

Au-Decorated Ti_3C_2 MXene Sensor for Enhanced Detection of Gaseous Toxins (CO , COCl_2 , H_2S , NH_3 , NO_2): A DFT Study

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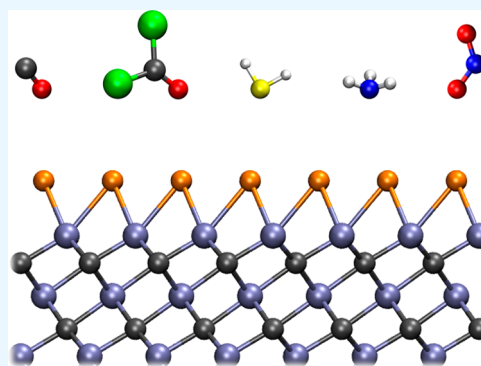
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ABSTRACT: The rising level of toxic gases in the environment poses a high demand for efficient gas sensing materials. MXenes, an emerging class of two-dimensional (2D) materials, have gained significant interest in this area for having an active-site rich structure, tunable surface properties, and remarkable stability. Herein, an extensive density functional theory (DFT) study is conducted to investigate the sensing properties of pristine and Au-functionalized Ti_3C_2 MXene for five toxic gas molecules: CO , COCl_2 , H_2S , NH_3 , and NO_2 . Pristine Ti_3C_2 displays high affinity for CO , H_2S , and NH_3 , as assessed by density of states and a large binding energy, resulting in the chemisorption of these gas molecules providing a relatively large recovery time. In contrast, Au-functionalized Ti_3C_2 is able to sense all five toxins which are physisorbed on it, as indicated by lower adsorption energy and faster recovery time. As an example, the adsorption energy computed for CO is -0.14 eV and the resulting recovery time 0.21 ns. These results reveal that Au-functionalized Ti_3C_2 can serve as a highly efficient material for toxic gas sensing, particularly CO .



1. INTRODUCTION

Chemical gas sensors are extensively used in many industries, such as automotive, medical, process automation, safety, product and quality control, as well as pollution monitoring.¹ Due to the growing needs arising from environmental changes, scientists are searching for efficient sensing materials.² Two-dimensional (2D) materials are currently revolutionizing the gas-sensing industry.³ In addition to being cost- and energy-efficient, they provide exceptional sensitivity to a wide range of gases.⁴ Notably, 2D materials have a large surface-to-volume ratio, which equips them with an extended active surface area for gas-molecule interaction.⁵ This property enhances their sensing ability enabling them to detect up to parts-per-billion (ppb) pollutant concentrations.⁶ However, these materials have certain limitations as they require extremely lengthy recovery periods and have constraints on accuracy.⁷

Toxic and hazardous gases, including phosgene (COCl_2) and carbon monoxide (CO), are frequently utilized in manufacturing industries.⁸ Incomplete combustion of fuels results in the formation of severely harmful carbon monoxide. This gas is challenging to detect for having colorless, tasteless, and odorless properties.⁹ The primary sources of carbon monoxide emissions in the environment are industrial activities, automobile exhaust, and fossil fuel consumption.¹⁰ By binding with hemoglobin in blood, carbon monoxide poses significant life threats.¹¹ Upon inhalation, it enters the body and reaches the bloodstream, displacing oxygen and causing

the heart and brain to malfunction due to oxygen deprivation.¹² The most common signs of CO accumulation in the human body are fatigue and headache. Elevated CO levels cause brain damage, unconsciousness, and even death.¹³ Phosgene, a toxic gas, is a key ingredient in making polyurethanes.¹⁴ Despite being hazardous, it is also an important substrate in making isocyanates and insecticides.¹⁵ Phosgene poses a significant threat to human life due to its subtle smell and delayed symptom onset, making it difficult to detect. Although it can be identified at 0.4 ppm, this is still above the safe exposure limit for humans.¹⁶ Particularly, the toxicity of this gas arises from its destructive effect on the pulmonary alveoli proteins. By damaging the blood-air barrier, phosgene ultimately causes asphyxiation.¹⁷ Due to high toxicity of these gases, various nanomaterials (graphene,^{18,19} graphdiyne,²⁰ phosphorene,^{21–23} MoS_2 ,²⁴ and many more^{25–27}) for their sensing have been explored.

