


 Cite this: *RSC Adv.*, 2022, 12, 29805

# Optimized study of the annealing effect on the electrical and structural properties of HDLC thin-films

 Hari Shankar Biswas,<sup>a</sup> Jagannath Datta,<sup>b</sup> Prasenjit Mandal,<sup>c</sup> Sandeep Poddar,<sup>d</sup> Amit Kumar Kundu<sup>e</sup> and Indranil Saha<sup>e</sup>

The plasma-enhanced chemical vapor deposition (PECVD) technique has been utilized for the facile surface deposition of hydrogenated diamond-like carbon (HDLC) thin-films onto Si(100) substrates. The as-deposited film surface is homogenous, free of pinholes, and adheres to the substrate. Annealing of the synthesized HDLC surface in a vacuum was performed in the temperature range of 200 to 1000 °C. A host of instrumental techniques, *viz.* FTIR spectroscopy, AFM, STM, and EC-AFM, were successfully employed to detect the morphological transformation in the HDLC films upon annealing. EC-AFM studies show irreversible biased behavior after undergoing a surface redox couple reaction and morphological change. Raman spectroscopy was carried out along with STM and EC-AFM to determine the functional nature and conductivity of the annealed surface.

Received 5th October 2022

Accepted 5th October 2022

DOI: 10.1039/d2ra06255k

[rsc.li/rsc-advances](https://rsc.li/rsc-advances)

## 1. Introduction

For the advancement of basic solid-state material science, nanotechnology and nanoscience are largely considered major driving forces. The quality of our daily lives has improved extensively due to the discovery of novel nanomaterials, and the most important class of such nanomaterials is based on carbon. Hydrogenated diamond-like carbon (HDLC) is an artificially synthesized tetrahedral structured carbon-based nanomaterial. Its properties are similar to those of a pure diamond, which is a purely tetrahedral structure of carbon atoms. The synthesis of HDLC materials is tunable, and hence their properties are variable. The term diamond-like carbon (DLC) was first used by Aisenberg and Chabot for the ion-beam deposition of carbon thin-films.<sup>1</sup> They immensely helped in understanding the growth mechanisms, material characteristics, and feasibility of the industrial applications of DLC. A slightly modified form of DLC, hydrogenated DLC (HDLC), does not stand for a particular composition, but it is a combination of sp<sup>2</sup> and sp<sup>3</sup> hybridized crystalline and amorphous carbon thin-films whose properties effectively depend on the synthetic nature and various treatments applied. Moreover, the synthesis and design

characterization of a carbon-based nanocrystalline diamond can be tuned. We can change and improve the colour, hardness, sp<sup>2</sup>/sp<sup>3</sup> ratio, *etc.*, by varying the deposition time and flow rate of gas.<sup>2</sup> The film can be characterized using different techniques, such as scanning probe microscopy (AFM and STM), FTIR spectroscopy, Raman spectroscopy, and EC-AFM.

Owing to the unique properties of their inherent characteristics, *viz.*, a low coefficient of friction,<sup>3</sup> chemical inertness, appreciable hardness, enhanced current resistance, and chemical non-reactivity, DLC films have received widespread traction in a host of commercial applications, such as biocompatible coatings in cardiac medicines and orthopedics as well as protective hard coatings in contact lenses for tribological applications. There are various industrial apparatuses as well as some tools used in everyday life, such as scissors, knives, and safety razors, that are coated with HDLC to retain their sharpness for a long period; gears, drill bits, and weapons are also coated with DLC.<sup>4</sup>

The structure, hardness, and electrical characteristics of HDLC can be significantly altered by thermal annealing.<sup>5,6</sup> It can enhance the degree of graphitization and diminish the intrinsic compressive stress of HDLC.<sup>7,8</sup> The thermal annealing of HDLC films and subsequent bond changes result in the cleavage of C–H bonds, leading to hydrogen depreciation. All these phenomena lead to an increase in the sp<sup>2</sup> content of the film, and when the sp<sup>2</sup> clusters are finally enough for overlap, the band gap diminishes, which improves the conductive nature of the material. This creates a scientific interest in their electronic applications. The knowledge of their electronic properties, such as the carrier mobilities, change in work function, and temperature-dependent Hall effect due to graphitization,

<sup>a</sup>Department of Chemistry, Surendranath College, Kolkata 700 009, India. E-mail: harishankarb7@gmail.com

<sup>b</sup>Analytical Chemistry Division, BARC, Variable Energy Cyclotron Centre, Kolkata 700 064, India

<sup>c</sup>Department of Chemistry, Santipur College, Nadia, Pin: 741 404, West Bengal, India

