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# Low-Temperature Observation of the Excited-State Decay of Ruthenium-(Mono-2,2':6',2"-Terpyridine) lons with Innocent Ligands: DFT Modeling of an <sup>3</sup>MLCT-<sup>3</sup>MC Intersystem Crossing Pathway

Chi-Wei Yin,\* Ming-Kang Tsai,\* and Yuan Jang Chen\*



series of Ru-tpy complexes. The target  $[Ru(tpy)(AN)_3]^{2+}$  and  $[Ru(tpy)(bpm)_{-}]^{2+}$ (AN)]<sup>2+</sup> complexes were found to have low-emission quantum yields in lowtemperature observations. To better understand this phenomenon, density functional theory (DFT) calculations were performed to simulate the singlet ground state  $(S_0)$ ,  $T_{e'}$  and metal-centered excited states  $({}^{3}MC)$  of these complexes. The calculated energy barriers between Te and the low-lying <sup>3</sup>MC state for  $[Ru(tpy)(AN)_3]^{2+}$  and  $[Ru(tpy)(bpm)(AN)]^{2+}$  provided clear evidence in



support of their emitting state decay behavior. Developing a knowledge of the underlying photophysics of these Ru-tpy complexes will allow new complexes to be designed for use in photophysical and photochemical applications in the future.

## 1. INTRODUCTION

Metal complexes have been extensively studied and have been widely utilized in various aspects of research during the past 50 years.<sup>1-15</sup> It can be generalized from these studies that the relative energetics between triplet metal-to-ligand charge transfer emitting states  $({}^{3}MLCT, T_{e})$  and triplet metal-centered excited states (<sup>3</sup>MC, emission quencher) of ruthenium polypyridine complexes would determine the suitability for a specific photochemical application.<sup>16,17</sup> For example, the rate constant for an internal conversion from <sup>3</sup>MLCT to <sup>3</sup>MC,  $k_{ic}$  $({}^{3}MLCT \rightarrow {}^{3}MC$  state), for Ru-polypyridine complexes guenches the phosphorescence at room temperature (RT),<sup>18</sup> a characteristic that could be harmful to photodynamic therapy (PDT) but favorable for photochemotherapeutic (PCT) applications.<sup>19,20</sup>

The electronic excited states of transition-metal complexes were first demonstrated experimentally by Crosby during 1970-1975.<sup>7,21-23</sup> In these seminal studies, the behavior of excited states of  $[Ru(bpy)_3]^{2+}$  (bpy = 2,2'-bipyridine) was experimentally probed by the emission spectra at 77 K. Since then,  $[Ru(bpy)_3]^{2+}$  has become a classic model of ruthenium polypyridine complexes, and Ru polypyridyl analogues of this complex have been widely discussed in many studies.<sup>6,24-29</sup> In addition, based on their photophysical and photochemical character, this class of complexes has been utilized in a variety of applications, including in dye-sensitized solar cells (DSSCs),<sup>24,25,30-32</sup> photocatalysis,<sup>33,34</sup> PCT,<sup>35,36</sup> and PDT.<sup>37,38</sup>

Comparably,  $[Ru(tpy)_2]^{2+}$  is another classic model of a ruthenium polypyridine complex. Photophysical and photochemical properties of  $[Ru(tpy)_2]^{2+}$  (tpy = 2,2':6',2"-terpyridine) and its analogues have been reported in a large number of publications in the past 40 years.<sup>39–45</sup>  $[Ru(tpy)_2]^{2+}$  appears to have different structural characteristics in comparison to [Ru(bpy)<sub>3</sub>]<sup>2+</sup> with efficient excited-state quenching at room temperature. For example, the bite angle between Ru and tpy is about 157°, which is an unfavorable angle and causes the Ru-tpy moiety to have a weaker ligand field than that for a Ru-bpy moiety.<sup>46</sup> Consequently,  $[Ru(tpy)_2]^{2+}$  should have a more stable ligand-field state (metal-centered state, <sup>3</sup>MC state) than  $[Ru(bpy)_3]^{2+}$ , a characteristic that enhances the probability of an ic ( ${}^{3}MLCT \rightarrow {}^{3}MC$  state) and a decreased energy barrier  $({}^{3}MLCT \rightarrow {}^{3}MC \text{ state})$  for  $[Ru(tpy)_2]^{2+.44}$  This perspective is further supported by the observations of nonemission at RT for  $[Ru(tpy)_2]^{2+}$  and phosphorescence emission at RT for  $[Ru(bpy)_3]^{2+2,45,47-49}$ 

The energy gap  $({}^{3}MLCT \rightarrow {}^{3}MC$  state) is an important parameter that determines the probability of ic (<sup>3</sup>MLCT  $\rightarrow$ 

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<sup>3</sup>MC state). If the <sup>3</sup>MC state of a Ru-polypyridine complex is lower in energy than its <sup>3</sup>MLCT state, ic (<sup>3</sup>MLCT  $\rightarrow$  <sup>3</sup>MC state) can take place easily, being favorable for PCT. On the other hand, if a complex has a higher <sup>3</sup>MC state over the <sup>3</sup>MLCT state, the probability of ic (<sup>3</sup>MLCT  $\rightarrow$  <sup>3</sup>MC state) would decrease and would have more potential for PDT. The presence of the tpy ligand could result in a stable <sup>3</sup>MC state as mentioned above, and varying the second ligand may introduce the possibility of tuning the relative energetics between <sup>3</sup>MC and <sup>3</sup>MLCT states.

