# Tetraiododiborane(4) ( $\left.B_{2} I_{4}\right)$ is a Polymer Based on sp ${ }^{\mathbf{3}}$ Boron in the Solid State 

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#### Abstract

Herein we present the first solid-state structures of tetraiododiborane(4) ( $\left.B_{2} I_{4}\right)$, which was long believed to exist in all phases as discrete molecules with planar, tricoordinate boron atoms, like the lighter tetrahalodiboranes(4) $B_{2} F_{4}, B_{2} C l_{4}$, and $B_{2} B r_{4}$. Single-crystal $X$-ray diffraction, solid-state $N M R$, and IR measurements indicate that $B_{2} I_{4}$ in fact exists as two different polymeric forms in the solid state, both of which feature boron atoms in tetrahedral environments. DFT calculations are used to simulate the IR spectra of the solution and solid-state structures, and these are compared with the experimental spectra.


EEver since Stock's low-yielding synthesis of $\mathrm{B}_{2} \mathrm{Cl}_{4}$ in 1925, ${ }^{[1]}$ the structures of the tetrahalodiboranes $(4)\left(\mathrm{B}_{2} \mathrm{X}_{4} ; \mathrm{X}=\mathrm{F}, \mathrm{Cl}\right.$, $\mathrm{Br}, \mathrm{I})$ in their various phases (solid, liquid, gas) have been of fundamental interest to synthetic chemists. ${ }^{[2,3]}$ The structures of tetrahalodiboranes(4) were revisited around the turn of the millenium by theoretical groups, thanks to more reliable modern computational methods. ${ }^{[4]}$ However, despite the ubiquity of diborane(4) species as reagents in organic chemistry, ${ }^{[5]}$ experimental studies of tetrahalodiboranes(4) remain very rare, and have long been hampered by their difficult syntheses and/or instability.

The structures of Lewis acidic Group 13 compounds are often influenced by the vacant $p$ orbitals of the Group 13 atoms themselves, which can be stabilized in a number of different ways. Interactions of the empty p orbital of Group 13

[^0]"anes" $\left(\mathrm{EX}_{3} ; \mathrm{E}=\right.$ Group 13 element $)$ with nearby electron lone pairs form the basis of a number of polymeric species, ${ }^{[6]}$ and are common methods of stabilization adopted by the higher homologues of boranes (e.g. $\mathrm{AlBr}_{3}$ and $\mathrm{AlI}_{3}$ ). In contrast, corresponding bridging halide interactions are negligible for most haloboranes (e.g. $\mathrm{BX}_{3}$ and $\mathrm{B}_{2} \mathrm{X}_{4} ; \mathrm{X}=\mathrm{F}$, $\mathrm{Cl}, \mathrm{Br}),{ }^{[7]}$ except under high pressures. ${ }^{[8]}$ The interactions which are known to have the most influence on the conformations of tetrahalodiboranes(4) are hyperconjugation and $\pi$-donation from the halide substituents. $\mathrm{B}_{2} \mathrm{~F}_{4}$ was found to have a small rotational barrier between the planar $\left(D_{2 h}\right)$ and the perpendicular $\left(D_{2 d}\right)$ geometries; ${ }^{[4]}$ however, X-ray and electron diffraction techniques, as well as Raman and IR spectroscopy, have verified tetrafluorodiborane(4) that adopts the planar $\left(D_{2 h}\right)$ geometry in all phases (Figure 1). ${ }^{[2 \mathrm{a}, 4 \mathrm{c}, 9]} \mathrm{B}_{2} \mathrm{Cl}_{4}$ is known to be perpendicular $\left(D_{2 d}\right)$ in the gas and liquid phases, ${ }^{[2 c, e, 9 d, 10]}$ but to be planar $\left(D_{2 h}\right)$ in the solid phase. ${ }^{[2 \mathrm{~b}, 9 \mathrm{~d}, 10,11]}$ The heavier and sterically more demanding $\mathrm{B}_{2} \mathrm{Br}_{4}$ results in the perpendicular $\left(D_{2 d}\right)$ form in all phases. ${ }^{[2 f, \mathrm{j}]}$

