC) and phosphodiesterases (PDE), and thereby modulates the intraciliary levels of adenosine monophosphate (cAMP) and guanosine monophosphate (cGMP) in a  $\text{Ca}^{2+}$  dependent manner [65]. The cAMP- and cGMP-dependent kinases phosphorylate dynein proteins in the bases of cilia and thereby induce the mechanical cilia movement. The complete set of equations making up the proposed model is presented in the Methods section. Below we provide a number of new insights into the mechanism of cilia regulation via a detailed investigation of the properties of this model.

### The mechanism of $\text{Ca}^{2+}$ -dependent inhibition of $\text{Ca}^{2+}$ channels

A subset of intraciliary  $\text{Ca}^{2+}$  channels have been reported to operate in an intraciliary  $\text{Ca}^{2+}$  dependent manner and have been proposed as major regulators of ciliary beat [49-51]. It is established that  $\text{Ca}^{2+}$  current is not inhibited by the double pulse application of depolarization impulses under voltage clamp conditions in those situations when the first transmembrane potential shift is equal to the equilibrium  $\text{Ca}^{2+}$  potential (+120 mV) [66]. Further experimental evidence reveals that  $\text{Ca}^{2+}$  current inactivation kinetics are delayed when  $\text{Ca}^{2+}$  ions are partially replaced by  $\text{Ba}^{2+}$  ions [67-71]. Altogether these findings suggest that the channels are not inhibited directly by the depolarizing shift of transmembrane potential, but that instead their conductivity is dependent on the intraciliary  $\text{Ca}^{2+}$  concentration. Some decrease of the inward current amplitude (by approximately 25%) upon transmembrane potential shift into the  $\text{Ca}^{2+}$  equilibrium level can be explained by the fact that  $\text{K}^+$  currents can contribute to the overall current measurements. Here we consider the intraciliary  $\text{Ca}^{2+}$  concentration-dependent  $\text{Ca}^{2+}$  channel inhibition and employ the developed model to analyse two potential scenarios for the  $\text{Ca}^{2+}$  channel conductivity regulation. In one case,  $\text{Ca}^{2+}$  ions bind to the  $\text{Ca}^{2+}$  binding site on the channel and thereby inhibit the channel's conductivity by direct interaction. The other possibility is that the  $\text{Ca}^{2+}$  binding protein interacts with the  $\text{Ca}^{2+}$  ion first and then this complex binds to the channel and inhibits its conductivity. In both cases the conductivity dependence on transmembrane potential is assumed to be monotonic according to the experimental data [66].

### Direct $\text{Ca}^{2+}$ -dependent $\text{Ca}^{2+}$ channel conductivity inhibition

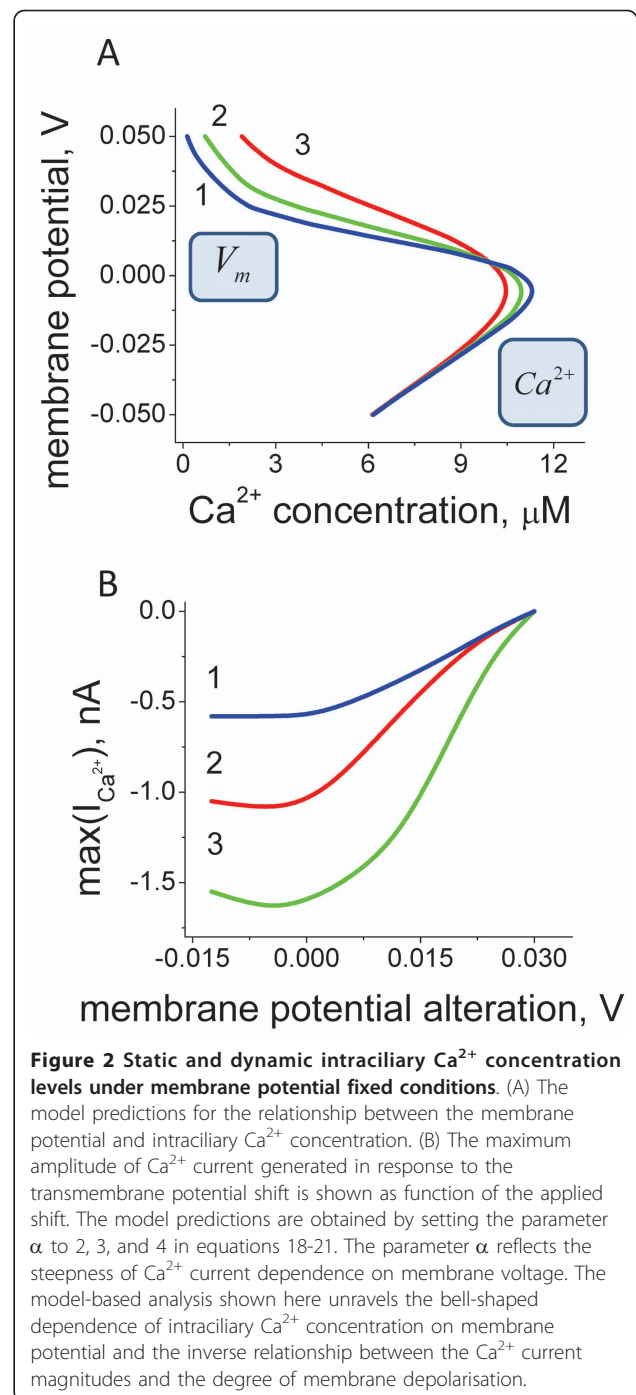
We first investigate a potential intraciliary  $\text{Ca}^{2+}$  regulatory mechanism via direct  $\text{Ca}^{2+}$  ion binding-dependent  $\text{Ca}^{2+}$  channel inhibition. The relationship between the  $\text{Ca}^{2+}$  channel conductivity and the transmembrane potential has been experimentally characterized by an early study in *Paramecium* species [66]. The equation

(18) in the Methods section approximates the experimentally established dependence. In the case of direct  $\text{Ca}^{2+}$ -dependent  $\text{Ca}^{2+}$  channel inhibition, one can show that the nullclines for non-dimensional  $\text{Ca}^{2+}$  concentration ( $\frac{d\text{Ca}^{2+}}{dt} = 0$ ) and for the number of open channels ( $\frac{dn}{dt} = 0$ ) from the system of differential equations (20) intersect at one stable point for all values in the physiological range of model parameters. The numerical solutions of the coupled differential equations (20) allow us to obtain the solutions for how the steady-state  $\text{Ca}^{2+}$  levels depend on the transmembrane potential. The model predictions for the steady-state ciliary  $\text{Ca}^{2+}$  channel conductivity dependence on the transmembrane potential under the voltage clamp conditions are shown on Figure 2A. The extracellular conditions are subject to constant change both in the case of ciliates as well as for multicellular organisms. The modifications in the external environment continuously shift the transmembrane potential. In order to estimate how the transmembrane potential alterations affect the inward  $\text{Ca}^{2+}$  current we derived the dependence for the  $\text{Ca}^{2+}$  ion flow (equation (21) in the Methods section). The experimental data-based (Figure 2A) [66] model for the inward  $\text{Ca}^{2+}$  current dependence on membrane potential (Figure 2B) predicts a significant reduction of the inward  $\text{Ca}^{2+}$  current amplitude as a function of membrane depolarization.

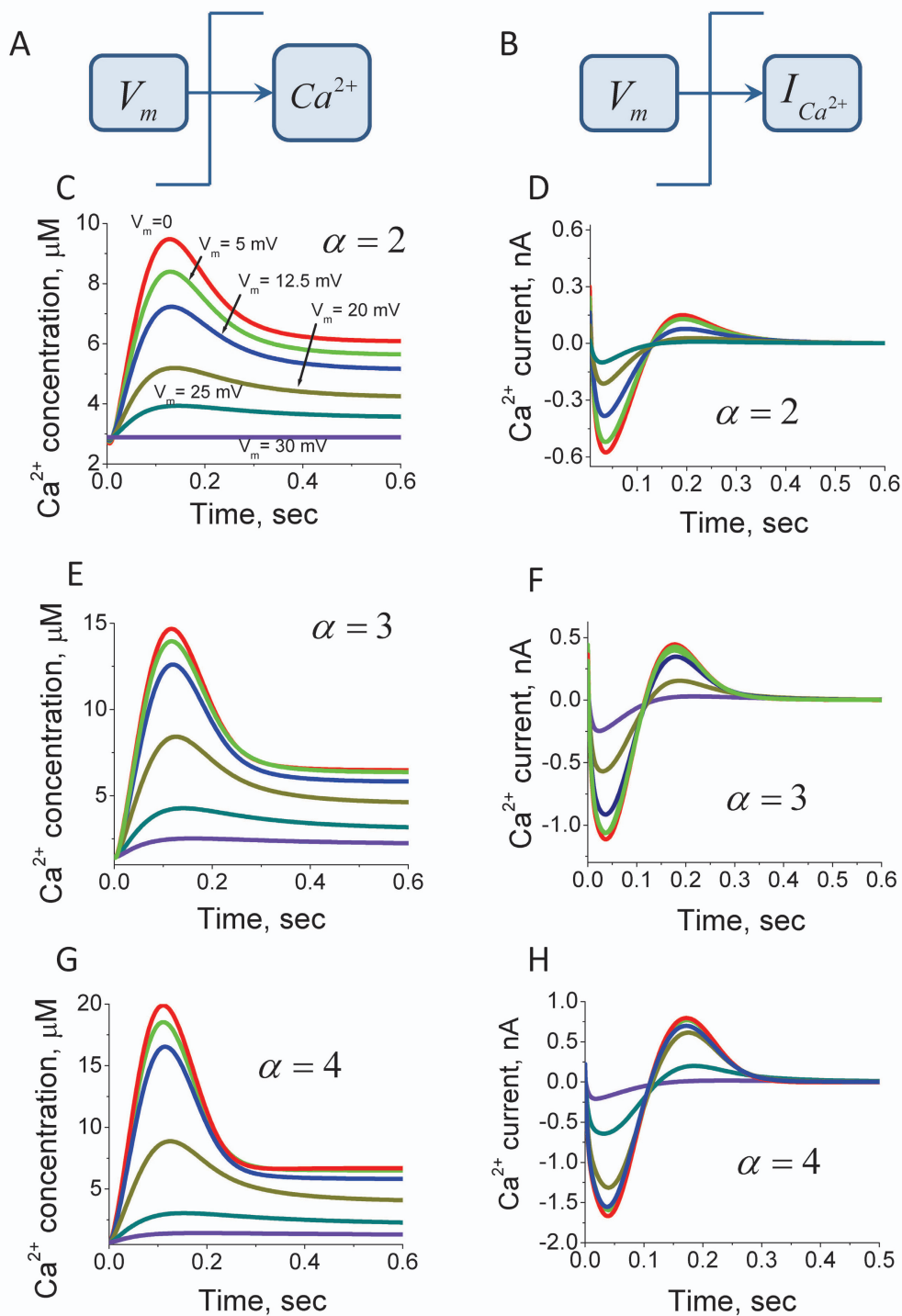
We next investigated the dynamic alterations of the intraciliary  $\text{Ca}^{2+}$  concentrations and the inward  $\text{Ca}^{2+}$  current in response to the transmembrane potential shifts. Equations (20) and (21) in the Methods section were used for quantitative estimations of the  $\text{Ca}^{2+}$  concentration and current responses to the normalised membrane potential shifts. The model predicts that the rapid depolarising alterations of the ciliary transmembrane potential leads to the generation of single  $\text{Ca}^{2+}$  spikes (Figure 3). Such responses can take place when  $\text{Ca}^{2+}$  channels are inhibited by  $\text{Ca}^{2+}$  ions and the channel's conductivity depends on the membrane potential in a monotonic manner (equation 18). The amplitude of those impulses depends on the steepness of the  $\text{Ca}^{2+}$  channels conductivity dependence on the membrane potential. The mechanism of the  $\text{Ca}^{2+}$  spike generation is mainly due to the delay of the  $\text{Ca}^{2+}$ -induced inhibition with respect to the  $\text{Ca}^{2+}$  conductivity alteration characteristic times. The described mechanism of the  $\text{Ca}^{2+}$  spike generation will only work if the  $\text{Ca}^{2+}$  currents significantly alter the  $\text{Ca}^{2+}$  concentration in the intraciliary compartments.

#### Indirect $\text{Ca}^{2+}$ channel conductivity regulation

In the previous section we considered  $\text{Ca}^{2+}$ -dependent  $\text{Ca}^{2+}$  channel regulation under the assumptions that



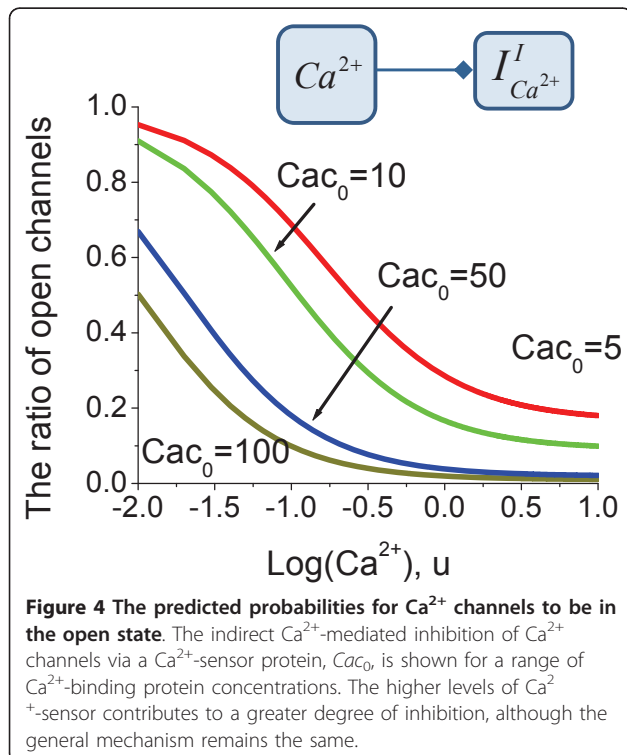
$\text{Ca}^{2+}$  channels have an intracellular  $\text{Ca}^{2+}$  binding site and  $\text{Ca}^{2+}$  ion binding closes the channels. However, several experimental studies have suggested that the conductivity of  $\text{Ca}^{2+}$  channels in cilia can also be regulated indirectly, via a  $\text{Ca}^{2+}$  binding protein. At present, there is no direct experimental evidence that explicitly favours either direct or indirect regulatory mechanism. We, therefore, investigated the second possibility for



**Figure 3 Systems model predictions for intraciliary  $Ca^{2+}$  concentration and current responses to the transmembrane potential shift.** The intraciliary  $Ca^{2+}$  concentrations (A) and  $Ca^{2+}$  currents (B) are calculated according to the model with direct  $Ca^{2+}$ -mediated ciliary  $Ca^{2+}$  channels inhibition in response to variable degree of transmembrane potential shift. The responses are colour coded according to the degree of applied transmembrane potential shift. The violet and red coloured lines represent dynamic responses obtained after the smallest and the largest depolarising shift of membrane potential, respectively. The non dimensional membrane potential values following voltage shift from initial  $\psi_0 = -1.2$  are shown in (C) and remain the same throughout the figure. The calculations suggest that ciliary membrane depolarisation induces an intraciliary  $Ca^{2+}$  spike over a wide physiological range of depolarising conditions (C, E and G), whereas the current generates a biphasic response (D, F and H). The calculations were carried out for three different levels of parameter  $\alpha$ , which reflects the steepness of the  $Ca^{2+}$  channels conductivity dependence on membrane potential.

indirect  $\text{Ca}^{2+}$ -dependent  $\text{Ca}^{2+}$  channels conductivity inhibition.

