



Article The Role of Quadruple Bonding in the Electron Transport through a Dimolybdenum Tetraacetate Molecule

Dmitry O. Arentov, Maxim R. Ryzhikov 🕒 and Svetlana G. Kozlova *🕒

Nikolaev Institute of Inorganic Chemistry SB RAS, Acad. Lavrentiev Ave. 3, 630090 Novosibirsk, Russia

* Correspondence: sgk@niic.nsc.ru

Abstract: A dimolybdenum tetraacetate ($Mo_2(O_2CCH_3)_4$) molecule is embedded between two electrodes formed by semi-infinite 1D monatomic chains of lithium, aluminum, and titanium atoms. Electron transport through the $Mo_2(O_2CCH_3)_4$ molecule is calculated. The role of quadrupole bonding in the transport properties of the studied systems is analyzed.

Keywords: DFT; NEGF; electron transport; Mo₂(O₂CCH₃)₄; d-orbitals

1. Introduction

Dimolybdenum tetraacetate Mo₂(O₂CCH₃)₄ (DMT) is a molecule whose electronic structure is characterized by a quadruple Mo-Mo bonding with the $\sigma^2 \pi^4 \delta^2$ configuration where eight electrons are 4d valence electrons of two molybdenum atoms [1–5]. Theoretical chemistry pays particular attention to the properties of δ bonding, which is weaker than σ and π bonds but nonetheless affects the electronic absorption spectra and the spatial structure of the molecules with quadruple bonds [6,7]. Since DMT molecules can form ordered structures on various surfaces, they are currently considered as promising molecular units for the design of nanoscale electronic devices [8–10]. The role of σ , π , and δ bonds in the electron transport through a single molecule can also be discussed, since these bonds are formed by different overlaps of d-orbital lobes.

Electron transport through a single molecule is commonly described using a method based on the nonequilibrium Green's function (NEGF) formalism [11–13]. The NEGF technique allows calculating the transmission function T(E) showing the sum probability for an electron with energy E to pass from the source to the drain through channels represented by the molecular orbitals (MOs) of the electrode—molecule—electrode system. T(E) depends both on the electronic structure of the electrodes and the molecule and on the interactions between them. In the case of weak contacts between the molecule and the electrodes, T(E) reflects the structure of energy levels of the MOs of the isolated molecule, and its shape is represented by a series of peaks (resonances) in the vicinity of MO energies [14]. Therefore, we assume that varying the distance between the contacts and the molecule that can be used to analyze the participation of σ , π , δ bonds in the electron transport.

In this work, we consider electron transport through a DMT molecule supplemented by two M atoms on both sides and embedded into a system of semi-infinite electrodes. The electrodes were represented by 1D chains of four M atoms (M = Li, Al, Ti) in the left (source) and right (drain) parts of the system, respectively. These s-, p-, and d-block metals are convenient models to study Fermi level shifts (E_F). The 1D chain is the simplest model of electron container needed to calculate transmission function in a molecular system. The central part of the system was represented in two ways (Figure 1): by four M atoms for monoatomic wires (M_6 – M_6) and by a DMT molecule embedded between M atoms of the monoatomic wires (M_6 –DMT– M_6). System M_6 – M_6 was used as a reference when studying transmission functions in systems (M_6 –DMT– M_6).



Citation: Arentov, D.O.; Ryzhikov, M.R.; Kozlova, S.G. The Role of Quadruple Bonding in the Electron Transport through a Dimolybdenum Tetraacetate Molecule. *Molecules* **2022**, 27, 6912. https://doi.org/10.3390/ molecules27206912

Academic Editor: Federico Totti

Received: 30 September 2022 Accepted: 13 October 2022 Published: 14 October 2022

Publisher's Note: MDPI stays neutral with regard to jurisdictional claims in published maps and institutional affiliations.



Copyright: © 2022 by the authors. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (https:// creativecommons.org/licenses/by/ 4.0/).



Figure 1. Model systems M_6 – M_6 (**a**) and M_6 –DMT– M_6 (**b**) (M = Li, Al, Ti).

