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Alkali Metal Cation Affinities of Neutral Maingroup-Element Hydrides across the Periodic Table

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Supporting Information

ABSTRACT: We have carried out an extensive quantum chemical exploration of gasphase alkali metal cation affinities (AMCAs) of archetypal neutral bases across the periodic system using relativistic density functional theory. One objective of this work is to provide an intrinsically consistent set of values of the 298 K AMCAs of all neutral maingroup-element hydrides XH_n of groups 15–18 along the periods 1–6. Our main purpose is to understand these trends in terms of the underlying bonding mechanism using Kohn–Sham molecular orbital theory together with a canonical energy decomposition analysis (EDA). We compare the trends in XH_n AMCAs with the trends in XH_n proton affinities (PAs). We also examine the differences between the trends in AMCAs of the neutral XH_n bases with those in the corresponding anionic XH_{n-1}⁻ bases. Furthermore, we analyze how the cation affinity of our neutral Lewis bases changes along the group-1 cations H⁺, Li⁺, Na⁺, K⁺, Rb⁺, and Cs⁺.

1. INTRODUCTION

Alkali metal cations are involved in many chemical and biological systems, such as in osmotic systems, electrolyte balances, ion channels, and electrochemistry.^{1–9} The thermodynamic affinity of Lewis bases for these cations, therefore, plays a significant role for predicting and understanding stability as well as reactivity in various molecular structures and chemical processes, for example, in ion-pair S_N^2 reactions.^{10–13} The alkali metal cation affinity (AMCA) is defined as the enthalpy change associated with heterolytic dissociation of the alkali cation (M⁺) complex of the neutral (XH_n) or anionic (XH_{n-1}⁻) Lewis base, as shown in eqs 1 and 2, respectively:

 $MXH_{n}^{+} \to M^{+} + XH_{n} \qquad \Delta H = AMCA \qquad (1)$

$$MXH_{n-1} \to M^{+} + XH_{n-1}^{-} \Delta H = AMCA$$
(2)

Despite the importance of this quantity, relatively little attention has been devoted to the AMCA if compared to, for example, the proton affinity (PA).^{14–19} Nevertheless, there were some theoretical^{20–30} and experimental^{31–40} attempts to better understand the trends and the features behind this quantity.

In our previous study,⁴¹ we found that the AMCAs of the anionic maingroup-element hydrides (XH_{n-1}^{-}) (eq 2) are significantly smaller than the corresponding proton affinities (PAs) and show similar although not identical trends, compared to PAs, if the Lewis-basic center X varies across

the periodic table. The reason for the smaller AMCA is mainly a much weaker HOMO–LUMO interaction between the Lewis base and the alkali metal cations if compared to the proton.⁴² This is due to the increase in the HOMO–LUMO gap and the decrease in the HOMO–LUMO overlap as the cation LUMO goes up in orbital energy and becomes more diffuse from proton 1s to alkali cation *ns* AO.

The present study extends our previous work in three ways: First, we shift our focus from the anionic Lewis bases XH_{n-1}^{-1} to the neutral maingroup-element hydrides XH_n . The main objective is to obtain a better understanding of the physical factors behind the trends in AMCA values across the periodic table based on a consistent set of accurate data in combination with detailed bonding analyses using Kohn–Sham molecular orbital (KS-MO) theory and a quantitative energy decomposition analysis (EDA). In addition to the AMCAs of all bases (ΔH_{298}) , we also report the associated 298 K entropies (ΔS_{298}) , provided as $-T\Delta S_{298}$) and 298 K Gibbs energies (ΔG_{298}) . In the second place, we wish to compare the AMCAs of the neutral bases $(XH_n)^{14-16}$ as well as the AMCAs of the corresponding anionic bases (XH_{n-1}^{-1}) .⁴¹ Third, we examine how the cation

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Table 1. Thermodynamic PA and AMCA Properties (in kcal mol⁻¹) for Neutral Maingroup-Element Hydrides at 298 K^a

