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## Research article

# Noise-induced network bursts and coherence in a calcium-mediated neural network

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# ABSTRACT

Noise-induced population bursting has been widely identified to play important roles in information processes. We construct a mathematical model for a random and sparse heterogeneous neural network where bursting can be induced from a resting state by a global stochastic stimulus. Importantly, the noise-induced bursting dynamics of this network are mediated by calcium conductance. We use two spectral measures to evaluate network coherence in the context of the network bursts, the spike trains of all neurons, and the individual bursts of all neurons. Our results show that the coherence of the network is optimized by an optimal level of the stochastic stimulus, which is known as coherence resonance (CR). We also demonstrate that the interplay of the calcium conductance and noise intensity can modify the degree of CR.

#### 1. Introduction

Bursting is one of the fundamental coding strategies for neuronal information processing and transmission in the brain [1, 2, 3, 4]. Its temporal pattern is characterized by the repetitive switch between a silent phase (with almost no spike emission) and an active phase (with two or more spikes with high firing rates). Network bursts (or population bursts) refer to synchronous or near-synchronous burst firing across a neural network [5]. The generation of burst firing is regulated by low-threshold calcium channels in various neuronal populations [6, 7]. Many calcium imaging studies have reported that such neuronal populations are relatively random, dispersed networks [8]. The synaptic connectivity between neurons is anti-correlated with their lateral distance [9, 10], thus, bursting networks mediated by calcium channels have a rather low connection probability.

Coherence is one of the most common measures used to quantify the correlation or synchronicity of the oscillatory patterns of neurons across a neural network [11]. The collective activities of neural networks are often influenced by a ubiquitous and often significant component—noise [12]. Generally, noise can be either local (independent and uncorrelated for each neuron in the network), or global (identical across the network) [13, 14]. It has long been shown that noise can play a constructive role to improve the performance of a dynamical sys-

tem through, for example, coherence resonance (CR). CR is a resonant mechanism where an appropriate amount of noise alone (i.e. without external periodic stimulus) drives a quiescent but excitable system to produce the most coherent oscillations [15, 16]. In neuronal dynamics, oscillations represent the neuronal spikes or bursts and the system represents a neuron or a network. CR has been observed in neural networks such as globally coupled networks [17, 18, 19, 20, 21, 22], randomly connected neural networks exhibiting single oscillations [23, 24], smallworld networks [25, 26], ring networks [27, 28], multiplex networks [29, 30, 31], and the influencer network of phase oscillators [32]. However, the stochastic dynamics of a calcium-mediated random and sparse heterogeneous bursting network have not been extensively investigated and the effect of CR in such a network remains elusive.

In this work, we consider a quiescent but excitable network mediated by calcium currents, where the synaptic connections between neurons are random and sparse. In order to reproduce the collective activities of such a network more realistically, spike-timing-dependent plasticity (STDP) is integrated into our network model. STDP is used to simulate the dynamic synapse between two communicating neurons, where the spike timing information of presynaptic and postsynaptic neurons is used to adjust the synaptic strength over time [33]. As a result, temporal information conveyed by spike timing is spatially distributed across a network [34]. Hence, three heterogeneous elements

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are incorporated into the proposed network model: diverse calcium conductance, dynamic synaptic currents due to random connectivity and STDP, and distinct local noise.

This study focuses on the influences of the first heterogeneous compartment, calcium conductance, on network dynamics. We first consider constant calcium conductance and show that a global stochastic stimulus can induce network bursts and increase coherence. In particular, our analysis demonstrates that at an optimal intensity of the global noise, the similarity relations between neurons are maximized, indicating the occurrence of CR. Secondly, by altering the values of the constant calcium conductance, we find that CR is robust and that the maximum degree of coherence increases when the calcium conductance approaches the excitation threshold. Finally, we explore the impact of heterogeneous calcium conductance on network coherence.

The remainder of this paper is organized as follows: Section 2 contains the description of the network model. Section 3.1 introduces the noise-induced bursting generated by this network. Sections 3.2-3.4focus on the CR of the network bursts, all spikes, and all bursts, respectively. Section 3.5 examines the effect of heterogeneous calcium conductance on network coherence. Discussion is given in Section 4.

#### 2. Mathematical model

We model a network of N = 100 neurons with random synaptic connections, which adapts the form of the reduced Morris-Lecar model with a linear slow subsystem [35].

$$C\frac{dv_i}{dt} = I_i - I_{Ca,i} - I_{K,i} - I_{L,i} + I_{loc,i} - I_{syn,i} + I_{glo}$$
(1)

$$\frac{dw_i}{dt} = \phi \lambda_w(v_i)(w_\infty(v_i) - w_i)$$
<sup>(2)</sup>

$$\frac{dI_i}{dt} = \epsilon(v_0 - v_i) \tag{3}$$

$$\frac{dg_i}{dt} = -\frac{g_i}{\tau_e} \tag{4}$$

where  $v_i$  is the membrane potential of the *i*th neuron for i = 1, 2, 3, ..., N.  $I_{Ca,i} = g_{Ca}m_{\infty}(v_i)(v_i - v_{Ca})$ ,  $I_{K,i} = g_kw_i(v_i - v_k)$ , and  $I_{L,i} = g_l(v_i - v_l)$  are the calcium, potassium, and leakage currents, respectively, with gating functions

$$m_{\infty}(v_{i}) = \frac{1}{2} \left( 1 + \tanh \frac{v_{i} - v_{1}}{v_{2}} \right),$$
  

$$w_{\infty}(v_{i}) = \frac{1}{2} \left( 1 + \tanh \frac{v_{i} - v_{3}}{v_{4}} \right),$$
  