2. Results and Discussion

2.1. Quadruple Bond in Mo₂(O₂CCH₃)₄ Molecule

The optimized interatomic distances of the DMT molecule are Mo-Mo = 2.134 Å, Mo-O = 2.111 Å, and \angle Mo-Mo-O = 91.62° (see also Tables S1 and S2). These parameters agree with the corresponding experimental values Mo-Mo = 2.093 Å, <Mo-O> = 2.119 Å, and < \angle Mo-Mo-O> = 91.8° [1]. Figure 2 shows MOs characterized by bonding (σ , π , δ) and antibonding (σ^* , π^* , δ^*) interactions of Mo 4d orbitals, in agreement with the calculation results reported previously for the electronic structure of the DMT molecule [4]. Therefore, we assume that these MOs can be used to estimate the role of σ/σ^* , π/π^* and δ/δ^* interactions in the electron transport through the molecule. The calculated HOMO—LUMO gap is 2.002 eV, in good agreement with the experimental value ~2.69 eV [7].



Figure 2. MOs (isosurface: 0.05 au) demonstrating bonding (σ , π , δ) and antibonding (σ^* , π^* , δ^*) interactions of Mo 4d orbitals for the optimized DMT structure. The digits in parentheses are the MO energies in eV.

2.2. Electron Transport through Systems M_6 -($Mo_2(O_2CCH_3)_4$)- M_6 (M = Li, Al, Ti)

The M-Mo distances were varied from the sums of their empirical atomic radii (d(Li-Mo) = 2.90 Å, d(Al-Mo) = 2.70 Å, d(Ti-Mo) = 2.85 Å [15]) to 6–7 Å with 0.05 Å step. It was established that the peaks of transmission functions ($T_{DMT}(E)$) of the M₆–DMT–M₆ system change their shape and energy positions. The peaks of $T_{DMT}(E)$ vanish above 6–7 Å.

The $T_{DMT}(E)$ at the most indicative distances (the shortest, the distance where the peaks narrowing sharply, the longest, and points between them) are given at SI (Figures S1–S3).

2.2.1. System Li₆-(Mo₂(O₂CCH₃)₄)-Li₆

Function $T_{wire}(E)$ for the Li wire has three plateaus (Figure 3). The embedding of a DMT molecule into a wire of lithium atoms decreases the $T_{DMT}(E)$ values compared to the $T_{wire}(E)$ values. The $T_{DMT}(E)$ for d(Li-Mo) = 2.90 Å has three humps in the energy range where function $T_{wire}(E)$ for the Li wire has three plateaus. The peaks of these humps are located at -2.45 eV, -0.92 eV, and 0.84 eV energies. Function $T_{DMT}(E)$ begins narrowing sharply at the d(Li-Mo) distance equal to 3.10 Å (Figure S1). The $T_{DMT}(E)$ function for d(Li-Mo) = 3.10 Å has peaks in the same regions at -3.43 eV, -3.28 eV, -2.48 eV, -0.97 eV, and -0.62 eV as well as some peaks above 0 eV.



Figure 3. Transmission functions $T_{DMT}(E)$ (green for d(Li-Mo) = 2.90 Å and blue for d(Li-Mo) = 3.10 Å) for the Li₆–DMT–Li₆ system and transmission function $T_{wire}(E)$ (brown) for the Li₆-Li₆ system. E_F denotes the Fermi level of electrodes. Energy positions of bonding (σ , π , δ) and antibonding (σ^* , π^* , δ^*) MOs are shown by vertical lines (black dots).

The LUMO and LUMO+5 (corresponding to δ^* and σ^* interactions, respectively) of the DMT molecule are the only orbitals that fall into non-zero value regions of both functions $T_{wire}(E)$ and $T_{DMT}(E)$. On the one side, the LUMO+1 and LUMO+2 are close to the dip (-2; -1.5 eV) of function $T_{wire}(E)$. On the other, these two orbitals are close to the $T_{DMT}(-1.5eV)$ peak at d(Li-Mo) = 3.1 Å. Thus, the participation of LUMO+1 and LUMO+2 (corresponding to two π^* interactions) in electron transport is ambiguous. The HOMO, HOMO-1, HOMO-2, and HOMO-3 are located beyond the transmission bands and apparently do not participate in conduction. As can be seen, the 1D chain can't be considered as a simple bulk electron container, since the transmission function of unperturbed by molecule 1D chain have some structure with humps, dips and plateau [16]. Thus, the DMT can provide channels for electrons only in the region where the transmission of the reference 1D chain of Li is non-zero.