		grou	p 15			grou	p 16			grou	ıp 17			gro	up 18	
period	base	ΔH	$-T\Delta S$	ΔG	base	ΔH	$-T\Delta S$	ΔG	base	ΔH	$-T\Delta S$	ΔG	base	ΔH	$-T\Delta S$	ΔG
							Prot	on Affiniti	ies							
P1													He	45.2	-5.9	39.3
P2	NH_3	203.4	-8.2	195.2	OH_2	164.7	-7.4	157.3	FH	117.6	-6.5	111.0	Ne	52.7	-5.6	47.0
P3	PH_3	185.8	-8.2	177.6	SH_2	170.6	-7.5	163.1	ClH	136.2	-6.4	129.8	Ar	93.9	-5.3	88.5
P4	AsH_3	176.7	-8.2	168.5	SeH_2	172.4	-7.5	164.9	BrH	141.9	-6.3	135.6	Kr	106.1	-5.2	100.9
P5	SbH_3	175.5	-8.1	167.4	${\rm TeH}_2$	178.5	-7.4	171.0	IH	151.9	-6.2	145.7	Xe	122.3	-5.1	117.2
P6	BiH ₃	161.4	-8.1	153.3	PoH_2	180.9	-7.4	173.5	AtH	156.1	-6.2	149.9	Rn	129.2	-5.0	124.2
							Lithium	Cation Af	finities							
P1													He	1.7	-4.6	-2.9
P2	NH_3	37.4	-7.1	30.2	OH_2	32.1	-6.8	25.4	FH	21.5	-5.8	15.7	Ne	2.7	-4.5	-1.8
P3	PH_3	23.2	-6.0	17.2	SH_2	22.2	-6.2	16.0	ClH	15.7	-5.4	10.3	Ar	6.7	-4.7	2.0
P4	AsH_3	19.6	-6.6	13.0	SeH_2	22.4	-6.2	16.2	BrH	16.2	-5.3	10.9	Kr	8.3	-4.6	3.7
P5	SbH_3	17.8	-6.4	11.4	${\rm TeH}_2$	23.0	-6.0	17.0	IH	17.3	-5.2	12.1	Xe	10.5	-4.6	6.0
P6	BiH ₃	11.4	-6.2	5.2	PoH_2	23.7	-6.1	17.6	AtH	18.1	-5.2	12.9	Rn	11.8	-4.5	7.3
							Sodium	Cation Af	finities							
P1													He	0.8	-3.9	-3.1
P2	NH_3	25.7	-6.8	18.9	OH_2	21.8	-6.4	15.4	FH	14.5	-5.3	9.3	Ne	1.2	-3.6	-2.4
P3	PH_3	15.4	-6.3	9.2	SH_2	14.7	-5.9	8.9	ClH	9.9	-5.0	4.9	Ar	3.5	-4.3	-0.7
P4	AsH_3	12.6	-6.2	6.4	SeH_2	15.1	-5.9	9.2	BrH	10.3	-5.0	5.3	Kr	4.6	-4.3	0.4
P5	SbH_3	11.2	-6.0	5.3	${\rm TeH}_2$	15.9	-5.7	10.2	IH	11.2	-4.9	6.3	Xe	6.2	-4.2	2.0
P6	BiH_3	6.1	-5.7	0.4	PoH_2	16.7	-5.8	10.9	AtH	11.9	-4.9	7.0	Rn	7.1	-4.2	2.9
							Potassium	Cation A	ffinities							
P1													He	0.4	-2.8	-2.4
P2	NH_3	17.8	-6.4	11.4	OH_2	15.7	-6.0	9.6	FH	10.7	-5.2	5.5	Ne	0.9	-4.1	-3.2
P3	PH_3	9.6	-5.7	3.9	SH_2	9.3	-5.3	4.0	ClH	6.3	-4.6	1.7	Ar	1.9	-3.9	-2.0
P4	AsH_3	7.3	-4.9	2.4	SeH_2	9.3	-5.4	3.9	BrH	6.3	-4.6	1.7	Kr	2.6	-4.0	-1.4
P5	SbH_3	5.9	-4.5	1.4	${\rm TeH}_2$	9.7	-5.2	4.5	IH	6.8	-4.4	2.4	Xe	3.5	-4.0	-0.4
P6	BiH_3	1.9	-4.4	-2.5	PoH_2	10.2	-5.3	4.8	AtH	7.2	-4.5	2.7	Rn	4.1	-3.9	0.2
							Rubidium	Cation A	ffinities							
P1													He	1.0	-4.4	-3.5
P2	NH_3	15.7	-6.2	9.5	OH_2	13.8	-5.9	7.9	FH	9.5	-5.1	4.4	Ne	0.9	-4.2	-3.3
P3	PH_3	8.2	-5.4	2.8	SH_2	8.0	-5.1	2.8	ClH	5.4	-4.4	1.0	Ar	1.5	-3.7	-2.2
P4	AsH_3	6.1	-5.3	0.8	SeH_2	7.9	-5.2	2.7	BrH	5.3	-4.4	0.9	Kr	2.1	-3.8	-1.7
P5	SbH_3	4.8	-5.0	-0.2	${\rm TeH}_2$	8.3	-4.9	3.4	IH	5.8	-4.2	1.6	Xe	2.9	-3.8	-0.9
P6	BiH_3	1.1	-4.7	-3.6	PoH_2	8.7	-5.1	3.6	AtH	6.1	-4.4	1.7	Rn	3.4	-3.8	-0.4
							Cesium	Cation Af	finities							
P1													He	0.9	-4.3	-3.4
P2	NH_3	13.7	-6.0	7.6	OH_2	12.1	-5.7	6.4	FH	8.3	-5.0	3.2	Ne	0.9	-4.1	-3.2
P3	PH_3	6.8	-5.1	1.7	SH_2	6.7	-4.9	1.8	ClH	4.5	-4.1	0.3	Ar	1.1	-3.3	-2.2
P4	AsH_3	4.9	-5.1	-0.1	SeH_2	6.6	-5.0	1.6	BrH	4.4	-4.2	0.2	Kr	1.6	-3.6	-2.0
P5	SbH_3	3.7	-4.7	-1.0	${\rm TeH}_2$	6.9	-4.6	2.3	IH	4.8	-4.0	0.8	Xe	2.3	-3.6	-1.3
P6	BiH ₃	0.4	-4.4	-4.0	PoH_2	7.3	-4.9	2.3	AtH	5.1	-4.2	0.9	Rn	2.8	-3.6	-0.8
^a Comp	uted at Z	ZORA-BI	986/QZ4	P//ZOR	A-BP86/	TZ2P fo	r the read	tion MX	${\rm H}_n^+ \to 1$	$M^+ + XI$	H _n at 298	.15 K ar	nd 1 atr	n.		