MXenes are a class of 2D materials composed of transition metal atoms and nonmetallic species like carbon or nitrogen as well as other elements including oxygen, sulfur, and chlorine,

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and characterized by high electrical conductivity, tunable surface chemistry, and excellent mechanical strength.²⁸ A detailed analysis of how MXenes interact with gas molecules has revealed the complex processes involved in their gas sensing activity.²⁹ This knowledge has led to the development of MXene-based highly sensitive gas detectors with great selectivity, which are relevant for several applications.³⁰ In particular, MXenes revolutionized gas sensors, making it possible to detect a broad spectrum of toxic gases, including nitrogen dioxide (NO₂),³¹ ammonia (NH₃),³² carbon monoxide (CO),^{33,34} COCl₂,³⁵ and hydrogen sulfide (H₂S).³⁶ Due to the unique properties of MXenes, researchers across the globe are focusing on unlocking their critical gas sensing capabilities, minimizing experimental risks.^{37,38}

Ab initio studies are widely used for their unparalleled atomic-level insights into sensing mechanisms and material properties.³⁹ Simulations have revealed that MXene surface terminations govern their interactions with gas molecules.⁴⁰ The synthesis of the titanium carbide (Ti₃C₂T_x) MXene opened a new frontier in materials research. Naguib et al.³⁸ established a comprehensive understanding of the unique two-dimensional material, achieved by integrating theoretical and experimental approaches. Extending this study, many researchers have directed their attention toward MXene-based material as gas sensors.^{41–43} MXenes with hydroxyl groups strongly react with NO₂ due to a strong chemical interaction of NO₂ with these functionalities. In contrast, fluorine-terminated MXenes show high sensitivity for NH₃, resulting from the Lewis acid–base interaction between NH₃ and fluorine.⁴⁴ To enhance the response toward NH₃ or other gases, different techniques, such as doping⁴⁵ and noble metal decoration⁴⁶ have been proposed.

In the present study based on density functional theory (DFT), pristine and Au-functionalized Ti₃C₂ MXenes are investigated as platforms for sensing toxic gas molecules such as CO, COCl₂, H₂S, NH₃, and NO₂. The optimized structures, electronic properties, and adsorption performance against the sample gases are evaluated together with the sensing performance of the material in terms of the recovery time.

2. METHODOLOGY

All DFT calculations presented in this work were carried out with the CP2K code⁴⁷ applying periodic boundary conditions to account for the in-plane periodicity of the systems. Spin-polarization was considered due to the presence of transition metal atoms. Double-zeta valence with polarization (DZVP) molecularly optimized basis sets⁴⁸ and Goedecker–Teter–Hutter pseudopotentials⁴⁹ were used. The generalized gradient approximation (GGA) as implemented in the Perdew–Burke–Ernzerhof (PBE) exchange–correlation functional was adopted⁵⁰ including the Grimme-D3 scheme to account for van der Waals interactions between MXenes and gas molecules.⁵¹ The plane-wave energy cutoff was set to 450 Ry and all the systems were relaxed below the force threshold of 0.01 eV/Å per atom. The Broyden–Fletcher–Goldfarb–Shanno (BFGS) algorithm was used for geometry optimization. Bader charges on gas molecules were computed with Critic2 code.⁵² Furthermore, a vacuum space of 20 Å was added perpendicular to the MXene basal plane to avoid artificial interactions between neighboring images. All the studied systems were visualized with VMD software package.⁵³

The adsorption energy (E_{ads}) of the gas molecule was calculated as

$$E_{\text{ads}} = E_{\text{sys}} - E_{\text{MXene}} - E_{\text{gas}} \quad (1)$$

where E_{sys} is the total energy of the gas@MXene system, E_{MXene} is the total energy of the Ti₃C₂/Au–Ti₃C₂ slab, and E_{gas} is the total energy of the isolated gas molecule.

The sensor recovery time (τ) was calculated as^{54,55}

$$\tau = \nu_0^{-1} e^{-E_{\text{ads}}/kT} \quad (2)$$

where ν_0 is the attempt frequency, k is the Boltzmann constant, and T is the temperature. The values of attempt frequency and temperature were set to 10¹² s^{−1} and 298 K, respectively.^{24,56}

The work function (ϕ) of the studied systems was evaluated as⁵⁷

$$\phi = E_{\text{inf}} - E_{\text{F}} \quad (3)$$

where E_{inf} is the electrostatic potential energy at infinity or vacuum and E_{F} is the Fermi energy.

