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Investigation of space division multiplexing in multimode step-index silica photonic crystal fibers

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ABSTRACT

The feasible distance is presented for space division multiplexed (SDM) transmission along multimode silica step-index photonic crystal fiber (SI PCF) by solving the time-independent power flow equation (TI PFE). These distances for two and three spatially multiplexed channels were determined to depend on mode coupling, fiber structural parameters, and launch beam width in order to keep crosstalk in two- and three-channel modulation to a maximum of 20% of the peak signal strength. We found that the length of the fiber at which an SDM can be realized increases with the size of the air-holes in the cladding (higher NA). When a wide launch excites more guiding modes, these lengths become shorter. Such knowledge is valuable for the use of multimode silica SI PCFs in communications.

1. Introduction

An increase of optical fiber networks capacity is being driven by the demand for more data bandwidth. Multiplexing of optical data in space, time, polarization, wavelength and phase, may be able to address the strong demand for capacity growth. Both high-power and low-power fiber optics systems can use multiplexing [1,2]. SDM is a technology that allows multiple optical signals to be transmitted over separate spatial paths within a single-core or multicore fibers. The modes which are transmitted over these separate spatial paths are used to transmit separate data streams within the fiber, effectively increasing the capacity of the fiber. In each of these approaches, the individual modes within the fiber have to be carefully controlled and manipulated to ensure that they can be used to transmit data effectively. One of the key advantages of SDM fiber optic systems is that they can provide a significant increase in capacity without the need for additional fibers or complex network infrastructure. This can help to reduce costs and simplify network management, making it easier to deploy and maintain high-speed communication networks. In the past ten years, SDM has received a lot of interest for the capacity development in fiber optics communications [1-8]. Each SDM channel in the case of SDM is given a radially distributed spatial location inside the fiber. The incident angle and strength of coupling between a propagating modes define the position of each channel within the fiber. While all subsequent channels characterized with an incident angle of $\theta_0 > 0^\circ$ seem as concentric rings, the central channel which is with an incident angle of $\theta_0 = 0^\circ$ parallel to the fiber axis actually looks like a disk. The

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optical fiber system's bandwidth is increased since each ring and the central disk constitute its own individual spatially modulated optical channel (Fig. 1). However, the intended SDM technique is hampered by random coupling between the channels. With increasing distance from the input end, this interference becomes more and more pronounced. It results from macro- and micro-bending as well as impurities and other flaws that are trapped into the fiber during the production process. Power is transferred between neighboring channels by these irregularities. Referred to as "mode coupling", the outcome is a crosstalk between the SDM channels. With each successive imperfection in the fiber that is encountered, effects spread outward from the input end as more and more of the launch light gets connected to nearby channels. It is unquestionably crucial for the operation of SDM to characterize such performance-limiting effects of mode coupling that result in crosstalk in an SDM scheme. By solving the TI PFE, we accomplish it in this study for the multimode silica SI PCF with a solid-core.

The refractive-index (RI) change in an optical fiber has historically been accomplished by chemical doping. Another method [9,10] includes fabricating a micro-structured pattern of minuscule air-holes in the cladding along the PCFs' length (Fig. 2(a)). Since the air-hole pattern reduces the cladding's effective RI, the fiber can act as a light guide. The RI profile of such a fiber is shown in Fig. 2(b). Commercially available simulation software packages are not adequate for multimode PCFs due to hundreds or even thousands of propagation modes. In this paper, the TI PFE is numerically solved to address this deficiency. In our knowledge for the first time in this work we investigate an SDM capability of multimode silica SI PCF.

2. SI PCF design

The SI PCF's cladding is covered with air-holes that are typically homogeneous in size and arranged in a form of triangular lattice. The pitch (Λ) or the size (*d*) of these air-holes can be changed to get the required effective RI, numerical aperture and dispersion (Fig. 2). Despite having uniform material qualities throughout, the middle part of the fiber without holes has the highest RI n_0 . Air-holes in the cladding lower this index's effective value n_1 , and larger or more sparsely distributed air-holes further lower it.