affinities of our neutral Lewis bases XH_n change along the group-1 cations H⁺, Li⁺, Na⁺, K⁺, Rb⁺, and Cs⁺.

2. METHODS

2.1. Basis Sets. All calculations were performed with the Amsterdam Density Functional (ADF) program.^{43,44} The numerical integration is performed by using a procedure developed by te Velde et al.^{45,46}

Molecular orbitals (MOs) were expanded using two large, uncontracted sets of Slater-type orbitals (STOs): TZ2P for geometry optimization and vibrational analysis and QZ4P for single-point energy calculations.⁴⁷ The TZ2P basis set is of triple- ζ quality, augmented by two sets of polarization functions (d and f on heavy atoms; 2p and 3d sets on H). The QZ4P basis, which contains additional diffuse functions, is of quadruple- ζ quality, augmented by four sets of polarization functions (two 3d and two 4f sets on C, N, and O; two 2p and two 3d sets on H). Core electrons (e.g., 1s for second-period, 1s2s2p for third-period, 1s2s2p3s3p for fourth-period, 1s2s2p3s3p3d4s4p for fifth-period, and 1s2s2p3s3p3d4s4p4d for sixth-period atoms) were treated by the frozen core approximation.⁴⁷ An auxiliary set of s, p, d, f, and g Slater-type orbitals was used to fit the molecular density and to represent the Coulomb and exchange potentials accurately in each selfconsistent field (SCF) cycle.

2.2. Density Functional. Energies and gradients were calculated using the local density approximation (LDA: Slater⁴⁸ exchange and VWN⁴⁹ correlation) with gradient corrections due to Becke (exchange) and Perdew (correlation) added self-consistently.^{50–52} This is the BP86 density func-



Figure 1. Alkali metal cation affinities AMCAs at 298 K of the neutral maingroup-element hydrides XH_n of groups 15–18 and periods 1–6 (P1–P6) and the corresponding proton affinities PAs, computed at ZORA-BP86/QZ4P//ZORA-BP86/TZ2P.

tional, which is one of the three best DFT functionals for the accuracy of geometries^{14–16,53} with an estimated unsigned error of 0.009 Å in combination with the TZ2P basis set. In a previous study^{14–16} on the proton affinities of anionic species, we compared the energies of a range of other DFT functionals, to estimate the influence of the choice of DFT functional. These functionals included the local density approximation (LDA), generalized gradient approximation (GGA), *meta*-GGA, and hybrid functionals. Scalar relativistic corrections were included self-consistently using the zeroth order regular approximation (ZORA).⁵⁴ Spin–orbit coupling effects were neglected because they are small for closed-shell systems as they occur in this investigation.

Geometries, vibrational frequencies, and thermodynamic corrections have been computed using the TZ2P basis set: ZORA-BP86/TZ2P level. All electronic energies have been

computed in a single-point fashion using the QZ4P basis set, based on the ZORA-BP86/TZ2P geometries: ZORA-BP86/QZ4P//ZORA-BP86/TZ2P. The bonding analyses have been carried out at the ZORA-BP86/TZ2P level of theory. All equilibrium geometries (see the Supporting Information) are verified by vibrational analyses to be (local) minima on the potential energy surface (zero imaginary frequencies).

2.3. Thermochemistry. Enthalpies at 298.15 K and 1 atm (ΔH_{298}) were calculated from electronic bond energies (ΔE) at ZORA-BP86/QZ4P//ZORA-BP86/TZ2P and vibrational frequencies at ZORA-BP86/TZ2P using standard thermochemistry relations for an ideal gas, according to eq 3:^{55,56}

$$\Delta H_{298} = \Delta E + \Delta E_{\text{trans},298} + \Delta E_{\text{rot},298} + \Delta E_{\text{vib},0} + \Delta (\Delta E_{\text{vib},0})_{298} + \Delta (pV)$$
(3)



Figure 2. Proton affinities (PAs) and lithium cation affinities (LiCAs) at 298 K of the neutral (XH_n) maingroup-element hydrides of groups 15–18 and periods 1–6 (P1–P6) and the corresponding proton affinities (PAs) and lithium cation affinities of the anionic (XH_{n-1}^{-}) maingroup-element hydrides of groups 14–17 and periods 2–6 (P2–P6), computed at ZORA-BP86/QZ4P//ZORA-BP86/TZ2P.