In this study, we attempt to control the relative energetics between <sup>3</sup>MLCT and <sup>3</sup>MC through adjusting the ligands of Rutpy complexes, being characterized by joint experimental and theoretical approaches. A series of Ru-tpy complexes were synthesized, and the ground and excited states were comprehensively characterized. The phosphorescence emission energies of these Ru-tpy complexes were measured as 13 000- $18\,000 \text{ cm}^{-1} (555-769 \text{ nm})$  at 77 K in the presence of amine (NH<sub>3</sub>), acetonitrile (AN), and bis(pyrazolyl)methane (bpm) ligands. Furthermore, the nonradiative decay rate constant  $(k_{nr})$ of the  $T_e$  state at room temperature may contain the  $k_{ic}$  in the ruthenium polypyridine complex system. However, the  $k_{nr}$  of a classical Ru-bpy complex can be inferred to be much greater than  $k_{ic}$  under low-temperature conditions at 77 K, as it complies with the energy gap law.<sup>3,50,51</sup> This consequently suggests that a series of classical Ru-bpy complexes could be used as a reference to compare with other series of complexes. In the present study, the series of classical Ru-bpy complexes (1)-(5) (the structures of these complexes are shown in Figure S1A1)<sup>51</sup> were also utilized to highlight the special properties of Ru-tpy complexes (a)-(e) (see Scheme 1).







### 2. EXPERIMENTAL SECTION

**2.1. Materials and Synthesis of Compounds.** In this work, five Ru-tpy complexes (a)–(e) were prepared, and the skeletal structures of these target Ru-tpy complexes are shown in Scheme 1. The  $[Ru(tpy)_2]^{2+}$  (a) and  $[Ru(tpy)(AN)_3]^{2+}$  (c) were prepared by literature procedures.<sup>41,52</sup> Details concerning the starting material and methods used for the synthesis of the new complexes  $[Ru(tpy)(NH_3)_3]^{2+}$  (b),  $[Ru(tpy)(bpm)(NH_3)]^{2+}$  (d), and  $[Ru(tpy)(bpm)(AN)]^{2+}$  (e) are presented in Section S1. The NMR (<sup>1</sup>H and <sup>13</sup>C) spectra and elemental analyses for these complexes are shown in Section S1.

**2.2. Instrumentation.** Details concerning the electrochemistry, absorption spectra, 77 K emission, 77 K emission lifetime, and 77 K emission quantum yield ( $\Phi_{(em)}$ ) measurements are presented in Section S1. On the other hand, the 77 K condition would decrease the probability of another quenching state and allow focusing on probing the <sup>3</sup>MLCT emission.

**2.3. Computational Details.** The optimizations of the geometry for the singlet ground state  $(S_0)$  and the triplet excited states ( $T_e$  and  ${}^{3}MC$ ) for all of the Ru-tpy complexes were characterized at the density functional theory (DFT) level with the B3PW91 functional  ${}^{53-56}$  and the Stuttgart/Dresden effective core potential (SDDall) basis set  ${}^{57-59}$  in the Gaussian 16 package.<sup>60</sup> To confirm the minimum structures, harmonic frequency analyses were carried out. The theoretical electronic absorption spectra were simulated by time-dependent DFT calculations with the IEF-PCM solvation model (acetonitrile).<sup>61-63</sup> The natural transition orbitals (NTOs)<sup>64</sup> for the selected excited states for the Ru-tpy complexes were calculated to provide a better vision of the excited states for TD-DFT (see Figure S3C1A–C5C).

Additionally, the B3LYP,  $^{65,66}$  PBE0,  $^{67}$  CAM-B3LYP,  $^{68}$  and LC-wPBE<sup>69</sup> functionals were benchmarked with B3PW91 for describing the low-energy <sup>1</sup>MLCT electronic absorption envelopes of  $[Ru(tpy)_2]^{2+}$  (see Figure S3A1). It appeared that data collected using the B3PW91 functional are in better agreement with the experimental absorption peaks, being consistent with other early reports and our previous studies.<sup>3,50,51,70</sup> Therefore, the B3PW91 functional was finally utilized to simulate the ground and excited states of the Ru-tpy complexes.

### 3. RESULTS

**3.1. Electrochemical and Spectroscopic Properties of the Complexes.** *3.1.1. Electrochemistry.* Cyclic voltammetry (CV) was performed at room temperature for each complex in acetonitrile solutions containing 0.1 M n-TBAH, and the measured potentials are listed in Table S1. All of the series of Rutpy complexes showed reversible Ru-based oxidation and reversible ligand-based reduction. The cyclic voltammogram for  $[Ru(tpy)_2]^{2+}$  displayed two  $tpy^{0/1-}$  oxidation couples at -1.27 and -1.51 V vs Ag/AgCl, and the other Ru-tpy complexes exhibited only one  $tpy^{0/1-}$  oxidation couple in the negative potential (see Figure S2A1).