3. RESULTS AND DISCUSSION

3.1. Structural and Adsorption Properties. The structure of the Ti₃C₂ MXene is characterized by 5 alternating

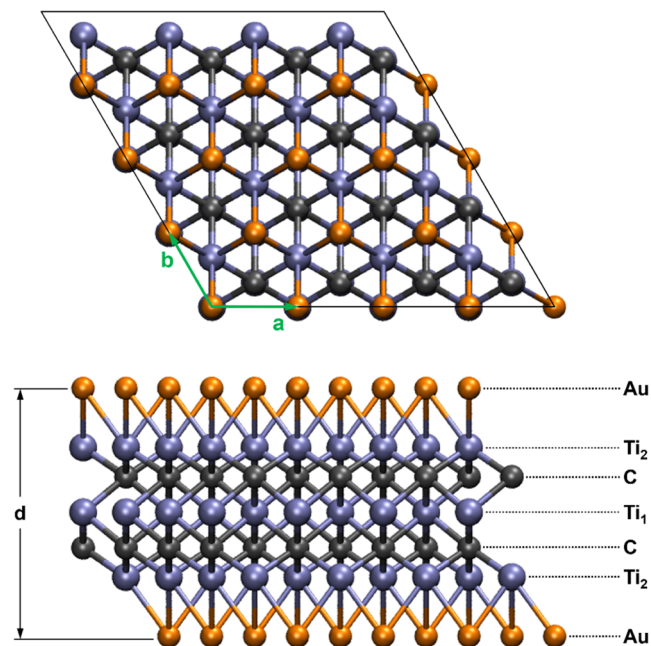
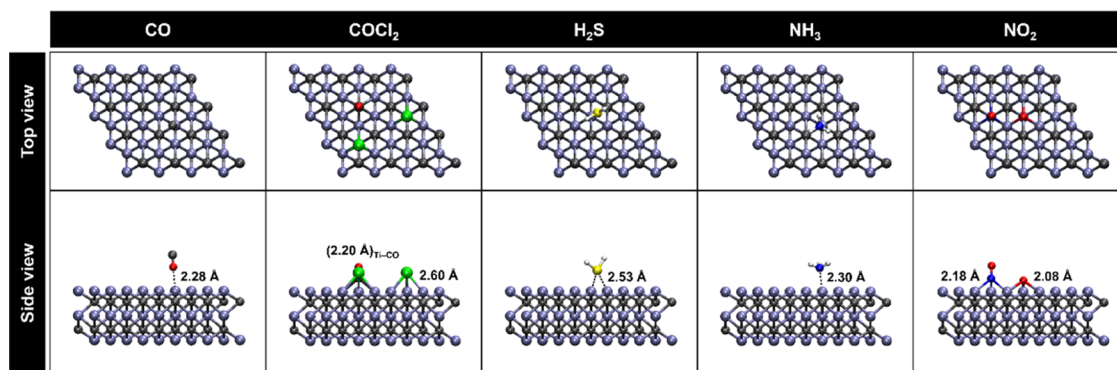
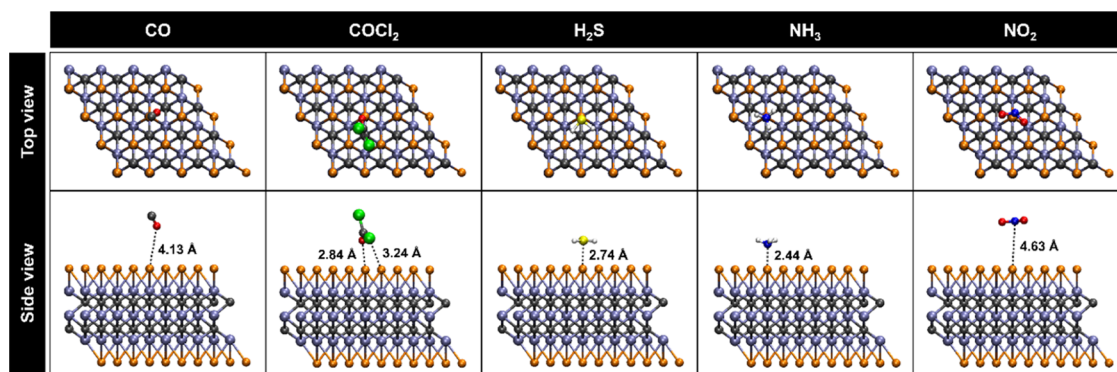


Figure 1. Top and side views of the optimized structure of the 4 × 4 supercell of the Au-decorated MXene, where the overall layer thickness d is indicated. Lattice parameters are shown with green arrows.

layers of C and Ti atoms with the latter at the interface with vacuo. The Au-functionalization of this material is achieved by bonding a layer of gold to the outermost layers of Ti atoms, see Figure 1, where the 4 × 4 supercell adopted in the simulations of gas adsorption is depicted. While the hexagonal symmetry of the Ti₃C₂ unit cell remains unchanged after functionalization with gold atoms, the in-plane lattice constant of the unit cell of Ti₃C₂ is increased by 0.16% with Au-decoration and the resulting slab thickness of the surface-modified MXene is 1.92 times that of the pristine one, see Table 1. In addition, Au atoms at hexagonal sites show triple coordination with Ti atoms with a bond length of 2.77 Å. The Ti₁–C bond length

Table 1. Optimized Lattice Constant (a), Slab Thickness (d), and Bond Lengths of Pristine and Au-Decorated Ti_3C_2 MXene

system	a (Å)	d (Å)	bond lengths (Å)		
			$\text{Ti}_1\text{--C}$	$\text{Ti}_2\text{--C}$	$\text{Ti}_2\text{--Au}$
Ti_3C_2	3.07	4.63	2.20	2.03	
$\text{Au--Ti}_3\text{C}_2$	3.08	8.91	2.19	2.07	2.77

