



Article Initial Decomposition Mechanism of 3-Nitro-1,2,4-triazol-5-one (NTO) under Shock Loading: ReaxFF Parameterization and Molecular Dynamic Study

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Citation: Du, L.; Jin, S.; Nie, P.; She, C.; Wang, J. Initial Decomposition Mechanism of 3-Nitro-1,2,4-triazol-5-one (NTO) under Shock Loading: ReaxFF Parameterization and Molecular Dynamic Study. *Molecules* **2021**, *26*, 4808. https://doi.org/10.3390/ molecules26164808

Academic Editors: Aurora Costales and Fernando Cortés-Guzmán

Received: 16 June 2021 Accepted: 29 July 2021 Published: 9 August 2021

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Copyright: © 2021 by the authors. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (https:// creativecommons.org/licenses/by/ 4.0/). Abstract: We report a reactive molecular dynamic (ReaxFF-MD) study using the newly parameterized ReaxFF-lg reactive force field to explore the initial decomposition mechanism of 3-Nitro-1,2,4-triazol-5-one (NTO) under shock loading (shock velocity >6 km/s). The new ReaxFF-lg parameters were trained from massive quantum mechanics data and experimental values, especially including the bond dissociation curves, valence angle bending curves, dihedral angle torsion curves, and unimolecular decomposition paths of 3-Nitro-1,2,4-triazol-5-one (NTO), 1,3,5-Trinitro-1,3,5-triazine (RDX), and 1,1-Diamino-2,2-dinitroethylene (FOX-7). The simulation results were obtained by analyzing the ReaxFF dynamic trajectories, which predicted the most frequent chain reactions that occurred before NTO decomposition was the unimolecular NTO merged into clusters ($(C_2H_2O_3N_4)_n$). Then, the NTO dissociated from $(C_2H_2O_3N_4)_n$ and started to decompose. In addition, the paths of NO₂ elimination and skeleton heterocycle cleavage were considered as the dominant initial decomposition mechanisms of NTO. A small amount of NTO dissociation was triggered by the intermolecular hydrogen transfer, instead of the intramolecular one. For α -NTO, the calculated equation of state was in excellent agreement with the experimental data. Moreover, the discontinuity slope of the shockparticle velocity equation was presented at a shock velocity of 4 km/s. However, the slope of the shock-particle velocity equation for β -NTO showed no discontinuity in the shock wave velocity range of 3–11 km/s. These studies showed that MD by using a suitable ReaxFF-lg parameter set, could provided detailed atomistic information to explain the shock-induced complex reaction mechanisms of energetic materials. With the ReaxFF-MD coupling MSST method and a cheap computational cost, one could also obtain the deformation behaviors and equation of states for energetic materials under conditions of extreme pressure.

Keywords: 3-Nitro-1,2,4-triazol-5-one (NTO); ReaxFF-lg parameterization; reactive molecular dynamic; shock Hugoniot state; shock-induced initial decomposition mechanisms

1. Introduction

The decomposition mechanism of 3-Nitro-1,2,4-triazol-5-one (NTO) [1] under different conditions obviously influences the safety performance of NTO-based insensitive ammunition (IM) [2–5]. The thermal-induced decomposition kinetic process and burning characteristics of NTO [6–11] have been studied via various experimental techniques. In addition, the theoretical works concerning the electron structures and possible decomposition paths of NTO [12–14] were also reported. However, the atom-scale understanding about the shock-induced chemistry reaction of NTO still remain unclear.

Due to the progress of simulation technology, non-equilibrium ReaxFF reactive molecular dynamics [15–17] provided the high possibility to investigate the shock-induced

decomposition reaction of NTO at the molecular level. ReaxFF [15], developed by van Duin, is a bond-order-dependent reactive force field based on massive quantum mechanics (QM) calculation, which enables the accurate simulation of energetic materials under a realistic condition with a cheap computation. The shock-induced decomposition mechanisms of typical energetic materials (such as RDX [17], HMX [18], TATB [18], PETN [19], and TNT [20]) revealed by the ReaxFF molecular dynamics were in good agreement with the data from the experimental observations. To date, unfortunately, the poor transferability of the existing ReaxFF force fields [17,18,21,22] has led to no suitable ReaxFF potentials for simulating the thermal- and shock-induced chemistry reaction processes of NTO, hexanitrohexaazaisowurtzitane (CL-20), and dihydroxylammonium 5, 5'-bistetrazole-1, 1'-diolate (TKX-50). Therefore, developing a new ReaxFF parameter set to describe the physical and chemical properties of NTO, CL-20, and TKX-50 is a basic work for the exploration of the atom-scale decomposition chemistry of those energetic molecules.