$$\lambda_{w}(v_{i}) = \frac{1}{3} \cosh \frac{v_{i} - v_{3}}{2v_{4}}.$$

 $w_i$  is a gating variable of  $I_{K,i}$  with  $\phi$  as the scaling rate of channel opening.  $I_i$  is the linear feedback input current with feedback coefficient  $\epsilon$ .  $g_i$  is the time-varying conductance of the synaptic current with time constant  $\tau_e$ .  $g_{Ca}$ ,  $g_k$ , and  $g_l$  are the maximum conductance of the calcium, potassium, and leakage currents, with corresponding reversal potentials  $v_{Ca}$ ,  $v_k$ , and  $v_l$ .  $I_{loc,i} = D_1\xi_i$  and  $I_{glo} = D_2\eta$  represent the local intrinsic noise (unique for each neuron) and global external stochastic stimulus (same for all neurons), respectively, where  $\xi_i$  and  $\eta$  are independent Gaussian white noise with mean 0 and variance 1, and  $D_1$  and  $D_2$  are scaling parameters for the local and global noise intensities, respectively. The parameter values of this model are listed in Table 1.

This network is randomly connected with a probability of 15%. That is, the connection of each pair of neurons is randomly and independently realized with a probability of 15%, which is a realistic assumption for a calcium-sensitive neural network based on experimental study in [10]. An example synaptic connectivity map is presented in Fig. 1A. The synaptic current for the *i*th neuron,  $I_{syn,i}$ , in Equ. (1), is averaged by the number of incoming connections from pre-synaptic

#### Table 1. Parameter values.

Parameter	Value	Unit	Parameter	Value	Unit
$v_0$	-20	mV	v <sub>I</sub>	-50	mV
$v_1$	-1	mV	$g_k$	1.2	mS
$v_2$	15	mV	$g_l$	0.6	mS
$v_3$	10	mV	$\phi$	1	1/ms
$v_4$	5	mV	e	0.001	mS/ms
$v_{Ca}$	90	mV	С	1	μF
$v_k$	-100	mV	w <sub>e</sub>	0.03	mS
$\Omega_d$	4	1/s	U <sub>e</sub>	20	mV
$\Omega_f$	4	1/s	$g_{Ca}$	$0.63\sim0.646$	mS
$ au_e$	0.55	ms			

neurons. Namely,  $I_{syn,i} = \frac{1}{N_{pre}} \sum_{\substack{j=1 \ j \neq i}}^{N} p_{ji} g_i (v_j - v_e)$  where  $N_{pre}$  represents the number of incoming synapses to neuron *i*;  $p_{ji} = 1$  if the *j*th neuron (pre-synaptic) and *i*th neuron (post-synaptic) are connected and 0 otherwise; and  $\sum_{i=1}^{N} \sum_{j=1}^{N} p_{ji} = 15\% \times N \times N = 1500$ . The synapses are

modeled by a phenomenological model with a STDP mechanism [36, 37], which is an important feature for synaptic memory formation and removal. In particular, synapse release is defined by the product of the two variables,  $x_s$  and  $u_s$ , which represent the fractions of neurotransmitter available and docked for release, respectively. Between action potentials,  $u_s$  and  $x_s$  follow the dynamics

$$\frac{du_s}{dt} = -\Omega_f u_s,$$
$$\frac{dx_s}{dt} = \Omega_d (1 - x_s).$$

Whenever a pre-synaptic action potential arrives at a post-synaptic cell, the excitatory conductance increases according to  $g_i \leftarrow g_i + w_e u_s x_s$ , where  $w_e$  is the synaptic weight.

To simulate our network model, we used the *Brian2* package in Python and the Euler-Maruyama method. We then exported the simulation results (i.e., the spike times, global stochastic stimuli, and membrane potentials) and proceeded with our analysis using MATLAB. For example, MATLAB was used to compute the peri-stimulus time histogram (PSTH), power spectrum density (PSD), and signal to noise ratios (SNRs). 50 trials were used to average the PSDs and SNRs in Figs. 2–8. The network connectivity is different in different trials but with a constant connection probability of 15%. In Fig. 8, a different set of uniformly distributed random numbers was generated and used to model different  $g_{Ca}$  values in the network in different trials.

# 3. Results

#### 3.1. Noise-induced bursting

This study examines the stochastic dynamics of an excitable neural network, whose population activity is characterized by bursts when a global stochastic stimulus is applied. Its corresponding deterministic network (where  $D_1 = D_2 = 0$ ) is quiescent for  $g_{Ca} < 0.648$  and exhibits periodic bursts of three or more spikes when  $g_{Ca} \ge 0.648$ . The proposed network rests in the excitable regime (i.e.  $g_{Ca} < 0.648$ ), where bursting spikes are stimulated by local intrinsic noise,  $D_1\xi_i$ , or external stochastic input,  $D_2\eta$ . Here the network has two heterogeneous components: the local noise  $(I_{loc,i})$  and synaptic currents  $(I_{syn,i})$ , due to random connections between neurons and STDP). When  $g_{Ca}$  is slightly lower than the excitation threshold (0.648), bursting spikes can be evoked by local noise alone, as shown by one voltage segment in Fig. 1B with  $g_{Ca} = 0.646$ ,  $D_1 = 0.007$ , and  $D_2 = 0$ . Our simulation shows, by taking  $g_{Ca} = 0.646$ ,  $D_1 = 0.007$ , and  $D_2 = 0$ , bursts occur for all 50 trials with a low average occurrence rate of population bursts (0.64 bursts/second). When  $g_{Ca}$  is much lower than 0.648, local noise alone can not evoke bursting behavior (see the voltage segment with  $g_{Ca} = 0.642$ ,  $D_1 = 0.007$ , and  $D_2 = 0$  in Fig. 1B).