Figure 2. Optimized structures of the studied toxins at the pristine Ti_3C_2 surface.Figure 3. Optimized configurations of the studied toxins at the $\text{Au--Ti}_3\text{C}_2$ surface with the minimal distances from the substrates indicated by dotted lines in each plot.Table 2. Adsorption Energies (E_{ads}) of the Studied Toxic Gases and the Sensor Recovery Time (τ) of Pristine and Au-Decorated Ti_3C_2

system	E_{ads} (eV)	τ (s)
$\text{CO@Ti}_3\text{C}_2$	−0.33	3.53×10^{-07}
$\text{COCl}_2\text{@Ti}_3\text{C}_2$		
$\text{H}_2\text{S@Ti}_3\text{C}_2$	−1.40	4.64×10^{11}
$\text{NH}_3\text{@Ti}_3\text{C}_2$	−1.19	1.29×10^{08}
$\text{NO}_2\text{@Ti}_3\text{C}_2$		
$\text{CO@Au--Ti}_3\text{C}_2$	−0.14	2.11×10^{-10}
$\text{COCl}_2\text{@Au--Ti}_3\text{C}_2$	−0.54	0.0016
$\text{H}_2\text{S@Au--Ti}_3\text{C}_2$	−0.71	1.07
$\text{NH}_3\text{@Au--Ti}_3\text{C}_2$	−0.66	0.13
$\text{NO}_2\text{@Au--Ti}_3\text{C}_2$	−0.68	0.32

decreases by 0.41% in Au-functionalized Ti_3C_2 whereas the $\text{Ti}_2\text{--C}$ bond length is relaxed by 1.57%.

The adsorption of CO, COCl_2 , H_2S , NH_3 , and NO_2 is first simulated on the pristine Ti_3C_2 surface, see Figure 2. COCl_2 and NO_2 chemically react with the surface and are therefore not considered in the subsequent analysis. The CO molecule is bound to the surface Ti atom in a vertical configuration forming a $\text{Ti}_2\text{--O}$ bond with length 2.28 Å; the corresponding adsorption energy is −0.33 eV. The adsorption of H_2S takes

Table 3. Performance Comparison of Various Gas Sensing Materials

material	gas	E_{ads} (eV)	τ (s)	references
$\text{Au--Ti}_3\text{C}_2$	CO	−0.14	2.11×10^{-10}	this study
In_2O_3		−0.52	6.17×10^{-04}	58
Co-g-GaN		−1.22	4.10×10^{08}	59
$\text{Au--Ti}_3\text{C}_2$	COCl_2	−0.54	0.0016	this study
Si-embedded MoS_2		−1.23	5.77×10^{08}	24
Al-Pc		−0.86	3.47×10^{02}	60
$\text{Au--Ti}_3\text{C}_2$	H_2S	−0.71	1.07	this study
In_2O_3		−2.35	5.30×10^{27}	58
Mn-doped BP		−0.85	2.70	61
$\text{Au--Ti}_3\text{C}_2$	NH_3	−0.66	0.13	this study
Rh-InSe		−1.11	5.64×10^{06}	62
Co-g-GaN		−0.68	0.30	59
$\text{Au--Ti}_3\text{C}_2$	NO_2	−0.68	0.32	this study
Rh-InSe		−3.09	1.52×10^{40}	62
Co-g-GaN		−3.67	1.00×10^{50}	59

place through the coordination of the S atom with two adjacent Ti atoms at a distance of 2.53 Å from each of them and with an energy release of −1.40 eV. The optimized geometry of NH_3 adsorbed on Ti_3C_2 shows the molecule lying on top of the surface with the N atom facing Ti and forming a

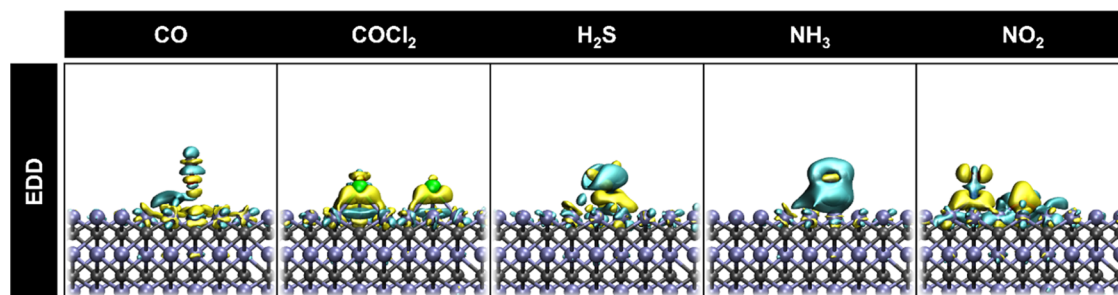


Figure 4. Electron density difference (EDD) between pristine Ti_3C_2 and the studied gases. Yellow (cyan) isosurfaces represent charge accumulation (depletion). Isovalues set to $0.0008 \text{ e}/\text{\AA}^3$.