The model for  $\text{Ca}^{2+}$  ion interactions followed by the interactions with the  $\text{Ca}^{2+}$  channels (described in the Methods section) is developed in line with a previously suggested modelling methodology for  $\text{Ca}^{2+}$ -CaM interactions [63,64]. The model predictions for the number of open channels as a function of intraciliary  $\text{Ca}^{2+}$  concentration are shown on Figure 4 as derived by equation (25) in the Methods section. Given that the exact nature of the  $\text{Ca}^{2+}$  binding protein acting as a mediator between  $\text{Ca}^{2+}$  ions and  $\text{Ca}^{2+}$  channels is not established, equation (25) has been solved for different physiologically possible ratios of total number of channels to the dissociation constant for the  $\text{Ca}^{2+}$ -binding protein interactions with  $\text{Ca}^{2+}$  channels. According to the derived models, the comparison of direct versus indirect  $\text{Ca}^{2+}$  channels inhibition can be carried out by setting the non dimensional  $\text{Ca}^{2+}$  concentration to zero ( $u = 0$  in equation (25)). In this case, the predictions of equation (25) for indirect  $\text{Ca}^{2+}$  channels conductivity inhibition almost coincide with the model for direct  $\text{Ca}^{2+}$ -dependent inhibition. In both models most of the  $\text{Ca}^{2+}$  channels are open in the lower range of  $\text{Ca}^{2+}$  concentrations. However, the model for indirect  $\text{Ca}^{2+}$  channels inhibition predicts that the number of open  $\text{Ca}^{2+}$  channels would equal  $\frac{1}{\text{Cac}_0 + 1}$  when intraciliary  $\text{Ca}^{2+}$  reaches

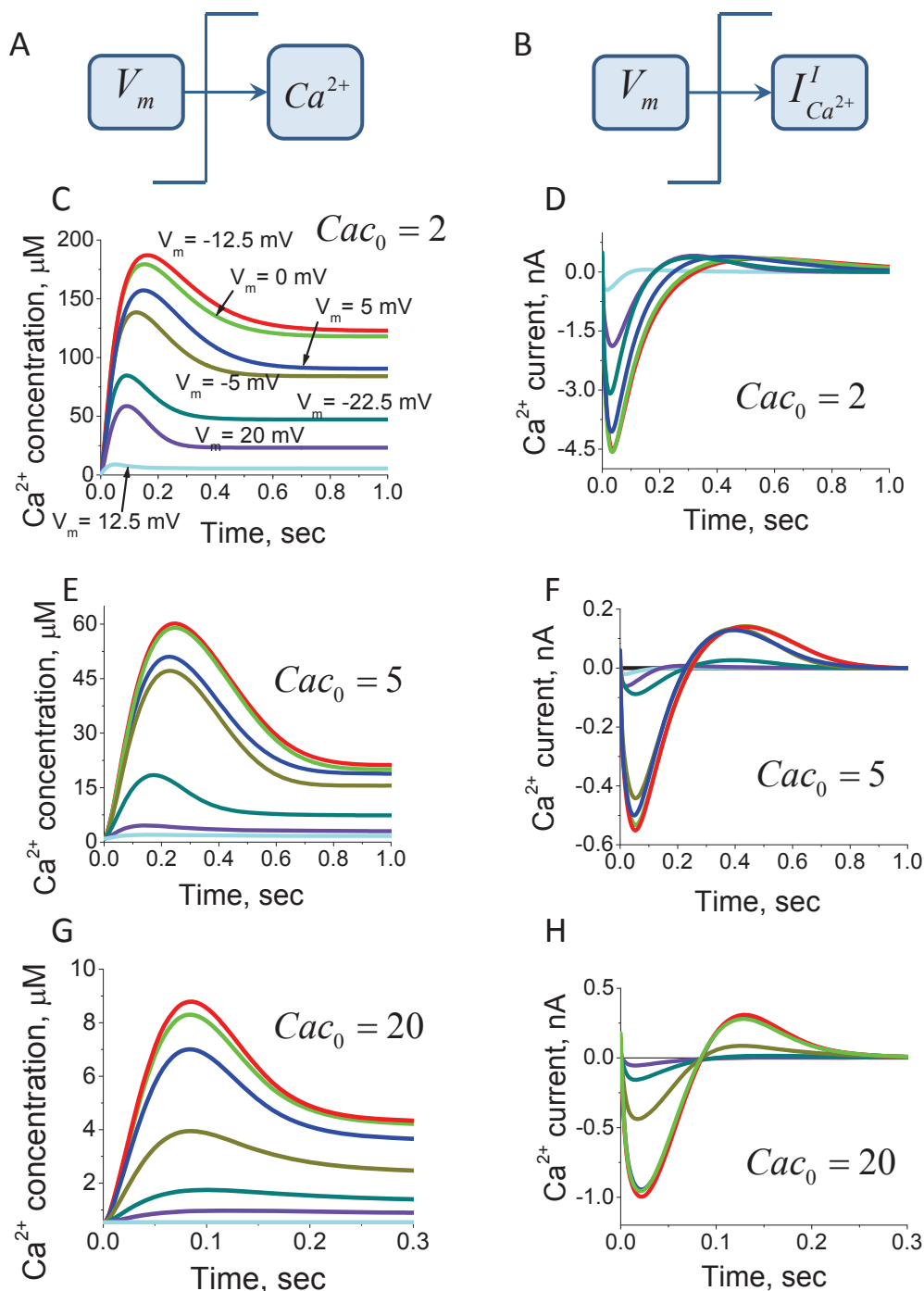


high concentrations. In other words, in the case of the indirect mechanism of inhibition, high  $\text{Ca}^{2+}$  does not inhibit the channels completely. The number of remaining channels in the open state would depend on the concentration of the regulatory  $\text{Ca}^{2+}$  binding protein.

We next analysed the dynamics of the ciliary  $\text{Ca}^{2+}$  channels inhibition in response to a change in intraciliary  $\text{Ca}^{2+}$  concentration and in the transmembrane potential. As mentioned earlier, the  $\text{Ca}^{2+}$  channel conductivity is not inhibited by the membrane potential, but rather has a monotonic dependence on the transmembrane potential difference as shown on Figure 2B. We found that the characteristic time  $\tau_{\text{Ca}^{2+}}$  of  $\text{Ca}^{2+}$  channel alterations in response to step changes in  $\text{Ca}^{2+}$  is inversely proportional to the total number of channels. The model predictions for intraciliary  $\text{Ca}^{2+}$  concentration and inward  $\text{Ca}^{2+}$  current in response to depolarising transmembrane potential changes from  $V_0$  to  $V_1$ , are shown on Figure 5. According to the derived equation (30) in the Methods section, the inward  $\text{Ca}^{2+}$  current changes with a characteristic time  $\tau_V$ , which can only be estimated by considering the  $\text{K}^+$  current contribution. In the first instance we only consider active and passive  $\text{Ca}^{2+}$  transport (equation (31) in Methods). We neglected by the kinetics for the channels of active transport due to the assumption that the kinetics of alterations of active transport are much faster than the characteristic alteration times of passive transport (equation (32) for the  $\text{Ca}^{2+}$  currents). The model predictions for the steady-state dependence of transmembrane potential on intraciliary  $\text{Ca}^{2+}$  concentration and the dynamic  $\text{Ca}^{2+}$  current amplitude dependence on the membrane potential are shown in Figure 6A and 6B, respectively.

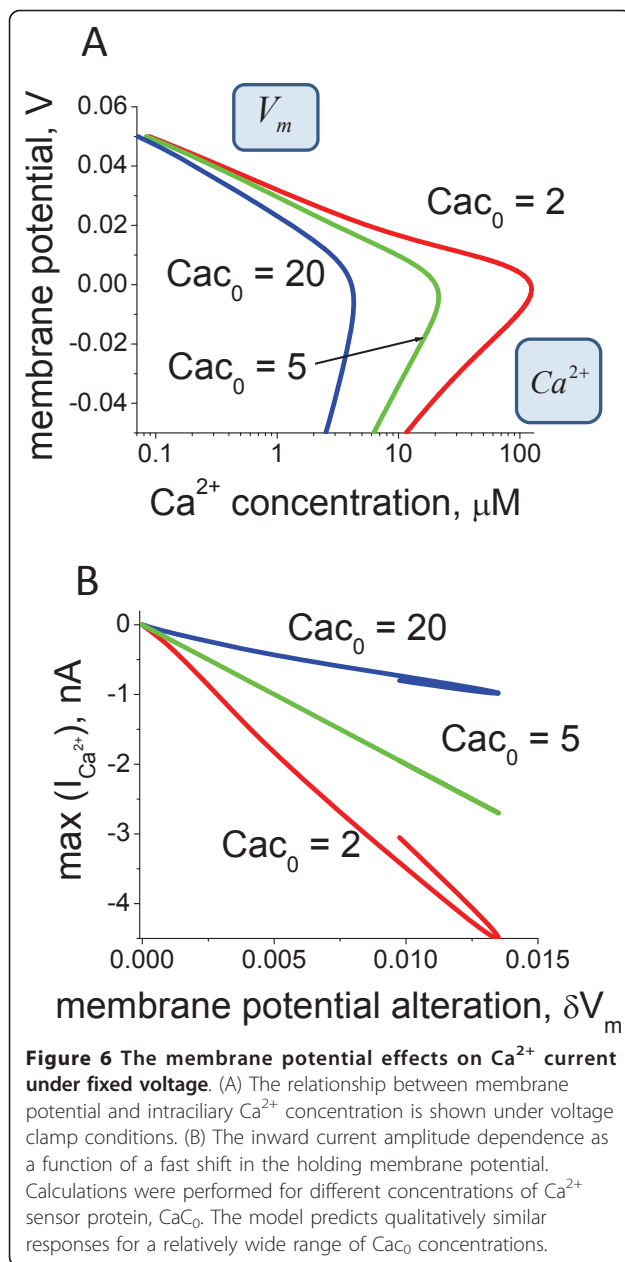
The cilia's external environment is subject to constant change and can significantly affect the behavioural responses of ciliates and modulate ciliary beating in multicellular organisms. In order to account for the effects of extraciliary  $\text{Ca}^{2+}$  variations we estimated the amplitudes of intraciliary  $\text{Ca}^{2+}$  spike generation under different external  $\text{Ca}^{2+}$  concentrations and variable transmembrane potentials. Figure 7A shows that the increase of the  $\text{Ca}^{2+}$  concentration in the external solution increases the amplitude of the generated intraciliary  $\text{Ca}^{2+}$  spike. The amplitude of the spikes goes to zero when the membrane potential equals the equilibrium potential for  $\text{Ca}^{2+}$  ions. The increase of membrane potential decreases the amplitude of the  $\text{Ca}^{2+}$  current (Figure 7B). This effect is due to the increase of the steady-state intracilia  $\text{Ca}^{2+}$  level and  $\text{Ca}^{2+}$ -dependent inhibition of the  $\text{Ca}^{2+}$  channels.

We noted earlier that there is a  $\text{Ca}^{2+}$  current in the cilia which transfers ions from the cilia into the cellular compartments. This current can be described by