The addition of a global stochastic stimulus,  $D_2\eta$ , increases both the burst-generation probability and the occurrence rate of population



**Fig. 1.** (A) An example synaptic connectivity map for a network of N = 100 neurons with a random connection probability of 15%. The connectivity varies for different trials. (B) Example voltage traces of one neuron in this network in response to a global stochastic stimulus,  $D_2\eta$ . This global stochastic stimulus (top row) has three different intensities  $D_2 = 0,0.03$ , and 0.15 for every 1500 ms segment. The intensity of local intrinsic noise,  $D_1$ , is constant over the three 1500 ms segments ( $D_1 = 0.007$ ). Note,  $g_{Ca} = 0.648$  is the excitation threshold between the quiescent and bursting regimes of the deterministic state of this network (where  $D_1 = D_2 = 0$ ). Thus, two calcium conductance values are taken (middle row:  $g_{Ca} = 0.646$  and bottom row:  $g_{Ca} = 0.642$ ) to illustrate the dynamical change of voltage traces.

bursts. As seen in the voltage traces of Fig. 1B, when the stochastic stimulus changes from a weak level ( $D_2 = 0.03$ ) to a relatively higher level ( $D_2 = 0.15$ ), the burst rate increases. Meanwhile, for a larger  $D_2$ , the number of spikes in a single burst event becomes more random. For example, Fig. 2B (middle row) shows that most bursts have 3 spikes when  $D_2 = 0.03$ , whereas the number of spikes within one burst ranges from 2 to 5 when  $D_2 = 0.15$ . Another observation is that an increase in  $D_2$  causes a higher voltage fluctuation on the slow silent phase in between consecutive burst events (i.e., the hyperpolarization stage of action potentials where voltage is around -27 mV to -16 mV), as illustrated in Fig. 1B, where  $D_2 = 0.15$ . The fluctuations in the slow silent phase can be used to determine the intensity level of global stochastic input, for example,  $D_2 = 0.03$  and 0.15 correspond to weak and intermediate levels, respectively.

#### 3.2. Coherence resonance (CR) of network bursts

To study network dynamics-and by extension network coherencewe subject our network to weak, intermediate, and strong levels of the global stochastic stimulus,  $D_2\eta$ . The change in network bursting can be visualized by raster plots and peri-stimulus time histograms (PSTH), as shown in Fig. 2A-2C. A raster plot is a collection of the spike times of individual neurons in a network, where each black dot in the raster plot represents a spike. The PSTH (blue curves in Fig. 2A-2C) summarizes the number of spikes from all neurons across the network at a certain time. Therefore, it records the timing of network bursts, and the height and width of the PSTH peaks indicate the synchrony of individual spikes within a burst. When a rather weak stimulus is applied, the network produces fewer population bursts. For example, there are 4 population bursts over a 1600 ms period with  $D_2 = 0.1$ , as shown in Fig. 2A. The height and the width of some of the PSTH peaks are relatively short and wide, respectively, which reflects a relatively low spike-to-spike synchronization within such burst events (see the 2nd burst in Fig. 2A). When the stimulus is increased to an intermediate level (e.g.  $D_2 = 0.225$  in Fig. 2B), bursting activity becomes more frequent and the PSTH has higher and narrower peaks; spiking events are tightly contained in bursts and the network becomes highly synchronized. Note,  $D_2 = 0.225$  is chosen because it is the optimal stimulus intensity of CR (as in Fig. 4A). However, when the stimulus is further increased to stronger levels (e.g.  $D_2 = 0.4$  in Fig. 2C), network bursts become more frequent but get less synchronized. The sharp peaks of the PSTH become broader, indicating that noise starts to dictate the dynamics of the network.

Such temporal changes in network dynamics may also be viewed by the power spectral density (PSD) in the frequency domain. PSTHs are used to calculate the PSD and the average PSD over 50 trials is shown in the inset of Fig. 2D. The black, red, and green PSD curves are labeled by letters A, B, and C, and correspond to Fig. 2A (weak stimulus case), Fig. 2B (intermediate/optimal stimulus case), and Fig. 2C (strong stimulus case), respectively. Three major features of a PSD are often considered. The first is the central frequency (or called the resonant frequency), which is the frequency location of the highest PSD point and the reciprocal of the average inter-burst interval (IBI). The central frequency increases with  $D_2$ , which agrees with the network dynamics in the time domain (Fig. 2A-2C) where the IBI is smaller with the increment of  $D_2$ . The central frequency is also positively correlated with the burst rate because a shorter IBI implies a higher burst rate. The second and third PSD features considered are the height and half-width of PSD peaks. As expected, the optimal stochastic stimulus results in the most pronounced PSD peak (red curve in Fig. 2D inset, with the largest height and the smallest half-width) as opposed to the cases of weak and strong stochastic stimulus (black and green curves in Fig. 2D). This is caused by higher and narrower PSTH peaks as shown in Fig. 2B.

The collective activities observed in Fig. 2A-2C and the inset of Fig. 2D indicate that the intensity of the global stochastic stimulus plays an important role in modifying the coherence of our network. To measure coherence more concretely, we use the signal-to-noise ratio (SNR) measure [15, 16],

$$\alpha = h_p (\Delta \omega / \omega_p)^{-1}, \tag{5}$$

where  $h_p$  and  $\omega_p$  denote the height and central frequency of the PSD peak, respectively, and  $\Delta\omega$  denotes the width of the PSD peak at halfmaximal power. For this network,  $\omega_p$  has a rather slight change as  $D_2$  increases (see Fig. 2D inset), so the ratio between  $h_p$  and  $\Delta\omega$  dominates  $\alpha$ . As discussed, the PSD peaks are most pronounced (i.e., large  $h_p$  and small  $\Delta\omega$ ) at the intermediate stimulus values as opposed to the weak and strong stimulus values. Therefore the SNR- $\alpha$  should peak at intermediate levels of the stochastic stimulus. The SNR- $\alpha$  curve is presented in Fig. 2D. One sees that for a weak stimulus values, and then decreases and tends toward zero for strong stimulus values. This is the characteristic pattern of CR [15, 16] and the peak of the SNR curve corresponds to the maximum degree of network coherence. The intensity of a stochastic stimulus that maximizes SNR is called the optimal intensity, and for  $g_{Ca} = 0.64$  in Fig. 2D, the optimal intensity is  $D_2 = 0.225$ .