Here, $\Delta E_{\text{trans,298}}$, $\Delta E_{\text{rot,298}}$, and $\Delta E_{\text{vib,0}}$ are the differences between the reactant (i.e., MXH_n^+ , the base-cation complex) and products (i.e., $M^+ + \text{XH}_n$, the cation and the neutral base) in translational, rotational, and zero-point vibrational energy, respectively. $\Delta (\Delta E_{\text{vib,0}})_{298}$ is the change in the vibrational energy difference as one goes from 0 to 298.15 K. The vibrational energy corrections are based on our frequency calculations. The molar work term $\Delta(pV)$ is $(\Delta n)RT$; $\Delta n = +1$ for one reactant MXH_n⁺ dissociating into two products M⁺ and XH_n. Thermal corrections for the electronic energy are neglected. The change of the Gibbs energy (ΔG) in the gas phase is calculated for 298.15 K and 1 atm (eq 4).

$$\Delta G_{298} = \Delta H_{298} - T \Delta S_{298} \tag{4}$$

2.4. Bond-Energy Decomposition Analysis. The bonding analyses have been carried out at the ZORA-BP86/TZ2P level of theory. The overall bond energy ΔE_{bond} (which corresponds to $-\Delta E$ in eq 3) between cation M⁺ and base XH_n is made up of two major components:⁵⁷⁻⁶⁰

$$\Delta E_{\text{bond}} = \Delta E_{\text{strain}} + \Delta E_{\text{int}} \tag{5}$$

Here, the strain energy ΔE_{strain} is the amount of energy required to deform the separate base from its equilibrium structure to the geometry that they acquire in the overall complex MXH_n⁺. The interaction energy ΔE_{int} corresponds to

the actual energy change when the geometrically deformed base combines with the cation to form the overall complex.

The interaction ΔE_{int} between the deformed reactants is further decomposed into three physically meaningful terms, in the conceptual framework provided by the Kohn–Sham molecular orbital (KS-MO) model (eq 6).^{43,57–60}

$$\Delta E_{\rm int} = \Delta V_{\rm elstat} + \Delta E_{\rm Pauli} + \Delta E_{\rm oi} \tag{6}$$

The ΔV_{elstat} term corresponds to the classical electrostatic interaction between unperturbed charge distributions $\rho_{\text{A}}(r) + \rho_{\text{B}}(r)$ of the deformed fragments A and B and is usually attractive. The Pauli repulsion ΔE_{Pauli} comprises the destabilizing interactions between occupied orbitals (more precisely, between same-spin orbitals) and is responsible for any steric repulsion. The orbital interaction ΔE_{oi} accounts for electronpair bonding, charge transfer (interaction between occupied orbitals on one fragment with unoccupied orbitals of the other fragment, including HOMO–LUMO interactions), and polarization (empty–occupied orbitals mixing on one fragment due to the presence of another fragment).

3. RESULTS AND DISCUSSION

3.1. AMCAs and PAs of Neutral Maingroup-Element Hydrides. Our ZORA-BP86/QZ4P//ZORA-BP86/TZ2P computed alkali metal cation affinities (AMCAs) and proton

Article



Figure 3. Energy decomposition analysis of proton affinity (PA) and lithium cation affinity (LiCA) energies ΔE of neutral bases XH_n, computed at ZORA-BP86/TZ2P.

affinities (PAs) at 298 K (ΔH), the corresponding entropies ΔS (provided as $-T\Delta S$ values), and free energies ΔG of all neutral maingroup-element hydrides of group 15–18 and periods 1–6 are summarized in Table 1 and Figure 1.

The AMCAs, but also the PAs, of the neutral maingroupelement hydrides XH_n are considerably weaker than those of the corresponding anionic Lewis bases XH_{n-1}^{-} . This weakening is due to the fact that dissociation of the complex in the latter is associated with charge separation and thus substantial electrostatic attraction (eq 2), whereas no charge separation occurs in the former (eq 1). Consequently, the AMCAs of the anionic bases benefit from a substantially more stabilizing electrostatic attraction than the neutral bases. This and other features behind AMCA trends are discussed in more detail, later on, in the section on the bonding mechanism.

The trend in AMCA along the various neutral maingroupelement hydrides XH_n shows characteristic patterns which are similar for each of the alkali metal cations (*vide infra*). Along the alkali metal cations, down group 1, the AMCA of a Lewis base XH_n , in general, systematically weakens. Furthermore, the XH_n AMCAs are substantially smaller than the corresponding PAs by about 150–200 kcal mol⁻¹ (see Table 1), similar to our earlier finding for the anionic Lewis bases XH_{n-1}^{-41} Our bonding analyses reveal that these differences mainly stem from weaker orbital interactions of XH_n with the alkali metal cations as compared to those with the proton (*vide infra*).

