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Reactive co-sputter deposition of Ta-doped tungsten oxide thin films for water splitting application

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This study aimed to investigate the structural, optical, and electronic properties of WO₂ thin films modified by Ta-doping, considering their potential application in photoelectrochemical (PEC) water splitting. Due to its unique physical and chemical properties, WO₃ films have been commonly suggested as a promising photoanode for hydrogen production. However, the wide bandgap and unsuitable band edge positions of WO₂ limit its PEC efficiency. Doping have been extensively applied as an effective strategy for bandgap engineering. Here, post-annealed WO, films with different concentrations of Ta dopant were synthesized via reactive magnetron co-sputtering, while DC and RF sputtering powers were varied with the aim of achieving the desired properties. EDX analysis showed that Ta atoms were doped into WO_3 in the range of 0-3.93 at%. As evident from SEM and AFM images, the surface morphology was significantly affected by increasing Ta doping, the formation of a granular structure with well-defined boundaries and increasing surface roughness (1.79-47.94 nm). XRD patterns confirmed that the incorporation of Ta atoms into a monoclinic WO₂ improved the crystallinity, especially in the (002) direction. Most importantly, a decrease in the average transparency (92.82-74.27%), an increase in visible absorption, a red shift of the fundamental absorption edge corresponding to a favorable drop in the optical bandgap energy (3.07-2.61 eV) were found with increasing Ta concentration. Notably, the substitution of W6+ ions with Ta dopant (0-3.93 at%) led to an upward shift in the valence band maximum (3.62-3.31 eV) and a downward shift in the conduction band minimum (0.55-0.70 eV). The WO₃ photoanode doped with 3.93 at% Ta exhibited the maximum photocurrent density of 0.65 mA/cm² (at 1 V vs. Ag/AgCl) under simulated sunlight. Furthermore, WO₂ photoanode doped with 3.93 at% Ta showed excellent photoresponsivity and slow electron-hole recombination. The obtained results predict the potential of Ta-doping coupled with post-annealing to optimize the structural and optoelectronic properties of sputtered WO₃ thin films as photoanode for use in efficient PEC water splitting.

Keywords WO₂, Ta-doping, Photoelectrochemical water splitting, Bandgap, Co-sputtering, Photocurrent

Nowadays, rising global energy consumption and the depletion of fossil fuels have become great challenges for humanity1. In addition, the combustion of fossil fuels leads to important environmental damages such as water and air pollution, global warming due to increased CO, concentration in the atmosphere, and the greenhouse effect^{2,3}. Therefore, it is necessary to find renewable, sustainable, and clean energy systems to meet future energy needs along with environmental protection^{4,5}. Recently, hydrogen (H₂) has attracted considerable attention as a ultimate clean and renewable energy source with the potential to be used in fuel cells and chemical industries without emitting toxic gases (CO_x, NO_x, SO_x, C_xH_y)⁵⁻⁷. Extraction of hydrogen from water and sunlight through photoelectrochemical (PEC) water splitting has been recognized as an efficient solution for solar energy utilization^{8,9}. The interesting discovery of the PEC water splitting was first reported in 1972 by Fujishima and Honda using a titanium dioxide photoanode¹⁰. Recently, the identification, development, and exploitation of ideal photoanode materials with a suitable designs to achieve better water splitting activity have been investigated $^{11-\hat{15}}$. To realize the most efficient PEC water splitting, semiconductors must possess essential properties, including adequate visible light absorption, effective charge separation and transportation, enhanced photocarrier collection, efficient water reduction and oxidation catalysis, and well stability in acid¹⁶. Numerous amount of research has been devoted to the use of metal oxide semiconductors such as TiO_2^{17} , ZnO^{18} , NiO^{19} , α -Fe₂O₃²⁰, Cu_2O^{21} , Cu/CuO^{22} , CuO-La₂O₃²³, CeO_2^{24} , SnO_2^{25} , $BiVO_4^{26}$, $NaTaO_3^{27}$, $KTaO_3^{28}$, $SrTiO_3^{29}$, and $CuFeO_2^{30}$, LaFeO₃³¹, and $CuBi_2O_4^{32}$ in PEC water splitting. Among these metal oxides, n-type semiconductor tungsten oxide (WO₃) has attracted considerable attention due to its potential applications in optoelectronic

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devices, gas sensors, optical detectors, electrochromic devices, solar cells, supercapacitors, high-temperature diodes, photocatalysts, MIS-based devices, hydrophilic coatings, schottky barrier diodes, etc.33-37 Furthermore, tungsten oxide is well-known as the promising electrode material for efficient solar-to-H2 conversion because of several advantages including, availability, non-toxicity, photosensivity, tunable bandgap, fast carrier transport, medium hole diffusion length, excellent photocorrosion resistance in aqueous solutions, high chemical stability, and inexpensive³⁸⁻⁴¹. The first successful realization of solar water splitting on tungsten oxide photoanode was reported by Hodes el al. 42. After that, many researchers have focused on the synthesis of tungsten oxide with favorable optical, electrical, and mechanical properties. Tungsten oxide is more preferred than ZnO and TiO, for PEC water splitting⁴³. According to reports, WO₃ can absorb almost 12% of the solar spectrum, which provides the maximum efficiency (~4.8%) of solar-to-hydrogen conversion in PEC systems⁴⁴. Despite these features, the implementation of highly efficient water splitting reaction by tungsten oxide faces major limitations. First, the bandgap of tungsten oxide is not narrow enough to absorb a significant part of sunlight⁴⁵. On the other hand, the position of the conduction band edge is not higher according to the redox potential of H⁺/H₂, so the excited electrons in the conduction band are not able to convert protons into hydrogen⁴⁶. Consequently, band edge engineering is required to narrow the bandgap and alter the band levels of WO₃⁴⁷. Numerous methods have been proposed for band edge engineering, such as metal doping, morphological control, composite material production, and surface sensitization⁴⁸. Among these, doping with metal ions is suggested as the most effective approach to tune the optical and electrical properties of metal oxides 48 . Recently, doping of WO₃ with metal ions such as Mg⁴⁹, Ti 50,51 , Zr⁵², Ta⁵³, Mo⁵⁴, Fe⁵⁵, Ni⁵⁶, Cu⁵⁷, Al⁵⁸, Ga⁵⁹, Sn⁶⁰, Bi⁶¹, Te⁶², and Gd⁶³ has been extensively investigated. Since the radius of Ta ions is close to that of tungsten, they are properly incorporated into the tungsten lattice. As a result, it assumed that the doping of Ta metal ions into tungsten oxide can easily reduce the optical band gap and improve the photoconversion efficiency⁶⁴. There are various techniques to produce metalion-doped WO₃ thin films including, hydrothermal⁶⁵, sol-gel⁶⁶, electrodeposition⁶⁷, spray pyrolysis⁶⁸, electron beam evaporation⁶⁹, thermal evaporation⁷⁰, and magnetron sputtering⁷¹. In the present work, the reactive magnetron co-sputtering method is applied for the deposition of Ta-doped tungsten oxide thin films with various Ta content. This method provides film deposition with uniform thickness, high purity, homogeneous surface, better adhesion to the substrate, large-area scale, and desired dopant concentration at room temperature⁷²⁻⁷ Different DC and RF sputtering powers have been applied to W and Ta metal targets, respectively. The variation of the sputtering powers enables us to adjust the composition (Ta/W ratio) of the samples to achieve suitable optoelectronic properties for applications in water splitting. Here, the annealing process is necessary to form crystalline WO, for better PEC activity. Based on earlier studies, annealing conditions and temperature were selected⁷⁶. The effect of Ta-doping and post-annealing on the structural, morphological, optical, and electronic characteristics of sputtered WO₃ thin films was analyzed in detail.