The heaviest tetrahalodiborane (4), $\mathrm{B}_{2} \mathrm{I}_{4}$, is perhaps the least well-understood of this particularly underexplored family of compounds. Haubold predicted $\mathrm{B}_{2} \mathrm{I}_{4}$ to be perpendicular in all phases, based on IR spectra of solid material as a nujol mull and on assumption of structural similarity to $\mathrm{B}_{2} \mathrm{Br}_{4}$ (Figure 1). ${ }^{[12]}$ It should be noted that the IR spectra also showed contamination with boron triiodide and further decomposition products which may have arisen from the nujol trituration process. To our knowledge, the only other spectroscopic data reported for $\mathrm{B}_{2} \mathrm{I}_{4}$ is its solution ${ }^{11} \mathrm{~B}$ NMR resonance $(\delta=70 \mathrm{ppm}) .{ }^{[12,13]}$

In recent years, interest in the reactivity of tetrahalodiboranes(4) has almost exclusively come from our laboratories and from the group of Kinjo, where $\mathrm{B}_{2} \mathrm{Br}_{4}$ has been used for the construction of molecules with boronboron multiple bonds such as diborenes and diborynes, ${ }^{[3,14]}$ as well as precursors to bis(boratabenzene) ligands. ${ }^{[15]}$ We have recently broadened our interests in tetrahalodiboranes(4), developing simple solution-phase techniques to prepare the remaining $\mathrm{B}_{2} \mathrm{X}_{4}$ species (i.e. $\mathrm{X}=\mathrm{F}, \mathrm{Cl}, \mathrm{I}$ ) from $\mathrm{B}_{2} \mathrm{Br}_{4}{ }^{[13]}$ Our subsequent explorations of the reactivity of $\mathrm{B}_{2} \mathrm{I}_{4}$ have demonstrated that this reactive molecule differs significantly from the lighter tetrahalodiboranes(4) in terms of its reactivity with low-valent metal complexes, ${ }^{[16]}$ and that while (like $\mathrm{B}_{2} \mathrm{Br}_{4}$ ) it can be used as a precursor for the synthesis of dihalodiborenes, the products showed dramatic differences from their bromo analogues. ${ }^{[17]} \mathrm{B}_{2} \mathrm{I}_{4}$ has also since been subjected to twofold quaternization of its boron atoms in the form of double addition of Lewis bases (providing adducts of
solid, liquid, gas
This work:
solid-state structures of $\mathrm{B}_{\mathbf{2}} \mathbf{I}_{\mathbf{4}}$


Figure 1. Geometries of the tetrahalodiboranes(4) in their various physical states (top) and the infinite polymeric solid-state forms of $\mathrm{B}_{2} \mathrm{I}_{4}$ presented herein (bottom). Note that dotted B-I interactions are used to highlight each $\mathrm{B}_{2} \mathrm{I}_{4}$ unit for clarity, and do not imply weaker bonds or longer B-I distances.
the form $\left.\mathrm{B}_{2} \mathrm{I}_{4} \mathrm{~L}_{2}\right)^{[18]}$ and halide anions (providing the hexaiododiborate dianion, $\left.\left[\mathrm{B}_{2} \mathrm{I}_{6}\right]^{2-}\right) .{ }^{[19]}$

Despite the high synthetic potential of $\mathrm{B}_{2} \mathrm{I}_{4}$ in main-group, transition-metal, organic, and materials chemistry, there remains almost no experimental evidence for its structure, and no solid-state structures have been reported. Herein we report two solid-state structures of $\mathrm{B}_{2} \mathrm{I}_{4}$, showing the compound to have (at least) two crystal forms, both of which are polymers based on intermolecular B-I interactions. These polymeric structures are drastically different from those of the lighter tetrahalodiboranes(4), which exist as discrete molecules with planar, $\mathrm{sp}^{2}$-hybridized boron atoms in the solid state. These findings are supported by IR and solid-state NMR spectroscopy and computational studies.

Treatment of $\mathrm{B}_{2} \mathrm{Br}_{4}$ or $\mathrm{B}_{2} \mathrm{Cl}_{4}$ with 1.33 equiv of $\mathrm{BI}_{3}$ leads to the formation of solid $\mathrm{B}_{2} \mathrm{I}_{4}$, which can be easily isolated by removal of the $\mathrm{BCl}_{3}$ or $\mathrm{BBr}_{3}$ byproducts by vacuum distillation. ${ }^{[13]}$ Sublimation at $1 \times 10^{-3} \mathrm{mbar} / 45^{\circ} \mathrm{C}$ yielded colorless crystals of tetraiododiborane(4), which were used for the spectroscopic studies presented herein. It should be noted that the amorphous material obtained prior to sublimation provided effectively identical NMR and IR data. The IR spectrum of $B_{2} I_{4}$ in toluene solution matches that reported by


Figure 2. Solution and solid-state IR spectra of $B_{2} I_{4}$.