Gibbs energies ΔG_{298} show the same trends as the corresponding enthalpic AMCA values (see Table 1). The reason is that the corresponding reaction entropies yield a relatively small and/or little varying contribution $-T\Delta S_{298}$ of -8 to -3 kcal mol⁻¹ along the entire set of model Lewis bases. Note that the AMCAs of most of the noble gases have negative Gibbs energies associated with their AMCAs. This means that these complexes are thermodynamically not stable and would dissociate spontaneously. For example, whereas ΔG_{298} for dissociating a sodium cation from radon amounts to +2.9 kcal mol⁻¹ (i.e., the complex is thermodynamically stable), the corresponding value for helium is -3.1 kcal mol⁻¹, indicating that dissociation of HeNa⁺ occurs spontaneously (see Table 1).

Along the second and third periods, the AMCA (XH_n) decreases from group 15 to 18. A major change occurs in the affinity when one descends down the groups. The neutral AMCAs (XH_n) show an inversion in the trend down a group going from group 15 to group 18. Thus, descending group 15, the AMCA *decreases*, but descending group 18, it *increases*. The same trend in the affinity is also observed in the PAs of the neutral maingroup-element hydrides (XH_n) , while it differs significantly from those previously found for the anionic conjugate bases (XH_{n-1}^{-}) of the maingroup-element hydrides where the AMCA, as well as PA, of the anionic bases (XH_{n-1}^{-}) always *decreases* down a group.

However, there is an interesting analogy between the AMCA, as well as PA, trends of XH_n and XH_{n-1} , which can be recognized in Figure 2. For both bases, the kink in the affinity trend along a period occurs after the step from the tricoordinate base (group 15 for XH_n and group 14 for XH_{n-1}^{-}) to the dicoordinate base (group 16 for XH_n and group 15 for XH_{n-1}^{-}). This kink is more pronounced for bases with a heavier protophilic center, that is, as we go from the third period down to the sixth period. In a previous study on proton affinities,¹⁴⁻¹⁶ this phenomenon has been ascribed to the valence ns electrons on the protophilic center "X" of the base becoming increasingly inactive down a period because of the relativistic stabilization of the ns AO. The sudden increase in cation affinity from a trivalent to a bivalent base is associated with an active np-type lone pair (which is always at higher energy than the *ns* electron pair) becoming available in the latter.

3.2. Bonding Mechanism: Variation of the Neutral Base. Our heterolytic $(M^+)-(XH_n)$ bonding analyses have been carried out at ZORA-BP86/TZ2P and comprise two complementary approaches: (i) quantitative analysis of the Kohn–Sham orbital interaction mechanism and (ii) the associated bond energy decomposition (see Figure 3). In the discussion, we focus on the lithium cation affinity (LiCA) due to the similarity in trends between this cation and the rest of the alkali metal cations. Furthermore, we compare this LiCA of the neutral bases XH_n with the corresponding PA and with the AMCAs of the anionic maingroup-element hydrides (XH_{n-1}^-) . Detailed numerical results from the analyses of all alkali metal cation (AMCA) as well as proton (PA) affinities can be found in the Supporting Information (see Tables S1–S6).

The trend in cation affinity ΔH is determined by that in the electronic cation affinity energy ΔE associated with reaction 1. Note that, for the bonding analysis, we use the bond energy $\Delta E_{\text{bond}} = -\Delta E$, that is, the energy change associated with bond formation $M^+ + XH_n \rightarrow MXH_n^+$. The main contributor to $\Delta E_{\text{bond}} = \Delta E_{\text{strain}} + \Delta E_{\text{intv}}$ in turn, is the interaction energy

 $\Delta E_{\rm int}$ between the two fragments which determines the overall trend in stability. ΔE_{bond} follows the same trend as ΔE_{int} because the relatively small strain energy ΔE_{strain} does not affect this trend in interaction. The reason is that, for n = 0-2, the bases XH_a do not have to deform much when forming the complex with the cation. An exception is constituted by the three-coordinate bases XH_n of group 15 which are sterically more crowded and undergo a slight, yet significant change in pyramidality as they bind to the cation.⁶¹ The ΔE_{strain} values in this group are in the order of 18 kcal mol⁻¹ for PAs and 2 kcal mol⁻¹ for LiCAs. This significant decrease in the ΔE_{strain} values going from the proton to the lithium cation can be ascribed to the weaker X-Li interaction which affects the XH₃ fragment to a lesser extent. In any case, as stated before, this strain effect is too small to change the overall trend in relative stability that is set by ΔE_{int} . A similar situation was previously found for the anionic conjugate bases (XH3-) of the maingroup-element hydrides of group 14.41

The interaction ΔE_{int} behind the cation affinities originates from a combination of three phenomena in the bonding mechanism: (i) electrostatic attraction ΔV_{elstat} which is weak for neutral XH_n, as compared to the situation of the anionic XH_n, due to the absence of charge separation in the former; (ii) orbital interaction ΔE_{oi} which is significantly stronger for H⁺ than Li⁺; and (iii) Pauli repulsion with the core AOs of the alkali cations which lessens all AMCAs and thus also the lithium cation affinity (LiCA), whereas no Pauli repulsion occurs in the PAs, as the proton has no core electrons. We recall that the AMCAs of the heavier alkali cations behave similarly to those of the lithium cation, with the understanding that they are even further weakened with respect to the corresponding PAs.