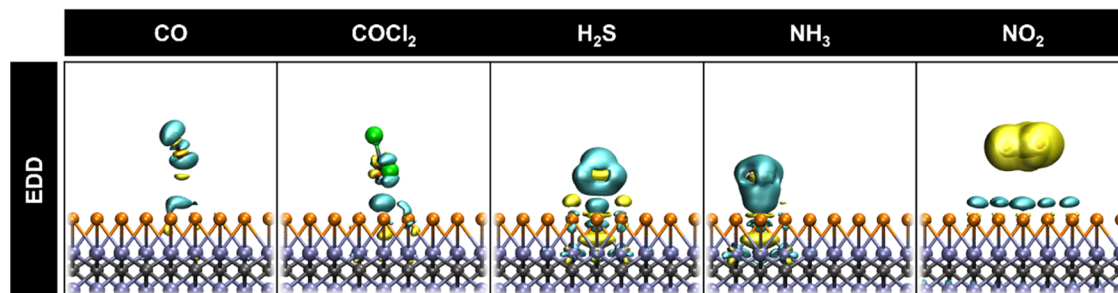


Figure 5. Electron density difference (EDD) between $\text{Au-Ti}_3\text{C}_2$ and the studied gas molecules. Yellow (cyan) isosurfaces represent charge accumulation (depletion). Isovalues set to $0.0008 \text{ e}/\text{\AA}^3$.

$\text{Ti}_2\text{-N}$ bond of length 2.30 \AA with a binding energy of -1.19 eV .

Moving on to the study of gas adsorption on Au-functionalized Ti_3C_2 (see Figure 3), we find physisorption in all considered cases, with comparatively large bond distances and weak interactions between the adsorbents and the underlying Au atoms. A slightly tilted configuration of CO with respect to the surface normal appears with the O pointing to the closest Au atom at a distance of 4.13 \AA : an adsorption energy of -0.14 eV is obtained in this case. Both values indicate a rather weak interaction between adsorbent and substrate. The adsorption of COCl_2 involves the formation of an Au–O bond with length 2.84 \AA while the Au–Cl interaction is markedly ionic with the two atoms at a distance of 3.24 \AA . The energy release obtained for this system is -0.54 eV . The binding of H_2S on Au-functionalized Ti_3C_2 occurs with the molecule in a horizontal arrangement: the Au–S bond has a length of 2.74 \AA while the separation between H and the closest Au atom is 3.10 \AA . An adsorption energy of -0.71 eV is obtained in this case. NH_3 and NO_2 are bound to the surface through Au–N interactions with energy releases of -0.66 and -0.68 eV , respectively. Despite these similar values of adsorption energy, the behavior of the two molecules on the Au-decorated Ti_3C_2 MXene is markedly different. In fact, NH_3 adsorbs at a distance of 2.44 \AA between N and Au while NO_2 is much further away from the substrate (N–Au distance: 4.63 \AA). We ascribe this behavior to the repulsion exerted to the MXene by the strongly electron-withdrawing NO_2 moiety. Finally, it can be observed that the adsorption of gases is energetically favorable in all cases which is quantified based on negative adsorption energies.

By comparing the adsorption energies obtained for CO, H_2S , and NH_3 on Ti_3C_2 and on its Au-functionalized counterpart, we notice that the computed values are larger for the pristine MXene, see Table 2. This evidence suggests that the gold

capping layer screens substrate–adsorbate interactions and it is further supported by their larger separations. We can use the computed values of adsorption energies to assess the performance of the MXene as a sensor for the considered toxic gas molecules, by evaluating the recovery by means of eq 2. The obtained values are listed in Table 2. Among the unreacted gas molecules on pristine Ti_3C_2 , CO shows the fastest recovery time of $0.35 \mu\text{s}$, making it an ideal analyte for pristine Ti_3C_2 -based gas sensors. In contrast, the recovery times calculated for H_2S and NH_3 can be considered infinitely long, indicating that sensing these two gas molecules is not possible on pristine Ti_3C_2 . In fact, the substrate cannot return its pristine status in a time window compatible with applications. On the other hand, the recovery times obtained for the toxins at $\text{Au-Ti}_3\text{C}_2$ are very reasonable in terms of sensor performance, ranging from 0.21 ns for CO adsorption to about one second for H_2S . Intermediate results are obtained for phosgene (1.60 ms), NH_3 (130 ms), and NO_2 (320 ms). These results indicate that gold functionalization of Ti_3C_2 can improve the detection of gaseous toxins compared to pristine material. Moreover, a comparison of performance parameters for various gas sensors is given in Table 3, indicating the superiority of the studied material over others reported in literature.

3.2. Charge Density Difference Analysis. After the assessment of the adsorption properties of Ti_3C_2 and its Au-functionalized counterparts, it is relevant to analyze the electronic properties of the composite systems. We start by considering the electron density difference between the electron density of the composite system and the one of its constituents. The visualization of this quantity helps understand the charge transfer or redistribution between the interacting systems and collect information about the nature of the chemical interactions in the complex. The charge transfer between pristine Ti_3C_2 and the unreacted gases is