3.1.2. UV–Vis Absorption Spectroscopy. The room-temperature (RT) absorption spectra for the Ru-tpy complexes (a)– (e) are shown in Figure 1, and the low-temperature absorption spectra for the Ru-tpy complexes (a)–(e) are shown in Figure S2B3. In Figure 1, the spectra of  $[Ru(tpy)_2]^{2+}$  and  $[Ru(tpy)-(AN)_3]^{2+}$  were in agreement with previously reported spectra.<sup>41,44</sup> For all of the complexes (a)–(e), the low-energy absorption bands were dominated by MLCT absorption and were assigned as the  $Ru(d\pi) \rightarrow tpy (\pi^*)$  CT band. The absorbance maximum ( $\varepsilon_{max}$ ) of two Ru-tpy moieties for complex (a),  $[Ru(tpy)_2]^{2+}$  (black curve in the lower panel of Figure 1), was about 15 800 M<sup>-1</sup> cm<sup>-1</sup> (at 21 000 cm<sup>-1</sup>), a value that is about 3 times larger than the  $\varepsilon_{max}$  of one Ru-tpy moiety for the target complexes (b)–(e) (about 4000–5000 M<sup>-1</sup> cm<sup>-1</sup> as shown in Figure 1).

On the other hand, in Figure 1, when comparing the energies of the absorption maximum  $(h\nu_{max(abs)})$  of <sup>1</sup>MLCT absorption of  $[Ru(tpy)(NH_3)_3]^{2+}$  (b, red) and  $[Ru(tpy)(AN)_3]^{2+}$  (c, blue), it could be observed that the  $h\nu_{max(abs)}$  of  $[Ru(tpy)(AN)_3]^{2+}$  is higher than that of  $[Ru(tpy)(NH_3)_3]^{2+}$  (see Table 1), a



**Figure 1.** Top panel: 298 K absorption spectra of target complexes (Rutpy complexes)  $[Ru(tpy)(NH_3)_3]^{2+}$  (b, red),  $[Ru(tpy)(AN)_3]^{2+}$  (c, blue),  $[Ru(tpy)(bpm)(NH_3)]^{2+}$  (d, dark cyan), and  $[Ru(tpy)(bpm)-(AN)]^{2+}$  (e, magenta) in acetonitrile; lower panel: 298 K absorption spectra of target complexes  $[Ru(tpy)_2]^{2+}$  (a, black).

difference that can be attributed to the significant disparity of the ligand field for the  $\pi$ -withdrawing AN ligand and the  $\sigma$ -only donor NH<sub>3</sub> ligand. A similar case can be observed when comparing the  $h\nu_{max(abs)}$  of  $[Ru(tpy)(bpm)(NH_3)]^{2+}$  (d, dark cyan) and  $[Ru(tpy)(bpm)(AN)]^{2+}$  (e, magenta).

3.1.3. 77 K Emission Spectra. The 77 K emission from the  $\operatorname{Ru}(d\pi) \rightarrow \operatorname{bpy}(\pi^*)$  <sup>3</sup>MLCT state was observed for the Ru-bpy complexes (1), (2), and (4),<sup>3,71</sup> and from the  $\operatorname{Ru}(d\pi) \rightarrow$ tpy( $\pi^*$ ) <sup>3</sup>MLCT state, the emission was observed for Ru-tpy complexes (a)–(e),<sup>40,49</sup> respectively; see Figure 2. The band shapes of the 77 K emission for the Ru-bpy and Ru-tpy complexes consist of a fundamental band  $(I_{0'-0})$  and vibronic sidebands  $(I_{0'-i})$  and  $j \neq 0$ .<sup>72</sup> For example, the 77 K emission spectra of  $[Ru(bpy)_3]^{2+}$  (1, gray) shows a fundamental band at about 16 500-18 000 cm<sup>-1</sup> and a vibronic sideband at about 15 000-17 000 cm<sup>-1</sup> (sum of first-order vibronic components,  $I_{0'-1}$ ). Basically, first-order vibronic sidebands originate from some vibrational modes (distortion modes). According to the literature, the assignment of intense vibrational modes for  $[Ru(bpy)_3]^{2+}$  is in the inter-ring of the bpy ligand. The relative intensities of the first-order vibrational mode contributions were obtained from time-dependent resonance-Raman (TD-rR) data.<sup>72,73</sup> Unfortunately, TD-rR data is not easy to obtain and no TD-rR data for the Ru-tpy complexes was found in the literature, so we were unable to perform an analysis of the contribution of the first-order vibronic mode for the Ru-tpy complexes. However, we were still able to observe that the intensity of the vibronic sideband in the emission spectra of the Ru-tpy complexes was lower than that for the Ru-bpy complexes with a similar emission energy  $(h\nu_{em})$  region. For example, in Figure 2, the intensity of the vibronic sideband for complex (b) is smaller than that for complex (1) with similar emission energies. On the other hand, the intensities of the vibronic sideband in the <sup>3</sup>MLCT emissions of the class of Ru-bpy complexes decrease with decreasing energy of the <sup>3</sup>MLCT emission,<sup>51,72,74</sup> a trend that was also observed in Ru-tpy complexes. Moreover, the emission spectra of [Ru(tpy)(bpm)-(AN)<sup>2+</sup> contain a broad and special vibronic sideband. As mentioned above, each of the Ru-tpy complexes should have its intrinsic Raman vibronic contribution to the emission intensity.