The electrostatic attraction ΔV_{elstat} of the neutral bases XH_n is relatively weak toward both lithium cations and protons if compared to the situation of the anionic bases XH_{n-1} in which case a strong Coulomb attraction between oppositely charged fragments occurs. This is the main difference between the cation affinities of the neutral and anionic Lewis bases.⁴ The exact trend in electrostatic interaction depends in an intricate manner on the shape and mutual penetration of the fragment charge distributions.⁶²⁻⁶⁴ Therefore, significant deviations of Coulomb's law $q_1 \cdot q_2 / r_{12}$ for two point charges occur. Still, one can observe a weakening in the values of $\Delta V_{\rm elstat}$ as the equilibrium bond length increases down the groups. This weakening of the AMCA down a group is also computed for the values of the electrostatic attraction of the other alkali metal cations. Likewise, essentially the same trend of a weakening in $\Delta V_{
m elstat}$ occurs if the cation itself varies down group 1, i.e., along M⁺ = Li⁺, Na⁺, K⁺, Rb⁺, and Cs⁺ (see Tables S2–S6 in the Supporting Information).

As mentioned above, for the PA, the orbital interaction ΔE_{oi} becomes the dominant component in the interaction going from anionic (XH_{n-1}^{-}) to neutral (XH_n) bases of the maingroup-element hydrides. For both cases, anionic as well as neutral bases, it is the main responsible interaction term causing the weakening in the affinity going from PA to LiCA. As can be seen in Figure 3, the orbital interaction ΔE_{oi} of the lithium cation is about 150 kcal mol⁻¹ weaker than the corresponding one of the proton. This strong weakening in ΔE_{oi} , going from H⁺ to Li⁺, is caused by the increase in the HOMO–LUMO gap (the energy of the LUMO increases drastically as one goes from the proton –13.6 eV to the lithium metal cation –6.9 eV) and the decrease in the HOMO–



Figure 4. Energy decomposition analysis of the H^+ -SH₂ and Li^+ -SH₂ interaction ΔE_{int} as a function of the bond distance *d*, computed at ZORA-BP86/TZ2P.



Figure 5. Alkali metal cation affinities (AMCAs) at 298 K of neutral maingroup-element hydrides of groups 15, 16, 17, and 18 (XH₃, XH₂, XH, X) and periods 1-6 (P1–P6), as a function of the cation, computed at ZORA-BP86/QZ4P//ZORA-BP86/TZ2P.

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Table 2. Linear Correlation Function of AMCA with PA and of AMCA with LICA Values (in kcal mol⁻¹) for Anionic and Neutral Maingroup-Element Hydrides, Together with the Correlation Coefficient r^2 and Standard Deviation SD^{*a*}

AMCA vs PA	r^2	SD	AMCA vs LiCA	r^2	SD
		Anionic Bas	ses (XH_{n-1}^{-})		
LiCA = 0.517 PA - 33.597	0.597	12.2			
NaCA = 0.410 PA - 13.255	0.710	7.5	NaCA = 0.718 LiCA + 24.335	0.976	2.1
KCA = 0.377 PA - 18.855	0.552	9.7	KCA = 0.757 LiCA + 1.532	0.996	0.9
RbCA = 0.368 PA - 19.841	0.541	9.7	RbCA = 0.746 LiCA - 0.965	0.995	1.0
CsCA = 0.377 PA - 26.389	0.495	10.9	CsCA = 0.795 LiCA - 11.815	0.989	1.6
		Neutral B	ases (XH _n)		
LiCA = 0.177 PA - 8.107	0.712	4.9			
NaCA = 0.125 PA - 6.667	0.683	3.7	NaCA = 0.719 LiCA - 1.179	0.996	0.4
KCA = 0.079 PA - 4.274	0.531	3.2	KCA = 0.505 LiCA - 1.718	0.964	0.9
RbCA = 0.066 PA - 3.427	0.480	3.0	RbCA = 0.440 LiCA - 1.581	0.942	1.0
CsCA = 0.055 PA - 2.842	0.431	2.7	CsCA = 0.382 LiCA - 1.563	0.916	1.0
Computed at ZORA-BP86/OZ4P//Z	ZORA-BP86/TZ	2P at 298.15 K a	and 1 atm. See Figure 6 for graphical r	epresentations of	the correlated

"Computed at ZORA-BP86/QZ4P//ZORA-BP86/TZ2P at 298.15 K and 1 atm. See Figure 6 for graphical representations of the correlated values.

LUMO bond overlap as a result of the diffuse nature of the alkali metal cation ns LUMOs (see also ref 42). However, both H⁺ and Li⁺ show a common trend in the orbital interaction. The orbital interaction becomes in general less stabilizing along a period and more stabilizing down a group.