Haubold, which is consistent with a perpendicular $D_{2 d}$ conformation (Figure 2). However, the IR spectrum of the solid material recrystallized from toluene at $-30^{\circ} \mathrm{C}$ did not show the typical bands expected for either the $D_{2 d}$ or $D_{2 h}$ forms of the molecule (vide infra), thereby indicating another structural motif.

In order to gain a better picture of the solid-state structure of $\mathrm{B}_{2} \mathrm{I}_{4},{ }^{11} \mathrm{~B}$ RSHE/MAS (RSHE = rotor-synchronized HahnEcho, MAS = magic angle spinning) NMR spectra were recorded for the amorphous and crystalline material. In both spectra a broad resonance at $\delta=-16 \mathrm{ppm}$ (FWHH $=$ 2940 Hz ) was detected, which is far upfield from the corresponding resonance observed in the solution ${ }^{11} \mathrm{~B}$ NMR spectrum ( $\delta=70 \mathrm{ppm}$ ) and suggests that the boron centers in the compound are quaternized in the solid state. Comparison of the chemical shift of this resonance with the solution ${ }^{11} \mathrm{~B}$ NMR resonances of Lewis base adducts $\left[\mathrm{B}_{2} \mathrm{I}_{4}\left(\mathrm{SMe}_{2}\right)_{2}\right]$ $(\delta=-20 \mathrm{ppm}),{ }^{[13]}\left[\mathrm{B}_{2} \mathrm{I}_{4}\left(\mathrm{PMe}_{3}\right)_{2}\right](\delta=-27 \mathrm{ppm}),{ }^{[17]}$ and $\left[\mathrm{B}_{2} \mathrm{I}_{4}-\right.$ $\left.\left(\mathrm{PCy}_{3}\right)_{2}\right] \quad(\delta=-27 \mathrm{ppm})^{[17]}$ support this assumption. Additional sharp minor signals detected in the ${ }^{11}$ B RSHE/MAS NMR spectra can be assigned to liquid $\mathrm{B}_{2} \mathrm{I}_{4}(\delta=70 \mathrm{ppm})$, liquid $\mathrm{BI}_{3}(\delta=-5 \mathrm{ppm})$ and further liquid decomposition products ( $\delta=56,19 \mathrm{ppm}$ ) (see the Supporting Information). This mixture can be explained by partial melting of the sample caused by pressure arising from MAS conditions (rotation at 15 kHz ), and presumably subsequent decomposition of $\mathrm{B}_{2} \mathrm{I}_{4}$. Experiments conducted with lower rotation rates ( $7-11 \mathrm{kHz}$ ) led to less decomposition of the sensitive compound, but did not eliminate the melting or the formation of the decomposition products.

Confirmation of the nonplanar boron atoms of $\mathrm{B}_{2} \mathrm{I}_{4}$ in the solid state was unequivocally provided by single-crystal X-ray diffraction of three separate samples. Single crystals grown from saturated toluene solutions at $-30^{\circ} \mathrm{C}$ provided one crystal form of the compound ( $\alpha$ form; Figure 3, top), while crystallization at ambient temperature or by sublimation (40$45^{\circ} \mathrm{C},<1 \times 10^{-3} \mathrm{mbar}$ ) provided a second form ( $\beta$ form; Figure 3, bottom). In both isomers the $\mathrm{B}_{2} \mathrm{I}_{4}$ units are connected via bridging iodides, which bridge boron atoms in a symmetrical fashion such that the two B-I distances of each


Figure 3. Linear ( $\alpha$ ) and helical ( $\beta$ ) solid-state forms of $\mathrm{B}_{2} \mathrm{I}_{4}$ as derived from single-crystal X-ray diffraction. Note that dotted $\mathrm{B}-\mathrm{I}$ interactions are used to highlight each $\mathrm{B}_{2} \mathrm{I}_{4}$ unit for clarity, and do not imply weaker bonds or longer B-I distances.
bridging iodide are effectively identical, resulting in infinite $1 D$ polymers. In the $\alpha$ form, dimers of $B_{2} I_{4}$ units form through three bridging iodides, leading to a norbornane-like [2.2.1] bicyclic structure formed from two five-membered $\mathrm{B}_{3} \mathrm{I}_{2}$ rings. Each dimer is connected to two other dimers in a linear fashion through two further $\mathrm{B}-\mathrm{I}$ bonds. The $\beta$ form is also made up of $\mathrm{B}_{2} \mathrm{I}_{4}$ dimers. In each dimer, two $\mathrm{B}_{2} \mathrm{I}_{4}$ units are connected through two fused monocyclic $\mathrm{B}_{3} \mathrm{I}_{2}$ rings, forming an edge-shared [3.3.0] bicyclic structure. On the supramolecular level, the $\alpha$ form adopts a roughly linear structure, while the $\beta$ form has a distinct helical superstructure with a pitch of ca. 10.7 Å.