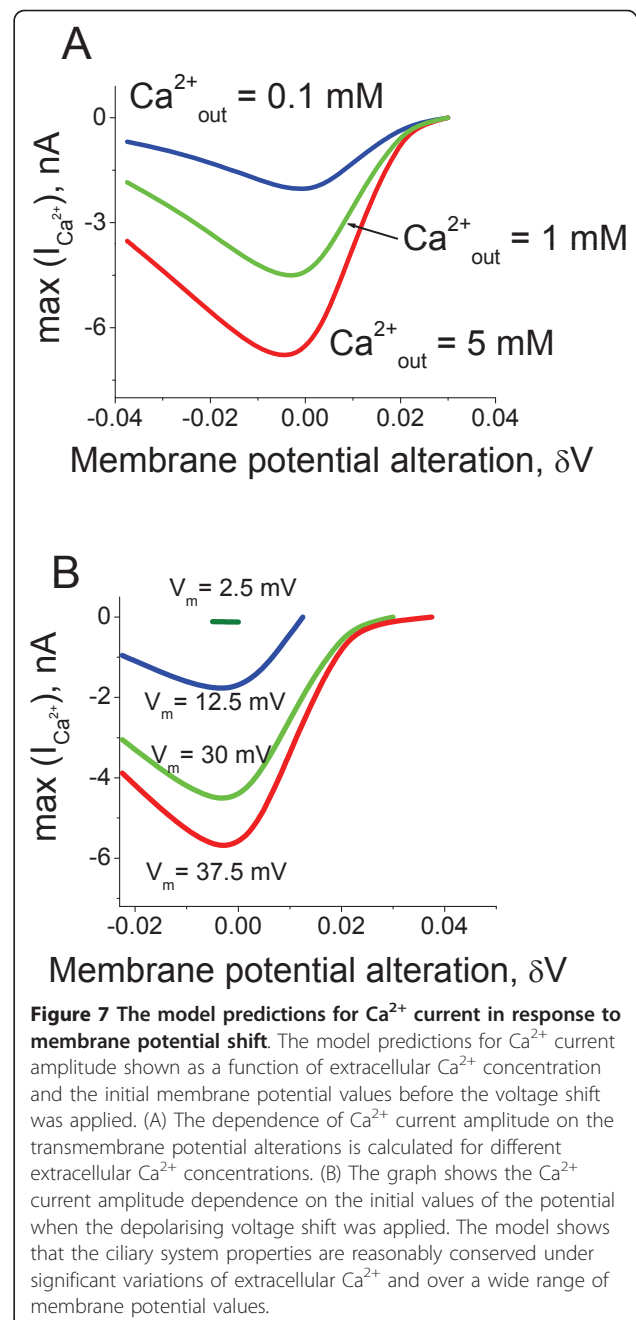


**Figure 5** The responses of  $\text{Ca}^{2+}$  current and concentration to the potential shift via indirect  $\text{Ca}^{2+}$  inhibition. The intracellular  $\text{Ca}^{2+}$  concentration (A) and  $\text{Ca}^{2+}$  current (B) responses to membrane depolarisation are computed under the assumption of indirect  $\text{Ca}^{2+}$  channel inhibition for different concentrations of  $\text{Ca}^{2+}$ -sensor protein,  $\text{Ca}c_0$ . The model predicts that in response to fast membrane depolarisation, the ciliary coupled  $\text{Ca}^{2+}$  and membrane potential system generates an intracellular  $\text{Ca}^{2+}$  concentration spike similar to the one observed under the direct  $\text{Ca}^{2+}$ -inhibition assumptions. The non dimensional membrane potential values following voltage shift from the initial  $\psi_0 = -1.2$  are shown in (C) and are the same throughout the figure. The major difference between direct (Figure 3) and indirect mechanisms of inhibition occurs in the steady-state levels of  $\text{Ca}^{2+}$  concentration that takes place after the transitional process. The comparison of intracellular  $\text{Ca}^{2+}$  concentration responses for increasing levels of the  $\text{Ca}^{2+}$  sensor protein (C, E and G) predicts that intracellular  $\text{Ca}^{2+}$  concentrations following the voltage shift are inversely dependent on the  $\text{Ca}^{2+}$  sensor concentration,  $\text{Ca}c_0$ . Different concentrations of  $\text{Ca}^{2+}$  sensor protein do not affect the membrane depolarisation-induced  $\text{Ca}^{2+}$  current (D, F and H).





equation (11) in Methods. The contribution of cilium-to-cell body current to the intracellular  $\text{Ca}^{2+}$  concentration dynamics was evaluated experimentally in [72,73]. It was shown that under depolarized membrane potential conditions the contribution of this current is very small and the intracellular  $\text{Ca}^{2+}$  is mainly pumped out of the cilia into the extracellular space by the active  $\text{Ca}^{2+}$  transport. According to other observations,  $\text{Ca}^{2+}$  current from cilia into the cellular compartment can be larger than the current generated by the active  $\text{Ca}^{2+}$  transport. In order to investigate the role and contribution of the cilia-to-cell compartment current, we introduced its contribution to the intracellular  $\text{Ca}^{2+}$  concentration



dynamics (equation (34)). We performed qualitative analysis of the  $\text{Ca}^{2+}$  concentration alterations in the cilia in the presence of the cilium-to-cell current and compared the  $\text{Ca}^{2+}$  dynamics with the case when this current was not present. We found that although the cilium-to-cell body current influences the intracellular  $\text{Ca}^{2+}$  concentration levels, it does not change the dynamics qualitatively when the membrane potential is depolarized and fixed.

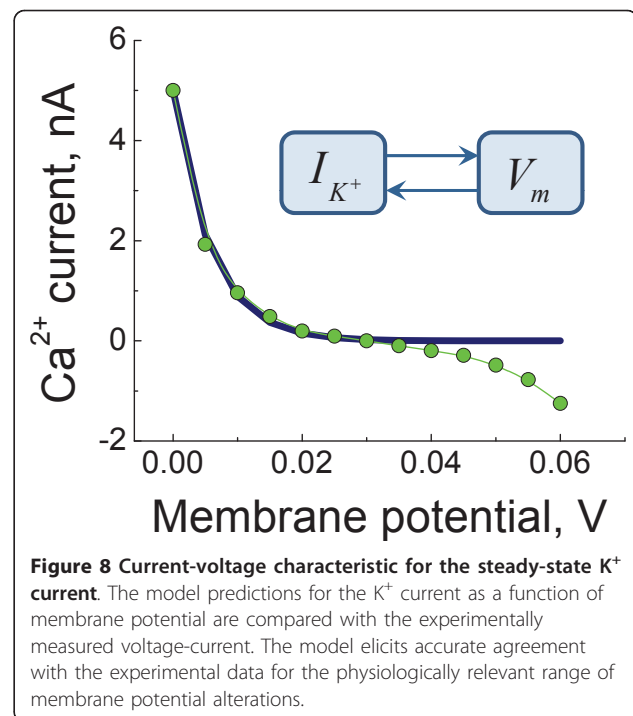
Our findings suggest that the cilium represents an excitable system with unique properties. The  $\text{Ca}^{2+}$ -dependent inhibition of  $\text{Ca}^{2+}$  channels inhibition

allows for the generation of single impulses of variable amplitude proportional to the degree of membrane depolarisation caused by variations in the external concentrations of ions. This system is able to generate a single spike despite unpredictable variations of ionic concentrations in the environment and is, therefore, very robust to alterations in the external conditions. Another interesting aspect of the ciliary excitation is the ability of the system to generate regulatory intraciliary  $\text{Ca}^{2+}$  impulses proportional to the degree of membrane depolarisation (Figures 3 and 5). This property can allow cells to sense and “automatically” respond to alterations in their environment.

### The contribution of $\text{K}^+$ currents

In the previous section, we analysed the dynamic properties of the intraciliary  $\text{Ca}^{2+}$  system under voltage clamp conditions. Several lines of evidence suggest that  $\text{K}^+$  currents contribute to the currents registered in cilia under voltage clamped conditions. The existence of  $\text{K}^+$  currents in cilia is supported by a number of experimental studies. The experimental data shows that the measured current is not equal to zero when the membrane potential equals the equilibrium membrane potential for  $\text{Ca}^{2+}$  ions. Instead, the current equals zero when membrane potential is about 10 mV while the equilibrium potential for  $\text{Ca}^{2+}$  ions equals 120 mV [66]. This observation suggests that both  $\text{Ca}^{2+}$  and  $\text{K}^+$  currents contribute to the overall current measured at early stages of current registration under voltage clamp, and therefore both currents need to be taken into the consideration in order to advance understanding of the mechanisms involved in ciliary regulation. At the same time, it has so far been impossible to register  $\text{Ca}^{2+}$  currents by inhibiting the  $\text{K}^+$  contribution. Various compounds can only partially block the  $\text{K}^+$  current when applied from inside of the membrane. Ciliary  $\text{K}^+$  currents have also been measured separately from  $\text{Ca}^{2+}$  currents.

In order to account for the contribution of  $\text{K}^+$  currents to the regulation of intraciliary  $\text{Ca}^{2+}$  concentration, we developed a model for the regulation of  $\text{K}^+$  currents by membrane potential (equation (35) in Methods). The dependence of  $\text{K}^+$  conductivity on the membrane potential is described by equation (36). Figure 8 shows the measured experimental values and the approximating curve calculated according to equation (36). The current on Figure 8 is normalized to 1 when membrane potential equals 0. In order to approximate the current, we only used the experimental values obtained under the membrane depolarized conditions, and this experimental data is approximated by equation (36). The predictions for intraciliary  $\text{Ca}^{2+}$  dynamics and  $\text{Ca}^{2+}$  impulse amplitude in response to the shift in membrane potential are shown in Figure 9A and 9B, respectively. A comparison

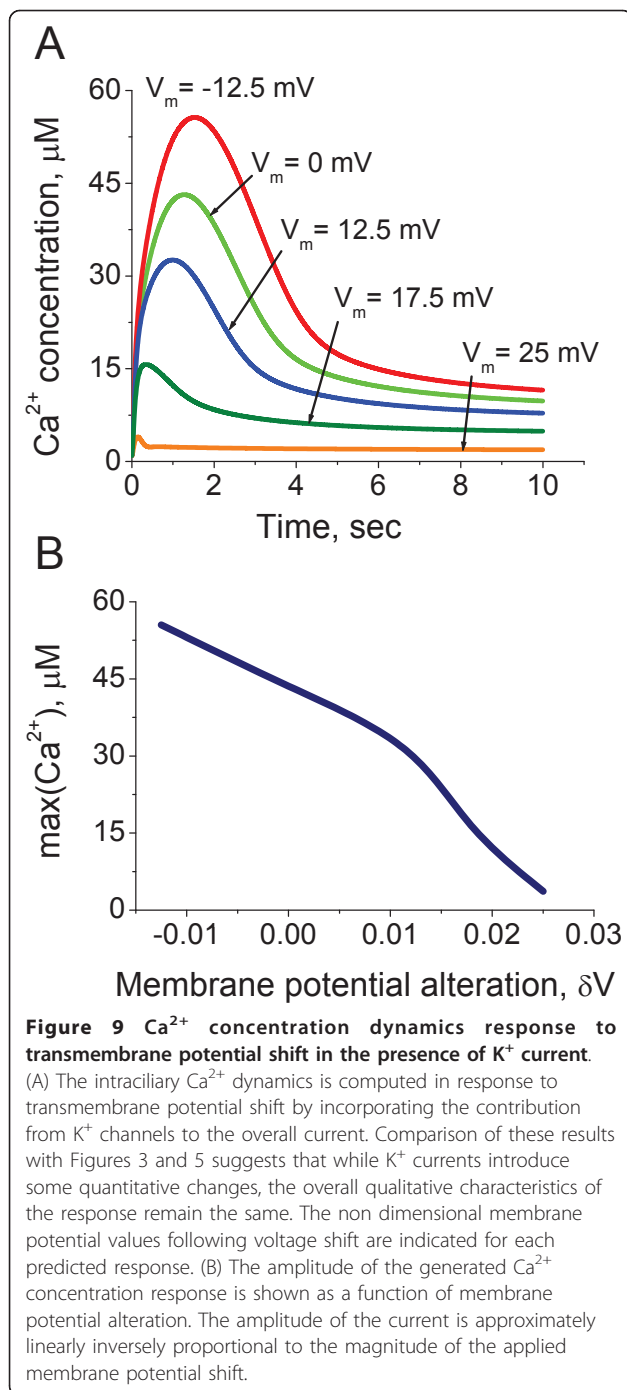


of the membrane potential shift-induced intraciliary  $\text{Ca}^{2+}$  spike generation in the presence of the  $\text{K}^+$  currents (Figure 9A) with the model when  $\text{K}^+$  channels are not incorporated (Figure 3 and 5), suggests that while the  $\text{K}^+$  currents change the quantitative values of the  $\text{Ca}^{2+}$  concentration dynamics, the general shape of the  $\text{Ca}^{2+}$  response remains the same.

### The transmembrane potential dynamics in the absence of voltage clamp

In the previous sections we investigated the mechanisms of the transmembrane potential shift-dependent  $\text{Ca}^{2+}$  spike generation under voltage clamp conditions. However,  $\text{Ca}^{2+}$  currents themselves can alter the membrane potential. Here we incorporate the membrane potential dependence on  $\text{Ca}^{2+}$  currents and investigate the membrane potential dynamics in the absence of voltage clamp (equations (40) and (41)). The non dimensional  $\text{Ca}^{2+}$  concentration and membrane potential are described by equation (42).

The predictions for the individual  $\text{Ca}^{2+}$  and  $\text{K}^+$  current responses to various membrane potential shifts are shown in Figure 10. One can note that the dynamics of the responses significantly differs between the ones obtained under voltage clamp conditions and the situation when the membrane potential is not fixed, but dependent on the currents (Figures 3, 5 and 9). Voltage current characteristics for  $\text{Ca}^{2+}$ ,  $\text{K}^+$  and full currents calculated using the current amplitudes from the currents dynamics are shown on Figure 10H. One can note that



the current is an almost exponentially growing function of membrane potential.

The monotonic dependence of  $\text{Ca}^{2+}$  current on transmembrane potential, and simultaneous  $\text{Ca}^{2+}$ -dependent inhibition of  $\text{Ca}^{2+}$  channels, represents a classical problem of two interconnected variables: intraciliary  $\text{Ca}^{2+}$  and membrane potential. In this system, increasing  $\text{Ca}^{2+}$  current with transmembrane potential depolarisation represents a positive feedback loop mechanism, whereas