3.1.4. Comparison of the 77 K Emission Quantum Yield for the Target Complexes. Basically, the 77 K  $k_{nr}$  of the typical Rubpy complexes was not influenced by the  $k_{ic}$  (emitting <sup>3</sup>MLCT state to <sup>3</sup>MC state) and the 77 K emission quantum yields (QYs) for the typical Ru-bpy complexes (red squares, codes (1)–(5); see Figure S1A1), and this decreased systematically with decreasing  $h_{vave(em)}$  (this trend is not rigorously considered to be fundamental; see the red dashed line in Figure 3).<sup>3,51</sup> On the other hand, the 77 K quantum yields (QY) for the Ru-tpy complexes (c) and (e) with an AN ligand are lower than expected in the given emission energy, as indicated by the blue

Table 1. Ambient Absorption, 77 K Emission, Emission Decay Constants, and Emission Yield of the Ru-tpy Complexes<sup>a</sup>

			77 K emission <sup><math>b</math></sup>	77 K lifetime <sup>b</sup>			77 K quantum yield <sup>b</sup>
code	complexes	$h\nu_{\max(abs)}, cm^{-1}/10^3, 298 K$ (87 K) <sup>a</sup>	${{\rm h}}  u_{ m max(em)},{ m cm}^{-1}/10^3 \ ({ m h}  u_{ m ave(em)},{ m cm}^{-1}/10^3)^c$	$k_{ m obs} \mu s^{-1} ( au_{ m mean},  au_{ m mean})^{d,e}$	$\mu s^{-1} f$	$\mu s^{-1} f$	$\Phi_{\rm em} \times 10^4$
(a)	$[Ru(tpy)_2]^{2+}$	21.03 (20.92)	16.55 (15.70)	0.09 (11.17)	0.030	0.059	$3400 \pm 700$
(b)	$[Ru(tpy)(NH_3)_3]^{2+}$	18.50 (17.48)	13.01 (12.50)	2.63 (0.38)	0.025	2.607	95 ± 15
(c)	$[Ru(tpy)(AN)_3]^{2+}$	22.96 (22.62)	18.07 (16.86)	0.06 (16.2) (60%) 0.58 (1.72) (40%)			2090 ± 200
(d)	$ \begin{matrix} [{\rm Ru(tpy)(bpm)} \\ ({\rm NH_3}) \end{matrix} \end{matrix} ]^{2+} $	21.51 (21.46)	14.47 (13.69)	0.78 (1.28)	0.031	0.75	399 ± 50
(e)	[Ru(tpy)(bpm) (AN)] <sup>2+</sup>	23.09 (23.36)	15.96 (15.13)	0.18 (5.71)	0.015	0.16	$840\pm260$

<sup>*a*</sup>In acetonitrile. <sup>*b*</sup>In butyronitrile. <sup>*c*</sup> $\nu_{ave} \approx \int \nu_m I_m d\nu_m / \int I_m d\nu_m$ . <sup>*d*</sup>Mean emission lifetime ( $\tau$ ). <sup>*e*</sup>Mean excited-state decay rate constant,  $k_{obs}=1/\tau$ . <sup>*f*</sup> $k_r$  is the radiative rate constant, and  $k_{nr}$  is the nonradiative rate constant,  $k_{obs} = k_r + k_{nr}$ .



**Figure 2.** Top panel: 77 K emission spectra of all of the Ru-tpy complexes in butyronitrile solvent glasses:  $[Ru(tpy)_2]^{2+}$  (a, black),  $[Ru(tpy)(NH_3)_3]^{2+}$  (b, red),  $[Ru(tpy)(AN)_3]^{2+}$  (c, blue),  $[Ru(tpy)(bpm)(NH_3)]^{2+}$  (d, dark cyan), and  $[Ru(tpy)(bpm)(AN)]^{2+}$  (e, magenta); lower panel: 77 K emission spectra of target complexes  $[Ru(tpy)_2]^{2+}$  (a, black) and reference complexes  $[Ru(bpy)_3]^{2+}$  (1, gray),  $[Ru(bpy)_2(en)]^{2+}$  (2, dash dot), and  $[Ru(bpy)(en)_2]^{2+}$  (4, black dash).

dashed line for the 77 K QY of the Ru-tpy complexes in Figure 3 (see black arrow in Figure 3). On the basis of the Ru-tpy complexes, it is possible that a <sup>3</sup>MC state exists that could cause emission quenching. We assume therefore that the unusually low 77 K quantum yield for complexes (c) and (e) that could exist as <sup>3</sup>MC excited states would reduce the value of the quantum yield at 77 K. This conclusion is consistent with our DFT calculations that have been discussed in detail in Section 4.1. As mentioned above, in the case of the Ru-tpy complexes, it is easier to control the probability of ic (<sup>3</sup>MLCT  $\rightarrow$  <sup>3</sup>MC) compared to the Ru-bpy complexes, a result that causes the Ru-tpy complex to have more flexibility for different applications.

3.1.5. 77 K Emission Lifetimes. To gain insights into the photophysical behavior of the Ru complexes, we probed the emission decay lifetime ( $\tau_{em}$ ) for Ru-tpy complexes. As shown in Table 1, all of the Ru-tpy complexes are marked by their <sup>3</sup>MLCT



**Figure 3.** Comparison of  $\ln(\Phi)$  vs  $h\nu_{ave(em)}$  for typical Ru-bpy complexes (red square) and target Ru-tpy complexes:  $[Ru(tpy)_2]^{2+}$  (a, blue square),  $[Ru(tpy)(NH_3)_3]^{2+}$  (b, blue square),  $[Ru(tpy)(AN)_3]^{2+}$  (c, green square),  $[Ru(tpy)(bpm)(NH_3)]^{2+}$  (d, blue square), and  $[Ru(tpy)(bpm)(AN)]^{2+}$  (e, green square). The red-dashed least-square line indicates typical Ru-bpy complexes<sup>3</sup> (1–5, red square) and the blue-dashed least-square line indicates target Ru-tpy complexes (a, b, d, blue square); the trend for these data is not considered to be rigorously fundamental.