#### 3.3. The network coherence in terms of all spikes

PSTH is a collective quantity describing population bursts and does not accurately capture the fast dynamics of intra-burst spikes (i.e., individual spikes within a burst), and as a result, its PSD (and coherence measure SNR- $\alpha$ ) covers only the low-frequency range of 0-15 Hz. To



**Fig. 2.** (A)-(C) The spike raster plots (black dots) and the peri-stimulus time histogram (PSTH, blue curves) for  $D_2 = 0.1$ , 0.225 and 0.4 (representing weak, intermediate, and strong stimulus levels, respectively) when  $g_{Ca} = 0.64$ . The vertical axis corresponds to the index of a neuron in the network. Each dot indicates that one neuron generates a spike at the time corresponding to the horizontal (time) axis. The bin-width of the PSTH is 20 ms. (D) Signal-to-noise ratio (SNR)  $\alpha$  v.s. the intensity of stochastic stimulus  $D_2$ . SNR- $\alpha$  is defined in Equ. (5). Letters A, B and C correspond to  $D_2$  values as in panels (A)-(C), which represent three noise intensities of global stimulus: weak; optimal; and strong. Inset of (D): the power spectral density (PSD) of PSTH v.s. frequency for  $D_2 = 0.1$ , 0.225 and 0.4. PSDs are computed based on PSTH and averaged over 50 trials. Letters A-C beside the PSD curves indicate their corresponding example raster plots and PSTHs in panels (A)-(C).

study network coherence, we must consider larger bandwidths of frequency to account for both bursts and intra-burst dynamics (i.e., both fast and slow dynamics). Therefore, we analyzed the spike trains of individual neurons in the network. The histogram of inter-spike intervals (ISIs) of spike trains shows a bimodal distribution (two separate and independent peaks): one peak is located at shorter ISIs corresponding to the fast intra-burst dynamics, and the other one is at longer ISIs due to the slow dynamics of burst events. As shown in Fig. 3A, with  $g_{Ca} = 0.64$ and  $D_2 = 0.05$ , the majority of intra-burst ISIs are concentrated around 10 ms and inter-burst ISIs are located on the interval from 345 ms to 415 ms.  $D_2 = 0.05$  is chosen because it is the optimal stimulus intensity of CR as shown later in Fig. 4A. This temporal feature of spike trains can also be illustrated by the PSD as shown in Fig. 3B. The PSD peak at the low frequency (around 2.5 Hz) corresponds to the longer inter-burst ISIs, and the PSD peak at the high-frequency (around 100 Hz) results from the shorter intra-burst ISIs, whereas the PSDs of PSTH in the inset of Fig. 2D do not have a peak at a frequency range higher than 15 Hz.

To get a more precise insight on the coherence of all spikes (from both low- and high-frequency ranges) in one network, we use an inputoutput SNR measure [38]: the ratio between the power of spike trains (output) and the power of the global stochastic stimulus (input). This SNR measure is commonly used to select the best recording location through spike sorting, and also to assess the reliability of neural information transmission [39]. To distinguish from the first SNR measure in Equ. (5), we denote the second SNR measure as  $\beta$ . That is,

$$\beta = \frac{P_{output}}{P_{input}} = \frac{1}{N} \sum_{i=1}^{N} \frac{P_{ST,i}}{P_N},\tag{6}$$

where  $P_{ST,i}$  is the power of the spike train of the *i*th neuron, and  $P_N$  is the power of global stochastic stimulus,  $D_2\eta$ .  $D_2\eta$  is identical for all neurons in the network, and thus (6) can be equivalently written as

$$\beta = \frac{\frac{1}{N} \sum_{i=1}^{N} P_{ST,i}}{P_N}.$$

We then introduce the re-scaled PSD which helps us to understand how  $P_{ST,i}$  and  $P_N$  affect SNR- $\beta$ . The re-scaled PSD of spike trains is defined by the averaged PSD over N spike trains in a network divided by  $P_N$ , i.e.

$$\tilde{S}(f) = \frac{\frac{1}{N} \sum_{i=1}^{N} S_i(f)}{P_N},$$
(7)

where *f* is the frequency,  $S_i(f)$  is the PSD of the spike train generated by the *i*th neuron, and  $\tilde{S}(f)$  is the re-scaled PSD.  $P_N$  is proportional to  $D_2$  because  $\eta$  takes the form of white noise, which has a constant PSD. The re-scaled PSDs for three  $D_2$  values (0.03, 0.05, and 0.08) are demonstrated in the inset of Fig. 4A. Similar to the PSD in Fig. 3B, the re-scaled PSDs have peaks at both low- and high-frequency ranges. In particular, the re-scaled PSD corresponding to  $D_2 = 0.05$  (red curve in Fig. 4A inset) is higher than the other two PSD curves; consequently, SNR- $\beta$  is expected to be larger at  $D_2 = 0.05$ .