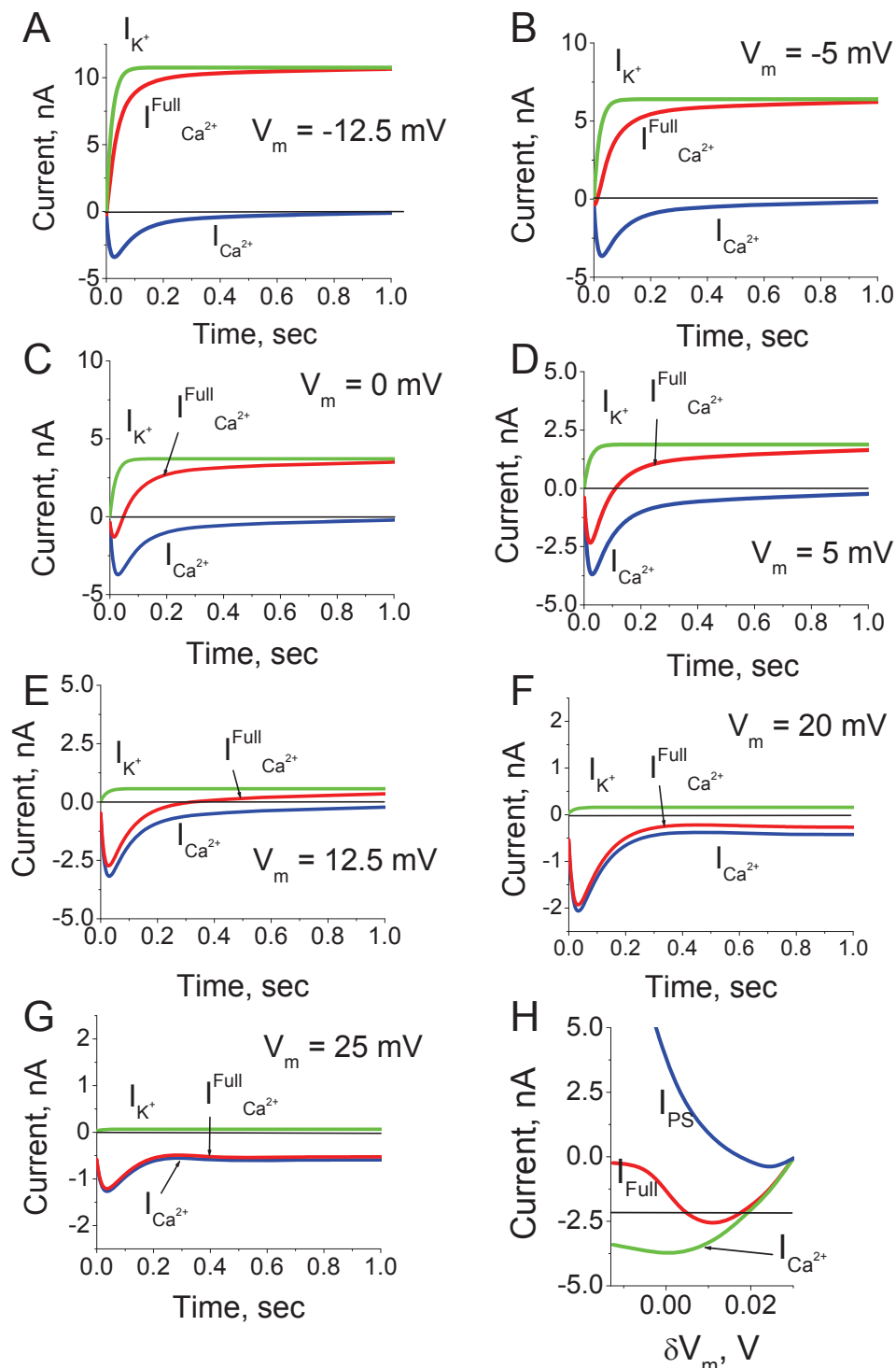
the intraciliary  $\text{Ca}^{2+}$  concentration-dependent  $\text{Ca}^{2+}$  channels inhibition represents a negative feedback loop. We, therefore, sought to investigate the range of potential dynamical properties of the ciliary system emerging from the coupling of  $\text{Ca}^{2+}$  current and membrane potential described by equations (42).

Figure 11 shows the  $\text{Ca}^{2+}$  current and membrane potential dynamics and the phase diagrams for an increasing range of inward current. We found that the inward current into the cilium can modify the dynamic properties of the  $\text{Ca}^{2+}$ -membrane potential system. In all cases, the null cline  $\frac{dV}{dt} = 0$  represents the N-shaped curve. In the physiological range of non-dimensional membrane potential  $V$  (50 mV, -0.025 mV) and intraciliary  $\text{Ca}^{2+}$  ( $u$ ) from 0.04 to 40  $\mu\text{M}$ , the null cline  $\frac{d\text{Ca}^{2+}}{dt} = 0$  shows a monotonic growth.

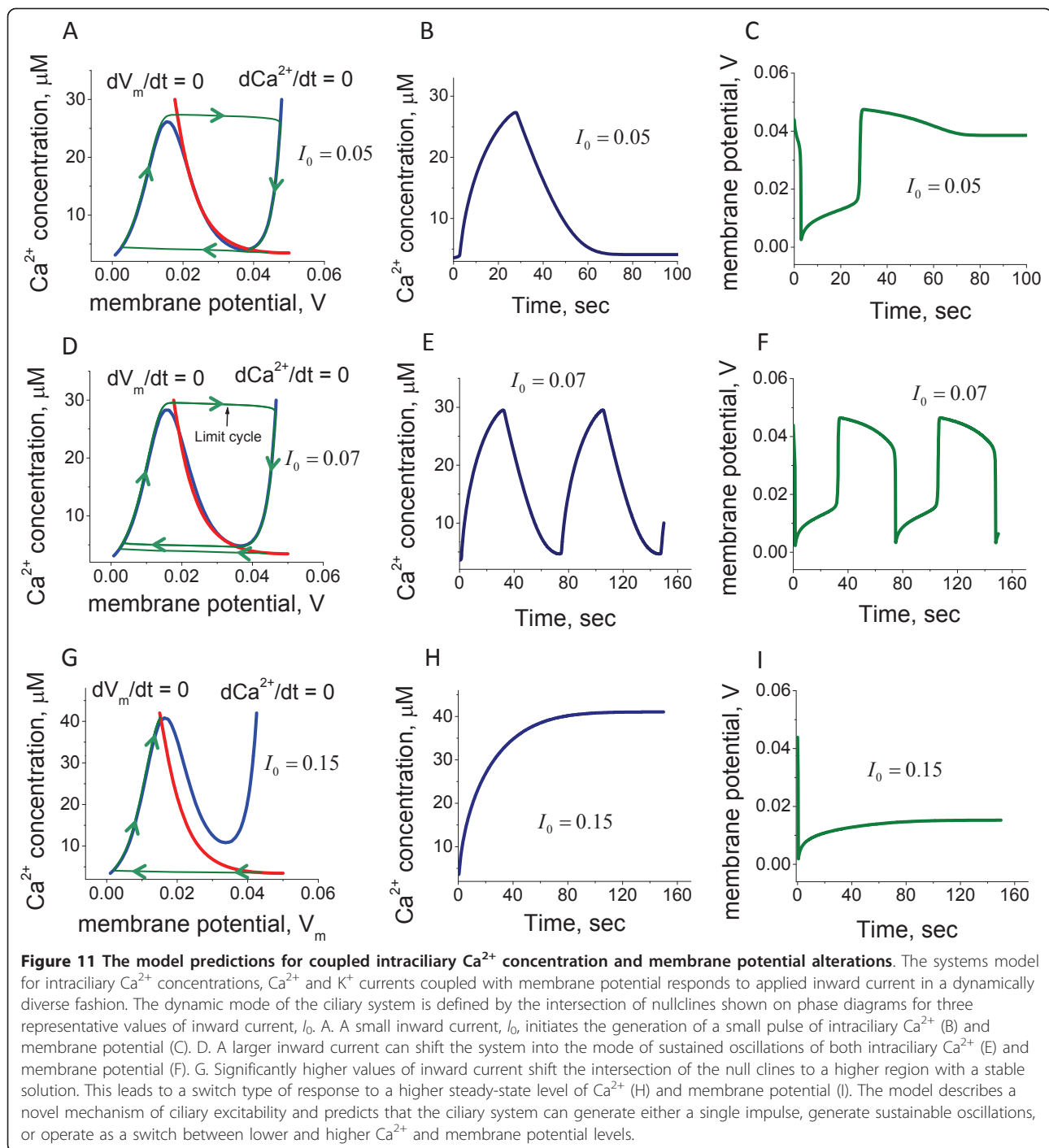
One can clearly see that there is significantly different response for different values of the inward current. When the influx of the ions is relatively small, the  $\frac{d\text{Ca}^{2+}}{dt} = 0$  null cline intersects the  $\frac{dV}{dt} = 0$  null cline in the left descending area (Figure 11A); such a null cline crossing results in a stable solution. In this case the system responds by the generation of a single impulse of both intraciliary  $\text{Ca}^{2+}$  concentration and the membrane potential followed by a return to homeostatic levels (Figure 11A, B and 11C). Further increasing the current causes the null cline  $\frac{d\text{Ca}^{2+}}{dt} = 0$  to intersect with the null

cline  $\frac{dV}{dt} = 0$  in the middle region of the ascending area, leading to an unstable solution with a limit cycle formed around the area that represents the oscillations. (Figure 11D, E and 11F). However, further increase of the current causes the  $\frac{d\text{Ca}^{2+}}{dt} = 0$  null cline to intersect with the null cline  $\frac{dV}{dt} = 0$  in the right descending area,

resulting in a stable solution with a slight increase of the homeostatic  $\text{Ca}^{2+}$  and membrane potential levels (Figure 11G, H and 11I). The key conclusion from this analysis is that the external ionic conditions can initiate essentially different dynamic properties of the system regulating ciliary movement. One of the key factors that affect the ciliary beat cycle is the level of intraciliary  $\text{Ca}^{2+}$ . Our findings suggest that in response to the external conditions, there are several possibilities for intraciliary  $\text{Ca}^{2+}$  upregulation. The system can generate a single spike (Figure 11B) of variable amplitude (data not shown), permanently increase  $\text{Ca}^{2+}$  in a dynamic fashion and maintain the high intraciliary levels (Figure 11E), or operate in a monostable multivibrator mode (cilia can



**Figure 10 The dynamics of current alterations in the absence of voltage clamp.** The  $Ca^{2+}$  and  $K^+$  current dynamics were calculated in response to the membrane potential shift in the absence of voltage holding conditions (A-G). The membrane potential was depolarised from the normalised value  $V_0 = -30$  mV to -12.5 mV (A) and 25 mV (G), which represent the smallest and the largest change shift, respectively. The currents elicited by the membrane potential alteration within the range are shown in (B)-(F).  $I_{Ca^{2+}}$ ,  $I_{K^+}$  and  $I_{Full}$  represent calcium, potassium and full currents, respectively. The comparison with the  $Ca^{2+}$  current responses to membrane potential shift obtained under the voltage clamped conditions (Figure 3, 5 and 9) reveals significant differences in the dynamics. (H) Voltage current characteristic is calculated using the current amplitudes from the currents dynamics for the full ( $I_{Full}$ ) and  $Ca^{2+}$  ( $I_{Ca^{2+}}$ ) currents. The steady-state voltage current relationship is calculated according to the stationary current values, ( $I_{PS}$ ).



generate a  $\text{Ca}^{2+}$  spike in response to any alteration of membrane potential) (Figure 11H). These three possibilities can be associated with the different modes of ciliary beat observed in human cilia as well as in various ciliates.

The dynamic properties of excitable systems with two interdependent variables are reasonably well understood at a theoretical level. In the present case,  $\text{Ca}^{2+}$  and

membrane potential represent the slow and fast variables, respectively. This study, therefore, establishes that the dynamic properties of ciliary systems, where the  $\text{Ca}^{2+}$  and  $\text{K}^+$  channel conductivities represent monotonic function of membrane potential and the  $\text{Ca}^{2+}$  channels conductivity inversely depends on intraciliary  $\text{Ca}^{2+}$  concentration, are comparable with the properties of excitable systems based on the “N-shape” dependence of the

$\text{Na}^{2+}$  channel conductivity on membrane potential [74]. At the same time, it is essential to note that the mechanism of excitation described in motile cilia is different from the “classical” one described in most excitable cells and systems that involve  $\text{IP}_3$   $\text{Ca}^{2+}$  channels [75,76].

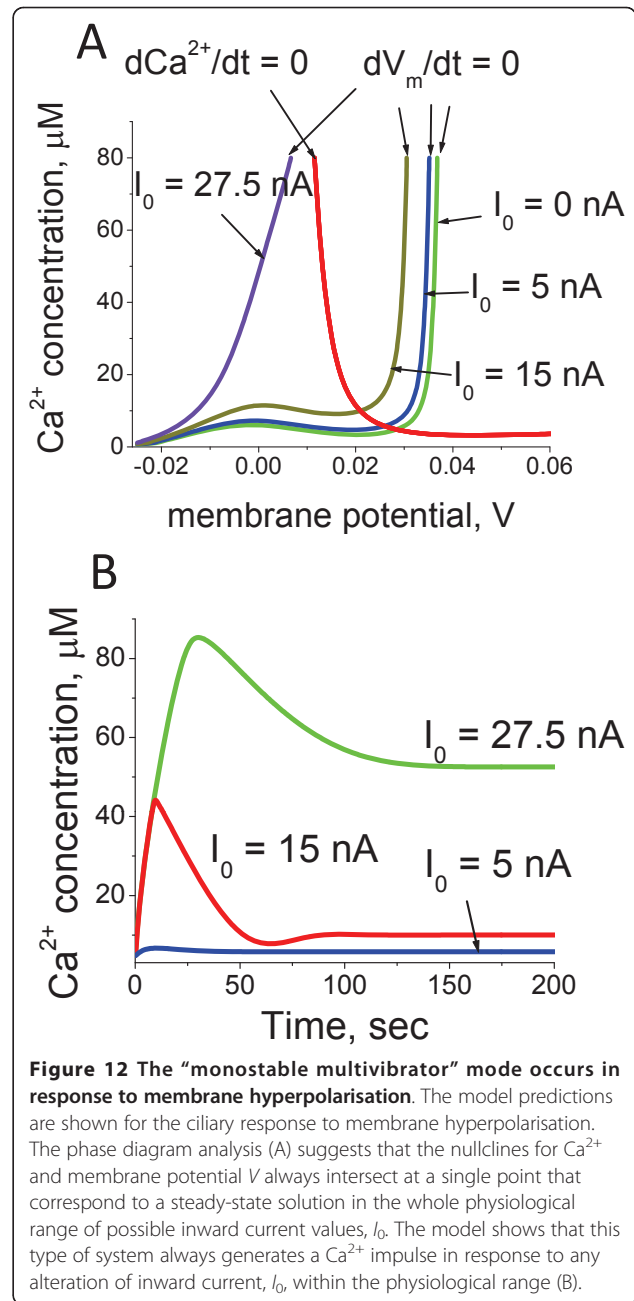
**The membrane hyperpolarisation-dependent currents modulate the excitatory properties of the ciliary system**

The ciliary transmembrane potential can shift in two directions. In the previous section we investigated the intraciliary  $\text{Ca}^{2+}$  responses caused by membrane depolarisation. Here we assess the implications of the membrane hyperpolarisation which has been shown to activate the current from cilia into the cell body [77,78]. We introduced the corresponding term into our model for the  $\text{Ca}^{2+}$  ions movement via the membrane as a function of the corresponding membrane potential shift (equation 43). By assuming the potential independent mechanism for  $\text{Ca}^{2+}$  and  $\text{K}^+$  ion expulsion, the system of intraciliary  $\text{Ca}^{2+}$  and membrane potential is derived as shown in equation (46) in the Methods section.