The $B-B$ bond lengths of the two solid-state forms of $B_{2} I_{4}$ are almost identical ( $\alpha$ form: 1.695(6) $\AA$; $\beta$ form: $1.68(2) \AA$ ), although the high experimental error in the latter precludes more detailed analysis. As expected, the B-I distances (2.301(4)-2.336(4) $\AA$ ) involving bridging iodides are significantly longer than those of terminal iodides ( $\alpha$ form: 2.184(4) $-2.193(4) \AA$; $\beta$ form: $2.189(9)-2.180(9) \AA)$.

Theoretical simulations were performed to establish a correlation between the structural and spectral properties of $\mathrm{B}_{2} \mathrm{I}_{4}$ in solution and in its two solid-state forms. Firstly, geometry optimization confirmed that the minimum-energy structure of molecular $\mathrm{B}_{2} \mathrm{I}_{4}$ in toluene solution possesses $D_{2 d}$ symmetry. Secondly, the calculated IR spectrum of an isolated $\mathrm{B}_{2} \mathrm{I}_{4}$ molecule in the $D_{2 d}$ conformation is in very good agreement with the experimental solution IR spectrum (Figure 4) and with the assignment of vibrations made by Haubold, ${ }^{[12]}$ confirming the $D_{2 d}$-symmetric geometry of $\mathrm{B}_{2} \mathrm{I}_{4}$ in solution. The theoretical IR spectrum represents a weighted


Figure 4. Comparison of the experimental solution-state (toluene) IR spectrum of $B_{2} I_{4}$ with the simulated IR spectrum of isolated $B_{2} I_{4}$ molecules ( $D_{2 d}$ symmetry). The most intense vibrational modes are presented as insets.
sum of vibrational spectra of $\mathrm{B}_{2} \mathrm{I}_{4}$ species with different isotopic compositions corresponding to a relative abundance of ${ }^{10} \mathrm{~B}$ and ${ }^{11} \mathrm{~B}$ isotopes in a ratio of 20:80 (see the Supporting Information for individual spectra of the isotopomers).

Simulations of the vibrational spectra of both solid-state structures ( $\alpha$ and $\beta$ forms) were also performed. The calculated spectra of the two crystal forms possess noticeably distinct features, as shown in Figure 5a (green line) and 5b (blue line). The three most intense vibrational modes of the


Figure 5. $\mathrm{a}, \mathrm{b}$ ) Comparison of the experimental (solid-state) spectrum of $\mathrm{B}_{2} \mathrm{I}_{4}$ with the simulated IR spectra of the linear (a) and helical (b) forms of $\mathrm{B}_{2} \mathrm{I}_{4}$.c) Depiction of the vibration modes corresponding to bands $A, B$, and $C$ in (b).
helical $\beta$ structure represent deformation (A), rotation (B), and shifting (C) of the boron skeleton, as shown in Figure 5c (see Movies $\mathrm{A}, \mathrm{B}$, and C provided in the Supporting Information). The experimental solid-state IR spectrum of $\mathrm{B}_{2} \mathrm{I}_{4}$ (Figure $5 \mathrm{a}, \mathrm{b}$, black line) is poorly fit by the calculated IR of the $\alpha$ form (Figure 5 a), but is fit very well by that of the $\beta$ form (Figure 5b). Although the crystals used for the IR spectrum were grown at $-30^{\circ} \mathrm{C}$ from toluene, and thus should correspond to the $\alpha$ form, a change to the $\beta$ form likely takes place when the sample is ground and pressed into a pellet for the measurement. This hypothesis is in line with the pressure sensitivity of $\mathrm{B}_{2} \mathrm{I}_{4}$ observed during the fast rotation of the sample when we measured the solid-state MAS NMR spectra. It should also be noted that the amorphous material obtained from solvent evaporation (Figure 3) provides an essentially identical solid-state IR spectrum. The moderately intense vibrational band at ca. $670 \mathrm{~cm}^{-1}$, which is completely absent in the simulated IR spectra, can be attributed to a minor fraction of molecular $\mathrm{B}_{2} \mathrm{I}_{4}$ in the $D_{2 d}$ form, as established in solutions of the compound.