Methodology

Thin films of Ta₂O₅, WO₃, and Ta-doped WO₃ were synthesized through reactive magnetron co-sputtering technique under Ar/O₂ atmosphere at room temperature using a MECA-2000 system, as shown in Fig. 1. Tungsten (W) and Tantalum (Ta) metal targets (99.99% purity, 6 mm thickness, and 76.2 mm diameter) were attached to DC and RF supplied magnetrons, respectively. Glass and quartz substrates (1×1 cm²) were successively cleaned in an ultrasonic bath according to the standard protocol by acetone, ethanol, and deionized water for 15 min and then dried with N2 gas. To ensure the uniformity of film deposition, the substrates were placed in a holder rotating at 30 rpm. The substrates were fixed at a distance of 15 cm from the target. The synthesis chamber was evacuated using rotary and turbomolecular pumps to reach a pressure of 4.9×10^{-5} mbar. The targets were pre-sputtered to remove surface contaminations in a pure argon plasma for 10 min. The flow of Ar gas was kept at 25 sccm, simultaneously, O, was introduced the chamber as reactive gas with the flow of 6.25 sccm. The flow of gases was controlled by mass flow controllers (MKS-Instruments) with an accuracy of 0.01 sccm. The working pressure during the deposition was carefully tuned at 7×10^{-3} mbar by a stepper motor-controlled high-vac gate valve (VAT PM-5). The DC and RF sputtering powers were varied to control the concentration of W and Ta (at%) in the film composition. Sputtering conditions were mentioned in Table 1. Variable RF power (10-60 W) was applied to the tantalum target in order to incorporate different concentrations of Ta atoms into the tungsten oxide structure to achieve the desired properties. The thickness of the films was determined by vibrating quartz method (MAXTEK, TM-350) about 200 nm. Afterward, the asdeposited samples were annealed at 550 °C for 120 min in a furnace with a heating/cooling ramp (rest) of 8 °C/ min in the air ambient. Energy dispersive X-ray spectroscopy (EDX, Tescan-Mira III) was employed to detect the elemental composition of the prepared thin films. The surface morphology of the samples was observed by a field emission scanning electron microscope (FESEM, Tescan-Mira III). The surface topography and roughness of the films were investigated using an atomic force microscope (AFM, Nanosurf mobile S). X-ray diffraction (XRD, Siemens D500) was employed to investigate the crystalline structure of the prepared samples. UV-Vis spectra were collected by spectrophotometer (GENESYS-5) in the wavelength range of 200-1100 nm. From the recorded data, optical parameters such as fundamental absorption edge and band gap energy were obtained. The photoelectrochemical performance of the prepared samples were measured using a conventional three-electrode setup under 1.5 G illumination of a 100 W Xe lamp. For PEC studies, undoped and Ta-doped WO₃ thin films were depositied on indium tin oxide (ITO) (200 nm) coated quartz substrates, which acted as working electrode. The area exposed to light was set to 1 cm² for all samples. A platinum wire was applied as a counter electrode, and a saturated Ag/AgCl was used as a reference electrode. Photocurrent measurements were conducted using a linear sweep voltammograms (LSV) in the potential range of -0.3 to 1.8 V (vs. Ag/AgCl) with a scan rate of 20 mV/s at room temperature in a 0.2 M Na₂SO₄ electrolyte. The dependency of photocurrent on irradiationtime were evaluate using a two-electrode setup under on/off light illumination at 0 V bias. The measurements were conducted directly between the the pt wire and the samples.

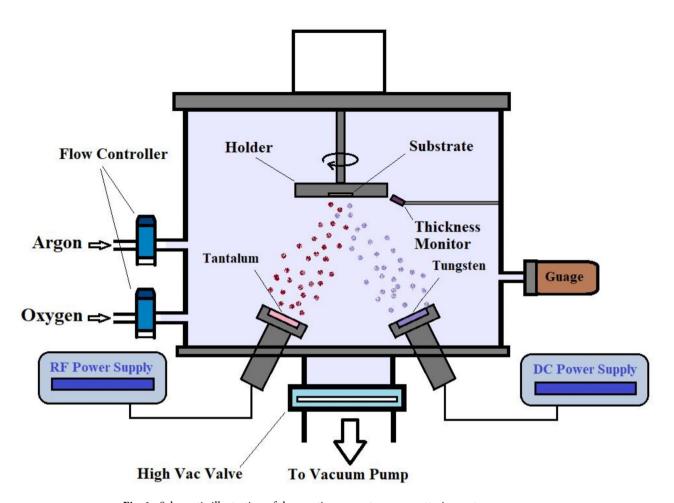


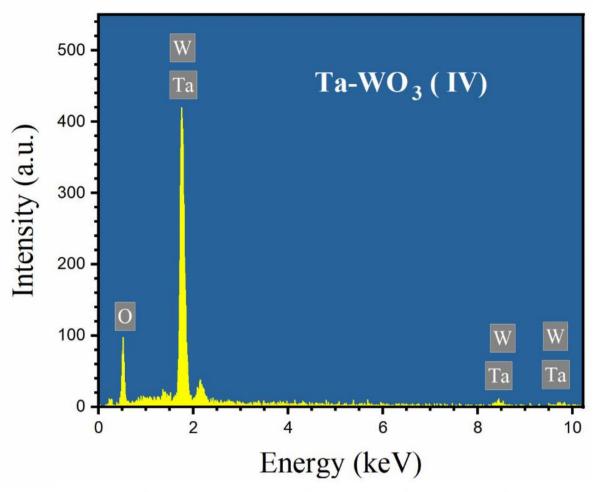
Fig. 1. Schematic illustration of the reactive magnetron co-sputtering system.