The SNR- $\beta$  as a function of  $D_2$  is illustrated in Fig. 4A for  $g_{Ca} = 0.64$  and it also shows a characteristic pattern of CR, with a maximum at



**Fig. 3.** (A) The inter-spike interval histogram (ISIH) when  $g_{Ca} = 0.64$  and  $D_2 = 0.05$ . In order to clearly show both intraburst spikes and burst events, different horizontal- and vertical- scales are used on the left and right parts of (A). (B) PSD v.s. frequency when  $g_{Ca} = 0.64$  and  $D_2 = 0.05$ . Note, a logarithmic scale on the horizontal axis is used to show both intraburst spike and burst events clearly. Panel (B) is the representation of panel (A) in the frequency domain.



**Fig. 4.** (A) SNR- $\beta$  calculated by the power spectra of spike trains v.s.  $D_2$  with  $g_{Ca} = 0.64$ . The SNR curve reaches a peak (red solid circle) at  $D_2 = 0.05$ . Inset of (A): The re-scaled PSDs of spike trains in the frequency domain. The re-scaled PSD is defined in Equ. (7). (B)-(D) show three example spike raster plots for the three  $D_2$  values as labeled in (A), and other parameter values are as same as in Fig. 2.



**Fig. 5.** (A) Average burst rate produced by each neuron in a network v.s.  $D_2$  for  $g_{Ca} = 0.64$ . Inset of (A): an example binary burst train, where the burst onset is defined by the firing time of the first spike of a burst event in the corresponding voltage trace. (B) SNR- $\beta$  curve calculated by the power spectra of the burst trains for  $g_{Ca} = 0.64$ , and other parameter values are the same as in Fig. 2. The largest SNR is marked by a solid dot and corresponds to  $D_2 = 0.05$ .

 $D_2 = 0.05$ . Comparing these two coherence measures (SNR- $\alpha$  in Fig. 2D and SNR- $\beta$  in Fig. 4A), two major differences are observed: (a) the noise intensity range is [0.03, 1] for SNR- $\alpha$ , but CR measured by SNR- $\beta$  occurs over the weak intensity range of [0.001, 0.3]; and (b) the optimal intensity is 0.225 in Fig. 2D but it is 0.05 in Fig. 4A. The differences above are due to the contrasting focuses of the two SNR functions, SNR- $\alpha$  in Equ. (5) and SNR- $\beta$  in Equ. (6). Although they both evaluate CR, SNR- $\alpha$  characterizes the similarity of the frequency content of neuronal oscillations (i.e. bursts across the network here), while SNR- $\beta$  focus on the reliability of the neuronal responses (i.e. all spikes across network here) to the input (stochastic stimulus) over time.

The example raster plots presented in Fig. 4B-4D demonstrate the spatio-temporal patterns of firing for this neural network with respect to the increment of  $D_2$  from 0.03 to 0.08. When subjected to a very weak level of stimulus (e.g.  $D_2 = 0.03$  in Fig. 4B) bursting is induced at a very low rate (around 0.6 bursts/second) with very low synchronization. When the stimulus is increased slightly (e.g.  $D_2 = 0.05$  in Fig. 4C), both population bursts and individual spikes become more frequent. The network becomes synchronized, spiking events are tightly grouped in bursts, and the SNR- $\beta$  reaches its peak as predicted by the analysis of re-scaled PSDs. This also correlates with the sharp peaks in the histogram of the ISIs (Fig. 3A). When the network is subjected to higher—but still relatively weak—levels of stimulus (e.g.  $D_2 = 0.08$  in Fig. 4D) the stochastic stimulus becomes overpowering, and some network bursts start to lose synchrony. The observations above are in line with the observations for SNR- $\alpha$  in Fig. 2, despite the difference in  $D_2$ values.

#### 3.4. The network coherence in terms of individual bursts

The PSTH estimates the timing of population bursts across a network. In order to accurately record the onset (i.e. the occurrence times) of the bursts produced by individual neurons, we mark a burst event by its initial spike time. Thus, a burst train can be formed for each neuron. A burst train is a binary sequence, which takes the value 1 at the initial spike of all bursting events and 0 otherwise, as demonstrated in the inset of Fig. 5A. The occurrence of bursts is identified using the dynamic burst threshold method in [40].

As shown in Fig. 5A, the average burst rate produced by each neuron in a network increases with respect to  $D_2$ , which is in line with the voltage time series of a single neuron (Fig. 1B). To quantify the coherence of the burst trains, we use the input-output SNR measure similar to Equ. (6), that is,

$$\beta = \frac{1}{N} \sum_{i=1}^{N} \frac{P_{BT,i}}{P_N},$$
(8)

where  $P_{BT,i}$  denotes the power of the burst train generated by the *i*th neuron. Fig. 5B illustrates the change of SNR on different  $D_2$  for

 $g_{Ca} = 0.64$  (i.e. the same parameter values as Fig. 4A). The SNR- $\beta$  curve in Fig. 5B shows the characteristics of CR: it rapidly increases for very weak levels of the global stochastic stimulus ( $D_2 < 0.05$  in Fig. 5B), reaches a peak value at  $D_2 = 0.05$  in Fig. 5B, and then decreases for larger stimulus intensities ( $D_2 > 0.05$  in Fig. 5B). The SNR- $\beta$  curves calculated from burst trains and spikes trains share the same optimal noise intensity, but it is different from the optimal intensity presented in SNR- $\alpha$  calculated from PSTH (Fig. 2D) because a higher  $D_2$  leads to a lower re-scaled PSD curve (Fig. 4A inset) and consequently SNR- $\beta$  is continuously decreasing for  $D_2 > 0.05$ .