2.2.2. System Al₆-(Mo₂(O₂CCH₃)₄)-Al₆

The $T_{DMT}(E)$ function of the Al₆–DMT–Al₆ system for d(Al-Mo) = 2.70 Å has close to unity values over the largest part of the calculated energy range (Figure 4). The only exception is the energy range from -2.7 eV to -0.27 eV, where $T_{DMT}(E)$ is close to 3.

Function $T_{DMT}(E)$ deviates from $T_{wire}(E)$ in the energy region from -7.8 eV to -6.9 eV and above the Fermi level (-3.7 eV) up to 0.3 eV (Figure 4). The transmission function became more resolved at d(Al-Mo) = 4.30 Å, still keeping most of features of short distance $T_{DMT}(E)$. Function $T_{DMT}(E)$ for d(Al-Mo) = 4.30 Å is represented by peaks at -6.69 eV, -5.87 eV, -1.38 eV, and -0.36 eV. The DMT molecule suppresses electron transport compared to the reference Al wire, similar to the system with Li.



Figure 4. Transmission functions $T_{DMT}(E)$ (green for d(Al-Mo) = 2.70 Å and blue for d(Al-Mo) = 4.30 Å) for the Al₆–DMT–Al₆ system and transmission function $T_{wire}(E)$ (brown) for the Al₆-Al₆ system. E_F denotes the Fermi level of electrodes. Energy positions of bonding (σ , π , δ) and antibonding (σ^* , π^* , δ^*) MOs are shown by vertical lines (black dots).

The energy regions of functions $T_{wire}(E)$ and $T_{DMT}(E)$ (for d(Al-Mo) = 2.70 Å) contain DMT MOs exhibiting bonding (σ , π , δ) and antibonding (σ^* , π^* , δ^*) interactions. At d(Al-Mo) = 4.30 Å, HOMO-3 (σ -bond), HOMO-2 and HOMO-1 (two π -bonds), LUMO+1 and LUMO+2 (two π^* -interactions), LUMO+5 (σ^* -interaction) can contribute to $T_{DMT}(E)$. The δ/δ^* -interactions (HOMO and LUMO) don't participate in transmission at this distance.

2.2.3. System Ti₆-(Mo₂(O₂CCH₃)₄)-Ti₆

Function $T_{DMT}(E)$ is smaller than $T_{wire}(E)$ over the largest part of the calculated energy range, meaning that the embedded molecule significantly affects the electron transport through the Ti wire (Figure 5). The energy ranges of functions $T_{wire}(E)$ and $T_{DMT}(E)$ (for d(Ti-Mo) = 2.85 Å) contain HOMO, LUMO, LUMO+1, LUMO+2, and LUMO+5. For d(Ti-Mo) = 4.80 Å, the region of non-negligible $T_{DMT}(E)$ values contains only LUMO+1 and LUMO+2 (two π^* -interactions).



Figure 5. Transmission functions $T_{DMT}(E)$ (green for d(Ti-Mo) = 2.85 Å and blue for d(Ti-Mo) = 4.80 Å) for the Ti₆–DMT–Ti₆ system and transmission function $T_{wire}(E)$ (brown) for the Ti₆-Ti₆ system. E_F denotes the Fermi level of electrodes. Energy positions of bonding (σ , π , δ) and antibonding (σ^* , π^* , δ^*) MOs are shown by vertical lines (black dots).

2.2.4. Discussion

The Mo-Mo distance in DMT (2.13 Å) is smaller than interatomic distances in 1D chains (2.86 Å–3.04 Å). The DMT molecule can be viewed as a defect disturbing the electronic structure of the monatomic wire. In addition, the embedded DMT molecule can be considered as a kind of filter that blocks some of the electron transport channels originally present in the 1D monatomic chains. This idea can explain the fact that the transmission functions of systems with an embedded DMT molecule are smaller than those of the wires and that not all MOs exhibiting bonding (σ , π , δ) and antibonding (σ ^{*}, π ^{*}, δ ^{*}) interactions of Mo 4d orbitals can participate in the electron transport through the DMT molecule.