In this study, we performed extensive QM calculations to parameterize a new ReaxFFlg force field for 3-Nitro-1,2,4-triazol-5-one (NTO). Then, we used the new ReaxFF-lg parameters to model the shock-induced decomposition process of NTO crystal. The shock Hugoniot states and kinetic decomposition mechanism in the shock velocity range of 6–11 km/s were studied to obtain a complete map of the shock-induced initial decomposition mechanisms of NTO under extreme conditions.

2. Method

2.1. ReaxFF Parameterization

The general ReaxFF potential functions partition the total energy into the contributions as:

$$E_{Reax} = E_{bond} + E_{lp} + E_{over} + E_{under} + E_{val} + E_{pen} + E_{coa} + E_{c2} + E_{triple} + E_{tors} + E_{conj} + E_{H-bond} + E_{vdW}$$
(1)
+ $E_{coulomb}$

The low-gradient model (E_{lg}) [21] is the long-range dispersion-correction for E_{Reax} to improve the descriptions of molecule crystalline:

$$E_{Reaxlg} = E_{Reax} + E_{lg} \tag{2}$$

The critical difference between ReaxFF and unreactive force fields is that the bond order (BO'_{ij}) decides the connectivity for two atoms, instead of a fixed bond. The BO'_{ij} (Equation (3)) is updated in every iteration and allows for the breaking and recreation of chemical bonds during the molecular dynamics. Therefore, all covalent interactions, such as E_{bond} and E_{val} , are expressed in the terms of the corrected BO'_{ij} .

$$BO'_{ij} = BO^{\sigma}_{ij} + BO^{\pi}_{ij} + BO^{\pi\pi}_{ij} = e^{[P_{bo1}(\frac{r_{ij}}{r_0^{\sigma}})^{p_{bo2}}]} + e^{[P_{bo3}(\frac{r_{ij}}{r_0^{\pi}})^{p_{bo4}}]} + e^{[P_{bo5}(\frac{r_{ij}}{r_0^{\pi\pi}})^{p_{bo6}}]}$$
(3)

To consider the Coulomb interactions between all atom pairs in the system, the atomic charges are updated using the Electron Equilibration Method (EEM) [23] approach at each iteration step. A similar method, called the QEq [24] charge derivation approach, is also suitable for calculating the atomic charge distribution during the dynamics. The atomic charge calculation in ReaxFF is the most time-consuming step, so the calculation cost of ReaxFF-MD is about two orders of magnitude higher than that of the classic non-reaction force field. However, the computation costs of ReaxFF are still far less than the ab-initio molecular dynamics.

The parameters of the ReaxFF are usually fitted from massive QM data and experimental values. In this study, the parametrization workflow of a new force field consisted of the following steps (Figure 1):



Figure 1. Workflow of the ReaxFF parameterization.

(1) Generation of the reference QM data (training set). In this study, the DFT methods with Grimme dispersion correction [25] were used to obtain the bond dissociation curves, valence angle bending curves, dihedral angle torsion curves, and unimolecular decomposition paths [13,26,27] of NTO, RDX, and FOX-7, respectively. Other important properties, such as the geometries (bond length, valence angles, and dihedral angles), interaction energy of dimers by hydrogen bonds and vdW interaction, dissociation potential energy surfaces of small molecules, cohesive energy, and crystal equations of state were also added to the training set. The B3LYP-D3BJ and M062X-D3 functional [28] were used to optimize the minimum points and transition state structures and predict the reaction energies. Moreover, the dimers' interaction energy and crystal equation of state were evaluated by the PWPB95-D3BJ approach [29] and PBEsol [30] with TS dispersion correction [31], respectively. Finally, the training set contained over 40 chemical reactions and over 1700 molecules for characterizing the atomic interactions under various environments. For detailed information about the training data, please see the Supplementary Materials.