excited-state lifetimes with one-exponential decay fitting, except for complex (c),  $[Ru(tpy)(AN)_3]^{2+}$ . The lifetime of  $[Ru(tpy)-(AN)_3]^{2+}$  presents two-exponential decay (16.2  $\mu$ s, 60%; and 1.72  $\mu$ s, 40%), and the longer one could be characterized by the relaxation process Ru-tpy <sup>3</sup>MLCT  $\rightarrow$  ground state (S<sub>0</sub>); this data, which is similar to the lifetime of  $[Ru(tpy)_2]^{2+}$  in the given region, and the short-lifetime component could be attributed to the process <sup>3</sup>MLCT  $\rightarrow$  <sup>3</sup>MC. This behavior has been discussed in Section 4.

On the other hand, the function of the observed rate constant  $(k_{obs})$  from the emission decay lifetime  $(\tau_{em})$  could be given by<sup>1</sup>

$$k_{\rm obs} = k_{\rm r} + k_{\rm nr} + k_{\rm q} \cdots$$
 (1)

where  $k_r$  is a radiative rate constant,  $k_{nr}$  is a nonradiative rate constant, and  $k_q$  is the quenching reaction rate constant. Additionally, if no  $k_q$  mechanism is operative, the ideal model for emitting state decay, eq 1, would exclude  $k_q$  and can be given by<sup>45,75</sup>

$$(\tau_{\rm em})^{-1} = k_{\rm obs} = k_{\rm r} + k_{\rm nr} ({\rm at}\ 77\ {\rm K})$$
 (2)

According to the formalism eq 1 and eq 2, the  $k_{\rm nr}$  values can be obtained from emission quantum yield ( $\Phi_{\rm em}$ ) and one exponential lifetime data ( $\tau_{\rm em}$ ),  $\Phi_{\rm em} = k_{\rm r}/k_{\rm obs}$ , for the Ru-tpy complexes (a), (b), (d), and (e). On the other hand, the behavior of nonradiative decay ( $k_{\rm nr}$ ) for these four target complexes can be compared with the reference complexes (1)~(5) in the next section.

3.1.6. Comparison of Nonradiative Decay  $(k_{nr})$  for Target and Reference Complexes. As the nonradiative decay  $(k_{nr})$  at

77 K plays a vital rule in the excited-state decay kinetics, the nonradia formalism, eq. 3 and eq. 4, based on a single distortion mode

formalism, eq 3 and eq 4, based on a single distortion mode  $(h\nu_k)$  provides a method to probe the  $k_{nr}$  for different chromophores.<sup>76,77</sup>

$$k_{\rm nr} \approx C_{\rm nr} \exp\left[-A\left(\frac{h\nu_{\rm ave}}{h\nu_{\rm k}}\right)\right]$$
 (3)

$$\ln(k_{\rm nr}) \approx \ln C_{\rm nr} - \left[A \left(\frac{h\nu_{\rm ave}}{h\nu_{\rm k}}\right)\right] \tag{4}$$

where  $C_{\rm nr} \approx (H_{\rm eg}^{\ 2})[8\pi^3/(h^3\nu_{\rm k}h\nu_{\rm ave})]^{1/2}$  and  $A \approx \ln(h\nu_{\rm ave}/\lambda_{\rm k}) - 1$ . The  $\nu_{\rm k}$  frequency is assumed to be a harmonic distortion mode,  $H_{\rm eg}$  is a function of spin–orbit coupling for a phosphorescence-emitting state, and  $\lambda_{\rm k}$  is obtained from the Huang–Rhys parameters for this  $\nu_{\rm k}$  mode.<sup>78,79</sup> We use the average emission energy ( $h\nu_{\rm ave(em)}$ , Table 1) for Ru-bpy and Ru-tpy phosphorescence chromophores in eq 4 to fit the red and blue linear plots, respectively; see Figure 4. However, the



**Figure 4.** Low-temperature  $\ln(k_{nr})$  vs  $h\nu_{ave(em)}$  for Ru-bpy complexes (red square) and Ru-tpy complexes (blue squares):  $[Ru(tpy)_2]^{2+}$  (a),  $[Ru(tpy)(NH_3)_3]^{2+}$  (b),  $[Ru(tpy)(bpm)(NH_3)]^{2+}$  (d), and  $[Ru(tpy)-(bpm)(AN)]^{2+}$  (e). The respective slope, intercept, and  $R^2$  for least-squares trend lines are as follows: -1.11, 29.5, and 0.99 (red line for Ru-bpy complexes)<sup>3</sup>; -1.15, 29.2, and 0.99 (blue line for Ru-tpy complexes).