This is a direct result of the trend in lone-pair orbital energies of the neutral bases XH_n of the maingroup-element hydrides. Along the periods, the lone-pair orbitals of the protophilic atom become more compact and stable: -6.3, -9.4, -13.6, to -13.6 eV for NH₃, OH₂, FH, and Ne, respectively. On the contrary, down the groups, these HOMO orbitals become less stable and go up in energy, especially down group 18, as the principal quantum number increases from n = 1 until n = 6: -15.6, -13.6, -10.3, -9.3, -8.3, and -7.8 eV along He, Ne, Ar, Kr, Xe, and Rn, respectively. On the other hand, down group 15, the orbital energies of the bases change relatively little: -6.3, -6.9, -7.0, -6.7, to -6.9 eV along NH₃, PH₃, AsH₃, SbH₃, and BiH₃, respectively. This phenomenon causes directly the inversion in the trends where the affinity decreases down the group 15 and it increases down group 18.

Furthermore, the orbital interaction ΔE_{oi} of the neutral maingroup-element hydrides XH_n with either H⁺ or Li⁺ (see Tables S1 and S2) is weaker than that of the corresponding anionic bases XH_{n-1}^{-,41} The reason is that the latter, i.e., XH_{n-1}^{-,} have higher-energy orbitals due to the net negative electrostatic potential that the electrons experience.

A characteristic of the proton–base complexes is that the electrostatic interaction becomes repulsive. The reason is the absence of the repulsive Pauli term in the interactions with a proton. In general, around the equilibrium distance, this Pauli repulsion is the main counteracting term against the attractive $\Delta V_{\rm elstat}$ and $\Delta E_{\rm oi}$ components in the bonding mechanism.^{63,64} Its absence in the case of complexes with protons leads to a shorter equilibrium H–[X]⁺ bond length. At these short distances, $\Delta V_{\rm elstat}$ becomes repulsive due to the nuclear repulsion which starts to overtake and dominates all other terms in $\Delta V_{\rm elstat}$ and prevents the bond distance from becoming 0 (see Figure 4).⁵⁷

3.3. Variation of the Alkali Metal Cations. The AMCAs of the neutral bases XH_n decrease down the alkali group (see Figure 5), similar to the previously studied AMCA of the anionic bases $(XH_{n-1}^{-})^{.41}$ This decrease is caused by a combination of two electronic mechanisms. First, the

involvement of the low-lying 2p AO of the lithium cation in the bonding contributes an extra stabilization, compared to sodium and heavier alkali cations. The reason is that, down group 1, the valence np AO goes up in energy (from -4.9 to -4.2 to -3.7 to -3.6 to -3.3 eV along Li⁺, Na⁺, K⁺, Rb⁺, and Cs⁺). Consequently, the HOMO–LUMO gap with the base increases and the associated stabilization becomes unimportant.

Second, the weakening in the AMCAs down the alkali cation group, that is, from Na⁺ to Cs⁺, is mainly caused by the increase in the HOMO-LUMO gap associated with the increase in energy of the alkali cation ns acceptor AO from $-7.1 \text{ eV} (\text{Na}^+)$ to $-6.0 \text{ eV} (\text{K}^+)$ to $-5.9 \text{ eV} (\text{Rb}^+)$ to -5.5 eV(Cs⁺). This translates into less stabilization coming from the associated orbital interaction. Furthermore, the HOMO-LUMO bond overlap decreases because the alkali cation valence ns AO becomes slightly more diffuse down this alkali cation group. For example, the overlap values between alkali metal cation ns and NH₃ $2a_1$ at their equilibrium distances amount to 0.286, 0.251, 0.196, 0.181, and 0.166 for Li⁺, Na⁺, K⁺, Rb⁺, and Cs⁺, respectively (overlap values not shown in a table). This order in overlap values is also found if we take consistently the same M-NH₃ distance for all M. Thus, if we chose, for example, an M-N distance of 2.75 Å (i.e., the equilibrium distance in the case of M = K), these overlap values still decrease as 0.213, 0.206, 0.196, 0.186, and 0.165 along Li⁺, Na⁺, K⁺, Rb⁺, and Cs⁺, respectively. Both trends, in the HOMO-LUMO gap and bond overlap, agree with earlier findings of Geerlings et al., who interpreted them in terms of the hardness of the Lewis acids and bases.⁶⁵⁻⁶⁷ We recall that the cation affinity increases down group 18 for all alkali metal cations. Note that, in absolute terms, this increase is lighter for the heavier alkali metal cations (see Figure 5). The reason for the smaller variation in cation affinity in the latter case is simply the aforementioned weaker overall affinity, originating from the higher-energy LUMO and thus larger HOMO-LUMO gap, of the heavier alkali cations.