2.1. Time-independent power flow equation

We investigate power flow inside the SI PCF using Gloge's TI PFE equation [11]:

$$\frac{\partial P(\theta, z)}{\partial z} = -\alpha(\theta)P(\theta, z) + \frac{D}{\theta}\frac{\partial}{\partial\theta}\left(\theta\frac{\partial P(\theta, z)}{\partial\theta}\right)$$
(1)

The term *P* (θ ,*z*) is the optical power as a function of the angle of propagation θ and propagation distance *z* and *D* is the coupling coefficient [11–14]. Since mode coupling is unaffected by the modal attenuation, equation (1) reads:



Fig. 1. SDM schematics with three channels.



(b)

Fig. 2. (a) Cross section of multimode SI PCF with a solid-core and rings of air-holes in the cladding, where Λ is the hole-to-hole spacing (pitch), *d* is diameter of air-holes in the cladding. (b) The referent multimode SI PCF's RI profile.

$$\frac{\partial P(\theta, z)}{\partial z} = \frac{D}{\theta} \frac{\partial P(\theta, z)}{\partial \theta} + D \frac{\partial^2 P(\theta, z)}{\partial \theta^2}$$
(2)

The explicit finite-difference method (EFDM) is used to solve equation (2) for the optical power $P(\theta, z)$ [14]. At launch, the following Gaussian input power distribution was assumed:

$$P(\theta, z) = \frac{1}{\sigma\sqrt{2\pi}} \exp\left[-\frac{(\theta - \theta_0)^2}{2\sigma^2}\right]$$
(3)

where θ_0 is the mean value for the launch angle and σ is standard deviation.

3. Numerical results

In order to investigate the capacity of SDM for this multimode silica SI PCF, we were able to calculate $P(\theta, z)$ along the fiber by employing the TI PFE. This has been done for various widths of Gaussian input power distributions and cladding's air-hole diameters *d*. For PCFs shown in Fig. 2(a), the effective parameter *V* can be utilized to calculate the effective RI of the cladding n_{fsm} :

$$V = \frac{2\pi}{\lambda} a_{eff} \sqrt{n_0^2 - n_{fsm}^2}$$
(4)

where n_0 is the RI of the core, $n_1 \equiv n_{fsm}$ is the effective RI of the cladding and $a_{eff} = \Lambda/\sqrt{3}$ [15]. Using the effective parameter V, equation (4) may be used to determine n_{fsm} [15]:

$$V\left(\frac{\lambda}{\Lambda},\frac{d}{\Lambda}\right) = A_1 + \frac{A_2}{1 + A_3 \exp(A_4\lambda/\Lambda)}$$
(5)

with the fitting parameters A_i (i = 1 to 4):

$$A_i = a_{i0} + a_{i1} \left(\frac{d}{\Lambda}\right)^{b_{i1}} + a_{i2} \left(\frac{d}{\Lambda}\right)^{b_{i2}} + a_{i3} \left(\frac{d}{\Lambda}\right)^{b_{i3}} \tag{6}$$

where the coefficients a_{i0} to a_{i3} and b_{i1} to b_{i3} (i = 1 to 4) are given in Table 1. The RI n_{fsm} , relative RI difference Δ and critical angles θ_m for varied air-holes diameter of the cladding d, pitch $\Lambda = 3 \mu m$ and $\lambda = 850$ nm, are summarized in Table 2.