lifetime of this complex (c) can be fitted by two exponential decays, thus making it impossible to show in Figure 4. In Figure 4, the fits within the Ru-bpy and Ru-tpy complexes reveal that each polypyridine (bpy and tpy) can be described by a single line (red and blue lines), indicating that both MLCT chromophores are in reasonable agreement with the energy gap law.<sup>76</sup> The h $\nu_{ave}$  scale in Figure 4 is related to the  $h\nu_{ave}/h\nu_k$  scale of a target active distortion mode,  $\nu_k$ , and is associated with the harmonic model whose energy approximates that of the emitting state,  $jh\nu_k \approx h\nu_{ave}$ . On the other hand, the slopes of the red and blue lines for two different chromophore complexes are similar, and this resemblance indicates that the trends for the overall results of

nonradiative decay for these two different chromophore series have a certain degree of similarity based on the low temperature observation. Since the multiatom molecular-type chromophores can be considered to be multidistortion modes of the photophysical relaxation from an emitting state, attempts were made to find the Huang-Rhys parameters by time-dependent resonance-Raman (TD-rR) observations for many compound systems.<sup>78</sup> However, rare Ru-diimine complexes can be obtained with Huang-Rhys parameters by TD-rR observations, such as the  $[Ru(bpy)_3]^{2+}$  (15  $\nu_k$  modes)<sup>73</sup> and  $[Ru(NH_3)_4(bpy)]^{2+}$  (11  $\nu_{\rm k}$  modes).<sup>79</sup> The interpretation of the TD-rR results must be rigorous and it is difficult to compare the intrinsic terms ( $H_{eg}$  and  $\nu_{\rm k}$ ) in C<sub>nr</sub> amplitudes between the fitted parameters of leastsquare lines from Ru-bpy and Ru-tpy chromophores, as shown in Figure 4. Based on Figure 4, the only conclusion is that the low-temperature  $k_{nr}$  behaviors of the target Ru-tpy (exclude complex c) and the typical Ru-bpy chromophores are similar, and the low-temperature  $k_{nr}$  relaxation of the typical Ru-bpy chromophore can be considered to be "pure  $k_{nr}$  relaxation" for  $^{3}MLCT \rightarrow S_{0}.^{3}$ 

**3.2. Computational Results.** 3.2.1. Calculated Absorption Spectra. The correlation between  $\Delta E_{1/2}$  and the calculated  $S_1$  transition (using  $S_0$ -optimized structures) appeared to maintain a linear relationship, as shown in Figure 5, where  $\Delta E_{1/2}$  denotes the difference of the first oxidation  $E_{1/2(1st)}$  (Ru<sup>III/II</sup>) and reduction  $E_{1/2(1st)}$  (tpy<sup>0/1–</sup>) from the electrochemical measurements.

In Figure 6, the most notable electronic transitions, corresponding to the observed strong absorption peaks in the  $15\ 000-25\ 000\ cm^{-1}$  region for complexes (a), (b), and (d), can be assigned to <sup>1</sup>MLCT absorption. The detailed electronic structural information representing these experimental observed peaks were characterized by the significant oscillator strengths



**Figure 5.** Correlation of the calculated (B3PW91) S<sub>1</sub> transition using S<sub>0</sub> minimum geometry and electrochemical data  $F\Delta E_{1/2}$  for the [Ru-(tpy)<sub>2</sub>]<sup>2+</sup> (a), [Ru(tpy)(NH<sub>3</sub>)<sub>3</sub>]<sup>2+</sup> (b), [Ru(tpy)(AN)<sub>3</sub>]<sup>2+</sup> (c), [Ru-(tpy)(bpm)(NH<sub>3</sub>)]<sup>2+</sup> (d), and [Ru(tpy)(bpm)(AN)]<sup>2+</sup> (e) complexes. The linear dashed line was fitted with the slope of 1.02 and the intercept of -2209.



**Figure 6.** Comparison of the observed absorption (black curves) and DFT-calculated absorptions (red curves) for B3PW91 functionals in the left panels (red bars for  $S_n$  with oscillator strengths in the logarithm scale, log(f)), and their selected NTOs (right panel) of target complexes  $[Ru(tpy)_2]^{2+}$  (a, panel A),  $[Ru(tpy)(NH_3)_3]^{2+}$  (b, panel B), and  $[Ru(tpy)(NH_3)_3]^{2+}$  (d, panel C).

that were obtained by the TD-DFT calculations. The corresponding natural transition orbitals of these identified theoretical transitions (see Figure 6) appear to be localized on the Ru and tpy ligand for the donor and acceptor orbitals, respectively.

3.2.2. NTOs and Spin Density of the <sup>3</sup>MLCT Emitting State and <sup>3</sup>MC States for Complexes (a), (c), and (e). The minimum structures of the <sup>3</sup>MLCT and <sup>3</sup>MC states were estimated using the ground-state triplet optimizations, where the partial liganddissociated geometries were adopted for the starting point of locating the <sup>3</sup>MC minimum structures, as has been demonstrated by Soupart et al.<sup>2</sup> in 2018. Additionally, to reconfirm the electronic structures of <sup>3</sup>MLCT and <sup>3</sup>MC minimum structures, the NTO of the first spin-forbidden S<sub>1</sub>-to-T<sub>1</sub> excitation was exanimated by TD-DFT calculations with the aforementioned triplet minimum structures (see Figure S3D1). Because the bite angle of the tpy ligand is about 157°, the tpy ligand could bring a relatively unfavorable ligand field to the d-electrons of the metal. That could consequently destabilize the <sup>3</sup>MLCT state and result in an efficient downhill ic (<sup>3</sup>MLCT  $\rightarrow$  <sup>3</sup>MC). However, not all of the synthesized Ru-tpy complexes appeared to have <sup>3</sup>MC states lower than <sup>3</sup>MLCT states, as indicated by the success of <sup>3</sup>MC triplet optimization calculations. The optimized <sup>3</sup>MC structures were only identified in the cases of the [Ru(tpy)<sub>2</sub>]<sup>2+</sup>, [Ru(tpy)(AN)<sub>3</sub>]<sup>2+</sup>, and [Ru(tpy)(bpm)(AN)]<sup>2+</sup> complexes, however, and could not be achieved in the cases of [Ru(tpy)-(NH<sub>3</sub>)<sub>3</sub>]<sup>2+</sup> and [Ru(tpy)(bpm)(NH<sub>3</sub>)]<sup>2+</sup> complexes (see Figure