#### 3.5. The effect of calcium conductance on CR

As demonstrated in subsection 3.1, neural dynamics change with the calcium conductance,  $g_{Ca}$ . In the preceding subsections we have considered  $g_{Ca} = 0.64$ . To study the effects of  $g_{Ca}$  on the network coherence, we take various  $g_{Ca}$  values in the excitable regime (i.e.  $g_{Ca} < 0.648$ ) and use SNR- $\beta$  measure calculated from both spike trains (Equ. (6)) and burst trains (Equ. (8)). A series of SNR- $\beta$  optimization curves, corresponding to four parameter values of  $g_{Ca}$  (0.638, 0.64, 0.642, and 0.645), are plotted for both the spike trains (Fig. 6A) and burst trains (Fig. 6B). All of the SNR curves display the characteristic pattern of CR, that is, they increase, reach a peak value, and then decrease towards zero. The SNR curve corresponding to a larger  $g_{Ca}$  is above the SNR curve corresponding to a smaller  $g_{Ca}$  for  $D_2 \in [0.005, 0.15]$ , indicating that the coherence degree is enhanced by increasing the  $g_{Ca}$  value over the weak noise intensity range. Moreover, the maximum degree of coherence (i.e. the height of the SNR peaks) and the corresponding optimal intensities  $(D_2)$  vary across different  $g_{Ca}$  values.

To capture  $g_{Ca}$ -dependent change in the height of SNRs, we computed the maximum SNR values for seven  $g_{Ca}$  values in the excitable regime, ranging from 0.63 to 0.645, and plotted them in Fig. 7A for spike trains and Fig. 7B for burst trains. Four  $g_{Ca}$  values in Fig. 6A are part of these seven values. As would be expected, when  $g_{Ca}$  is increased (i.e., closer to the excitation threshold), the peak value of SNR increases nearly exponentially. This implies that a higher network coherence is expected at the optimal intensity as the excitable system approaches the excitation threshold.

We also computed the optimal intensities of the global stochastic stimulus,  $D_2$ , for these seven  $g_{Ca}$  values. We found that the SNRs calculated from the power spectra of spike trains and burst trains share the same optimal intensity for one  $g_{Ca}$ ; they are plotted as a function of  $g_{Ca}$  in Fig. 7C. With increased  $g_{Ca}$ , the optimal intensity of the SNR decreases. This suggests that when the excitable system is closer to the excitation threshold, a smaller noise intensity is enough to drive the network to its best possible coherence.

As CR degree is sensitive to a slight change of  $g_{Ca}$ , in order to further study the effects of  $g_{Ca}$  on network coherence, we add a third heterogeneous component to this network. We let  $g_{Ca}$  be a uniform random



**Fig. 6.** (A) SNR- $\beta$  calculated by the power spectra of the spike trains v.s. the intensity of stochastic stimulus,  $D_2$ . (B) SNR- $\beta$  calculated from the burst trains v.s.  $D_2$ . For both panels,  $g_{Ca} = 0.638$ , 0.64, 0.642 and 0.645, and they are in the excitable regime. In order to clearly show all SNR curves the y-axis has different scales on the top and bottom in both panels.



**Fig. 7.** (A) The maximum degree of SNR- $\beta$  calculated by the power spectra of the spike trains v.s.  $g_{Ca}$ . (B) The maximum degree of SNR- $\beta$  from the burst trains v.s.  $g_{Ca}$ . (C) The optimal intensities of global stochastic stimulus ( $D_2$ ) v.s.  $g_{Ca}$ . For all three panels,  $g_{Ca} = 0.63$ , 0.632, 0.635, 0.638, 0.64, 0.642, and 0.645; all of which are in the excitable regime.

variable ranging from 0.63 to 0.645 (i.e.  $g_{Ca} \sim U(0.63, 0.645)),$  so that each neuron in our network may have a different  $g_{C_{\alpha}}$  value. SNR- $\beta$  measure is used to evaluate the network coherence, and we calculate the SNR- $\beta$  using the power spectra of both spike and burst trains (see Eqs. (6) and (8)). In Fig. 8, one sees that the SNR-for both spike and burst trains—sharply increases from  $D_2 = 0.005$  to  $D_2 = 0.03$ , reaching a peak at approximately  $D_2 = 0.03$  and then for larger  $D_2$  tends towards 0. In other words, the network displays a resonant behavior and the optimal stimulus intensity (e.g.  $D_2 = 0.03$  here) would induce the best coherence (with maximal SNR  $\approx 0.028$  for spike trains and 0.008 for burst trains). Compared to the results in Figs. 6–7, this  $g_{Ca}$ -varied network has similar peak SNR values and optimal intensities to the  $g_{Ca}$ -fixed network with  $g_{Ca} = 0.0642$  where SNR  $\approx 0.0253$  for spike trains and 0.0077 for burst trains, and  $D_2 = 0.04$ , although the mean  $g_{Ca}$  value here is 0.6375. This is because the CR degree increases nearly exponentially with increased  $g_{Ca}$  as shown in Fig. 7A-7B.

#### 4. Discussion

Many numerical studies on the stochastic dynamics of neural networks employ homogeneous networks [20, 22, 41], globally connected networks [17, 18, 19, 21, 22]. We consider a heterogeneous network mediated by calcium channels; also, it has random and sparse synaptic connections and produces bursts when stimulated by external input. This study provides different views to evaluate the coherence of such a network. In terms of the regularity of the network output, the coherence information extracted from the population bursts indicates that the resonant coherence occurs at the intermediate noise intensity, which agrees with previous studies of globally connected bursting networks (e.g., [19, 20]). This coherence information is helpful for understanding the influences of global noise on collective dynamics (e.g., population bursts) and the potential physiological functions of a neural network. If network performance is evaluated based on the efficiency of the network response to the stochastic input, a weak stochastic input can stimulate both the coherence of all spikes across the network and the correlation of neuronal bursts to reach the optimal level. This implies that the improvement of neuronal communication within a network can be achieved using weak noise.