Thin film	DC power (W)	RF power (W)	Deposition rate (A/S)
Ta ₂ O ₅	0	100	0.1
WO ₃	100	0	1.1
Ta-WO ₃ (I)	90	10	1
Ta-WO ₃ (II)	70	30	0.9
Ta-WO ₃ (III)	60	40	0.7
Ta-WO ₃ (IV)	40	60	0.3

Table 1. Sputtering deposition conditions.

Results and discussion Morphological analysis

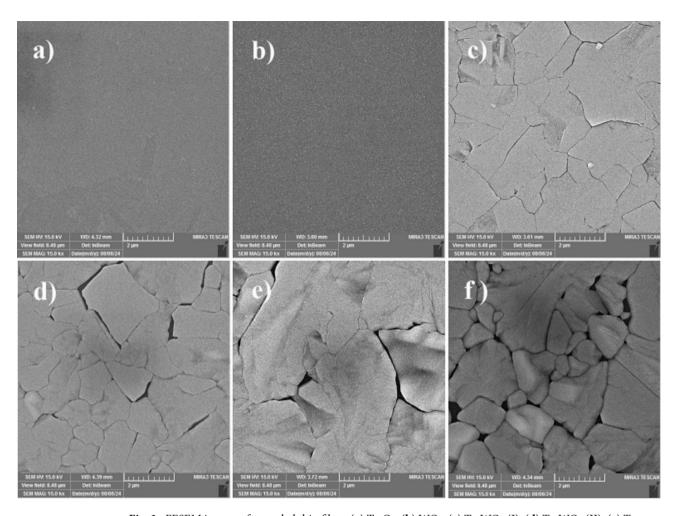
The compositional analysis of the prepared samples was performed using EDX spectroscopy. Figure 2 depicts the detailed EDX spectrum for the Ta-doped WO₃(IV) sample, indicating the existence of W, Ta, and O elements as reported⁷⁷. The elemental quantification of Ta₂O₅, WO₃, and Ta-doped WO₃ thin films annealed at 550 °C is presented in the inset table of Fig. 2. According to the obtained data, the incorporation of Ta in the doped WO₃ films is evident. The concentration of Ta in doped WO₃ films increases from 0.63 to 3.93% as atomic percentage with increasing RF sputtering power (10 to 60 W). Figure 3 shows surface FESEM micrographs of Ta₂O₅, WO₃, and Ta-doped WO₃ thin films annealed at 550 °C. The surface of undoped Ta₂O₅ and WO₃ films appeared to be smooth, dense, crack-free, and uniform, as shown in Fig. 3a,b. Morphological changes with cracking were significant for Ta-doped WO₃ films after annealing at 550 °C, Fig. c–f. This can be attributed to the improved crystallinity after Ta doping, as observed by XRD patterns. It was found that cracks in the surface morphology of Ta-doped thin films became more prominent with increasing Ta concentration. With further increase in Ta concentration to 3.93 at%, a granular structure with well-defined boundaries and close-packed grains was observed. It is concluded that the surface morphology of Ta-doped WO₃ thin films was significantly affected by the annealing process compared to pure Ta₂O₅ and WO₃. Further morphological characterization



Sample	W (At%)	Ta (At%)
Ta ₂ O ₅	0	29.78
WO ₃	13.56	0
Ta-WO ₃ (I)	14.33	0.63
Ta-WO ₃ (II)	10.73	2.06
Ta-WO ₃ (III)	9.35	3.27
Ta-WO ₃ (IV)	8.46	3.93

Fig. 2. EDX spectrum of Ta-doped WO $_3$ (IV) thin film annealed at 550 °C (top) and atomic percentage of identified elements from the EDX spectra (down).

of the prepared samples was performed by AFM analysis. Figure 4 shows two-dimensional (2D) and three-dimensional (3D) AFM images of ${\rm Ta_2O_5}$, WO $_3$, and Ta-doped WO $_3$ thin films annealed at 550 °C. The surface roughness of the prepared samples was measured in the term of root-mean-square ($\sigma_{\rm rms}$) roughness, as presented in Table 2. It was found that undoped ${\rm Ta_2O_5}$ and WO $_3$ thin films have uniform and smooth surfaces with root-mean-squared roughness ($\sigma_{\rm rms}$) values of 1.79 nm and 1.98 nm, respectively. Cracked surface morphology with high roughness ($\sigma_{\rm rms}$ ~ 14.71–47.94 nm) was found was observed for Ta-doped WO $_3$ thin films, as confirmed by SEM images. Figure 5 displays the increasing trend of the $\sigma_{\rm rms}$ value with increasing Ta concentration. This result suggests that Ta doping combined with annealing can form rough surfaces that are preferred for water splitting application 78. Memar et al. reported that increasing the surface roughness can contribute to increasing the electrolyte/electrode contact area, which leads to the enhancement of photocurrent density. On the other hand, a rough surface allows a large portion of sunlight to be captured due to multiple scattering 79. Similarly, new



reports suggested that the surface roughening of the photoanode films serves as an efficient method to optimize the photoelectrochemical response⁸⁰.

Structural analysis

The crystallographic properties of the synthesized thin films were investigated using XRD measurements. The XRD paterns of ${\rm Ta_2O_5}$, ${\rm WO_3}$, and ${\rm Ta\text{-}doped\ WO_3}$ thin films annealed at 550 °C are shown in Fig. 6. For the ${\rm Ta_2O_5}$ film, the absence of any diffraction peak confirmed its amorphous nature. The crystallization resistance was expected for the annealed ${\rm Ta_2O_5}$ films in the temperature range of 300–700 °C as reported in the literature ⁸¹. For undoped and ${\rm Ta\text{-}doped\ WO_3}$ films, the detected XRD diffractions were consistent with the monoclinic phase of ${\rm WO_3}$ (JCPDS No: 00-033-1387). The XRD peaks related to ${\rm Ta\text{-}W\text{-}O}$ or oxides of ${\rm Ta\ were}$ not found for ${\rm Ta\text{-}doped\ WO_3}$ films. The well-resolved peaks at 23.23°, 24.41°, 34.1°, 41.88°, and 56.07° correspond to the (002), (200), (112), (022), (222), and (124) planes, respectively ^{82,83}. The presence of more peaks in the XRD pattern of ${\rm Ta\text{-}doped\ WO_3}$ compared to undoped ${\rm WO_3}$ indicates the formation of better crystalline ${\rm WO_3}$ films after ${\rm Ta\ doping}$. This result suggests that the transition to a crystalline phase occurs more easily for combined oxide of ${\rm Wa\ nd}$ Ta than for pure ${\rm Ta_2O_5}$ and ${\rm WO_3}$. The strongest peak (200) of ${\rm Ta\text{-}doped\ WO_3}$ films was analyzed to determine crystallite size by Debye–Scherrer relation ⁴¹:

$$D = 0.9\lambda/\beta\cos\theta \tag{1}$$

where λ is the X-ray wavelength (1.5418 A°), β is FWHM of (200) peak, and θ is the diffraction angle. With increasing RF sputtering power (10–60 W), a increase in crystallite size of Ta-doped WO $_3$ was found from 73.81 to 110.5 nm, as shown in Table 2. According to previous reports, the crystal plane orientation of WO $_3$ films significantly affects the water splitting efficiency. Li et al. reported that WO $_3$ doped with different concentrations (0–20 wt%) of Mg after sintering at 950 °C showed different crystallization behavior, which affected their photoelectrochemical performance. For example, a new phase (MgWO $_4$) was found for 10 wt% Mg-doped WO $_3$, which was neither MgO nor WO $_3$ 49 . On the other hand, it has been reported that the incorporation of Ta or Ti ions into the WO $_3$ crystal lattice prepared by the hydrothermal method prevents the formation

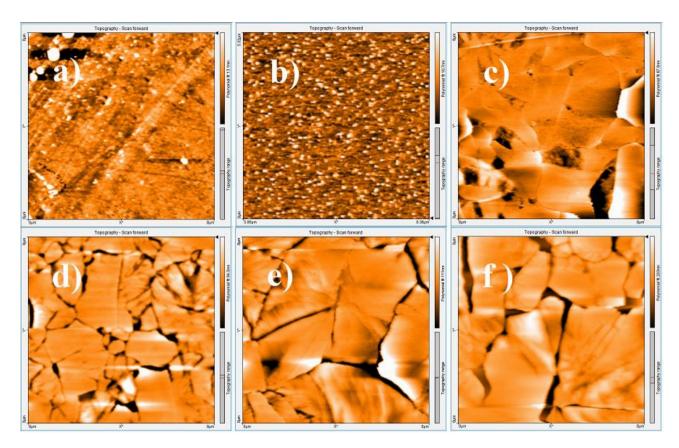


Fig. 4. AFM images of annealed thin films: (a) Ta_2O_5 , (b) WO_3 , (c) $Ta-WO_3$ (I), (d) $Ta-WO_3$ (II), (e) $Ta-WO_3$ (III), and (f) $Ta-WO_3$ (IV).

Thin film	σ _{rms} (nm)	Crystallite size (nm)	T _{av} (%)
Ta ₂ O ₅	1.79	_	92.82
WO ₃	1.98	_	90.01
Ta-WO ₃ (I)	14.71	73.81	88.06
Ta-WO ₃ (II)	16.04	74.86	84.80
Ta-WO ₃ (III)	30.42	55.40	74.00
Ta-WO ₃ (IV)	47.94	110.5	74.27

Table 2. Surface roughness, crystallite size, and average transmittance for sputter deposited thin films.

of a pure monoclinic phase after annealing at 500 °C, and multiple phases (hexagonal and monoclinic) are usually produced 64,84 . In the present study, the formation of a pure monoclinic phase, which is preferred for the water splitting reaction, was successful for Ta-doped WO₃ thin films deposited by magnetron co-sputtering after annealing at 500 °C. Recently, Wang et al. reported that WO₃ films grown with enriched the (002) plane by hydrothermal method enable more efficient water splitting reaction with remarkable photocurrent density 85 . In the present case, with increasing Ta concentration above 3.27 at%, a peak at 23.23° corresponding to (002) plane was noticed. In conclusion, Ta doping into WO₃ by promoting the growth of (002) direction can be helpful to enhance the water splitting performance. It is clear that increasing Ta consentration into WO₃ lattice leads to a gradual shift in the XRD peak positions of (222) and (124) planes towards higher 20 value. This is attributed to the substitution of some W⁶⁺ ions by Ta atoms in the WO₃ crystal structure. When the radius of the substituting material is different from the host material, the occurrence of compressive strain in the crystal leads to the reduction of lattice parameters 60,84 .

Optical properties

The bandgap nature and optical parameters of the synthesized thin films were examined by UV–Vis spectroscopic analysis in the wavelength range of 200–1100 nm. Figure 7 represents the transmittance spectra of ${\rm Ta}_2{\rm O}_5$, WO $_3$, and Ta-doped WO $_3$ thin films annealed at 550 °C. The existence of oscillatory modes in the spectra is due to the interference phenomenon caused by multiple reflections of light at the film-substrate interface ⁸⁶. The comparative analysis of the spectra exhibits that the optical behavior of Ta-doped WO $_3$ thin films strongly

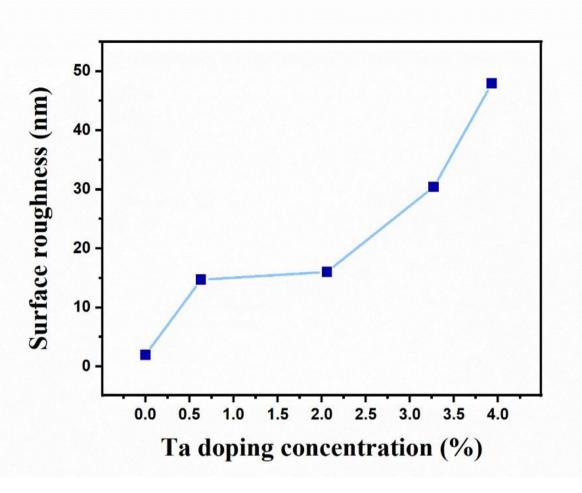


Fig. 5. variation of surface roughness as a function of Ta doping concentration.

depends on the Ta concentration. It is clear that the transparency of WO₃ films in the visible region decreases rapidly with increasing Ta concentration. The average transmittance (T_{av}) was calculated in the visible range ($\lambda = 380-780$ nm) using the following relation⁸⁷:

$$T_{av} = \frac{\int V(\lambda) T(\lambda) d\lambda}{\int V(\lambda) d\lambda}$$
 (2)