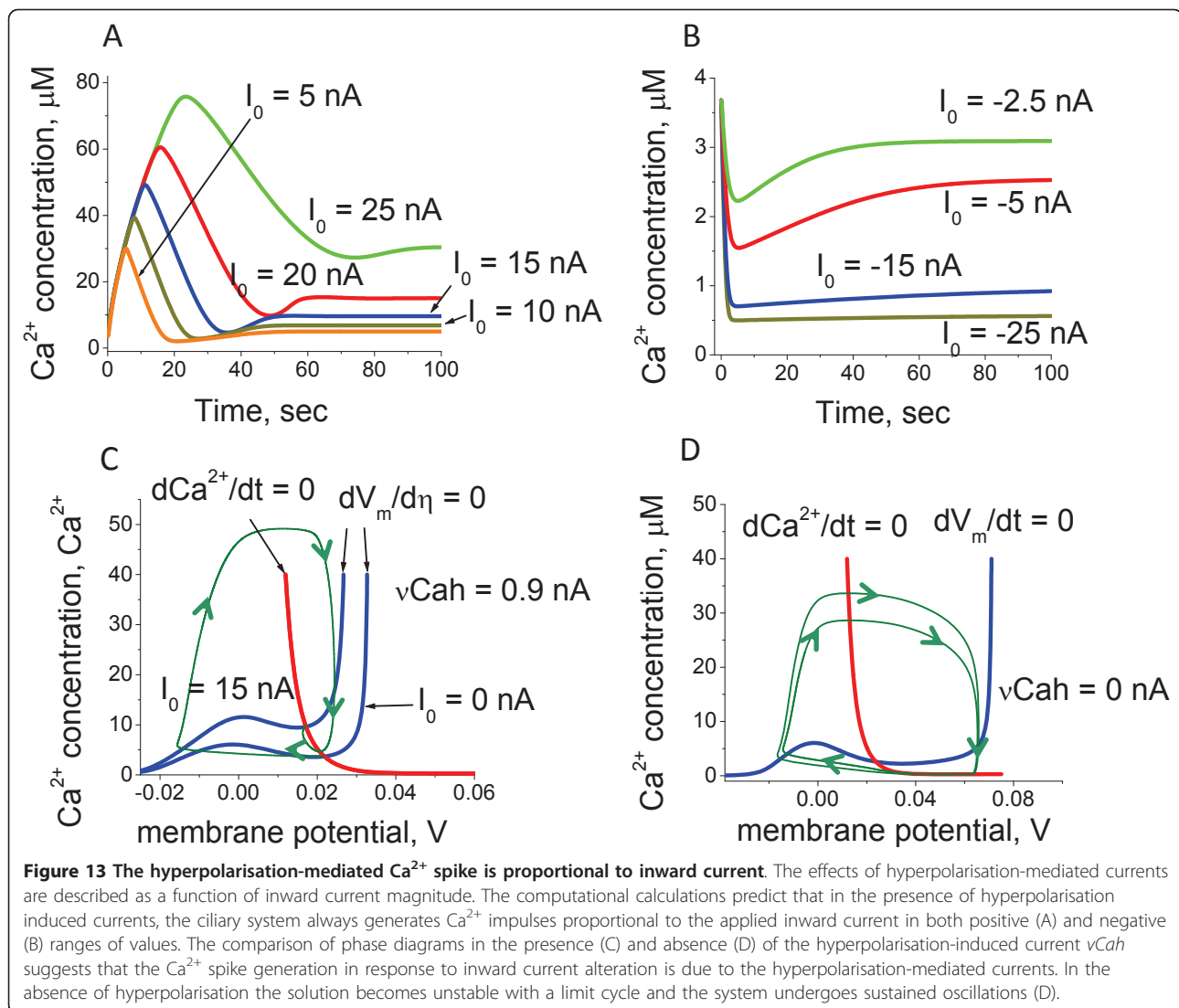
The model predictions for intraciliary  $\text{Ca}^{2+}$  and membrane potential dynamics in the absence of the hyperpolarisation-induced current are shown in Figure 11D, E and 11F. In this case the system remains in the mode of steady oscillations. We found that such oscillatory mode can be significantly modulated by the hyperpolarisation-induced and the external conditions-dependent current. Figure 12A shows the phase diagram in the presence of the hyperpolarising current. In contrast to the situation described on Figure 11 when the hyperpolarising current was absent, the ciliary system generates a single spike in response to alteration of external ionic concentrations over the whole physiological range of the inward current (Figure 12B and 12C). Under certain combinations of the hyperpolarisation-induced currents and the  $\text{Ca}^{2+}$ -dependent  $\text{K}^+$  currents the phase diagram modifies in a manner so that the system responds by generating a single spike in response to alteration of the inward current of any magnitude (Figure 13). This indicates that the ciliary system under membrane hyperpolarising conditions can become a monostable multivibrator.

**The role of cilia-to body  $\text{Ca}^{2+}$  current under membrane hyperpolarisation**

Intraciliary  $\text{Ca}^{2+}$  concentration has been experimentally estimated to be approximately one order of magnitude higher in comparison with the intracellular levels. The  $\text{Ca}^{2+}$  current generated by the ion flow from the ciliary compartment into the cell has been reported by a number of groups [28,52,53], but the role this current plays in the regulation of the ciliary beat remains unclear. In order to address this question, we introduced the term



for this current into the equation that describes the intraciliary  $\text{Ca}^{2+}$  concentration (equation (47) in the Methods section). In the absence of direct measurements of the dependence of conductivity on membrane potential we set the conductivity to increase in response to hyperpolarisation of membrane potential. A summary of the model responses for all channel conductivities as a function of membrane potential is shown in Figure 14. By numerically solving the coupled equations for intraciliary  $\text{Ca}^{2+}$  and membrane potential with the cilia to the cell body contribution, we found that this current



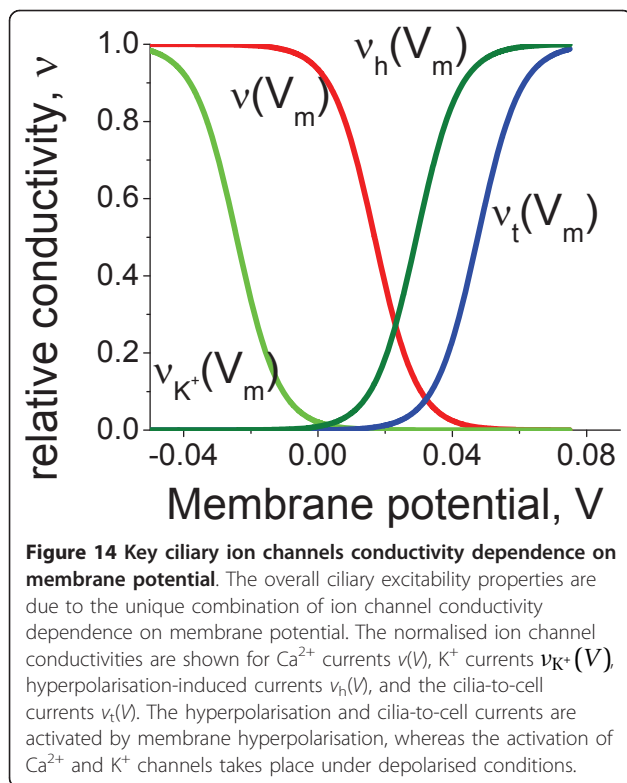
does not qualitatively change the general dynamic properties of the system.

Despite the lack of a noticeable contribution to the ciliary dynamic properties, this current requires a special consideration. Experimental studies have clearly demonstrated that intraciliary  $\text{Ca}^{2+}$  is significantly higher than intracellular  $\text{Ca}^{2+}$  concentration. At the same time, if the conductivity of protein structures governing the  $\text{Ca}^{2+}$  ions movement from cilia to the body is high, most of the intraciliary ions would move from cilia into the cell body in a very short time. A simple calculation suggests that if  $\text{Ca}^{2+}$  could freely flow from cilia into the body, the intraciliary concentration would become equal to the intracellular  $\text{Ca}^{2+}$  concentration in less than 100  $\mu\text{s}$  due to the difference in the volumes of the cell body and intraciliary compartments. Experimental measurements in ciliates show that the hyperpolarisation-induced backwards movements can last longer than 100  $\mu\text{seconds}$ . It

is also known that the avoidance reaction that requires long term elevation of intraciliary  $\text{Ca}^{2+}$  concentration can be observed in hyperpolarizing solutions. During all this time the intraciliary  $\text{Ca}^{2+}$  concentration can be several orders of magnitude higher than the intraciliary concentration. In this study, we have demonstrated that the steady-state  $\text{Ca}^{2+}$  current under the depolarized membrane potential conditions can only be reduced by the  $\text{Ca}^{2+}$ -dependent inhibition of  $\text{Ca}^{2+}$  channels. All these observations suggest that the  $\text{Ca}^{2+}$  removal from cilia to the cell body occurs in a membrane potential dependent manner.

#### The mechanism of $\text{Ca}^{2+}$ and cyclic nucleotide-dependent CBF regulation

In addition to intraciliary  $\text{Ca}^{2+}$  and  $\text{K}^+$  potassium levels being coupled with the membrane potential modulation, cyclic nucleotides contribute to the regulation of one of



the major ciliary beat parameters, frequency. Intraciliary  $\text{Ca}^{2+}$  levels activate a variety of adenylate cyclases (AC) and phosphodiesterases (PDE) that produce and hydrolyse cyclic nucleotides, respectively, and thereby modulate the intraciliary cAMP and cGMP levels. At the same time, cAMP and cGMP-dependent kinases phosphorylate dynein arms [45] in the bases of cilia and thereby induce the ciliary movement [79].

In a previous work we showed that the ciliary beat frequency can have a “double” bell shape dependence on  $\text{Ca}^{2+}$  concentration [65] due to the differential regulation of adenylate and guanylate cyclase isoforms in a  $\text{Ca}^{2+}$ -calmodulin (CaM) dependent manner [63,64]. One bell-shape was due to the cAMP production by a combination of AC and PDE, whereas another one was mediated by cGMP production and degradation. Our results proposed an explanation for seemingly conflicting experimental evidence suggesting that CBF can both decrease and increase with increasing  $\text{Ca}^{2+}$  concentration. Here we extend our previous analysis and describe the conditions when one or the other peak can be significantly reduced or even disappear. Figure 15A shows the model predictions for the ciliary beat frequency in comparison with the experimental data [65,80]. Recent studies demonstrated that hormones and pharmacological agents can regulate both function and structure of cilia by interfering with the cyclic nucleotide signalling pathways [81,82]. These findings

support the possibility for the development of novel therapeutic strategies for ciliary pathologies by modulating the dynamic mode of beating via ciliary membrane receptors [83].

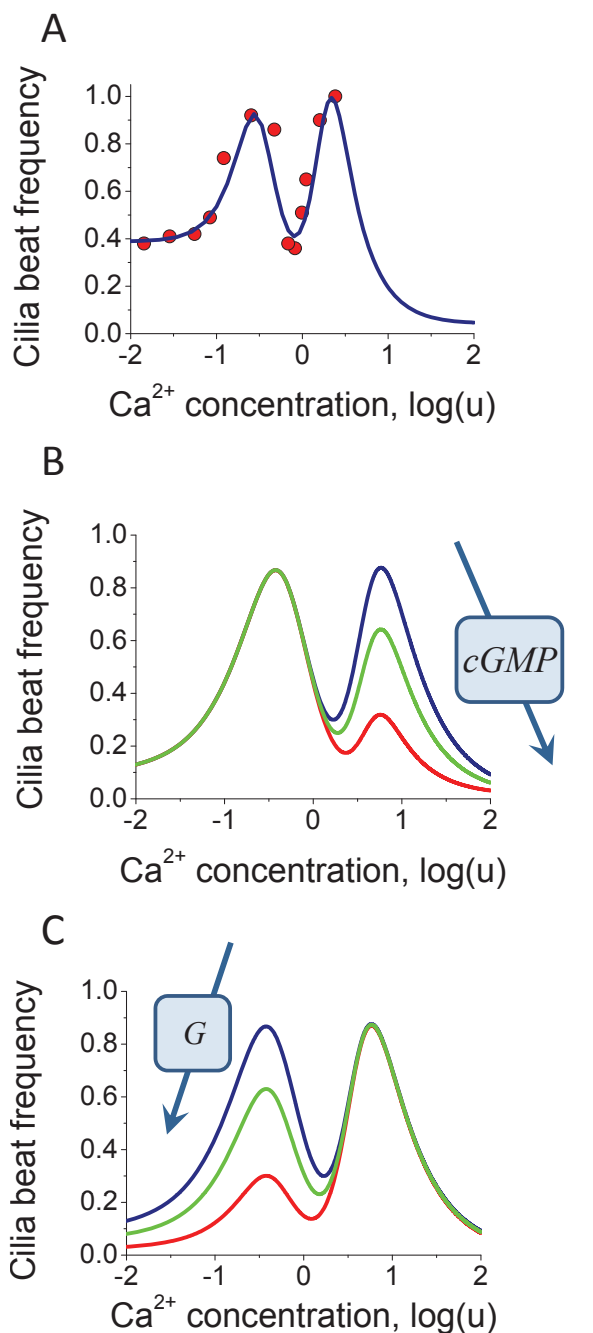
Figure 15B and 15C demonstrate that the “amplitude” of each peak can be significantly diminished if the activity of the AC or GC, respectively is modulated by a temporary or permanent, internal or external signal. Under such a scenario, CBF can only increase or decrease if it happens to be on one slope of the bell-shaped dependence. Therefore, according to our analysis, different organisms with the same underlying ciliary regulatory system can achieve all possible CBF regulatory modes as a function of  $\text{Ca}^{2+}$  concentration: the reverse bell-shaped dependence, if the “peak” values shown on Figure 15A occur at the lower and higher limits of the physiological range for  $\text{Ca}^{2+}$  concentration, the bell shape dependence that can be either cAMP and cGMP dependent, and either monotonic increase or decrease if the physiological range of  $\text{Ca}^{2+}$  concentrations occur at one of the slopes. Our model, therefore, describes the core  $\text{Ca}^{2+}$ -dependent regulatory mechanisms of cilia beat, but also provides an explanation for the differences observed between cilia in different single cell organisms as well as tissue specific differences. It also unravels the mechanism for how various stimuli modulate the rate of CBF by signalling via  $\text{Ca}^{2+}$ - and G-protein mediated pathways.