Using the EFDM [14], we solved the TI PFE (2) for the multimode SI PCF with silica core with $2a = 50 \mu m$, $n_0 = 1.45$ and $D = 2.3 \times$ 10^{-6} rad²/m [16]. In the calculations one can use the same value of D for conventional silica fibers and silica PCF since it is predominantly governed by fiber material. We examined scenarios with $d = 1 \mu m$ and $2 \mu m$ and Gaussian incident beams with (FWHM)_z = $_{0} = 1.3$ and 5°. Gaussian beams (3) were launched into multimode silica SI PCF with solid-core with d = 1 µm at their respective mean input angles of $\theta_0 = 0$ and 4° (two-channel SDM) and $\theta_0 = 0$, 2 and 4° (three-channel SDM). Gaussian beams were launched into multimode silica SI PCF with $d = 2 \mu m$ at their respective mean input angles of $\theta_0 = 0$ and 8° (two-channel SDM) and $\theta_0 = 0$, 4 and 8° (three-channel SDM). Figs. 3 and 4 demonstrate the normalized power distributions coming from the output fiber-ends in the case of two- and three-channel SDM, respectively, for such input into multimode silica SI PCF with $d = 2 \mu m$. As seen in Figs. 3(a) and 4(a), where the distributions are essentially unchanged from when the fibers were first launched and exhibit no evidence of signal crosstalk, such distributions remain almost unchanged within the first few tens of meters of fibers. Figs. 3(b) and 4(b) show that with increasing fibers' lengths, coupling between propagating modes results in the broadening of Gaussian distributions. As a result, there is some crosstalk between adjacent co-propagating optical channels because the tails of the various Gaussians overlap to an extent of about 20% of their maxima. In Figs. 3(c) and 4(c), the situation worsens significantly for longer fibers. The rings and the center dot in Fig. 1 that result from these distributions may still be recognizable, but they are severely blurred and morphed close together. This is due to the large crosstalk between neighboring co-propagating optical channels, which is evident as the individual distributions are more difficult to resolve from the overall image. Figs. 3(d) and 4(d) are related to circumstances at "coupling length" L_c (which is the length in a particular fiber at which the highest order launched modes shifted their distributions from ring into disk form). It should come as no surprise that the mode-distributions for the optical channels overlap significantly. The rings and the center dot seen in Fig. 1 are no longer recognizable in the images they produce, which have changed into overlapping disks. In this fiber, three-channel SDM is feasible up to lengths that are significantly shorter than those at which two-channel SDM can be accomplished, with the tolerable mixing threshold of 20% of the maximum intensity being used (Figs. 5 and 6). It should be noticed that these lengths z_{SDM} , with the acceptable degree of crosstalk between two and three spatially multiplexed channels, are significantly shorter than the corresponding lengths L_c (Table 3). It is clear that the length L_c increases as the diameter d of air holes increases (increase of NA). This is a result of higher-order modes participating more frequently in higher NA PCFs. We also demonstrate that the length at which EMD is created decreases with increasing launch beam width (Table 3). This is because a wide launch beam drives the EMD at shorter distances than a narrow launch beam does.

It should also be mentioned that in our earlier research, we successfully examined the SDM capabilities in standard multimode SI plastic optical fibers utilizing the TI PFE [17]. In comparison to the multimode silica SI PCF examined in this work, the transmission lengths for space division multiplexing in standard multimode SI plastic optical fibers are significantly shorter. This is dominantly a result of multimode plastic optical fibers having substantially stronger mode coupling than multimode silica optical fibers. For standard plastic optical fibers, a typical value of coupling coefficient is $D \sim 10^{-4} \text{ rad}^2/\text{m}$ [14,17]. The results reported in this work could be of great interest not only in fiber optic communication systems but also for various fiber optic sensory systems [18–20]. Finally, it should be mentioned that a similar modeling approach can be used for investigation of bandwidth in each SDM channel.

4. Conclusion

The transmission length limit for SDM in multimode silica SI PCF with solid core is investigated using the TI PFE. For the acceptable

Table 1	
Fitting coefficients in equation (6).	