Here, $V(\lambda)$ is the photopic luminous efficiency, representing the sensitivity of the standard observer for photometry⁸⁸. The obtained values of T_{av} for the prepared samples are listed in Table 2. Pure Ta_2O_5 films exhibit the highest T_{av} value in the visible region. It was clear that the optical transparency of doped films decreases noticeably with increasing Ta concentration, as shown in Fig. 11. This can be explained by two reasons: Firstly, the intra-band excitation of charge carriers in the conduction band states leads to increased absorption as a result of reduced transmission^{89,90}. Meanwhile, the reduction in T_{av} with increasing Ta concentration can be associated with the enhanced light scattering caused by the improvement of crystallinity and the increase in surface roughness as evident from structural and morphological studies (XRD and AFM results)^{91,92}. A sharp drop in the transmittance value between 300 and 400 nm corresponds to the fundamental absorption, indicating the semiconducting nature of the prepared films⁹³. Figure 8 represents the absorbance spectra of the prepared thin films. A careful observation of the spectra exhibits the enhancement of absorption in the visible region with increasing Ta concentration into the WO₃ lattice. The fundamental absorption edge corresponds to the direct transition of an electron from the top of valence band to the bottom of conduction band, which can be employed to find the nature (direct or indirect) and value of the bandgap⁹⁴. The optical absorption coefficient (α) is a necessary quantity to find the band gap energy of thin films, which can be determined using transmittance (T) and film thickness (d) as follows^{75,95}:

$$\alpha = \frac{1}{d} \ln \left(\frac{1}{T} \right) \tag{3}$$

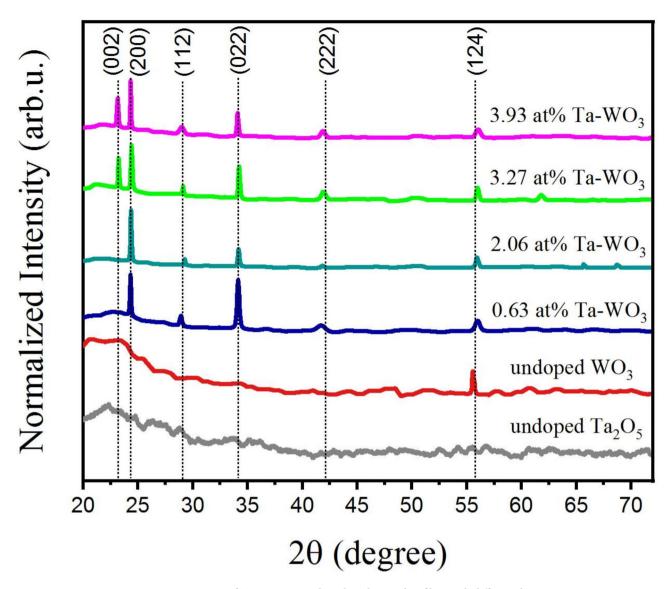


Fig. 6. XRD pattern of Ta₂O₅, WO₄, and Ta-doped WO₄ thin films with different dopant concentrations.

Figure 9 shows a sharp drop in α values in the UV spectral region and a relatively stabilized behavior in the visible light range. It is also found that the absorption coefficient of WO₃ films gradually increases with increasing Ta concentration, which results from enhanced absorption. A significant shift of the absorption edge towards longer wavelengths after Ta-doping leads to a reduction in the bandgap energy. The optical bandgap is an effective property to evaluate the photosensitivity of materials. The optical bandgap energy (E_g) of thin films can be extracted through the well-known Tauc relation ⁹⁶:

$$\alpha h v = B (h v - Eg)^m \tag{4}$$

Here, B is the band edge constant that depends on the crystalline quality, h is the Planck's constant, ν is the frequency of incident light, α is the absorption coefficient, and m is the power coefficient. The power coefficient is related to the nature of transition, which takes the values of 1/2, 2, 3/2, or 3 for direct allowed, indirect allowed, direct forbidden, and indirect forbidden electron transitions, respectively. Figure 10 depicts the Tauc plots for Ta_2O_5 , WO_3 , and Ta-doped WO_3 thin films annealed at 550 °C, considering the indirect allowed transition (n = 2). Extrapolation of the linear part of $(\alpha h\nu)^{1/2}$ until crossing the energy axis was used to estimate bandgap values⁹⁷. The extracted E_g values for the Ta_2O_5 , $Table V_3$, and Ta-doped $Table V_3$ thin films to be 4.00, 3.07, and 2.61–3.01 eV, respectively. Table 3 shows that the optical bandgaps of the prepared samples are very close to the results reported for these materials synthesized by other techniques. The decreasing trend of the optical bandgap of $Table V_3$ films from 3.07 to 2.61 eV with increasing $Talle V_3$ concentration (0–3.93%) is depicted in Fig. 11. The narrowing of the bandgap and the enhancement of visible light absorption after $Talle V_3$ and $Talle V_3$ and $Talle V_3$ respectively. Table 3 shows that the valence band $Talle V_3$ and $Talle V_3$ and $Talle V_3$ respectively. Table 3 shows that the optical bandgaps of the prepared samples are very close to the results reported for these materials synthesized by other techniques. The decreasing trend of the optical bandgap of $Talle V_3$ films from 3.07 to 2.61 eV with increasing $Talle V_3$ and $Talle V_3$ is depicted in Fig. 11. The narrowing of the bandgap and the enhancement of visible light absorption after $Talle V_3$ and conduction band (CB) of $Talle V_3$ and $Talle V_3$ reported for the section of localized impurity levels between the valence band (VB) and conduction band (CB) of $Talle V_3$ reported for the section $Talle V_3$ reporte

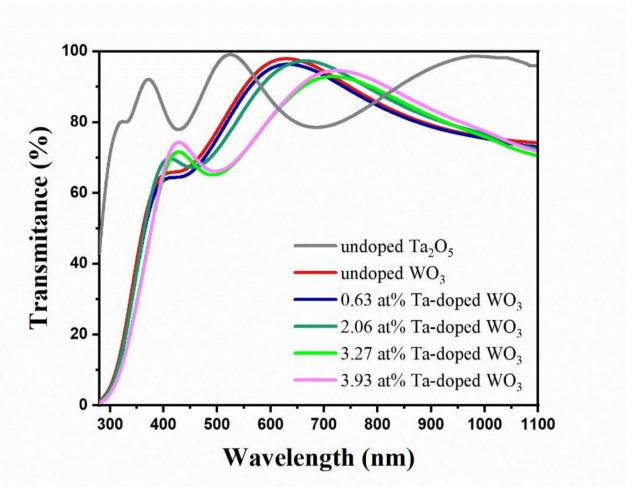


Fig. 7. Transmittance spectra of ${\rm Ta_2O_5}$, ${\rm WO_3}$, and ${\rm Ta\text{-}doped\ WO_3}$ thin films with different dopant concentrations.

impurity level to the CB^{91} . Several recent reports confirm the reduction of the bandgap energy of tungsten oxide after doping with metal ions 45,47,60,64,98. The bandgap value of 2.61 eV ensures better absorption of the incident light and enables the appropriate use of Ta-doped WO_3 as a photoanode in PEC water splitting.