An unequivocal rationale for the differences in solid-state structures between $\mathrm{B}_{2} \mathrm{I}_{4}$ and its lighter homologues is elusive, although the observed bridging of the iodide groups is indeed intuitive from consideration of the properties of the related monoboron halides $\mathrm{BX}_{3}(\mathrm{X}=\mathrm{F}$ to I$)$. The Lewis acidity of $\mathrm{BX}_{3}$ species is generally acknowledged to increase in the order $\mathrm{BF}_{3}<\mathrm{BCl}_{3}<\mathrm{BBr}_{3}<\mathrm{BI}_{3}$, although the reasons behind this have been the subject of extensive debate. ${ }^{[20]} \mathrm{A}$ theoretical study using the addition of two electrons to $\mathrm{BX}_{3}$ species (in order to model an extreme case of a Lewis base) found that the two-electron affinity of the B atom of $\mathrm{BX}_{3}$ species increases as the halide is changed from $F$ to $I$, while the energy required to pyramidalize the compounds also becomes lower. ${ }^{[20]}$ These two effects are likely also at play in the $\mathrm{B}_{2} \mathrm{X}_{4}$ series, overall making pyramidalization at boron and iodide bridging favorable in the solid-state structure of $\mathrm{B}_{2} \mathrm{I}_{4}$ but not its lighter homologues.

In order to investigate the bonding in the solid-state structures of $\mathrm{B}_{2} \mathrm{I}_{4}(\alpha)$ and $\mathrm{B}_{2} \mathrm{I}_{4}(\beta)$ in more depth, an analysis based on the quantum theory of atoms in molecules (QTAIM) ${ }^{[21]}$ was performed. The QTAIM representation of the helical $(\beta)$ structure is schematically shown in Figure 6. The electron density parameters at the bond critical points (BCPs) and the bonding analysis of the linear ( $\alpha$ ) structure are provided in the Supporting Information. The simulations reveal that, in both solid-state structures, the strength of the iodide-bridging bonds $\left(B-I_{B}\right.$, blue lines in Figure 6) is approximately $80 \%$ of that of the terminal $B-I$ bonds ( $B-I_{T}$, yellow lines), making them stable against cleavage. According to a reported theoretical study, ${ }^{[22]}$ the strengths of bridging halogen bonds relative to the terminal bonds increases from $F$ to I, which makes the formation of halogen bridges possible for $\mathrm{B}_{2} \mathrm{I}_{4}$ but not for the lighter homologues. Moreover, our QTAIM calculations have uncovered the presence of BCPs between the iodides in the $\left[\mathrm{B}_{2} \mathrm{I}_{4}\right]_{n}$ strands of both crystal forms, suggesting weak I $\cdots$ I interactions based on dispersion forces (magenta lines in Figure 6). These calculated intrastrand $\mathrm{I} \cdots \mathrm{I}$ interactions add a further stabilizing effect to


Figure 6. QTAIM representation of the solid-state structure of $\mathrm{B}_{2} \mathrm{I}_{4}(\beta)$. Semitransparent spheres denote the positions of boron (B), terminal iodide $\left(I_{T}\right)$, and bridging iodide $\left(I_{B}\right)$ atoms. The following electron density critical points are depicted: attractors (black), bond critical points (red), ring critical points (light blue), and local minima (gray).
a polymeric species that would not be possible with the lighter homologues $\mathrm{B}_{2} \mathrm{X}_{4}(\mathrm{X}=\mathrm{F}, \mathrm{Cl}, \mathrm{Br})$.

Herein we use IR spectroscopy, single-crystal X-ray diffraction, and DFT calculations to confirm the solution structure of $\mathrm{B}_{2} \mathrm{I}_{4}$ and establish its polymeric nature in the solid state. Two polymeric crystal forms are structurally authenticated, one of which possesses a distinct helical structure. Through computational simulation of the IR spectra of the two solid forms, the experimental solid IR spectrum of the samples is shown to correspond to the form adopted at higher temperatures. The massive discrepancies between the solidstate structures of $\mathrm{B}_{2} \mathrm{I}_{4}$ and the lighter homologues $\mathrm{B}_{2} \mathrm{X}_{4}(\mathrm{X}=$ $\mathrm{F}-\mathrm{Br})$ further underline the fundamental differences observed in the few available reactivity studies, ${ }^{[16-19]}$ and hint at more surprises to come in the chemistry of the tetrahalodiboranes(4).

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## Conflict of interest

The authors declare no conflict of interest.
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