## Discussion

We develop a new computational model for  $\text{Ca}^{2+}$  and membrane potential-dependent ciliary regulation that explains how different ciliary beating regimes are regulated. The model describes a novel mechanism of excitability based on the membrane potential-dependence of  $\text{Ca}^{2+}$  currents (Figure 2) and simultaneous intraciliary  $\text{Ca}^{2+}$ -concentration mediated inhibition of  $\text{Ca}^{2+}$  channels (Figure 4). Our analysis shows that motile cilia constitute an excitable system with a novel mechanism of excitability. The ciliary system is able to generate a  $\text{Ca}^{2+}$  spike in response to a wide range of transmembrane depolarisation (Figure 3, 5 and 9). The major difference in the ciliary excitation described here, with respect to classical excitation mechanisms, is that ciliary excitability is robust to a wide range of ionic variations in the environment.

The excitability mechanism of cells in evolutionary advanced organisms is based on a combination of the N-shaped dependence of the quick inward cationic current on the transmembrane potential and slow alterations of the  $\text{K}^+$  conductivity [84-87]. The ciliary voltage-current characteristic (Figure 10H) suggests several functional dynamic modes of operation: i) single impulse generation, ii) oscillator, iii) trigger (Figure 11), all





**Figure 15 Predictions for ciliary beat frequency modulation.** (A) The comparison of experimental data and model predictions for the CBF dependence on  $\text{Ca}^{2+}$  concentrations. The experimental data for CBF [80] is shown as circles whereas the model prediction is represented as continuous blue line [65]. (B) GC activity depends on a number of intracellular mediators. The model describes the correlation of the right “peak” with the GC activity levels. (C) The model predicts that extracellular signals via G-protein pathway can significantly modulate the left “peak” of the plot of CBF dependence on  $\text{Ca}^{2+}$  concentration. In both cases (B and C) the diminishing amplitude is shown as green and red lines. The model describes the distinct contribution of both internal and external signals to CBF modulation.

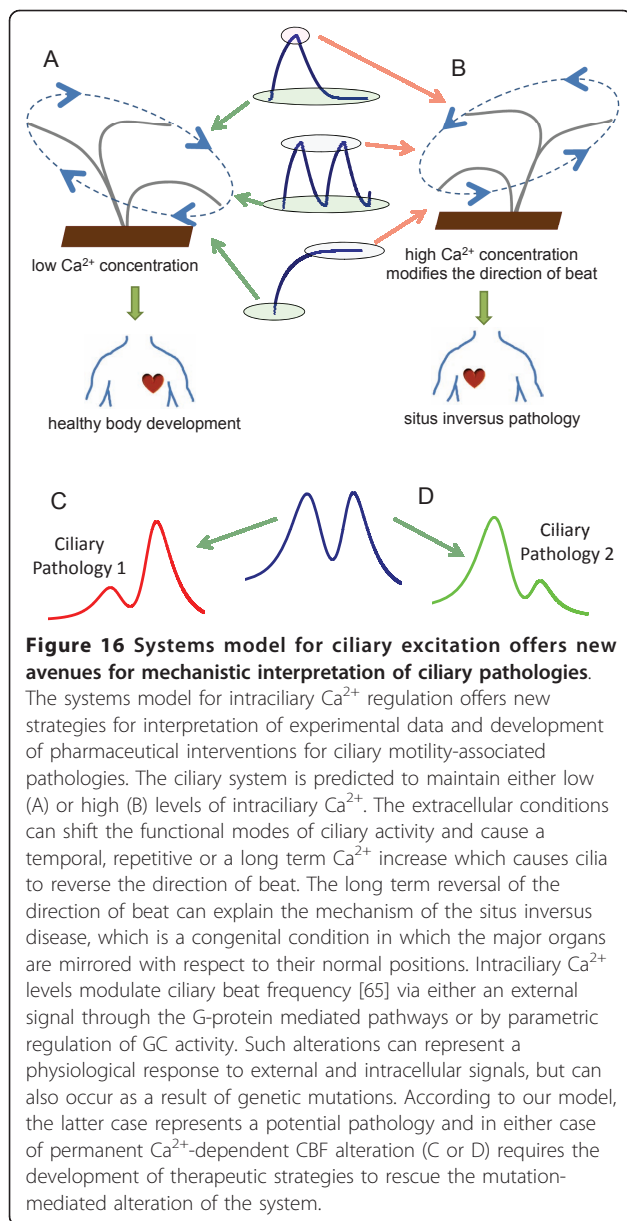
initiated by membrane depolarisation. At the same time, the hyperpolarisation-induced  $\text{Ca}^{2+}$  currents switch the system into the mode of a monostable multivibrator, when cilia can generate a  $\text{Ca}^{2+}$  spike in response to any alteration of membrane potential. The dynamics of such a system depends on the transmembrane potential. In other words, any alterations in the transmembrane potential (for example, initiated by variations of the external ion concentrations) switch functional performance of the system or make it non-excitable.

It was originally believed that  $\text{Ca}^{2+}$ , cAMP and cGMP each represent an independent pathway of ciliary regulation, however, there is by now a significant amount of evidence that strongly suggests that all three pathways are intimately interconnected [88]. It is well established that cAMP and cGMP are synthesized by AC isoforms and hydrolysed by PDEs in a  $\text{Ca}^{2+}$ -CaM-dependent manner. In this work we describe the mechanism of the cross talk between the three circuits and explain how CBF can be modulated via extra- and intraciliary pathways (Figure 15).

## Conclusions

### Therapeutic applications of systems model for intraciliary $\text{Ca}^{2+}$ regulation

Our detailed analysis of the effects of several  $\text{Ca}^{2+}$  and potassium currents and membrane potential on intraciliary  $\text{Ca}^{2+}$  levels offers a new way of interpreting ciliary motility associated pathologies. There are two main  $\text{Ca}^{2+}$ -mediated parameters that govern the motile function of cilia: the direction and the frequency of beat. Our model shows that there can be several dynamic regimes of intraciliary  $\text{Ca}^{2+}$  alterations during which the intraciliary  $\text{Ca}^{2+}$  concentration can be either at low or high levels, temporarily or for a significant period of time (Figure 16A). Experimental evidence suggests that high  $\text{Ca}^{2+}$  reverses the direction of cilia strike (Figure 16B), and modifies the frequency in a highly nonlinear manner (Figure 15) via synthesis and hydrolysis of cyclic nucleotides (Figure 15). In a previous work [89], we showed that genetic mutations that alter the dynamic properties of a system in a permanent manner can lead to disease. In this study, we propose a conceptually similar mechanism for the pathologies associated with ciliary motility that can take place in some human diseases [6]. Our study suggests that while CBF regulation in a multicellular organism can be modulated by a number of intracellular and extracellular factors (Figure 15), genetic mutations that directly or indirectly, affect the  $\text{Ca}^{2+}$ -mediated CBF dependence can dramatically impair the essential processes such as clearing function in airways, male fertility, or the determination of the left-right axis during development (Figure 16C and 16D). The treatment of the pathologies associated with this



mechanism would rely on the restoration of the original  $Ca^{2+}$ -dependent CBF dependence.

### Future perspective

At present there is limited understanding of the underlying biological mechanisms that govern ciliary motility. This study describes the modes of intraciliary  $Ca^{2+}$  dynamics in a highly detailed fashion. It shows the conditions that switch the system between the modes of  $Ca^{2+}$  spike generation, oscillatory dynamics and a trigger. The interdependent influences of  $Ca^{2+}$  and  $K^+$  currents, transmembrane potential and cyclic nucleotides modulate the

ciliary beat frequency and the direction of beat in a highly nonlinear manner. The further development of mathematical models of this system is still required to represent ciliary movements as a function of  $Ca^{2+}$  concentration and obtain the detailed understanding of ciliary motility which will be crucial for the development of new treatments for human diseases. While the core protein regulatory machinery involved in ciliary motility is very likely to be conserved, some variations in response to increased  $Ca^{2+}$  between single cell ciliates and mammalian cilia have been reported [90]. We would argue that those differences are not due to the change in the mechanisms of  $Ca^{2+}$ -dependent regulation but are rather caused by variations in the parameters of the regulatory circuits. The further investigation of single cell ciliates may allow a greater degree of characterisation of ciliary movement mechanisms, because in these systems alterations of ciliary motility translate into movement trajectories which can be easily observed.

## Methods

### Model Description

Figure 1 provides a schematic outline of the network regulating intraciliary  $Ca^{2+}$  concentration that is considered in our model. Intraciliary  $Ca^{2+}$  concentration is regulated by the currents of passive and active  $Ca^{2+}$  transport, as well as by  $Ca^{2+}$  leak into the extracellular space and into the cell body.

A basic mathematical model for intraciliary  $Ca^{2+}$  concentration and its relationship to transmembrane potential was proposed for the first time in [91]. A large number of recent experimental findings now allow the formulation of a more advanced model that includes the crucial aspects of the molecular mechanisms governing cilia movement. Below we describe the complete model for intraciliary  $Ca^{2+}$  regulation developed in this study.

The dynamics of intraciliary  $Ca^{2+}$  alteration are given by:

$$V_R \cdot \frac{d[Ca^{2+}]}{dt} = \frac{S_R}{z \cdot F} \cdot (I_{Ca^{2+}}^p + I_{Ca^{2+}}^a + I_{Ca^{2+}}^u) + I_{Ca^{2+}}^t + J([Ca^{2+}], [CaM_0]), \quad (1)$$

where  $V_R$  - is the cilium volume,  $S_R$  - is the cilium surface area, and  $I_{Ca^{2+}}^p$  and  $I_{Ca^{2+}}^a$  are the  $Ca^{2+}$  currents through the channels of passive and active  $Ca^{2+}$  transport, respectively.  $I_{Ca^{2+}}^t$  is the current from the cilium into the cell body.  $I_{Ca^{2+}}^u$  is the  $Ca^{2+}$  leakage current.  $J([Ca^{2+}], [CaM_0])$  is the function that encounters  $Ca^{2+}$  binding to and release from CaM, the main  $Ca^{2+}$  binding protein in cilia,  $z = 2$  is the  $Ca^{2+}$  ions charge, and  $F$  is the Faraday constant.

The dynamics of  $Ca^{2+}$  concentration alterations within the cilium are defined by the individual contribution of

the  $Ca^{2+}$  currents. The current via the channels of passive  $Ca^{2+}$  transport is given by:

$$I_{Ca^{2+}}^p = i_{Ca^{2+}}^p \cdot N_{Ca^{2+}}^p, \quad (2)$$

where  $N_{Ca^{2+}}^p$  is the density of functionally active  $Ca^{2+}$  channels.  $i_{Ca^{2+}}^p$  is the time averaged current via a single  $Ca^{2+}$  channel.  $Ca^{2+}$  channel activity is modulated externally by a number of metabolic pathways. Therefore, we only consider the pool of functionally active  $Ca^{2+}$  channels. The time-averaged current via a single  $Ca^{2+}$  channel is given by:

$$i_{Ca^{2+}}^p = \sum_i g_{Ca^{2+}}^i(V_m, Ca^{2+}) \cdot (V_m - E_{Ca^{2+}}), \quad (3)$$

where  $g_{Ca^{2+}}^i(V_m, Ca^{2+})$  is the conductivity of a single channel in the state  $i$  (in the most general state  $Ca^{2+}$  channels can have a number of states with different degrees of conductivity),  $E_{Ca^{2+}} = \left(\frac{R \cdot T}{2 \cdot F}\right) \cdot \ln\left(\frac{Ca_{out}^{2+}}{Ca_{in}^{2+}}\right)$  is the  $Ca^{2+}$  potential in the equilibrium,  $V_m$  is the transmembrane potential of the cilia membrane.

The  $Ca^{2+}$  leakage current is given by:

$$I_{Ca^{2+}}^u = g_{Ca^{2+}}^u \cdot (V_m - E_{Ca^{2+}}) \quad (4)$$

Assuming that the active  $Ca^{2+}$  transport system extrudes one  $Ca^{2+}$  ion per cycle, the current generated by the plasma membrane  $Ca^{2+}$  pump is given by:

$$I_{Ca^{2+}}^A = i_A \cdot N_{Ca^{2+}}^A, \quad (5)$$

where  $N_{Ca^{2+}}^A$  is the density of the plasma membrane  $Ca^{2+}$  pump protein complexes bound to one  $Ca^{2+}$  ion, and  $i_A$  is the time averaged  $Ca^{2+}$  current via a single  $Ca^{2+}$  channel. By introducing further assumptions that all the channels of active  $Ca^{2+}$  transport are saturated by ATP and that all bound  $Ca^{2+}$  molecules are released into the extracellular space, this current would be defined by the dynamics of the active transport channels bound to a  $Ca^{2+}$  ion:

$$\frac{dN_{Ca^{2+}}^A}{dt} = k_A^p \cdot [Ca_r^{2+}] \cdot (N_{Ca^{2+}}^{00} - N_{Ca^{2+}}^A) - (k_A^m + k_A^p) \cdot N_{Ca^{2+}}^A, \quad (6)$$

where  $N_{Ca^{2+}}^{00}$  is the density of the active  $Ca^{2+}$  transport channels,  $k_A^p$  and  $k_A^m$  are the association and dissociation constants for the  $Ca^{2+}$  ion interaction with the active  $Ca^{2+}$  transport channels, respectively.  $k_A^p$  is the constant that defines the  $Ca^{2+}$  ion transition from the bound state into the  $Ca^{2+}$  channel. By introducing new non-dimensional

variables:

$$\omega = \frac{N_{Ca^{2+}}^A}{N_{Ca^{2+}}^{00}}, \eta = n^m \cdot t, k_a = \frac{k_A^p \cdot K_{CaM}}{n^m}, k_b = \frac{k_A^m + k_A^p}{n^m}, u = \frac{Ca^{2+}}{K_{CaM}},$$

(where  $n^m$  is the dissociation constant of the protein regulating the passive  $Ca^{2+}$  transport channels), equation (6) takes the following form:

$$\frac{d\omega}{d\eta} = k_a \cdot u \cdot (1 - \omega) - k_b \cdot \omega, \quad (7)$$

The steady state-current through these active transport channels is given by:

$$I_{Ca^{2+}}^A = \tau \cdot \frac{u}{k_A + u}, \quad (8)$$

$$\text{where } k_A = \frac{k_A^m + k_A^p}{k_A^p \cdot K_{CaM}}, \tau = i_A \cdot N_{Ca^{2+}}^{00}, u = \frac{Ca^{2+}}{K_{CaM}}.$$

The model in [91] represented the leakage current from cilium into the cell body by the following expression:

$$i = \vartheta \cdot (Ca_r^{2+} - Ca_t^{2+}). \quad (9)$$

where  $\vartheta$  is the effective diffusion constant, and  $Ca_t^{2+}$  and  $Ca_r^{2+}$  are the intraciliary and intracellular  $Ca^{2+}$  concentrations, respectively. The Nernst equations allow a more accurate modelling of the leakage current from cilium into the cell body as follows:

$$I_{Ca^{2+}}^T = g_t(V_m) \cdot \left(V_{rt} - \frac{R \cdot T}{2 \cdot F} \cdot \ln\left[\frac{Ca_r^{2+}}{Ca_t^{2+}}\right]\right), \quad (10)$$

where  $g_t(V_m)$  is the overall conductivity of the cilium base area,  $V_{rt}$  is the difference of the potential between cell body and cilia,  $[Ca_r^{2+}]$  is the intraciliary  $Ca^{2+}$  concentration, and  $[Ca_t^{2+}]$  is the intracellular  $Ca^{2+}$  concentration. Equation (10) can then be represented in the following non dimensional form:

$$I_{Ca^{2+}}^T = \beta_1 \cdot \left(\psi_{rt} - 0.5 \cdot \ln\left(\frac{u}{u_t}\right)\right), \quad (11)$$

where

$$\beta_1 = \frac{g_t(V_m) \cdot R \cdot T}{F}, \psi_{rt} = \frac{V_{rt} \cdot F}{R \cdot T}, u = \frac{Ca_r^{2+}}{K_{CaM}}, u_t = \frac{Ca_t^{2+}}{K_{CaM}}.$$

In the following sections we derive the models and analyse the individual contributions of the different types of  $Ca^{2+}$  currents to the intraciliary  $Ca^{2+}$  homeostasis.