3. Theoretical Calculations

The geometry of the dimolybdenum tetraacetate (DMT) molecule was optimized by the ADF engine of the AMS2020 package [17,18] using the BP86 density functional [19] and the TZP all-electron basis set for all atoms [20]. All of the calculated vibrational frequencies were real, indicating that the optimized structure of the DMT molecule corresponds to the energy minimum (Table S1). Scalar relativistic effects were considered within the ZORA approach [21].

The electron transport was studied using the nonequilibrium Green's function (NEGF) formalism [11–13] as implemented in the BAND engine of the AMS2020 package [22,23], the BP86 density functional [24–26], and the ZORA [27] approach for scalar relativistic effects.

To reproduce the electronic properties of bulk metals most closely, the interatomic distances in lithium (d(Li-Li) = 3.040 Å), aluminum (d(Al-Al) = 2.864 Å), and titanium (d(Ti-Ti) = 2.951 Å) chains were chosen equal to the interatomic distances in the corresponding crystals [28–30]. The atoms of the studied systems were calculated using the TZP basis set [18] with the following frozen cores: Li(1s), Al(1s.2p), Ti(1s.2s.2p.3s.3p), Mo(1s.2s.2p.3s.3p.3d.4s.4p), O(1s), C(1s).

The GGA-BP86 density functional was chosen due to its time cost advantages in T(E) calculations [31]. For efficiency reasons, the radial part of the basis set functions is

multiplied by a Fermi—Dirac function controlled by two main parameters: distance from the center R = 7 Å and decay width Δ = 0.7 [23]. Function T(E) was calculated using the self-consistent field method in the energy range $-5 \text{ eV} < \text{E} - \text{E}_{\text{F}} < 5 \text{ eV}$ at a temperature of 316 K (corresponding to the kT energy of 0.001 Hartree).

4. Conclusions

It was shown on the example of M₆–DMT–M₆ systems (M = Li, Al, Ti) that an embedded DMT molecule can be considered as a defect decreasing the transmission function. It was revealed that the transmission function depends on d(Mo-M) distances. It was shown that increasing the d(Mo-M) values affects not only the shape of the transmission function but also the number of DMT MOs involved in the electron transport through the molecule. No dependence of the transmission function on different overlaps of d-orbital lobes of Mo was revealed. For M₆–DMT–M₆ (M = Al, Ti) systems, the highest values of transmission functions were found in the region of LUMO+1 and LUMO+2 characterized by the π^* interaction. However, these values can be explained not by the overlaps of d-orbital lobes but by the presence of two such MOs.

Although the results of this work may seem purely theoretical, they provide a first insight into the role of overlapping d-orbital lobes in the electron transport through a molecule. The relevance of the present study is substantiated by the success that has been achieved in experimental works devoted to the fabrication of 1D monatomic chains of metals [32]. The greatest progress in this field was achieved for gold wires [33]. For many other metals, nanowires with a diameter smaller than 20 nm were obtained, and the work continues to reduce their thickness still further [34]. In this regard, the transition to monatomic wires can be viewed as the next step towards the design of single molecule devices.

Supplementary Materials: The following supporting information can be downloaded at: https://www.mdpi.com/article/10.3390/molecules27206912/s1, Figure S1: Transmission functions in systems L_{i_6} -($Mo_2(O_2CCH_3)_4$)- L_{i_6} depending on distances d(Li-Mo); Figure S2: Transmission functions in systems A_{l_6} -($Mo_2(O_2CCH_3)_4$)- A_{l_6} depending on distances d(Al-Mo). Figure S3: Transmission functions in systems Ti₆-($Mo_2(O_2CCH_3)_4$)- A_{l_6} depending on distances d(Al-Mo). Table S1: Optimized coordinates (Å) of $Mo_2(O_2CCH_3)_4$)- T_{l_6} depending on distances d(Ti-Mo). Table S1: Optimized coordinates (Å) of $Mo_2(O_2CCH_3)_4$ molecule (BP86/TZP level); Table S2: Frequencies and their intensities of $Mo_2(O_2CCH_3)_4$ molecule.