(2) The weights (σ) assignment of the objective function between the ReaxFF prediction and QM data. The objective function has the form:

$$Error = \sum_{i=1}^{n} \left[\frac{x_{i,TS} - x_{i,ReaxFF}}{\sigma_i} \right]^2$$
(4)

where the sum runs over all the training set entries. Each difference between the reference property ($x_{i,TS}$) and the ReaxFF value ($x_{i,ReaxFF}$) is weighted individually via the σ_i . In this study, the properties of NTO were assigned smaller weights to highlight the importance during the force field optimization.

(3) The Covariance Matrix Adaption Evolution Strategy (CMA-ES) optimizer of AMS2020 [32], coupled with the RiPSOGM global optimization algorithm [33], was used to fit the new ReaxFF parameter set. At the beginning of the optimization, the ReaxFF parameters from reference [21] were chosen as a good initial force field.

(4) The crystal properties and reaction energy profiles were considered as the most important items for validating the new force field. Other properties, such as the enthalpies of sublimation and geometries, were also used for the evaluations.

2.2. Molecular Dynamic Simulations

The shock-induced decompositions of β -NTO at various shock wave velocities (6–11 km/s) were simulated by the multi-scale shock technique (MSST) [34], coupled with ReaxFF-MD. The single crystal of β -NTO, obtained from the Cambridge Crystallographic Data Centre (CCDC No. 166510), was expanded to a 6 × 6 × 6 supercell that used as the initial state of the 3D periodic box. In the simulation, a 1.5 ps isothermal isochoric (NVT ensemble, langevin thermostat, 50 fs damping parameter) MD was firstly conducted to relax the atomic positions of the initial state. Then, a shock wave along the lattice vector direction *a* was loaded to study the non-equilibrium evolution of the modelling system. For each simulation, the dynamic evolution time was 50 ps, and the integral step was set to 0.1 fs. The 0.3 bond's cutoff was used for the analysis of intermediates and products. All dynamic simulations were performed on the LAMMPS [35] package. For post-processing of the dynamic trajectory, the python scripts embedded in the MAPS [36] package were used.

3. Result and Discussion

3.1. Parameterization and Verification

3.1.1. New Reaxff-Lg Parameterization for NTO

Based on the massive reference QM data and experimental values, we obtained multiple ReaxFF-lg parameter sets of NTO from training. Then, one force field with the best performance was chosen as the new ReaxFF-lg potentials of NTO. This ReaxFF-lg parameters are presented in the Supplemental Materials.

3.1.2. Description of the Crystal Properties

To validate the chosen force field, we studied the equilibrium structure ($3 \times 3 \times 3$ supercell) of β -NTO by a 1 ps NVT-MD and subsequent 5 ps NPT-MD under standard conditions (298 k and 1 atm). The damping parameters were 50 fs for the thermostat and 500 fs for the barostat. After the 6 ps simulation, the calculated lattice constants were compared to the experimental data. An additional 25 ps NVT-MD under a lower temperature (5 k) was performed to calculate the heat of sublimation (ΔH_{sub}). In order to measure the differences between the equilibrium structures and initial geometries during this NVT-MD, we also calculated the average values of the root-mean-square deviations (RMSD) for the atomic position by considering all the dynamic trajectories. For comparison, the calculation results by MD using the original and new ReaxFF-lg parameters, are listed in Table 1.

Table 1. The lattice constants, density, heat of sublimation, and root-mean-square deviations of β -NTO.

Item		New ReaxFF-lg	Original ReaxFF-lg [21]	Experiment	
a		9.317	9.380	9.326	
b		5.445	5.482	5.450	
с		8.851	9.093	9.040	
α		90.0	90.0	90.0	
β		101.47	101.47	101.47	
γ		90.0	90.0	90.0	
density (g/cm ³)		1.926	1.885	1.918	
density error (%)		-0.417	-1.721	_	
ΔH_{sub} (kcal/mol)		25.62	39.89	23.42–28.65 [37]	
RMSD (Å)	crystal	0.204	0.660	-	
	molecule	0.063	0.247	-	

As shown in Table 1, the lattice constants and density were predicted by the original and new ReaxFF-lg force fields. Obviously, the results of the new ReaxFF-lg were closer to the experimental density. In addition, the error of ΔH_{sub} , predicted by the original ReaxFF-lg and obtained by the experiment, was more than 40% while it was reduced to a very low level using the new ReaxFF-lg. Moreover, the RMSD of β -NTO and single-molecule NTO,

calculated by the new ReaxFF-lg, showed lower values than that of the original ReaxFF-lg. This fact indicated that the inner molecule vibration of the crystal predicted by the new ReaxFF-lg MD was more similar to the original unit cell structure. Another noticeable phenomenon observed from the original ReaxFF-lg MD was that the hydrogen atoms in the nitrogen heterocycle ring were not on the same plane as the ring, exhibiting a bad reproduction of original single crystal structure. Fortunately, this phenomenon did not happen in the new ReaxFF-lg MD. For detailed information about this phenomenon, a few animations are provided in the Supplemental Materials for visualizing the differences.