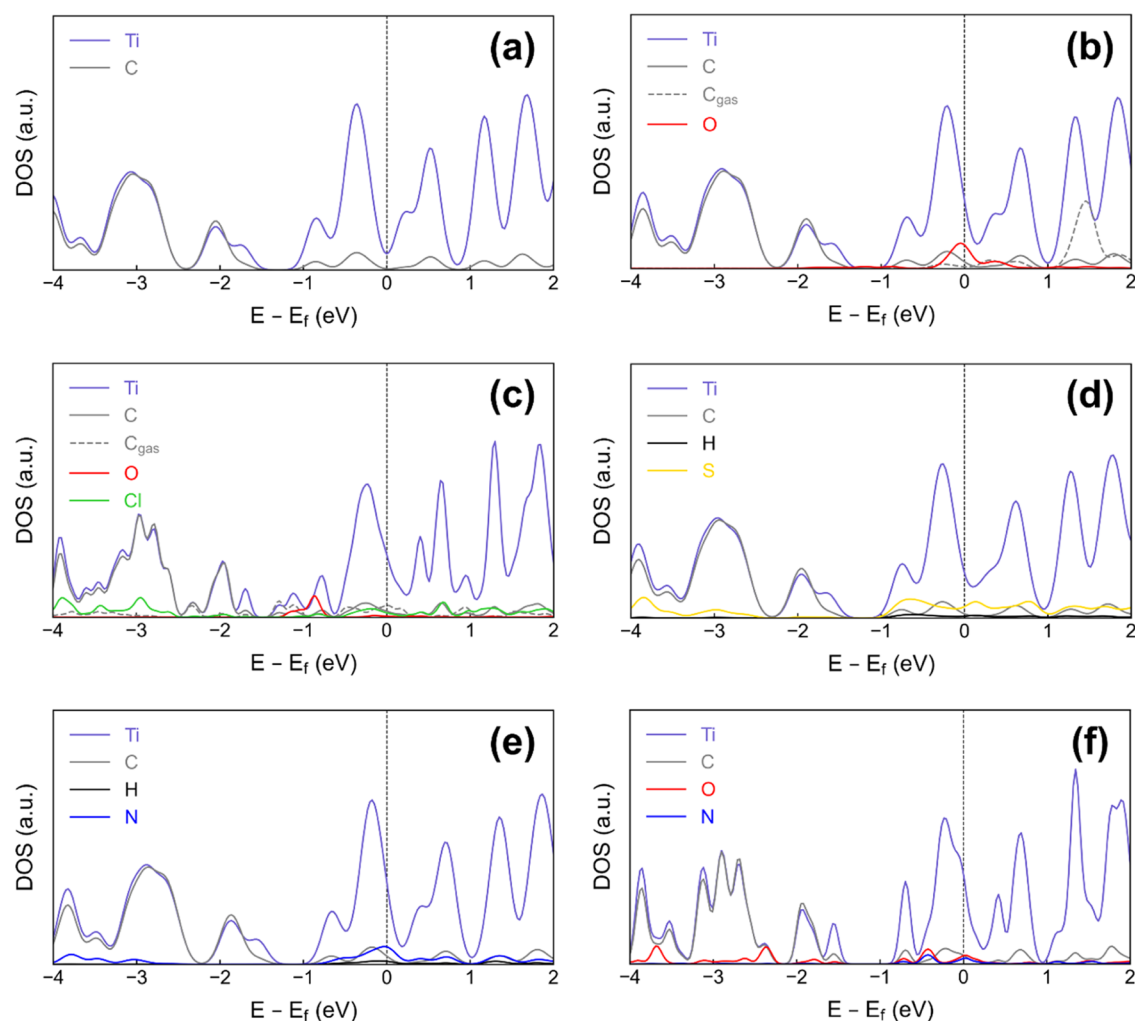


Figure 6. Projected density of states for (a) Ti_3C_2 , (b) $\text{CO}@ \text{Ti}_3\text{C}_2$, (c) $\text{COCl}_2@ \text{Ti}_3\text{C}_2$, (d) $\text{H}_2\text{S}@ \text{Ti}_3\text{C}_2$, (e) $\text{NH}_3@ \text{Ti}_3\text{C}_2$, and (f) $\text{NO}_2@ \text{Ti}_3\text{C}_2$. The Fermi energy (E_f) is set at zero.

illustrated in Figure 4. It can be observed that a small amount of electron density is distributed between O and Ti in the $\text{CO}@ \text{Ti}_3\text{C}_2$ system confirming the presence of the Ti–O chemical bond. In the $\text{H}_2\text{S}@ \text{Ti}_3\text{C}_2$ system, a relatively large electron accumulation between S and adjacent Ti atoms is obtained, indicating the formation of a Ti–S chemical bond. Similarly, in the $\text{NH}_3@ \text{Ti}_3\text{C}_2$ system, the Ti–N bond can be visualized since electron density is concentrated between N and Ti atoms. Furthermore, it can be noted that in all studied systems, electron density is withdrawn from the Ti atom. The low adsorption energy of $\text{CO}@ \text{Ti}_3\text{C}_2$ can be attributed to small charge transfer whereas the large electron density transfer in $\text{H}_2\text{S}@ \text{Ti}_3\text{C}_2$ and $\text{NH}_3@ \text{Ti}_3\text{C}_2$ is responsible for their relatively high binding energies.