Author Contributions: D.O.A.—Investigation, Visualization; M.R.R.—Writing—review & editing, methodology; S.G.K.—Writing—review & editing, conceptualization, supervision. All authors have read and agreed to the published version of the manuscript.

Funding: The reported study was funded by RSF project N 22-23-00245.

Institutional Review Board Statement: Not applicable.

Informed Consent Statement: Not applicable.

Data Availability Statement: The data presented in this study are available on request from the corresponding author.

Conflicts of Interest: There are no conflicts of interest to declare.

References

- Cotton, F.A.; Mester, Z.C.; Webb, T.R. Dimolybdenum tetraacetate. Acta Crystallogr. Sect. B Struct. Crystallogr. Cryst. Chem. 1974, 30, 2768–2770. [CrossRef]
- 2. Kelley, M.H.; Fink, M. The molecular structure of dimolybdenum tetra-acetate. J. Chem. Phys. 1982, 76, 1407–1416. [CrossRef]
- 3. Benard, M. A study of Hartree–Fock instabilities in Cr₂(O₂CH)₄ and Mo₂(O₂CH)₄. J. Chem. Phys. **1979**, 71, 2546–2556. [CrossRef]
- Blaudeau, J.P.; Pitzer, R.M. Ab initio studies of ligand effects on the metal-metal bond in dimolybdenum complexes. J. Phys. Chem. 1994, 98, 4575–4579. [CrossRef]
- 5. Cotton, F.A.; Murillo, C.A.; Walton, R.A. *Multiple Bonds between Metal Atoms*; Springer Science & Business Media: Berlin, Germany, 2005.
- Cotton, F.A. Spectroscopic and quantum theoretical studies of species with metal-to-metal bonds. J. Mol. Struct. 1980, 59, 97–108. [CrossRef]