### Model for intraciliary $Ca^{2+}$ -dependent $Ca^{2+}$ channel conductivity inhibition

In this model we assume that  $Ca^{2+}$  channels located within cilia have a  $Ca^{2+}$  binding site on the intracellular site of the channel. According to such a model,  $Ca^{2+}$  ion binding to that site mediates the channel's transition into the closed state with no conductivity. We further assume that the characteristic time for the transition

from the conductive to the non conductive states is much smaller than the characteristic  $Ca^{2+}$  alteration times. In that case, the dynamics of the  $Ca^{2+}$  channels transition into the closed state in response to  $Ca^{2+}$  increase is given by:

$$\frac{d[N]}{dt} = -n^p \cdot [N] \cdot [Ca^{2+}] + n^m \cdot ([N_0] - [N]), \quad (12)$$

where  $N$  is the number of channels in the open state,  $N_0$  is the total number of channels,  $n^p$  and  $n^m$  are the association and dissociation constants for the interaction of  $Ca^{2+}$  ions with the  $Ca^{2+}$  channels, respectively. The steady-state solution of equation (12) is given by:

$$[N] = [N_0] \cdot \frac{K_C}{K_C + [Ca^{2+}]}, \quad (13)$$

where  $K_C = \frac{n^m}{n^p}$ . In the non dimensional form this solution is given by:

$$n = \frac{k_C}{k_C + u}, \quad (14)$$

$$\text{where } n = \frac{[N]}{[N_0]}, k_C = \frac{K_C}{K_{CaM}}, u = \frac{Ca^{2+}}{K_{CaM}}.$$

When  $u = 0$ ,  $n = 1$  and for  $u = \infty$   $n = 0$ , the current via these channels equals:

$$I_{Ca^{2+}} = [N] \cdot g(V_m) \cdot (V_m - E_{Ca^{2+}}). \quad (15)$$

The alteration of the overall conductivity of the  $Ca^{2+}$  channels in the *Paramecium* cilia in response to the shift of transmembrane potential from  $V_0$  to  $V_1$  under the voltage clamped condition is given by:

$$g(V, t) = g(V_1) - (g(V_1) - g(V_0)) \cdot \exp\left(-\frac{t}{\tau_p}\right), \quad (16)$$

where  $\tau_p$  is the characteristic time of the transmembrane potential alteration from  $V_0$  to  $V_1$ .

The equation for the conductivity alterations (16) can be represented as follows:

$$g(V_m, t) = g_0 \cdot (v(V_1) - (v(V_1) - v(V_0)) \cdot \exp(-\frac{t}{\tau_p})), \quad (17)$$

where  $\psi = \frac{V \cdot F}{R \cdot T}$ ,  $v(\psi, t)$  is the  $Ca^{2+}$  channel conductivity dependence on the transmembrane potential and on time.

An early study investigated the  $Ca^{2+}$  channels' conductivity dependence on transmembrane potential  $\psi$  in *Paramecia* [66] and approximated it by the following equation:

$$v(\psi) = \frac{\exp(\alpha \cdot (\psi + d))}{\lambda + \exp(\alpha \cdot (\psi + d))}, \quad (18)$$

where  $\alpha$ ,  $d$  and  $\lambda$  are the parameter values that allow the best representation of the available experimental data. In this model, the steepness of the dependence of the conductivity on membrane potential is represented by the parameter  $\alpha$ .

For simplicity we assume that the elementary intraciliary volume with inward  $Ca^{2+}$  current does not contain any  $Ca^{2+}$  binding proteins and  $Ca^{2+}$  is extruded into the extracellular space by the active  $Ca^{2+}$  transport only. The  $Ca^{2+}$  binding to the  $Ca^{2+}$  channels is assumed to occur much faster than the characteristic times of  $Ca^{2+}$  channel inhibition. Under such assumptions, the intraciliary  $Ca^{2+}$  dynamics in response to the transmembrane potential shift under the voltage clamp is given by:

$$V_R \cdot \frac{d[Ca^{2+}]}{dt} = \frac{S_R}{z \cdot F} \cdot \left( -[N] \cdot (g(V_m) \cdot (V_m - E_{Ca^{2+}})) - \beta \cdot \frac{[Ca^{2+}]}{K_A + [Ca^{2+}]} \right). \quad (19)$$

In non dimensional form, equation (19) takes the following form:

$$\begin{cases} \frac{du}{d\eta} = s \cdot \left( -b \cdot n(\eta) \cdot (v(\psi_1) - (v(\psi_1) - v(\psi_0)) \cdot \exp(-\frac{\eta}{\tau_0})) \right) \cdot \\ \left( \psi_1 - 0.5 \cdot \ln\left(\frac{u_{out}}{u}\right) - \frac{u}{k_A + u} \right), \\ \frac{dn}{d\eta} = -k \cdot n \cdot u + (1 - n), \end{cases} \quad (20)$$

$$\text{where } u = \frac{Ca^{2+}}{K_{CaM}}, u_{out} = \frac{Ca_{out}^{2+}}{K_{CaM}},$$

The initial conditions for the equations (20) when  $\eta = 0$  and  $\psi = \psi_0$  are equal to  $u_0$  ( $\psi_0$ ) and  $n(u_0)$ , which can be obtained by numerical solution of these equations. The non dimensional  $Ca^{2+}$  current is equal to:

$$i_{Ca^{2+}} = -b \cdot n(\eta) \cdot \left( v(\psi_1) - (v(\psi_1) - v(\psi_0)) \cdot \exp\left(-\frac{\eta}{\tau_0}\right) \right) \cdot \left( \psi_1 - 0.5 \cdot \ln\left(\frac{u_{out}}{u}\right) - \frac{u}{k_A + u} \right). \quad (21)$$

### Indirect $Ca^{2+}$ channel conductivity regulation

The second model of intraciliary calcium regulation assumes the  $Ca^{2+}$  binding protein interacts with the  $Ca^{2+}$  channel and thereby causes the  $Ca^{2+}$  channels to close. By considering that the  $Ca^{2+}$  ion- $Ca^{2+}$  binding protein occurs on a faster time scale than the  $Ca^{2+}$  binding protein- $Ca^{2+}$  channels interaction, the steady-state solution for the  $Ca^{2+}$  binding protein in complex with a  $Ca^{2+}$  ion is given by:

$$[CaC] = [CaC_0] \cdot \frac{[Ca^{2+}]}{K_C + [Ca^{2+}]}, \quad (22)$$

where  $[CaC_0]$  is the total concentration of the  $Ca^{2+}$  binding protein and  $K_C = \frac{k^m}{k^p}$  is the equilibrium dissociation constant.

The  $Ca^{2+}$  binding protein interaction with the  $Ca^{2+}$  channels is given by:

$$\frac{d[C]}{dt} = -n^p \cdot [C] \cdot [CaC] + n^m \cdot ([C_0] - [C]), \quad (23)$$

where  $C$  is the number of channels in the open conductive state, and  $C_0$  is the total number of channels. The steady-state solution of equation (23) is given by:

$$[C] = [C_0] \cdot \frac{K_{CC}}{K_{CC} + [CaC]}, \quad (24)$$

where  $K_{CC} = \frac{n^m}{n^p}$  is the equilibrium dissociation constant for the  $Ca^{2+}$  binding protein- $Ca^{2+}$  channels interaction. By combining equations (22) and (24) one obtains the dependence of the open  $Ca^{2+}$  channels on the  $Ca^{2+}$  concentration:

$$c = \frac{k_C + u}{k_C + (cac_0 + 1) \cdot u'} \quad (25)$$

$$\text{where } c = \frac{[C]}{[C_0]}, u = \frac{Ca^{2+}}{K_{CaM}}, cac_0 = \frac{CaC_0}{K_{CC}}, k_C = \frac{K_C}{K_{CaM}},$$

Equation (23) can then be represented in the following non dimensional form:

$$\frac{dc}{d\eta} = -c \cdot cac_0 \cdot \frac{u}{k_C + u} + (1 - c), \quad (26)$$

where

$$\eta = n^m \cdot t, c = \frac{[C]}{[C_0]}, k_C = \frac{K_C}{K_{CaM}}, u = \frac{Ca_{2+}}{K_{CaM}}, cac_0 = \frac{[CaC_0]}{K_{CC}},$$

The solution of equation (26) that describes the dynamics of  $Ca^{2+}$  channels in the open conductive state in response to a  $Ca^{2+}$  shift from  $u_0$  to  $u_1$   $Ca^{2+}$  level is given by:

$$c(t) = c^\infty - (c^\infty - c^0) \cdot \exp\left(-\left(cac_0 \cdot \frac{u_1}{k_C + u_1} + 1\right) \cdot n^m \cdot t\right), \quad (27)$$

where

$$c^\infty = \frac{k_C + u_1}{k_C + (cac_0 + 1) \cdot u_1}, c^0 = \frac{k_C + u_0}{k_C + (cac_0 + 1) \cdot u_0}.$$

The solution (27) suggests that the number of open channels would change exponentially in response to a  $Ca^{2+}$  surge. The characteristic time for such an exponential change is given by:

$$\tau_{Ca^{2+}} = \frac{k_C + u}{(k_C + (cac_0 + 1) \cdot u) \cdot n^m}. \quad (28)$$

For cases when  $cac_0 \gg 1$  and  $u \gg 1$ , the characteristic time approximately equals  $\tau_{Ca^{2+}} \approx \frac{1}{[CaC_0] \cdot n^p}$ .

Under the assumption that the alteration of the transmembrane potential difference influences the  $Ca^{2+}$  channel conductivity, the channel conductivity as a function of  $Ca^{2+}$  concentration can be represented as follows:

$$g_{Ca^{2+}}(t) = [C_0] \cdot g_0 \cdot v(\psi) \cdot c(t). \quad (29)$$

When membrane potential changes from  $\psi_0$  to  $\psi_1$ , and the alteration of the intraciliary  $Ca^{2+}$  concentration is delayed, the conductivity changes from one value to another exponentially with the characteristic time  $\tau_V$ :

$$g(\psi, t) = g_0 \cdot (v(\psi_1) - (v(\psi_1) - v(\psi_0)) \cdot \exp(-t/\tau_V)). \quad (30)$$

Here we consider a simplified scenario when only passive and active  $Ca^{2+}$  transport channels are present. For such a model, the alteration of  $Ca^{2+}$  concentration in a cilium is given by:

$$V_R \cdot \frac{d[Ca^{2+}]}{dt} = \frac{S_R}{z \cdot F} \cdot (I_{Ca^{2+}}^P + I_{Ca^{2+}}^A). \quad (31)$$

By substituting formulas for the passive and active  $Ca^{2+}$  currents into equation (31), one obtains:

$$V_R \cdot \frac{d[Ca^{2+}]}{dt} = \frac{S_R}{z \cdot F} \cdot \left(-[C_0] \cdot c(t) \cdot g_0 \cdot v(V_m, t) \cdot (V_m - E_{Ca^{2+}}) - \beta \cdot \frac{[Ca^{2+}]}{K_A + [Ca^{2+}]}\right). \quad (32)$$

In this equation, we include the kinetics for the active  $Ca^{2+}$  channels due to the assumption that the dynamics of currents via the active  $Ca^{2+}$  channels is much faster than the dynamics of currents through the passive  $Ca^{2+}$  transport.

The non dimensional representation of Equations (32) and (23) is given by:

$$\begin{cases} \frac{du}{d\eta} = a \cdot \left(-b \cdot c(\eta, u) \cdot (v(\psi_1) - (v(\psi_1) - v(\psi_0)) \cdot \exp\left(-\frac{\eta}{\tau_0}\right))\right) \cdot \left(\psi_1 - 0.5 \cdot \ln\left(\frac{u_{out}}{u}\right)\right) - \frac{u}{k_A + u} \\ \frac{dc}{d\eta} = -c \cdot cac_0 \cdot \frac{u}{k_C + u} + (1 - c), \end{cases} \quad (33)$$

$$\text{where } u = \frac{Ca^{2+}}{K_{CaM}}, u_{out} = \frac{Ca_{out}^{2+}}{K_{CaM}},$$

By substituting the non dimensional representation of the  $Ca^{2+}$  current (21) into (33) the equation for the intraciliary  $Ca^{2+}$  dynamics (33) take the following form:

$$\begin{cases} \frac{du}{d\eta} = a \cdot \left(-b \cdot c(\eta, u) \cdot (v(\psi_1) - (v(\psi_1) - v(\psi_0)) \cdot \exp\left(-\frac{\eta}{\tau_0}\right))\right) \cdot \left(\psi_{st} - 0.5 \cdot \ln\left(\frac{u}{u_i}\right)\right) - \frac{u}{k_A + u} + o_V(\psi_1) \cdot \left(\psi_{st} - 0.5 \cdot \ln\left(\frac{u}{u_i}\right)\right) \\ \frac{dc}{d\eta} = -c \cdot cac_0 \cdot \frac{u}{k_C + u} + (1 - c), \end{cases} \quad (34)$$

$$\text{where } u = \frac{Ca^{2+}}{K_{CaM}}, u_{out} = \frac{Ca_{out}^{2+}}{K_{CaM}},$$

### Potassium current

Here we consider the mechanism of  $K^+$  current activation by membrane potential. The full  $K^+$  current in a cilium is given by:

$$I_{K^+} = N_{K^+} \cdot g_{K^+}^0 \cdot v_{K^+}(t, V_m) \cdot (V_m - E_{K^+}), \quad (35)$$

where  $N_{K^+}$  is the number of open  $K^+$  channels,  $g_{K^+}^0$  is the maximal conductivity, and  $E_{K^+}$  is the equilibrium  $K^+$  potential.