The electron density difference isosurfaces for the gas@Au– Ti_3C_2 systems are depicted in Figure 5. It can be noticed that all the investigated toxic gases interact with the Au– Ti_3C_2 surface via weak intermolecular interactions, since in all cases most of the electron density is situated around the atoms and not at the gas–MXene interface. The weak intermolecular forces between gases and Au-decorated MXene can be accountable for comparatively lower binding energies and larger bond distances.

3.3. Density of States and Work Function. Next, we examine the electronic structure of the studied systems with

the aid of the projected density of states, see in Figures 6 and 7. The metallic character of Ti_3C_2 is evident in Figure 6a where the populated region crossing the Fermi level is expectedly dominated by Ti-states. Interestingly, the density of states of pristine Ti_3C_2 slightly changes after adsorption. These variations mainly consist of an upshift of the electronic states around the Fermi level (Figure 6b–f). In addition, the electronic population of CO, H_2S , and NH_3 molecules around the Fermi level indicates strong interactions between the sensor and analyte. The smearing of electronic states of the H_2S molecule is likely due to the redistribution of the electron density around adjacent Ti atoms, see Figure 4.

The projected density of states of Au-functionalized Ti_3C_2 MXene [Figure 7(a)] indicates that the MXene remains metallic even in the presence of Au surface modification. The overlap of Au and Ti density of states around the Fermi level shows the Au–Ti chemical bond visualized in the inset, where the redistribution of electron density (yellow isosurface) causes the surface to become positively charged (cyan isosurface). The density of states of Au– Ti_3C_2 remains unchanged on gas adsorption, as shown in Figure 7(b–f). Moreover, in contrast to the gas@ Ti_3C_2 systems, a lower and more smeared density of states is obtained which for the gas molecules, testifying the previously discussed absence of strong chemical interactions between the Au-functionalized Ti_3C_2 substrate and the toxins.