**Figure 7.** Illustrations of DFT results with  $T_{e^3}^{3}MC_a$ , and  ${}^{3}MC_b$  coordinates for  $[Ru(tpy)(AN)_3]^{2+}$  (right panel) and  $[Ru(tpy)(bpm)(AN)]^{2+}$  (left panel). Mulliken orbital population analyses, reorganization energies ( $\lambda_{reg}$ ) of Ru, and tpy moieties for the NTOs of  $T_{e^3}^{3}MC_a$ , and  ${}^{3}MC_b$  excited states.

S3D2–D4). These results are in reasonable agreement with our experimental data (Figure 3).

On the other hand, two distinguishable <sup>3</sup>MC minimum structures, <sup>3</sup>MC<sub>a</sub> (equatorial plane) and <sup>3</sup>MC<sub>b</sub> (vertical axis), were identified for  $[\text{Ru}(\text{tpy})(\text{AN})_3]^{2+}$  and  $[\text{Ru}(\text{tpy})(\text{bpm})-(\text{AN})]^{2+}$ , respectively. For example, the <sup>3</sup>MC<sub>a</sub> denotes the elongation of the equatorial AN to Ru and <sup>3</sup>MC<sub>b</sub> denotes the case of axial ANs for  $[\text{Ru}(\text{tpy})(\text{AN})_3]^{2+}$  (see Figure 7). In Figure 7, the orbital compositions of the NTOs of <sup>3</sup>MC and <sup>3</sup>MLCT states for each complex show very different results, in that the donor and acceptor of <sup>3</sup>MLCT states are located on the Ru and the tpy ligand, respectively, but those of <sup>3</sup>MC states are mostly located on Ru. On the other hand, the reorganization energies ( $\lambda_{\text{reg}}$ ) of the <sup>3</sup>MC states for each complex are larger than those of the <sup>3</sup>MLCT state (see Figures 7 and S3F2).

Because of the consistency between the DFT calculations and the experimental results, these complete results could be used as a database to support and explain the photophysical behavior of Ru-tpy complexes in Section 4.

### 4. DISCUSSION

4.1. Consideration of the T<sub>e</sub> (<sup>3</sup>MLCT Emitting State) and <sup>3</sup>MC (Emission Quenching State) States for Observing 77 K Low-Emission Quantum Yields. As demonstrated earlier in Figure 3, the unusual 77 K low quantum yields with  $h\nu_{ave(em)}$  and two exponential emission lifetimes for the  $[Ru(tpy)(AN)_3]^{2+}$  complex suggested the presence of an accessible <sup>3</sup>MLCT  $\rightarrow$  <sup>3</sup>MC internal conversion process. As reported by Tsai et al.,<sup>70</sup> the elongation of the Ru–N coordination distance from the <sup>3</sup>MLCT (T<sub>e</sub>) structure could lead to the formation of an <sup>3</sup>MC minimum structure. In Figure 8, the ic (<sup>3</sup>MLCT  $\rightarrow$  <sup>3</sup>MC<sub>b</sub>) energy barriers from <sup>3</sup>MLCT to the lowest <sup>3</sup>MC state (<sup>3</sup>MC<sub>b</sub>) for complexes (c) and (e) were approximated using the constrained optimization procedure, in

which the distances between Ru and AN were internally scanned by optimizing other degrees of geometric freedom. As shown in Figure 8, the ic (<sup>3</sup>MLCT  $\rightarrow$  <sup>3</sup>MC) energy barrier for [Ru(tpy)- $(AN)_3$ <sup>2+</sup> is about 750 cm<sup>-1</sup> (see Figure S3E1A-E1C for the details), and the corresponding barrier for [Ru(tpy)(bpm)-(AN)<sup>2+</sup> is about 1700 cm<sup>-1</sup>. The ic (<sup>3</sup>MLCT  $\rightarrow$  <sup>3</sup>MC<sub>b</sub>) energy barrier for [Ru(tpy)(bpm)(AN)]<sup>2+</sup> is about 8.5 times larger than  $4k_{\rm B}T$  (about 200 cm<sup>-1</sup> at 77 K) at 77 K, while the ratio of  $[Ru(tpy)(AN)_3]^{2+}$  is less than 4, suggesting a more accessible reaction channel. If ic  $({}^{3}MLCT \rightarrow {}^{3}MC_{b})$  occurs, this assumes that the T<sub>e</sub> (<sup>3</sup>MLCT) lifetime should have two different fits for the decay rate constant  $(k_{obs1} = k_r + k_{nr} \text{ and } k_{obs2} = k_{ic})$  and the quantum yield for  $[Ru(tpy)(AN)_3]^{2+}$  should decrease with the population of ic ( ${}^{3}MLCT \rightarrow {}^{3}MC_{b}$ ). The current theoretical characterization findings are consistent with the two exponential lifetime data and low quantum yields of  $h\nu_{ave(em)}$  for [Ru(tpy)- $(AN)_3$ <sup>2+</sup> at 77 K.