0 1				
	i = 1	i = 2	i = 3	<i>i</i> = 4
a_{i0}	0.54808	0.71041	0.16904	-1.52736
a_{i1}	5.00401	9.73491	1.85765	1.06745
a_{i2}	-10.43248	47.41496	18.96849	1.93229
a_{i3}	8.22992	-437.50962	-42.4318	3.89
<i>b</i> _{<i>i</i>1}	5	1.8	1.7	-0.84
b _{i2}	7	7.32	10	1.02
<i>b</i> _{<i>i</i>3}	9	22.8	14	13.4

Table 2

Effective RI of the cladding n_1 (obtained using equations (5) and (6)), relative RI difference $\Delta = (n_0 - n_1)/n_0$, where $n_0 = 1.45$, and the critical angle θ_m for varied air-holes diameter d (wavelength: 850 nm).

d (μm)	1.0	2.0
<i>n</i> ₁	1.443717	1.423679
$\Delta = (n_0 - n_1)/n_0$	0.00433	0.01815
θ_m (deg)	5.34	10.93



Fig. 3. The evolution of the normalized output angular power distribution with fiber length for the case with $d = 2 \mu m$ calculated for two Gaussian beams (two-channel SDM) with input angles $\theta_0 = 0^\circ$ (red line) and 8° (blue line), with (FWHM)_{z = 0} = 1° for: (a) z = 30 m; (b) z = 280 m; (c) z = 900 m and (d) $z \equiv L_c = 1650$ m.

crosstalk between 2- and 3-spatially multiplexed channels, these distances were found to be dependent on the strength of mode coupling, the size of the air holes in the cladding and the width of the incident beam. An SDM can be realized at longer lengths because the EMD is established at longer fiber lengths with larger cladding's air-holes. These lengths shorten for a wide launch that stimulates additional guiding modes. As a comparison, the transmission lengths for SDM in multimode plastic optical fibers are significantly shorter. This is dominantly a result of multimode plastic optical fibers having substantially stronger mode coupling than multimode silica optical fibers. When employing a SDM using a multimode silica SI PCFs, results reported in this work should be taken into account.

Author contribution statement

Svetislav Savović: Conceived and designed the experiments; Performed the experiments; Analyzed and interpreted the data; Contributed reagents, materials, analysis tools or data; Wrote the paper.

Konstantinos Aidinis, Alexandar Djordjevich, and Rui Min: Analyzed and interpreted the data; Contributed reagents, materials, analysis tools or data; Wrote the paper.

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Fig. 4. The evolution of the normalized output angular power distribution with fiber length for the case with $d = 2 \mu m$ calculated for three Gaussian beams (three-channel SDM) with input angles $\theta_0 = 0^\circ$ (red line), 4° (green line) and 8° (blue line), with (FWHM)_{z = 0} = 1° for: (a) z = 30 m; (b) z = 80 m; (c) z = 400 m and (d) $z \equiv L_c = 1650$ m.



Fig. 5. Length z_{SDM} for two and three spatially multiplexed channels with minimal crosstalk in multimode silica SI PCF as a function of (FWHM)_z = $_0$ of the incident beam, for air-holes diameter in the cladding d = 1 µm.

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Data availability statement

Data will be made available on request.

Additional information

No additional information is available for this paper.



Fig. 6. Length z_{SDM} for two and three spatially multiplexed channels with minimal crosstalk in multimode silica SI PCF as a function of (FWHM)_z = 0 of the incident beam, for air-holes diameter in the cladding $d = 2 \mu m$.

Table 3

Minimal crosstalk length z_{SDM} for two and three spatially multiplexed channels and coupling length L_c for EMD in multimode silica SI PCF with d = 1 µm and d = 2 µm, for different (FWHM)_{z = 0} of the incident beam.

	$(FWHM)_{z = 0=10}$			$(FWHM)_{z = 0=3o}$			$(FWHM)_{z = 0=5o}$		
	zSDM [m] 2- channel	zSDM [m] 3- channel	<i>L</i> _c [m]	<i>zSDM</i> [m] 2- channel	<i>zSDM</i> [m] 3- channel	<i>L</i> _c [m]	zSDM [m] 2- channel	zSDM [m] 3- channel	<i>L_c</i> [m]
d = 1	85	35	500	75	30	410	62	25	300
$d=2 \ \mu m$ μm	280	80	1650	260	70	1520	250	60	1450

Declaration of competing interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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