Interestingly, down group 1, the alkali cation nd AOs rapidly descend in energy, especially from sodium (3.8 eV) to potassium (-3.7 eV), and begin to play the role of valence orbitals that can accept charge in donor-acceptor interactions. In that capacity, they participate as acceptor orbitals in the HOMO-LUMO interaction with the base instead of the ns and np AOs. Our finding of a reduced role for the 6s and 6p

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Figure 6. Correlation of alkali metal cation affinities (AMCAs) with proton affinities (PAs) and lithium metal cation affinities (LiCAs), computed at ZORA-BP86/QZ4P//ZORA-BP86/TZ2P (all values in kcal mol^{-1} ; see Table 2).

AOs and an enhanced role of the 6d AOs is in line with and augments the recent finding of Goesten et al.^{68,69} that the 6s AO has a significantly reduced importance as a valence orbital for the cesium cation in CsO_4^+ .

3.4. Correlation of the PAs and AMCAs. The linear functions for AMCA-PA and AMCA-LiCA correlations have been summarized in Table 2. Also, the correlation coefficient (r^2) and standard deviation (SD) have been provided for each correlation function with 20 data points for each alkali metal with the anionic bases and 21 data points with the corresponding neutral bases. The correlation coefficient values show, on the one hand, that there is a poor correlation between AMCAs and PAs of the anionic as well as of the corresponding neutral maingroup-element hydrides. The correlation coefficients are between 0.431 and 0.712 with relatively high standard deviations which vary from 2.7 to 12.2 kcal mol⁻¹, which are between 1.4 and 6.5% of the maximum cation affinity values. This poor correlation is due to the differences in the bonding mechanism between the proton and the alkali metal cations with the conjugate Lewis bases, which is

essentially derived from the difference in the orbital interaction. On the other hand, we find a satisfactory correlation between the anionic as well as neutral computed AMCA and the corresponding LiCA values with higher correlation coefficients between 0.916 and 0.996 with relatively low standard deviations which vary from 0.4 to 2.1 kcal mol⁻¹, which are between 0.3 and 1.3% of the maximum cation affinity values, as can be seen in Figure 6 and Table 2. The better quality of the latter correlation is attributed to the fact that the nature of the bonding mechanism is more similar among all AMCAs than between AMCAs and PAs. This agrees well with previous studies which also reported a good correlation between AMCAs and LiCAs.^{37,40,70,71}

4. CONCLUSIONS

Alkali metal cation affinities (AMCAs) of neutral maingroupelement hydrides XH_n in the gas phase are significantly smaller than the corresponding proton affinities (PAs) and the AMCAs of the anionic maingroup-element hydrides XH_{n-1} . These AMCAs show similar trends as the corresponding PAs,

as the Lewis-basic center X varies across the periodic table. Furthermore, AMCA values decrease along Li⁺, Na⁺, K⁺, Rb⁺, and Cs⁺. This follows from our quantum chemical analyses using relativistic density functional theory at ZORA-BP86/ QZ4P//ZORA-BP86/TZ2P.

The alkali metal cation affinity and the proton affinity decrease along neutral second-period maingroup-element hydrides NH₃, OH₂, FH, and Ne as valence 2p AOs of the protophilic atom become more compact and stable. This trend changes down the periodic table, that is, for the higher periods. The AMCA and PA of the maingroup-element hydrides XH_n decrease down group 15, while they increase down group 18. This is due to the more significant decrease of the HOMO-LUMO gap and, thus, the more significant stabilization in orbital interactions down group 18 than down group 15, in combination with the fact that the electrostatic attraction weakens in all cases down the periodic table as bond distances become longer. Thus, in the case of group 18, the significant enhancement of the orbital interactions overrules the weakening in electrostatic attraction and causes an overall increase in cation affinity. On the other hand, the trend in orbital interactions is too weak to overcome the weakening in electrostatic attraction down group 15, which results in the aforementioned decrease in affinity.

The AMCA and PA of neutral maingroup-element hydrides XH_n are weaker than the corresponding ones of the anionic maingroup-element hydrides XH_{n-1}^{-} , mainly, because the former go without whereas the latter go with charge separation upon dissociation. Furthermore, weaker orbital interactions and the presence of Pauli repulsion in alkali metal cation complexes are the main factors behind the fact that all AMCAs are weaker than the corresponding PAs.

The various AMCAs show similar trends with respect to variation in the bases. However, the AMCAs become smaller as the alkali cation varies down group 1 because, as the principal quantum number increases from n = 2 until 6, the alkali cation ns LUMO goes up in energy and becomes more diffuse, which leads to a weaker and longer bond toward the base. The similarity in bonding mechanisms among the AMCAs is reflected by excellent linear correlations between AMCA and LiCA. This linear relationship may be employed for accurate estimates of AMCAs based on quantum chemical data that need to be computed only for the LiCAs.

ASSOCIATED CONTENT

S Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acs.jpca.9b03814.

Energy decomposition analysis (EDA) of proton, lithium, sodium, potassium, rubidium, and cesium cation affinities and Cartesian coordinates, total energies E, and number of imaginary frequencies NIMAG of all equilibrium structures (PDF)

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Notes

The authors declare no competing financial interest.

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