Photoelectrochemical characterization

To evaluate the photoelectrochemical performance of the prepared electrodes, photocurrent-voltage measurements were conducted in a three-electrode setup by linear sweep voltammetry (LSV) in 0.2 M Na₂SO₄ electrolyte under one sun illumination. Figure 12 displays the LSV plots for undoped and 3.93 at% Ta-doped WO₃ thin films in the potential range of -0.4-1.6 V versus Ag/AgCl. It is clear that both undoped and Ta-doped electrodes exhibit negligible photocurrent responses under dark conditions. For undoped WO, photoanode, the photocurrent of ~0.19 mA/cm² was found at 1 V (vs. Ag/AgCl) under 1-sun irradiation. The maximum photocurrent obtained for undoped WO₂ in the present study is relatively less compared to the values reported in the literature^{60,84}. This is mainly associated with the material production process and the orientation of the crystal planes. According to previous studies, the water oxidation reaction occurs on the (002) plane more efficiently than on other planes of the monoclinic phase. Because the energy change required to realize the water oxidation reaction in this plane is 1.62 eV, which is much less than the energy of 1.49 eV required for the reaction on the (200) plane⁸⁵. Here, the undoped tungsten oxide films, despite forming a monoclinic crystalline phase after annealing, do not have the desired (002) crystal plane for the water splitting reaction, resulting in low photocurrent. More importantly, the substitution doping of Ta into WO₃ leads to a significant increase in photocurrent density. The samples prepared with different Ta dopant concentrations (0-3.93 at%) were analyzed in detail. The WO₃ film doped with 3.93 at% Ta exhibited the maximum photocurrent density of ~0.65 mA/ cm² at 1 V (vs. Ag/AgCl). This value was about 3.42 times greater than undoped tungsten oxide electrodes. In this study, the achived photocurrent values are in good agreement with the results reported for Ta-doped WO, synthesized using hydrothermal method by Kalanur et al.⁶⁴ This can be due to the growth of the (002) crystal plane and the decrease in the bandgap energy (absorption of a large part of the incident light)¹¹⁰. Morever, the photocurrent value mentioned above was found to be noticeably higher than result reported for Ta-doped WO.

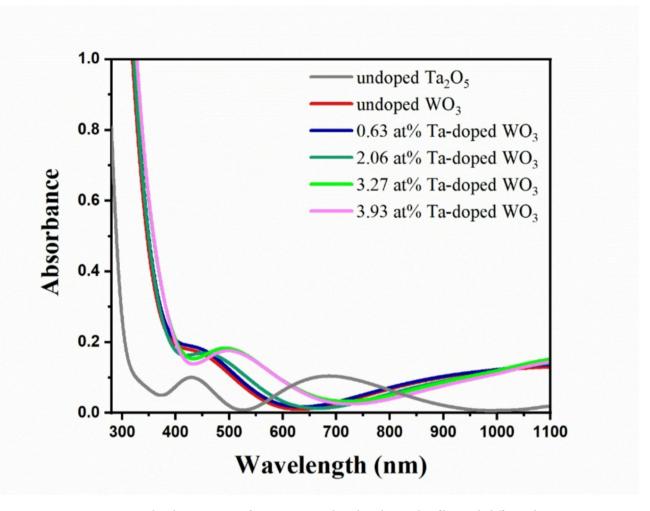


Fig. 8. Absorbance spectra of ${\rm Ta_2O_5}$, ${\rm WO_3}$, and ${\rm Ta\text{-}doped\ WO_3}$ thin films with different dopant concentrations.

prepared via spray pyrolysis by Enesca et al.¹¹¹ This results indicate the the success of the sputtering method to synthesize thin films with desired properties. Similarly, Mohamed et al. reported photocurrent enhancement in TiO, N,/TiO, bilayer and nitrogen-doped TiO, films prepared by reactive dc magnetron sputtering method for application in PEC efficient water splitting. This phenomenon is due to the bandgap narrowing effect caused by nitrogen incorporation 112,113. Comparatively, the photocurrent values obtained for sputter-deposited WO₃ and Ta-doped WO₃ films are listed in Table 4, along with previously reported results for WO₃ doped with various metals synthesized by different techniques. Overall water splitting activity was assessed using a two-electrode setup under on/off light illumination at 0 V bias. The measurements were carried out directly between the samples (working electrode) and the pt wire (counter electrode). Figure 13 depicts the dependency of photocurrent density on irradiation-time for undoped and Ta-doped WO₂ photoanodes. At the moment of illumination, a rapid increase in the photocurrent from zero to a certain value was observed for both undoped and doped samples. It is noticeable that the current density exhibited relatively stable trend during the period of light exposure and suddenly dropped to near zero at the moment of light extinction. Therefore, excellent photoresponsivity and slow electron-hole recombination are expected for WO₃ films doped with 3.93 at% Ta compared to undoped films¹¹⁴. On the other hand, the transient photocurrent (at 0 V) can be attributed to photogeneration of electron/hole pairs under light illumination. Finally, it can be concluded that the higher photocurrent of the Ta-doped electrodes is due to the increased electron/hole generation resulting from the narrowing of the optical bandgap.