We assume that the  $K^+$  channel conductance is dependent on the membrane potential (the opposite of the case of  $Ca^{2+}$  channel conductance). The  $K^+$  channel conductance dependence on the membrane potential is given by:

$$v_{K^+}(\psi) = \frac{\exp(\alpha_{K^+} \cdot (\psi + d_{K^+}))}{\lambda_{K^+} + \exp(\alpha_{K^+} \cdot (\psi + d_{K^+}))}. \quad (36)$$

We further assume that in response to the fast membrane potential alteration from  $\psi_0$  to  $\psi_1$ , the  $K^+$  conductivity is changing exponentially with a characteristic time  $\tau_{K^+}$ . In that case the conductivity dynamics over time are given by:

$$v_{K^+}(t) = v_{K^+}(\psi_1) - (v_{K^+}(\psi_1) - v_{K^+}(\psi_0)) \cdot \exp\left(\frac{-t}{\tau_{K^+}}\right). \quad (37)$$

The parameters for the  $K^+$  channels' conductivity dependence on the membrane potential can be estimated from experimental measurements of the fast  $K^+$  current (10 msec) as a function of membrane potential. In such a short time  $K^+$  current does not reach the maximum value and the  $Ca^{2+}$  current is almost equal to zero. Under such conditions, the equation for the  $K^+$  current is given by:

$$I_{K^+} = i_0 \cdot v_{K^+}(\psi_m) \cdot (\psi - \psi_{K^+}). \quad (38)$$

The full current measured under the voltage clamp conditions is then given by:

$$i_p = -b \cdot c(\eta, u) \cdot \left( v(\psi_1) - (v(\psi_1) - v(\psi_0)) \cdot \exp\left(\frac{-t}{\tau_0}\right) \right) \cdot \left( \psi - 0.5 \cdot \ln\left(\frac{u_{out}}{u}\right) \right) - b_1 \cdot \left( v_{K^+}(\psi_1) - (v_{K^+}(\psi_1) - v_{K^+}(\psi_0)) \cdot \exp\left(\frac{-t}{\tau_{K^+}}\right) \right) \cdot (\psi - \psi_{K^+}), \quad (39)$$

$$\text{where } b_1 = N_{K^+} \cdot \frac{g_{K^+}^0}{g_0}.$$

### The transmembrane potential dynamics

All the models described above were specifically developed under the voltage clamp conditions. However,  $Ca^{2+}$  as well as  $K^+$  currents can alter the membrane potential. In order to account for this effect, we considered the membrane potential dynamics in the absence of voltage clamp. Having established that the main contributors to the registered currents are the  $Ca^{2+}$  and  $K^+$

currents, we analyse the membrane potential dynamics as a function of  $Ca^{2+}$  and  $K^+$  current contributions:

$$C_m \cdot \frac{dV_m}{dt} = I_{Ca^{2+}} + I_{K^+}. \quad (40)$$

In the non dimensional form the equation for the membrane potential dynamics is given by:

$$\frac{d\psi}{d\eta} = \rho \cdot \left( \frac{-b \cdot (c(\eta, u) \cdot v(\eta, \psi) + v_{Ca^{2+}}^{st}) \cdot (\psi - 0.5 \cdot \ln\left(\frac{u_{out}}{u}\right))}{-b_1 \cdot (v_{K^+}(\eta, \psi) + v_{K^+}^{st}) \cdot (\psi - \psi_{K^+})} \right), \quad (41)$$

where  $\psi = \frac{V_m \cdot F}{R \cdot T}$ ,  $u = \frac{[Ca^{2+}]}{K_{CaM}}$ ,  $v_{Ca^{2+}}^{st}$  and  $v_{K^+}^{st}$  are the steady-state  $Ca^{2+}$  and  $K^+$  channel conductivities, respectively,  $\eta = n^m \cdot t$ ,  $\rho = \frac{g^0}{C_m \cdot n^m}$ .

In this equation, we take into consideration the contribution of the independent parts of the  $Ca^{2+}$  and  $K^+$  currents, conductivities  $v_{Ca^{2+}}^{st}$  and  $v_{K^+}^{st}$ , respectively, that do not depend on membrane potential or  $Ca^{2+}$  concentration. We also use the fact that  $Ca^{2+}$  and  $K^+$  channel conductivities change on a much faster time scale in comparison with membrane potential or  $Ca^{2+}$  concentration. Such an approximation allows us to employ the steady-state solutions for the channel conductivities as a function of membrane potential and  $Ca^{2+}$  concentration. The described assumptions and considerations lead to the following system of two nonlinear coupled equations for intraciliary  $Ca^{2+}$  concentration and membrane potential:

$$\begin{cases} \frac{du}{d\eta} = s \cdot \left( -b \cdot (c(u) \cdot v(\psi) + v_{Ca^{2+}}^{st}) \cdot \left( \psi - 0.5 \cdot \ln\left(\frac{u_{out}}{u}\right) \right) - \frac{u}{k_A + u} \right), \\ \frac{d\psi}{d\eta} = -\rho \cdot \left( \frac{-b \cdot (c(u) \cdot v(\psi) + v_{Ca^{2+}}^{st}) \cdot \left( \psi - 0.5 \cdot \ln\left(\frac{u_{out}}{u}\right) \right)}{-b_1 \cdot (v_{K^+}(\psi) + v_{K^+}^{st}) \cdot (\psi - \psi_{K^+}) + I_0} \right), \end{cases} \quad (42)$$

where  $I_0$  is the non dimensional inward current.

### Currents activated by the membrane hyperpolarisation

The membrane hyperpolarisation current is given by:

$$I_{Ca^{2+}}^{ht} = N_{Ca^{2+}}^{ht} \cdot g_{Ca^{2+}}^{ht}(\psi, t) \cdot (V_m - E_{Ca^{2+}}) \quad (43)$$

where  $N_{Ca^{2+}}^{ht}$  is the number of open, membrane hyperpolarisation-activated  $Ca^{2+}$  channels located on the cell body,  $g_{Ca^{2+}}^{ht}$  is the conductivity of the channel, and  $E_{Ca^{2+}}$  is the equilibrium potential for  $Ca^{2+}$  ions.

$$g_{Ca^{2+}}^{ht}(\psi, t) = g^0 \cdot v_h(\psi, t), \quad (44)$$

where

$$g^0 = \max(g_{Ca^{2+}}^{ht}(\psi, t)), v_h(\psi) = \frac{\exp(\alpha_h \cdot (\psi + d_h))}{\lambda_h + \exp(\alpha_h \cdot (\psi + d_h))}$$

By combining the contribution from different types of currents, one can derive the dynamics of the membrane

**Table 1 Parameter values employed in the systems model for the ciliary excitation**

Parameter	Value (dimensionless unless otherwise stated)	Figure No	Equation
$\alpha$	4	2A	18
$d$	0.4	2A	18
$\lambda$	0.5	2A	18
$b$	2	2B, 3, 5, 6, 7, 9, 10	20, 21, 34, 39
$k_A$	1	2B, 3, 5, 6, 7, 11	20, 34, 42
$u_{out}$	1000	2B, 3, 5, 6, 7	20, 21, 34
$\psi_0$	-1.2	2B, 3, 5, 6, 7, 8	20, 21, 34, 39
$V_0$	30 mV	2B, 3, 5, 6, 7, 8	20, 21, 34, 39
$\psi_1$	-1, -0.8, -0.5, -0.2, 0, 0.2, 0.5	2B, 3, 5, 6, 7, 8	20, 21, 34, 39
$V_1$	25, 20, 12.5, 5, 0, -5, -12.5 mV	2B, 3, 5, 6, 7, 8	20, 21, 34, 39
$s$	0.5	2B, 3, 5, 6, 7	20
$\tau_0$	0.02	2B, 3, 5, 6, 7, 9, 10	20, 34, 39
$k$	2	2B, 3, 5, 6, 7	20
$K_C$	1	4, 5, 6, 7	25
$CaC_0$	5, 10, 50, 100	4, 5, 6, 7	25, 34
$a$	4	7, 11	34, 42
$\alpha_{K^+}$	0.5	8, 10	36
$\lambda_{K^+}$	0.005	8, 10	36
$d_{K^+}$	0.5	8, 10	36
$b_1$	1	9	39
$b$	8	11, 12, 13, 14	42, 46
$caC_0$	20	11, 12, 13, 14	42
$v_{Ca^{2+}}^{st}$	0.01	11, 12, 13, 14	42, 46
$v_{K^+}^{st}$	0.01	11, 12, 13, 14	42, 46
$\rho$	10	11, 12, 13, 14	42, 46
$s$	0.5	11, 12, 13, 14	42, 46
$v_{Ca^h}$	0.9	12, 13, 14	46
$\lambda_h$	5	12, 13, 14	44
$d_h$	1	12, 13, 14	44
$\alpha_h$	4	12, 13, 14	44

potential alteration as follows:

$$C_m \cdot \frac{dV_m}{dt} = n_v \cdot \left( (-N_{Ca^{2+}}^0 \cdot (w_1 \cdot g_{Ca^{2+}}^1 + w_3 \cdot g_{Ca^{2+}}^3) + N_{Ca^{2+}}^0 \cdot g_{Ca^{2+}}^0) \cdot \left( V_m - \frac{R \cdot T}{2 \cdot F} \cdot \ln \left( \frac{u_{out}}{u_i} \right) \right) - i_A \cdot N_{Ca^{2+}}^{00} \cdot \omega + (N_{Ca^{2+}}^0 \cdot g_{Ca^{2+}}^1(t, V_m) + N_{Ca^{2+}}^0 \cdot g_{Ca^{2+}}^2(t, V_m) + N_{Ca^{2+}}^0 \cdot g_{Ca^{2+}}^3(V_m) \cdot v_1 + N_{Ca^{2+}}^0 \cdot g_{Ca^{2+}}^4(V_m) \cdot v_2) \cdot (V_m - E_K) - I_K^A \right) \quad (45)$$

The coupled system of differential equations for the  $Ca^{2+}$  and membrane potential dynamics can be further formulated as:

$$\begin{cases} \frac{du}{d\eta} = s \cdot \left( -b \cdot (c(u) \cdot v_{Ca^{2+}}(\psi) + v_{Ca^{2+}}^{st}) \cdot \left( \psi - 0.5 \cdot \ln \left( \frac{u_{out}}{u} \right) \right) - \frac{u}{k_A + u} \right), \\ \frac{d\psi}{d\eta} = -\rho \cdot \left( (-b \cdot (c(u) \cdot v(\psi) + v_{Ca^{2+}}^{st}) - v_{Ca^h} \cdot v_h(\psi)) \cdot \left( \psi - 0.5 \cdot \ln \left( \frac{u_{out}}{u} \right) \right) - b_1 \cdot (v_{K^+}^1(\psi) + v_{KCa} \cdot v_{K^+}^2(\psi, u) + v_{K^+}^{st}) \cdot (\psi - \psi_{K^+}) + I_0 \right), \end{cases} \quad (46)$$

where  $v_h(\psi)$  is the  $Ca^{2+}$  current contribution, activated by membrane depolarization, and  $v_{K^+}(\psi, u)$  is the  $Ca^{2+}$ -dependent  $K^+$  current contribution.

### Cilia-to body $Ca^{2+}$ current

The incorporation of the  $Ca^{2+}$  current from cilia to cell into the equations for the intraciliary  $Ca^{2+}$  and membrane potential system leads to an additional term being added to the equations (46):

$$\begin{cases} \frac{du}{d\eta} = s \cdot \left( \frac{-b \cdot (c(u) \cdot v(\psi) + v_{Ca^{2+}}^{st}) \cdot \left( \psi - 0.5 \cdot \ln \left( \frac{u_{out}}{u} \right) \right) - \frac{u}{k_A + u}}{-\frac{u}{k_A + u} + v_{Ca^h} \cdot v_h(\psi)} \cdot \left( \psi - 0.5 \cdot \ln \left( \frac{u_{out}}{u} \right) \right) \right), \\ \frac{d\psi}{d\eta} = -\rho \cdot \left( (-b \cdot (c(u) \cdot v(\psi) + v_{Ca^{2+}}^{st}) - v_{Ca^h} \cdot v_h(\psi)) \cdot \left( \psi - 0.5 \cdot \ln \left( \frac{u_{out}}{u} \right) \right) - b_1 \cdot (v_{K^+}(\psi) + v_{KCa} \cdot v_{K^+}^1(u, \psi) + v_{K^+}^{st}) \cdot (\psi - \psi_{K^+}) + I_0 \right), \end{cases} \quad (47)$$

In the general case the conductivity of the protein complexes governing the  $Ca^{2+}$  current from cilia to the cell body depends on the membrane potential:

**Table 2 The relationship between dimensional and non-dimensional quantities for Ca<sup>2+</sup> concentration and membrane potential**

Variables	Dimensional variables	Non-dimensional variables	Coefficient value
Calcium	Ca <sup>2+</sup> (M/L)	$u = \frac{Ca^{2+}}{K_{CaM}}$	$K_{CaM} = 4 \mu M$
Transmembrane potential	V <sub>m</sub>	$\psi_m = \frac{F \cdot V_m}{R \cdot T}$	$\frac{RT}{F} = -0.025 V$

$$v_{K^+}^{st} \quad (48)$$

All parameter values used in the above equations are given in Table 1. The relationship between dimensional and non-dimensional quantities for Ca<sup>2+</sup> concentration and membrane potential are given in Table 2.

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#### Authors' contributions

NVK, ANG, BG, MZQC, ID, CH and AA developed and implemented the project under the supervision of DGB, RNK, YU and NVV. XY, SKS and CMS contributed to the analysis of the model. All authors contributed to the writing of the final manuscript. All authors read and approved the final manuscript.

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