For investigating the prediction ability of the new ReaxFF-lg force field on the compression properties of the NTO crystals, the equation of states of α -NTO were calculated by ReaxFF-MD, coupled with the MSST methods (mentioned in Section 2.2). The calculated data are drawn in Figure 2, accompanied by the experimental data [38]. More shock Hugoniot data are provided in the Supplemental Materials.



Figure 2. Comparisons of the compression equation of state for α -NTO.

From Figure 2, the calculated results from the new force field are in good agreement with the experimental values, which means that the new ReaxFF-lg parameters accurately reproduced the shock compression properties of the α -NTO crystals. While the equation of state (EOS) data of α -NTO were not included in the training data, the new ReaxFF-lg parameters still successfully predicted the equation of states for α -NTO between 0–25 GPa with an excellent accuracy, indicating the good transferability of the new force field. Therefore, the new ReaxFF-lg parameters were suitable for simulating the deformation behavior of the NTO crystals during shock loading. However, the results from the original ReaxFF-lg MD overestimated the pressure of the Hugonito state, indicating that the original ReaxFF-lg almost failed to model the compression EOS of the NTO crystals.

3.1.3. Description of the Potential Energy Surface

The hydrogen transfer reaction and nitro group dissociation were vital for triggering the decomposition of the NTO molecule. In this paper, the bond dissociation energy profile of N-H and C-NO₂ single bond prediction by the QM and ReaxFF methods were considered for verification of the new ReaxFF-lg (shown in Figures 3 and 4).

From Figure 3, the prediction energy curves for the dissociation of the hydrogen atom attached to the ring by the new ReaxFF-lg obtained a good consistency with the QM data, indicating that the new ReaxFF-lg force fields could well reproduce the energies of the proton dissociation process. While the original ReaxFF-lg could predict the dissociation energy of one hydrogen atom, it could not describe the energy increases for another hydrogen atom dissociation when the N-H bond length exceeded 1.5 Å.



Figure 3. The dissociation energies for the two N-H single bonds at the skeleton ring of NTO (gray—carbon atom; white—hydrogen atom; red—oxygen atom; blue—nitrogen atom).



Figure 4. The dissociation energies of the C-N single bond during NO₂ and HONO elimination (gray—carbon atom; white—hydrogen atom; red—oxygen atom; blue—nitrogen atom).

From Figure 4, the original ReaxFF-lg failed to describe the NO₂ and HONO dissociation reaction energy profiles, thus poorly describing the initiation decomposition step of the NTO molecule. On the contrary, the predicted energy curves of the C-N bond broken by the new ReaxFF-lg were more approximate to the QM curves, indicating that the new ReaxFF-lg significantly improved the prediction accuracy of the energy changes in the NO₂ and HONO dissociation. It was also noticeable that the prediction energy at the C-N bond distances of 1.8–2.1 Å for the HONO dissociation by the new ReaxFF-lg were between the energy curves of the single and triplet states. Thus, we thought the fitting errors were acceptable for the HONO dissociation, and the fitting curves also gave a qualitatively correct energy profile.

To verify whether the ReaxFF-lg force field can accurately predict the valence angle bending and dihedral angle rotation energy barriers in the NTO molecule, the potential energy surfaces of the N-C-N, C-N-H, H-C-N-O, and N-C-N-O angles were calculated using the ReaxFF-lg and QM methods. These energy curves are shown in Figures 5 and 6.

As shown in Figures 5 and 6, the barriers of the N-C-N bond angle distortion and dihedral angles rotation calculated by the original ReaxFF-lg showed unacceptable errors, indicating that the original force field was completely unable to describe the angle potential energy surfaces of NTO. However, the energy curves obtained from the new ReaxFF-lg showed a good reproduction of the QM data. Thus, the new ReaxFF-lg could accurately predict the energy barriers of the valence angle distortion and dihedral angle rotation for the NTO molecule.