Band edge positions

To realize the water splitting reaction on electrode materials, the production of charge carriers and redox potential are crucial parameters 98 . In the present work, the band structure of the prepared films is schematically discussed to evaluate the possible mechanism of the PEC reaction. The positions of valence band minimum (VBM) and conduction band maximum (CBM) for undoped and Ta-doped WO $_{3}$ samples were obtained by the following equations 115 :

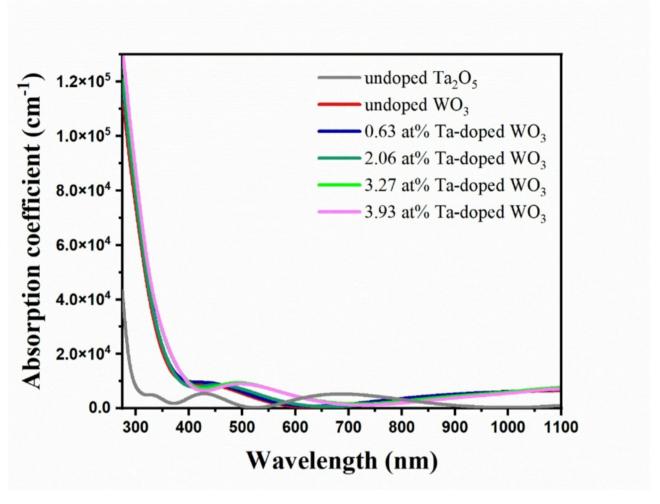


Fig. 9. Wavelength dependency of absorption coefficient for Ta₂O₅, WO₃, and Ta-doped WO₃ thin films with different dopant concentrations.

$$E_{CB} = X - E_{e} - \frac{1}{2}E_{g}$$
 (5)

$$E_{VB} = E_{CB} + E_{g} \tag{6}$$

Here, E_{CB} represents the conduction band edge, X is the determined Mulliken electronegativity, E_{e} is the energy of the free electrons on the hydrogen scale (~4.5 eV), E_{g} is the bandgap energy estimated from Tauc plot, and E_{VB} is the valence band edge. Considering the above equation, the calculated E_{CB}/E_{VB} values of undoped and Ta-doped WO $_{3}$ samples are displayed in Fig. 14. Considering that the redox potential of H^{+}/H_{2} is 0 eV, while the redox potential of $O_{2}/H_{2}O$ is 1.23 eV. Consequently, the water splitting reaction theoretically requires a minimum bandgap of 1.23 eV, which refers to light with a wavelength of about 1100 nm. From the Fig. 14, it can be seen that the replacement of W with Ta in WO_{3} leads to a downward shift in conduction band edge, an upward shift in valence band edge, and the significant reduction of the bandgap energy. This process occurs more noticeably with increasing Ta concentration.

Conclusion

In summary, this study reported the deposition of Ta-doped WO $_3$ thin films via reactive magnetron co-sputtering from W and Ta metal targets, and subsequent annealing at 550 °C. Tuning the elemental composition of Ta-doped WO $_3$ thin films by varying DC and RF sputtering powers allows narrowing of the band gap and altering of the band edge positions for use in PEC water splitting. By analyzing EDX results, it was concluded that the concentration of Ta dopant into WO $_3$ increased in the range of 0–3.93 at% with increasing RF sputtering power (0 to 60 W). From the FESEM and AFM micrographs, the growth of a granular structure and surface roughening were observed after Ta doping. Comparative studies of XRD patterns confirmed that the substitution of Ta atoms into the WO $_3$ structure during annealing improved the crystallinity, especially in the (002) direction. UV-Vis spectroscopic measurements showed that addition of Ta into WO $_3$ decreased the average transparency (from 92.82 to 74.27%), increased visible light absorption, and narrowed bandgap (from 3.07 to 2.61 eV). Linear sweep

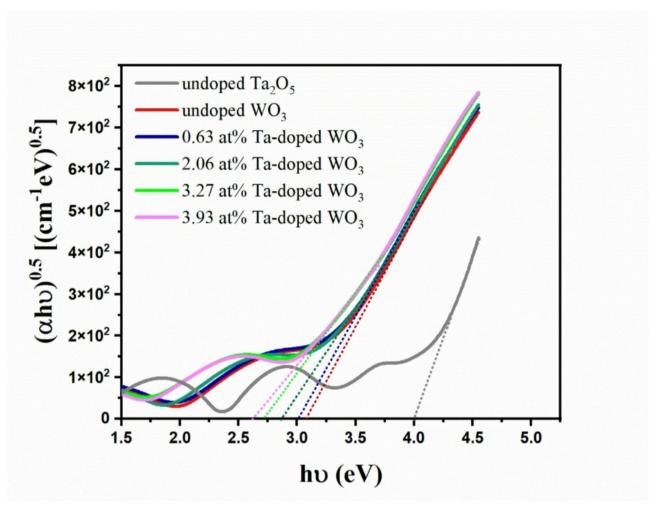


Fig. 10. Tauc plots for Ta₂O₅, WO₃, and Ta-doped WO₃ thin films with different dopant concentrations.

Thin film	E _g (eV)	Deposition techniques	References
	4.00	RF magnetron sputtering	This work
	4.18	Pulsed laser deposition	99
${\rm Ta_2O_5}$	4.20	Atomic layer deposition	100
	4.47	DC magnetron sputtering	101
	4.2	Ion-beam sputtering	102
WO ₃	3.07	DC magnetron sputtering	This work
	2.89-3.09	Impulse magnetron sputtering	103
	2.96-3.15	RF magnetron sputtering	104
	3.07	Pulsed dc magnetron sputtering	105
	2.90-3.18	DC magnetron sputtering	106
	3.00	Atomic layer deposition	82
	3.18	Chemical vapor deposition	107
	3.26	Evaporation-condensation	108
	3.20	Dual magnetron sputtering	109
Ta-doped WO ₃	3.01-2.61	Magnetron co-sputtering	This work
	2.61	Hydrothermal	64
	3	RF magnetron sputtering	77

 $\textbf{Table 3}. \ \ \text{Comparision of bandgap energy for Ta}_2\text{O}_5, \text{WO}_3, \text{and Ta-doped WO}_3 \text{ thin films}.$

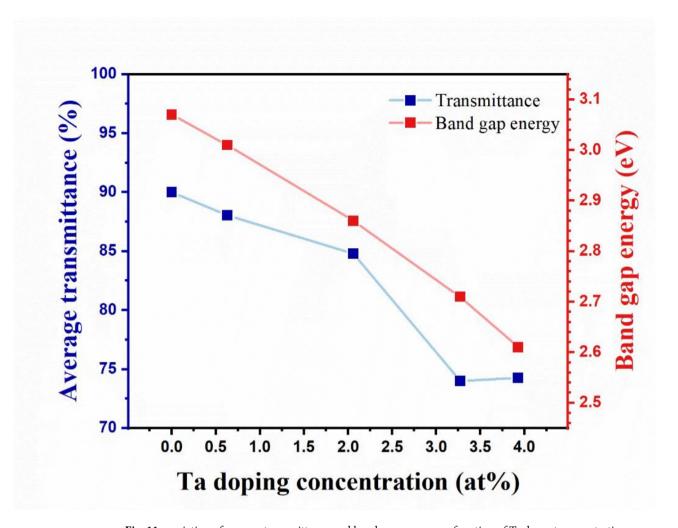


Fig. 11. variation of average transmittance and bandgap energy as a function of Ta dopant concentration.