- 7. Martin, D.S.; Newman, R.A.; Fanwick, P.E. Polarized electronic absorption spectra for dimolybdenum(II) tetraacetate. *Inorg. Chem.* **1979**, *18*, 2511–2520. [CrossRef]
- Aeschlimann, S.; Bauer, S.V.; Vogtland, M.; Stadtmüller, B.; Aeschlimann, M.; Floris, A.; Bechstein, R.; Kühnle, A. Creating a regular array of metal-complexing molecules on an insulator surface at room temperature. *Nat. Commun.* 2020, 11, 6424. [CrossRef]
- Aeschlimann, S.; Lyu, L.; Stadtmüller, B.; Aeschlimann, M.; Kühnle, A. Tailoring molecular island shapes: Influence of microscopic interaction on mesostructured. *Nano Res.* 2020, 13, 843–852. [CrossRef]
- Kollamana, J.; Wei, Z.; Lyu, L.; Zimmer, M.; Dietrich, F.; Eul, T.; Stöckl, J.; Maniraj, M.; Ponzoni, S.; Cinchetti, M.; et al. Control of Cooperativity through a Reversible Structural Phase Transition in MoMo-Methyl/Cu(111). *Adv. Funct. Mater.* 2017, 28, 1703544. [CrossRef]
- 11. Datta, S. Quantum Transport: Atom to Transistor; Cambridge University Press: Cambridge, UK, 2005; pp. 33–50.
- 12. Verzijl, C.J.O.; Thijssen, J.M. DFT-based molecular transport implementation in ADF/BAND. J. Phys. Chem. C 2012, 116, 24393–24412. [CrossRef]
- Li, R.; Zhang, J.; Hou, S.; Qian, Z.; Shen, Z.; Zhao, X.; Xue, Z. A corrected NEGF+DFT approach for calculating electronic transport through molecular devices: Filling bound states and patching the non-equilibrium integration. *Chem. Phys.* 2007, 336, 127–135. [CrossRef]
- 14. Lambert, C.J. Basic concepts of quantum interference and electron transport in single-molecule electronics. *Chem. Soc. Rev.* 2015, 44, 875–888. [CrossRef] [PubMed]
- 15. Slater, J.C. Atomic radii in crystals. J. Chem. Phys. 1964, 41, 3199–3204. [CrossRef]
- Verzijl, C.J.O.; Seldenthuis, J.S.; Thijssen, J.M.; Verzijl, C.J.O.; Seldenthuis, J.S.; Thijssen, J.M.; Verzijl, C.J.O.; Seldenthuis, J.S.; Thijssen, J.M.; Verzijl, C.J.O.; et al. Applicability of the wide-band limit in DFT-based molecular transport calculations. *J. Chem. Phys.* 2013, 138, 094102. [CrossRef]
- 17. Velde, G.T.; Bickelhaupt, F.M.; Baerends, E.J.; Fonseca Guerra, C.; van Gisbergen, S.J.A.; Snijders, J.G.; Ziegler, T. Chemistry with ADF. J. Comp. Chem. 2001, 22, 931–967. [CrossRef]
- ADF 2020.102, SCM, Theoretical Chemistry, Vrije Universiteit, Amsterdam, The Netherlands. Available online: http://www.scm. com (accessed on 30 September 2022).
- 19. Perdew, J.P.; Burke, K.; Ernzerhof, M. Generalized gradient approximation made simple. *Phys. Rev. Lett.* **1996**, *77*, 3865. [CrossRef] [PubMed]
- van Lenthe, E.; Baerends, E.J. Optimized Slater-type basis sets for the elements 1–118. J. Comput. Chem. 2003, 24, 1142–1156. [CrossRef] [PubMed]
- 21. van Lenthe, E.; Ehlers, A.; Baerends, E.-J. Geometry optimizations in the zero order regular approximation for relativistic effects. *J. Chem. Phys.* **1999**, *110*, 8943–8953. [CrossRef]
- 22. Velde, G.T.; Baerends, E.J. Precise density-functional method for periodic structures. Phys. Rev. B 1991, 44, 7888–7903. [CrossRef]
- 23. BAND 2020.102, SCM, Theoretical Chemistry, Vrije Universiteit, Amsterdam, The Netherlands. Available online: https://www.scm.com (accessed on 10 October 2022).
- 24. Becke, A.D. Density-functional exchange-energy approximation with correct asymptotic behavior. *Phys. Rev. A* **1988**, *38*, 3098. [CrossRef]
- Perdew, J.P. Density-functional approximation for the correlation energy of the inhomogeneous electron gas. *Phys. Rev. B* 1986, 33, 8822–8824. [CrossRef] [PubMed]
- 26. Perdew, J.P. Erratum: Density-functional approximation for the correlation energy of the inhomogeneous electron gas. *Phys. Rev.* B **1986**, 34, 7406. [CrossRef]
- 27. Philipsen, P.H.T.; Baerends, E.J. Relativistic calculations to assess the ability of the generalized gradient approximation to reproduce trends in cohesive properties of solids. *Phys. Rev. B* 2000, *61*, 1773–1778. [CrossRef]
- 28. Nadler, M.R.; Kempier, C.P. Crystallographic data 186. Lithium. Anal. Chem. 1959, 31, 2109. [CrossRef]
- 29. Wang, K.; Reeber, R.R. The perfect crystal, thermal vacancies and the thermal expansion coefficient of aluminium. *Philos. Mag. A* **2000**, *80*, 1629–1643. [CrossRef]
- 30. Pawar, R.R.; Deshpande, V.T. The anisotropy of the thermal expansion of α-titanium. Acta Cryst. A 1968, 24, 316–317. [CrossRef]
- Buehl, M.; Reimann, C.; Pantazis, D.A.; Bredow, T.; Neese, F. Geometries of Third-Row Transition-Metal Complexes from Density-Functional Theory. J. Chem. Theory Comput. 2008, 4, 1449–1459. [CrossRef]
- 32. Smit, R. H. M.; Untiedt, C.; Rubio-Bollinger, G.; Segers, R.C.; van Ruitenbeek, J.M. Observation of a parity oscillation in the conductance of atomic wires. *Phys. Rev. Lett.* **2003**, *91*, 076805. [CrossRef]
- Kizuka, T. Atomic configuration and mechanical and electrical properties of stable gold wires of single-atom width. *Phys. Rev. B* 2008, 77, 155401. [CrossRef]
- 34. Lu, Y.; Lou, J. Quantitative in-situ nanomechanical characterization of metallic nanowires. JOM 2011, 63, 35–42. [CrossRef]