Figure 5. Energy barriers for the distortion of the N-C-N and C-N-H bond angles in NTO (gray—carbon atom; white—hydrogen atom; red—oxygen atom; blue—nitrogen atom).



Figure 6. Rotational energy barriers for the H-C-N-O and N-C-N-O dihedral angles in NTO (gray—carbon atom; white—hydrogen atom; red—oxygen atom; blue—nitrogen atom).

3.2. Shock-Induced Chemistry of NTO

3.2.1. Shock Hugoniot State

For unreacted energetic materials, the shock wave velocity (U_s) and particle velocity (U_p) usually have the following linear relationship:

$$U_s = C_0 + SU_p \tag{5}$$

where C_0 and *S* are constants. Additionally, the pressures (*P*) of the shock compression are calculated by:

$$P = \rho_0 U_s U_p \tag{6}$$

where ρ_0 is the initial density.

Since the calculated equation of state for α -NTO in Section 2.2 was a verification of the new force field's good description ability of the shock compression properties of NTO crystals, we expanded the shock hugoniot data of α -NTO and calculated the U_s - U_p data of β -NTO. These values are listed in Table 2, and the pressure-volume curves are shown in Figure 7.

From Figure 7a, when the shock wave velocity is 4 km/s, there is an obvious inflection point at the curves of α -NTO, showing that the slope of equation 5 for α -NTO had different values after the inflection point. Engelke [39] reported that the energetic molecules containing aromatic rings showed discontinuities in the principal-shock Hugoniot. Our simulation results for α -NTO also showed the same characteristics. Thus, the data before and after the inflection point of α -NTO were fitted by the least squares method to obtain the corresponding C₀ and S. However, for β -NTO, the slope of equation 5 was no discontinuity

in the range of the shock wave velocity of 3–11 km/s, indicating that the particle velocity and the shock velocity maintain a good linear relationship.

$$\alpha$$
-NTO: $U_s = 1.5032 + 2.3005 U_p \ (2 \text{ km/s} \le U_s \le 4 \text{ km/s}, R^2 = 0.9912)$ (7)

$$\alpha$$
-NTO: $U_s = 2.8039 + 1.4136U_p (4 \text{ km/s} < U_s \le 11 \text{ km/s}, R^2 = 0.9966)$ (8)

$$\beta$$
-NTO: $U_s = 2.6449 + 1.4469 U_p (3 \text{ km/s} \le U_s \le 11 \text{ km/s}, R^2 = 0.9940)$ (9)

Table 2. The shock Hugoniot state of α -NTO and β -NTO.

U _s (km/s) —	U_p (km/s)		V/V ₀		P (GPa)	
	α-ΝΤΟ	β-ΝΤΟ	α-NTO	β-ΝΤΟ	α-NTO	β-ΝΤΟ
2	0.24	0.19	0.88	0.91	0.92	0.73
3	0.6	0.51	0.80	0.83	3.45	2.93
4	1.09	0.95	0.73	0.76	8.35	7.29
5	1.59	1.58	0.68	0.68	15.23	15.16
6	2.22	2.20	0.63	0.63	25.51	25.32
7	2.90	2.85	0.59	0.59	38.88	38.27
8	3.56	3.54	0.56	0.56	54.55	54.32
9	4.27	4.34	0.53	0.52	73.61	74.92
10	5.17	5.18	0.48	0.48	99.03	99.35
11	5.89	5.93	0.46	0.49	124.15	125.11

Note: V/V_0 is the relative volume.



Figure 7. Shock Huoniot state data of α -NTO and β -NTO. (**a**) shock velocity and particle velocity relation; (**b**) pressure via relative volume along Hugoniot line.

3.2.2. Initial Decomposition Mechanism

To understand the ideal detonation behavior of NTO, we mainly consider the evolution and reaction mechanism of NTO under strong shock waves (shock velocity >6 km/s). The massive reactions occurred during the ReaxFF-MD; therefore, we pointed out the cluster formation and initial decomposition paths of NTO, while neglecting the final product formation mechanism.