voltammetry data showed that the photocurrent density of WO $_3$ increased from 0.19 mA/cm² to 0.65 mA/cm² (at 1 V vs. Ag/AgCl) after doping with 3.93 at% tantalum. The incorporation of Ta ions into WO $_3$ lattice shifted the valence band up and the conduction band down. Performed studies suggest that sputtered WO $_3$ thin films can be modified by Ta-doping combined with post-annealing for application in efficient photoelectrochemical water splitting.

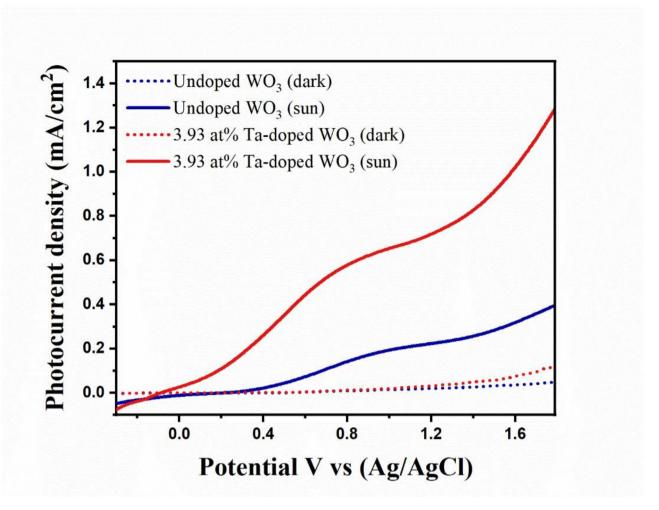


Fig. 12. Linear sweep voltammetry plots for undoped and Ta-doped WO₃.

Thin film	Photocurrent (mA/cm ²)	Measurement conditions	Deposition techniques	References
Ta-WO ₃	0.65	1 V (vs. Ag/AgCl)	Co-sputtering	This work
Ta-WO ₃	0.60	1.23 V (vs. RHE)	Hydrothermal	64
Gd-WO ₃	1.49	1 V (vs. Ag/AgCl)	Hydrothermal	63
Nb-WO ₃	0.414	1.23 V (vs. RHE)	Hydrothermal	76
Sn-WO ₃	0.427	1.23 V (vs. RHE)	Hydrothermal	60
Ti-WO ₃	1.139	1.23 V (vs. RHE)	Hydrothermal	84
V-WO ₃	0.201	1.23 V (vs. RHE)	Hydrothermal	47
Cr-WO ₃	0.024	1.23 V (vs. RHE)	Hydrothermal	47
Mn-WO ₃	0.295	1.23 V (vs. RHE)	Hydrothermal	47
Fe-WO ₃	0.442	1.23 V (vs. RHE)	Hydrothermal	47
Co-WO ₃	0.636	1.23 V (vs. RHE)	Hydrothermal	47
Ni-WO ₃	0.364	1.23 V (vs. RHE)	Hydrothermal	47
Cu-WO ₃	0.487	1.23 V (vs. RHE)	Hydrothermal	47
Zn-WO ₃	0.375	1.23 V (vs. RHE)	Hydrothermal	47

 $\textbf{Table 4}. \ \ \text{Comparision of photocurrent values for WO}_3 \ \text{thin films doped with various metals}.$

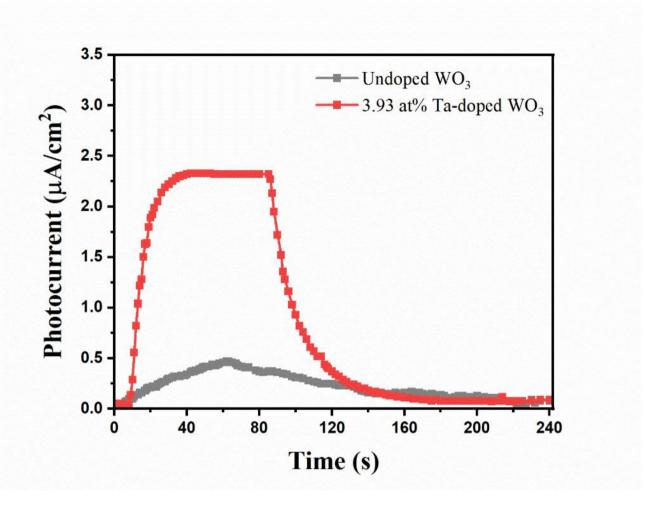


Fig. 13. Time dependence of photocurrent at 0 V for undoped and Ta-doped WO₃.

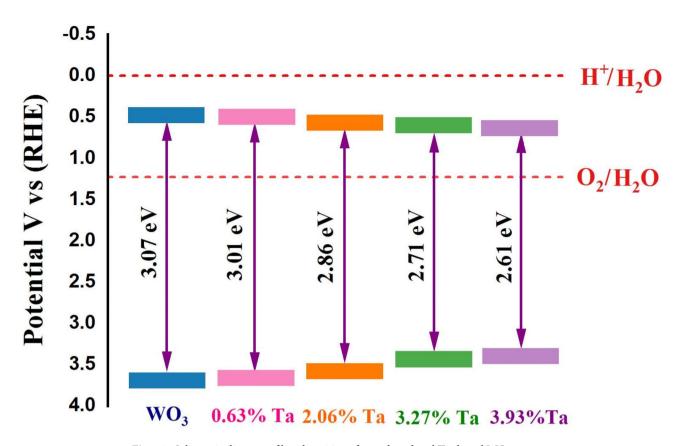


Fig. 14. Schematic diagram of band positions for undoped and Ta-doped WO₂.

Data availability

Data is provided within the manuscript or supplementary information files. The datasets used and/or analysed during the current study available from the corresponding author on reasonable request.

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Author contributions

H.Naghshara and S. Jafarpour are wrote the manuscript. H.Naghshara and S. Jafarpour prepared all of images. All authors reviewed the manuscript.

Declarations

Competing interests

The authors declare no competing interests.

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