The impact of the noise intensity, coupling strength, or the network topology on CR is often discussed (e.g. [28, 42]). Our study broadens the understanding of network coherence from a new viewpoint, i.e., how the interplay between calcium conductance and noise intensity affects CR. Calcium current has been found to regulate the excitation and resonance of individual neurons [43]. Our work demonstrates that the calcium current can also enhance network response to the stochastic stimuli: when the calcium conductance is closer to its excitation thresh-

0.03

0.02

0.01

SNR-B

# References

- [1] J.E. Lisman, Bursts as a unit of neural information: making unreliable synapses reliable, Trends Neurosci. 20 (1) (1997) 38-43
- [2] E.M. Izhikevich, N.S. Desai, E.C. Walcott, F.C. Hoppensteadt, Bursts as a unit of neural information: selective communication via resonance. Trends Neurosci, 26 (3) (2003) 161–167
- [3] F. Zeldenrust, W.J. Wadman, B. Englitz, Neural coding with bursts-current state and future perspectives, Front. Comput. Neurosci. 12 (2018) 48.
- [4] E. Williams, A. Payeur, A. Gidon, R. Naud, Neural burst codes disguised as rate codes, Sci. Rep. 11 (2021) 15910.
- T. Fardet, M. Ballandras, S. Bottani, S. Métens, P. Monceau, Understanding the generation of network bursts by adaptive oscillatory neurons, Front. Neurosci. 12 (2018) 41
- [6] S.M. Cain, T.P. Snutch, T-type calcium channels in burst-firing, network synchrony, and epilepsy, Biochim. Biophys. Acta, Biomembr. 1828 (7) (2013) 1572-1578.
- S.M. Joksimovic, P. Eggan, Y. Izumi, S.L. Joksimovic, V. Tesic, R.M. Dietz, J.E. Orfila, M.R. DiGruccio, P.S. Herson, V. Jevtovic-Todorovic, et al., The role of t-type calcium channels in the subiculum: to burst or not to burst?, J. Physiol. 595 (19) (2017) 6327-6348.
- [8] H. Takano, M. McCartney, P.I. Ortinski, C. Yue, M.E. Putt, D.A. Coulter, Deterministic and stochastic neuronal contributions to distinct synchronous CA3 network bursts, J. Neurosci. 32 (14) (2012) 4743-4754.
- [9] M.W. Reimann, A.-L. Horlemann, S. Ramaswamy, E.B. Muller, H. Markram, Morphological diversity strongly constrains synaptic connectivity and plasticity, Cereb. Cortex 27 (9) (2017) 4570-4585
- [10] B. Hellwig, A quantitative analysis of the local connectivity between pyramidal neurons in layers 2/3 of the rat visual cortex, Biol. Cybern. 82 (2) (2000) 111-121.
- [11] S.M. Bowyer, Coherence a measure of the brain networks: past and present, Neuropsychiatr. Electrophysiol. 2 (1) (2016) 1-12.
- [12] A.A. Faisal, L.P. Selen, D.M. Wolpert, Noise in the nervous system, Nat. Rev. Neurosci. 9 (4) (2008) 292–303
- [13] J.F. Lindner, B.K. Meadows, W.L. Ditto, M.E. Inchiosa, A.R. Bulsara, Array enhanced stochastic resonance and spatiotemporal synchronization, Phys. Rev. Lett. 75 (1995) 3-6.
- [14] J. Collins, C.C. Chow, T.T. Imhoff, Stochastic resonance without tuning, Nature 376 (6537) (1995) 236–238.
- [15] H. Gang, T. Ditzinger, C.-Z. Ning, H. Haken, Stochastic resonance without external periodic force, Phys. Rev. Lett. 71 (6) (1993) 807.
- [16] A.S. Pikovsky, J. Kurths, Coherence resonance in a noise-driven excitable system. Phys. Rev. Lett. 78 (5) (1997) 775.
- [17] A.V. Andreev, V.V. Makarov, A.E. Runnova, A.N. Pisarchik, A.E. Hramov, Coherence resonance in stimulated neuronal network, Chaos Solitons Fractals 106 (2018) 80-85
- [18] E. Yilmaz, M. Ozer, V. Baysal, M. Perc, Autapse-induced multiple coherence resonance in single neurons and neuronal networks, Sci. Rep. 6 (1) (2016) 1-14.
- [19] J.H. Kim, H.J. Lee, C.H. Min, K.J. Lee, Coherence resonance in bursting neural networks, Phys. Rev. E 92 (4) (2015) 042701.
- [20] S. Reinker, Y.-X. Li, R. Kuske, Noise-induced coherence and network oscillations in a reduced bursting model, Bull. Math. Biol. 68 (6) (2006) 1401–1427.
- [21] W.C. Stacey, D.M. Durand, Noise and coupling affect signal detection and bursting in a simulated physiological neural network, J. Neurophysiol. (2002).
- Y. Wang, D.T. Chik, Z. Wang, Coherence resonance and noise-induced synchronization in globally coupled Hodgkin-Huxley neurons, Phys. Rev. E 61 (1) (2000) 740.
- [23] H. Yu, L. Zhang, X. Guo, J. Wang, Y. Cao, J. Liu, Effect of inhibitory firing pattern on coherence resonance in random neural networks, Phys. A, Stat. Mech. Appl. 490 (2018) 1201-1210.
- [24] J. Pham, K. Pakdaman, J.-F. Vibert, Noise-induced coherent oscillations in randomly connected neural networks, Phys. Rev. E 58 (3) (1998) 3610.
- [25] X. Sun, M. Perc, Q. Lu, J. Kurths, Spatial coherence resonance on diffusive and small-world networks of Hodgkin-Huxley neurons, Chaos, Interdiscip. J. Nonlinear Sci. 18 (2) (2008) 023102.
- [26] Y.H. Zheng, O.S. Lu, Spatiotemporal patterns and chaotic burst synchronization in a small-world neuronal network, Phys. A, Stat. Mech. Appl. 387 (14) (2008) 3719-3728
- [27] C. Zheng, A. Pikovsky, Stochastic bursting in unidirectionally delay-coupled noisy excitable systems, Chaos, Interdiscip. J. Nonlinear Sci. 29 (4) (2019) 041103.
- [28] M. Masoliver, N. Malik, E. Schöll, A. Zakharova, Coherence resonance in a network of Fitzhugh-Nagumo systems: interplay of noise, time-delay, and topology, Chaos, Interdiscip. J. Nonlinear Sci. 27 (10) (2017) 101102.
- [29] M. Masoliver, C. Masoller, A. Zakharova, Control of coherence resonance in multiplex neural networks. Chaos Solitons Fractals 145 (2021) 110666.
- [30] M.E. Yamakou, J. Jost, Control of coherence resonance by self-induced stochastic resonance in a multiplex neural network, Phys. Rev. E 100 (2019) 022313.
- [31] N. Semenova, A. Zakharova, Weak multiplexing induces coherence resonance, Chaos, Interdiscip. J. Nonlinear Sci. 28 (5) (2018) 051104.