The  $[Ru(tpy)_2]^{2+}$  complex also had an optimized <sup>3</sup>MC state, but with a higher energy than T<sub>e</sub> (see Figure S3D2). This suggests that  $[Ru(tpy)_2]^{2+}$  would have an unfavorable or nonobservable ic  $({}^{3}MLCT \rightarrow {}^{3}MC_{b})$  at 77 K. Furthermore, the <sup>3</sup>MC excited state for [Ru(tpy)(NH<sub>3</sub>)<sub>3</sub>]<sup>2+</sup> and [Ru(tpy)- $(bpm)(NH_3)^{2+}$  could not be theoretically located using the same computational procedure. This is in good agreement with the experimental data for both complexes with their oneexponential decay lifetime for the  $k_{\rm obs}$  relaxation at 77 K. Due to the strong ligand stabilization effect of the amino ligands, the interplay of <sup>3</sup>MLCT and <sup>3</sup>MC states for complexes (b) and (d) can be schematically shown in Figure S3F1 (right), indicating a nonoptimizable <sup>3</sup>MC potential energy surface. Finally, according to these results, the Te relaxation processes for [Ru(tpy)- $(NH_3)_3^{2+}$  and  $[Ru(tpy)(bpm)(NH_3)]^{2+}$  do not have an ic  $({}^{3}MLCT \rightarrow {}^{3}MC)$  process at 77 K.



**Figure 8.** Transformation of the T<sub>e</sub> (<sup>3</sup>MLCT) state into an <sup>3</sup>MC<sub>b</sub> state for  $[\text{Ru}(\text{tpy})(\text{AN})_3]^{2+}$  (top) and  $[\text{Ru}(\text{tpy})(\text{bpm})(\text{AN})]^{2+}$  (lower). The Ru–N bond distances are elongated by *x* Å at each point starting from the T<sub>e</sub> state. (Ru–N(AN') for (c), x = 0.072; Ru–N(AN") for (c), x = 0.030; Ru–N(bpm) for (e), x = 0.018; and Ru–N(AN) for (e), x = 0.094).

### 5. CONCLUSIONS

A series of Ru-tpy complexes, including  $[Ru(tpy)_2]^{2+}$  (a),  $[Ru(tpy)(NH_3)_3]^{2+}$  (b),  $[Ru(tpy)(AN)_3]^{2+}$  (c),  $[Ru(tpy)-(bpm)(NH_3)]^{2+}$  (d), and  $[Ru(tpy)(bpm)(AN)]^{2+}$  (e), were synthesized and characterized. It is noteworthy that complexes (c) and (e) show low-emission quantum yields, a behavior that can be attributed to the presence of efficient ic (<sup>3</sup>MLCT  $\rightarrow$ <sup>3</sup>MC) pathways at 77 K. The low-lying <sup>3</sup>MC states for complexes (c) and (e) could favor PCT applications, while the low-lying <sup>3</sup>MLCT states for complexes (b) and (d) could favor PDT applications. The variation in the innocent ligands plays important roles in tuning the relative stability of the <sup>3</sup>MLCT or <sup>3</sup>MC minimum. Finally, this work demonstrates a complementary experimental and theoretical approach for the systematic characterization of these Ru-tpy complexes.

# ASSOCIATED CONTENT

# Supporting Information

The Supporting Information is available free of charge at https://pubs.acs.org/doi/10.1021/acsomega.3c01006.

S1, experimental section containing the materials (that were used in this work), syntheses, NMR spectra, and elemental analyses of the new Ru-tpy complexes, and the skeletal of reference and target complexes; S2, instrumentation section containing information on the instruments used, photophysical parameter, electrochemistry, low-temperature absorption, and emission spectra for the Ru-tpy complexes; S3, computational section containing the configurations of the HOMOs and LUMOs of TD-DFT-calculated low-energy  $S_n(S_0)$  excited states, orbital plots of these HOMOs and LUMOs at an S<sub>0</sub> geometry, comparisons of the RT-observed and TD-DFT-calculated absorption transitions, the TD-DFT parameters, NTO plots of  $S_n(S_0)$  and  $T_n$ , the orbital population the triplet transition energies, reorganization energies and the potential energy of the first triplet transition at each triplet geometry, and all of the spin densities of each triplet geometry for the Ru-tpy complexes; S3, detailed data for the energy gap of ic  $({}^{3}MLCT \rightarrow {}^{3}MC)$  for  $[Ru(tpy)(AN)_3]^{2+}(PDF)$ 

## AUTHOR INFORMATION

#### **Corresponding Authors**

- Chi-Wei Yin Department of Chemistry, Fu-Jen Catholic University, New Taipei City 24205, Taiwan, ROC;
  orcid.org/0000-0003-4972-5902; Phone: 886-2-2905-2491; Email: s26942421@gmail.com
- Ming-Kang Tsai Department of Chemistry, Fu-Jen Catholic University, New Taipei City 24205, Taiwan, ROC; Department of Chemistry, National Taiwan Normal University, Taipei 11677, Taiwan, ROC; orcid.org/0000-0001-9189-5572; Phone: 886-2-7734-6217; Email: mktsai@ntnu.edu.tw
- Yuan Jang Chen Department of Chemistry, Fu-Jen Catholic University, New Taipei City 24205, Taiwan, ROC;
  orcid.org/0000-0002-3134-8101; Phone: 886-2-2905-2490; Email: 054971@mail.fju.edu.tw

Complete contact information is available at: https://pubs.acs.org/10.1021/acsomega.3c01006

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#### Notes

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