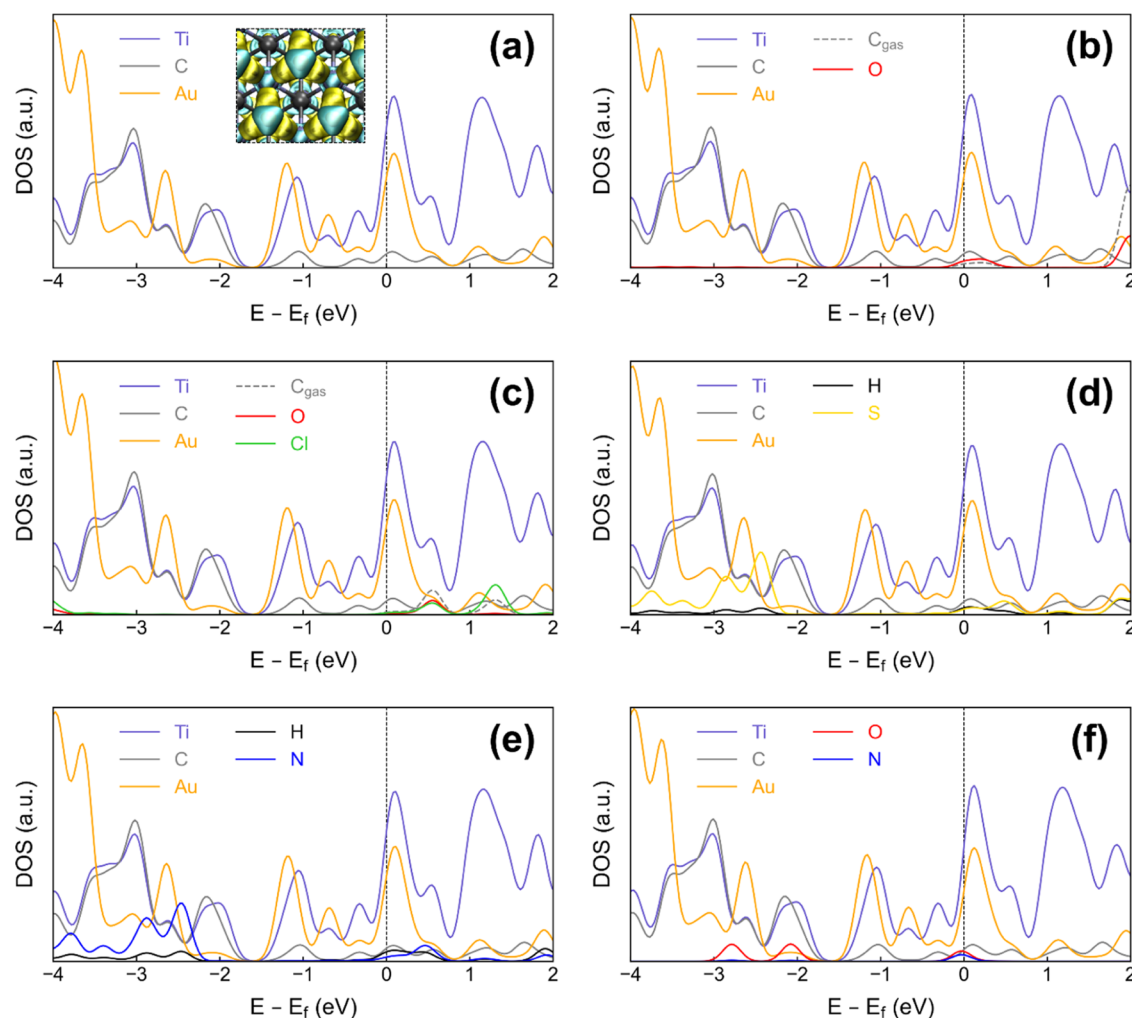


Figure 7. Projected density of states for (a) Au–Ti₃C₂ (inset shows EDD between Au and Ti₃C₂), (b) CO@Au–Ti₃C₂, (c) COCl₂@Au–Ti₃C₂, (d) H₂S@Au–Ti₃C₂, (e) NH₃@Au–Ti₃C₂, and (f) NO₂@Au–Ti₃C₂. The Fermi energy (E_f) is set to zero.

Table 4. Estimated Work Function (ϕ), Shift in Work Function ($\Delta\phi$), and Bader Charge (Q) on Gas Molecule for the Studied Systems

system	ϕ (eV)	$\Delta\phi$ (eV)	Q (e^-)
Ti ₃ C ₂	1.82		
CO@Ti ₃ C ₂	1.76	−0.06	0.17
H ₂ S@Ti ₃ C ₂	1.92	0.10	0.28
NH ₃ @Ti ₃ C ₂	1.84	0.02	0.02
Au–Ti ₃ C ₂	4.94		
CO@Au–Ti ₃ C ₂	5.01	0.07	0.01
COCl ₂ @Au–Ti ₃ C ₂	5.18	0.24	0.38
H ₂ S@Au–Ti ₃ C ₂	5.07	0.13	0.11
NH ₃ @Au–Ti ₃ C ₂	5.08	0.14	0.12
NO ₂ @Au–Ti ₃ C ₂	4.81	−0.13	0.15

The work functions and Bader charges for the considered systems are calculated to further assess the amount of interfacial charge transfer (Table 4). While pristine Ti₃C₂ is characterized by a work function of 1.82 eV, this value increases by more than 2.5 times in the presence of Au-functionalization. Comparing now the work function for each substrate with and without the gas molecules, we notice similar but not identical trends. Adsorption of H₂S and NH₃ leads to an increase in the work function of the pristine MXene by 100

and 20 meV, respectively. In contrast, interaction with CO lowers the work function by 60 meV. This behavior can be understood considering the larger polarity of the Ti–O bond formed between the MXene and this molecule, which withdraws electron from the substrate. A similar rationale can be used to interpret the results obtained for Au-functionalized Ti₃C₂. With all considered toxins except for NO₂, the work function increases compared to the one of the isolated gold-decorated MXene. The computed values range from 70 meV for CO@Au–Ti₃C₂ to 240 meV for COCl₂@Au–Ti₃C₂. In the H₂S@Ti₃C₂ system, a large increase in work function is obtained due to the binding of H₂S with two adjacent Ti atoms. Instead, the work function of NO₂@Au–Ti₃C₂ is 140 meV lower than the one of the isolated substrates.

4. CONCLUSIONS

We investigated Ti₃C₂ MXene (pristine and Au-functionalized) for sensing hazardous gases employing DFT calculations. Our findings revealed that the interaction of pristine Ti₃C₂ MXene with CO, H₂S, and NH₃ yields adsorption energies of −0.33, −1.40, and −1.19 eV, respectively. The resulting recovery time is 0.35 μ s for CO while it is significantly large with H₂S and NH₃. The analysis of the electronic properties showed that this material chemically reacts with these sample gases limiting its sensitivity as a

sensor. Au-functionalized Ti_3C_2 MXene, on the other hand, showed lower adsorption energy and reduced recovery time, which are characteristics of efficient and highly sensitive gas sensing materials. Notably, the Au-functionalized Ti_3C_2 MXene displayed exceptional sensing properties for CO with the lowest adsorption energy (-0.14 eV) and the shortest recovery time (0.21 ns) among studied sample gases. For COCl_2 , H_2S , NH_3 , and NO_2 , the adsorption energies are -0.54 , -0.71 , -0.66 , and -0.68 eV, respectively, resulting in recovery times ranging from 1.60 ms (with H_2S) to about 1 s (with COCl_2). The analysis of the electronic interactions of these composite materials revealed that the metallic nature of the MXene (both pristine and Au-functionalized) is unaltered on gas adsorption. The effect of the physisorbed toxins is to shift by a few hundreds of meV the electronic bands of the substrate, leading to corresponding variations in the calculated work function. This study shows that gold functionalization renders the MXene an effective detector of the sample gaseous toxins. In fact, the noble metal capping layer decreases the binding strength with the adsorbates and thus reduces the recovery time, making it suitable for sensing applications. The results presented in this work provide a basis for further theoretical and experimental investigations on Au-functionalized MXenes as promising materials for effective sensing of hazardous and toxic gases.

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Notes

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