0 0.05 0 0.1 0.15 D<sub>2</sub> (μA) Fig. 8. SNR- $\beta$  calculated by the power spectra of spike trains (solid line) and

 $\cap$ 

burst trains (dashed line) of a  $g_{Ca}$ -varied network of 100 neurons. A set of 100 uniformly distributed random numbers ranging from 0.63 to 0.645 was used to model different  $g_{Ca}$  values in this network. The largest SNR value for both the spike- and burst- trains occurs at  $D_2 = 0.03$ .

old, a smaller intensity of the stochastic stimulus is needed to induce the best coherence where a higher CR degree is achieved. This gives us hope for the experimental discovery of this noise-induced resonance effect in analyses of calcium-related brain response near the excitation threshold.

A rather small network connectivity probability is used here based on the experimental study [10]. A strong coupling usually promotes synchrony. This implies that if connection probability increases, the optimal  $D_2$  curve in Fig. 7C would shift downward and the optimal  $D_2$ in Fig. 8 would be smaller than 0.03; the maximum degree of SNR- $\beta$ in Fig. 7AB could increase. We assume that connection probability is constant for the sake of simplicity. However, the probability of connection can be a function of neuron location because synaptic connectivity between neurons is anti-correlated with their lateral distance [9, 10]. Therefore, the general effects of the connection probability function are an interesting topic requiring further investigation.

#### **Declarations**

#### Author contribution statement

Na Yu: Conceived and designed the experiments; Performed the experiments; Analyzed and interpreted the data; Contributed reagents, materials, analysis tools or data; Wrote the paper. Gurpreet Jagdev: Performed the experiments; Analyzed and interpreted the data; Contributed reagents, materials, analysis tools or data; Wrote the paper. Michelle Morgovsky: Performed the experiments; Analyzed and interpreted the data; Wrote the paper.

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#### Data availability statement

No data was used for the research described in the article.

### Declaration of interests statement

The authors declare no conflict of interest.

- [32] R. Tönjes, C.E. Fiore, T. Pereira, Coherence resonance in influencer networks, Nat. Commun. 12 (1) (2021) 1–8.
- [33] H. Shouval, S. Wang, G. Wittenberg, Spike timing dependent plasticity: a consequence of more fundamental learning rules, Front. Comput. Neurosci. 4 (2010) 19.
- [34] G.-q. Bi, M.-m. Poo, Synaptic modification by correlated activity: Hebb's postulate revisited, Annu. Rev. Neurosci. 24 (1) (2001) 139–166.
- [35] E.M. Izhikevich, Neural excitability, spiking and bursting, Int. J. Bifurc. Chaos 10 (06) (2000) 1171–1266.
- [36] H.M.G. Fuhrmann, I. Segev, M. Tsodyks, Coding of temporal information by activitydependent synapses, J. Neurophysiol. 87 (1) (2002) 140–148.
- [37] M. Stimberg, D.F. Goodman, R. Brette, M. De Pittà, Modeling neuron–glia interactions with the Brian 2 simulator, in: Computational Glioscience, Springer, 2019, pp. 471–505.
- [38] F. Rieke, D. Warland, Spikes: Exploring the Neural Code, MIT Press, Cambridge, Mass., 1999.

- [39] S.R. Schultz, Signal-to-noise ratio in neuroscience, Scholarpedia 2 (6) (2007) 2046.
- [40] J.V. Selinger, N.V. Kulagina, T.J. O'Shaughnessy, W. Ma, J.J. Pancrazio, Methods for characterizing interspike intervals and identifying bursts in neuronal activity, J. Neurosci. Methods 162 (1–2) (2007) 64–71.
- [41] H. Soula, C.C. Chow, Stochastic dynamics of a finite-size spiking neural network, Neural Comput. 19 (12) (2007) 3262–3292.
- [42] S.-Y. Kim, W. Lim, Noise-induced burst and spike synchronizations in an inhibitory small-world network of subthreshold bursting neurons, Cogn. Neurodyn. 9 (2) (2015) 179–200.
- [43] Y.Y.B. Hutcheon, R.M. Miura, E. Puil, Low-threshold calcium current and resonance in thalamic neurons: a model of frequency preference, J. Neurophysiol. 71 (2) (1994) 583–594.