(1) Cluster formation: the aggregation of the NTO molecules was caused by the strong compression during the shock wave front propagating across the NTO crystal. This was manifested as the NTO molecules merging into dimers, trimers, and even larger $(C_2H_2O_3N_4)_n$. Figure 8 illustrates the formation process of the NTO clusters.

$$C_2H_2O_3N_4 \longrightarrow C_4H_4O_6N_8 \longrightarrow C_6H_6O_9N_{12} \longrightarrow C_8H_8O_{12}N_{16} \longrightarrow C_{10}H_{10}O_{15}N_{20} \longrightarrow C_{12}H_{12}O_{18}N_{24}$$
(1)





$$C_{4}H_{4}O_{6}N_{8} + C_{6}H_{6}O_{9}N_{12} \longrightarrow C_{10}H_{10}O_{15}N_{20} \quad (2)$$

$$C_{4}H_{4}O_{6}N_{8} + C_{8}H_{8}O_{12}N_{16} \longrightarrow C_{12}H_{12}O_{18}N_{24} \quad (3)$$

$$2C_{4}H_{4}O_{6}N_{8} + 7C_{2}H_{2}O_{3}N_{4} \longrightarrow C_{20}H_{20}O_{30}N_{40} + C_{2}H_{2}O_{3}N_{4} \quad (4)$$

$$2C_{4}H_{4}O_{6}N_{8} + 7C_{2}H_{2}O_{3}N_{4} \longrightarrow C_{22}H_{22}O_{33}N_{44} \quad (5)$$

Figure 8. $(C_2H_2O_3N_4)_n$ formation.

The chemical reaction (1), listed in Figure 8, was the common chain reaction obtained from all simulations, of which the NTO dimer formation was the most frequently and earliest occurring reaction. From ReaxFF-MD, the mechanism of $(C_2H_2O_3N_4)_n$ formation was that the intermolecular NH . . . O hydrogen bonds were broken by shock compression, and then the oxygen atom of the nitro group attacked the hydrogen atom of the heterocyclic ring to produce the O-H single bond. Thus, the hydrogen atom was the bridge to form NTO clusters in the early evolution. In addition, $(C_2H_2O_3N_4)_n$ underwent a direct polymerization reaction (Reactions (4) and (5) in Figure 8) to form larger clusters when the shock wave velocity was higher than the Chapman-Jouguet detonation velocity of NTO (~8 km/s). Furthermore, the higher the shock wave velocity, the higher the degree of NTO polymerization.

(2) The initial decomposition paths of NTO: an interesting phenomenon was that almost all reactants were forming into $(C_2H_2O_3N_4)_n$ clusters. After that, the unimolecular NTO was dissociated from the $(C_2H_2O_3N_4)_n$ clusters and started to decompose along the paths shown in Figure 9. In addition, the Gibbs free energy changes of the reaction at high temperature and pressures (calculated by M062X-D3/def2-tzvp) are marked in Figure 9.



Figure 9. The initial decomposition paths of NTO.

From Figure 9, there were four decomposition paths for unimolecular NTO. The NO_2 elimination (path A) was found to be the early and frequently decomposition path, followed by heterocycle cleavage, thus forming the unstable intermediate (CH₂ON₂) and radicals CN. In addition, heterocycle cleavage (path B) and hydrogen transfer (paths C and D) both took place later and less frequently than path A because of their relative higher Gibbs free energy barriers. However, the occurring frequency of NTO decomposition along path B was far more than the intermolecular hydrogen transfer reactions.

Path A showed that NTO was first broken into the NO₂ and C₂H₂ON₃ ring, where the ring was split into CN and a three-atom nitrogen heterocycle (CH_2ON_2). Then, the $C_2H_2ON_2$ with a high activity could spontaneously undergo an intermolecular proton transfer to form CH₃ON₂ and CHON₂, and CHON₂ converted into CO and N₂. The CH₂ON₂ was also the product of the first decomposition step along path B and was accompanied by the O_2 N-CN formation. The intermediate O_2 N-CN could not only form NO₂ and CN by breaking the C-N bond, but also dissociate a nitrogen atom and rearrange to form O=CNO. The latter one spontaneously decomposed into CO and NO. The decomposition along path C was triggered by the hydrogen transfer towards carbonyl. The stable product C₂H₂O₃N₄ of path C broke the N-N bond of the heterocycle, then NO₂ elimination happened immediately, thus forming the $C_2H_3ON_3$. The decomposition paths A, B, and C all produced the same molecule, CH₃ON₂, and the C-NH₂ bond of this molecule easily lost its connection to produce OCNH and NH_2 . The reaction along path D was another different hydrogen transfer situation. It should have been the case that the $-NO_2$ group captured surrounding protons due to its strong electronegativity. However, the unstable intermediate, C₂H₂ON₃-HONO, automatically dissociated into HONO and C₂H₂ON₃, undergoing decomposition along path A at last.

Wang et al. explored the unimolecular decomposition mechanism of NTO by firstprinciple calculation [13]. They proposed that the hydrogen transfer at the nitrogen heterocycle and NO₂ dissociation were the initial decomposition steps of single NTO. However, our study did not find the intramolecular hydrogen transfer during MD; instead, we found the intermolecular hydrogen transfer paths (path C and D in Figure 9). Moreover, we found a new decomposition path: skeleton heterocycle cleavage (path B in Figure 9) under shock. On the other hand, the calculation results of Wang et al. [13] indicated that the energy barriers of hydrogen transfer were about 20 kcal/mol lower than that of NO₂ elimination. The opposite results that we acquired showed that the high temperature and high stress inside the NTO crystal significantly reduced the NO₂ elimination energy barrier and the NTO intramolecular stability. The spontaneous polymerization of NTO under a strong shock wave (shown in Figure 8) also inhibited the intramolecular hydrogen transfer reaction. Because of the contribution of the two parts, the NO₂ elimination and heterocycle cleavage became the primary decomposition pathways of NTO under shock wave.

4. Conclusions

A new ReaxFF-lg force field was customized to describe the properties of 3-Nitro-1,2,4triazol-5-one (NTO). This force field was trained from massive quantum mechanics data and experimental values, especially including the bond dissociation curves, valence angle bending curves, dihedral angle torsion curves, and unimolecular decomposition paths of NTO. Through verification, the new ReaxFF-lg parametrization was proved to be a suitable reactive force field for modelling the physical and chemistry properties of NTO.

Reactive molecular dynamic (ReaxFF-MD) simulations with the new ReaxFF-lg force field were performed to explore the initial decomposition mechanism of NTO under shock (shock velocity >6 km/s). These simulations showed that the most frequent chain reactions occurred before NTO decomposition was the single NTO molecules merged into dimers, trimers, and even large ($C_2H_2O_3N_4$)_n clusters.

Then, the NTO single molecule started to decompose after dissociating from the $(C_2H_2O_3N_4)_n$ clusters. The paths of NO₂ elimination and skeleton heterocycle cleavage were considered as the dominant initial decomposition mechanisms of NTO. A small amount of NTO decomposition was triggered by the intermolecular hydrogen transfer, instead of the intramolecular one. Moreover, the NTO crystals in α and β phase showed similar compression properties due to the similar molecular arrangement and smaller density differences between the two crystals.

These studies showed that MD, using a suitable ReaxFF-lg potential, provided detailed atomistic information to explain the shock-induced complex reaction mechanisms of energetic materials. By the ReaxFF-MD coupling MSST method, the deformation behaviors and equation of states of energetic materials under conditions of extreme pressure could be obtained with a cheaper computational cost.

Supplementary Materials: The following are available online. The Supplementary Materials produced the detailed information about the methods used in this study. The full training-data files can be requested by contacting the author. Figure S1: N-H bond distance during NVT-MD.; Figure S2: O=C-N-H bond distance during NVT-MD.; Table S1: Shock Hugoniot data of α -NTO.

Author Contributions: Data curation, P.N. and C.S.; Investigation, L.D.; Methodology, L.D.; Validation, S.J.; P.N.; C.S. and J.W.; Writing—original draft, L.D.; Writing—review and editing, L.D. and J.W. All authors have read and agreed to the published version of the manuscript.

Funding: The research was funded by the Fundamental Research Funds for the National Basic Research Program of China, giant number 2019-JCJQ-ZD-139-00.

Institutional Review Board Statement: Not applicable.

Informed Consent Statement: Not applicable.

Data Availability Statement: The full training-data files can be requested by contacting the author. Two email address are provided here: 3120185542@bit.edu.cn and dulixiaosong@163.com.

Acknowledgments: The authors would like to thank the Fundamental Research Funds for the National Basic Research Program of China (2019-JCJQ-ZD-139-00). The author also thanks David Furman for providing the flocky package.

Conflicts of Interest: The authors declare no conflict of interest.

Sample Availability: Samples of the compounds used in this study are available from the authors.

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