<sup>d</sup>Lincoln University College, Petaling Jaya-47301, Malaysia

<sup>e</sup>Department of Chemistry, Sripat Singh College, Jiaganj-742 123, Murshidabad, West Bengal, India



informs us of the possible applications of HDLC films in electronic devices. Some other applications of HDLC materials are in forming tools for deep drawing,<sup>9</sup> active sensors,<sup>10</sup> corrosion resistance,<sup>11</sup> biomedical implants and instruments,<sup>12</sup> surface coatings for coronary stents,<sup>13</sup> heart valves, orthopedic implants,<sup>14</sup> magnetic storage devices, and micro-electromechanical devices (MEMS).<sup>15</sup>

## 2. Experimental section

### 2.1 Materials

Potassium ferrocyanide ( $K_4[Fe(CN)_6] \cdot 3H_2O$ ; MW 368.32; 99.96%; Sisco Research Laboratories Pvt Ltd), helium gas (He; 99.999%; Hindustan Gases & Welding Com), methane ( $CH_4$ ; 99%; Praxair India Pvt. Ltd), and p-type mirror-polished Si(100) wafers were used in the experiments. Ultrapure water from the Milli-Q system (Millipore Co.) was used throughout the experiments.

### 2.2 Film deposition and annealing

The simple preparation of HDLC thin-films *via* the reduced pressure enhanced nucleation (BEN) method<sup>16</sup> at normal temperature in an asymmetrically capacitive RF (13.56 MHz) combined tool<sup>17</sup> involves the following two steps: (1) etching of the mirror-polished Si(100) substrate with a 20 mm diameter for 1200 s in pure hydrogen (flow rate  $\sim 500 \text{ cm}^3 \text{ min}^{-1}$ ) plasma, at a pressure of 0.0193 kPa, produced by a 30 W RF power-producing dc self-negative bias ( $\sim -200 \text{ V}$ ), to remove the oxide layer from the surface of Si (100); and (2) the *in situ* BEN process<sup>16</sup> using He (flow rate  $\sim 1500 \text{ cm}^3 \text{ min}^{-1}$ ) plasma produced by a 50 W RF power-producing dc self-negative bias ( $\sim -200 \text{ V}$ ), with  $H_2$  (flow rate  $\sim 500 \text{ cm}^3 \text{ min}^{-1}$ ) and  $CH_4$  (flow rate  $\sim 50 \text{ cm}^3 \text{ min}^{-1}$ ) gases at a total pressure of 0.0756 kPa and at a substrate temperature of  $\sim 287 \text{ K}$  for 1800 s deposition time. The as-prepared film<sup>18</sup> was annealed at constant temperatures of 200 °C, 400 °C, 700 °C, and 1000 °C for  $\sim 1800 \text{ s}$  at a high vacuum ( $\sim 1.3 \times 10^{-4} \text{ Pa}$ ) and cooled to  $\sim 27 \text{ }^\circ\text{C}$ .

### 2.3 Instrumental

Raman spectra were obtained on a confocal micro-Raman spectrometer (LabRAM HR Vis., Horiba Jobin Yvon S.A.S), which includes an 800 mm focal length Czerny–Turner-type spectrograph equipped with mirrors (reflective optics) with a high spectral resolution of  $0.1 \text{ cm}^{-1}$  per pixel at 488 nm with an  $1800 \text{ g mm}^{-1}$  gratin. Atomic force microscopy (AFM) topography of the sample was performed in the non-contact mode with a  $Si_3N_4$  cantilever using a multimode scanning probe microscope (Agilent AFM 5500 series system, USA) with a multipurpose small scanner with a low coherence laser (1 mW power, 670 nm wavelength ( $<50 \mu\text{m}$ ), scan range XY: 0–10  $\mu\text{m}$ ; Z: 0–2  $\mu\text{m}$ ; noise level: XY  $< 0.1 \text{ nm RMS}$ , Z  $< 0.02 \text{ nm RMS}$ ). Using scanning tunneling microscopy (STM) under the constant-tunneling-current mode with a pre-cut 80% platinum–20% iridium Agilent STM tip, the topography of the sample was obtained using a multimode scanning probe microscope (Agilent AFM 5500 series system, USA) with a multipurpose small

scanner with a scan range Z of 0–2 mm and a noise level Z of  $< 0.02 \text{ nm RMS}$ . In this instrument, a feedback loop maintained a constant tunneling current during scanning by vertically moving the scanner at each (X–Y) data point until reaching a given set-point value (0.5 nano Ampere). The computer stores the vertical position of the scanner at each (X–Y) datum to form a topographic image of the sample surface, that is, whether it is conductive or semiconducting. A potentiostat/galvanostat coupled with an AFM (EC-AFM, Agilent AFM 5500 series system, PicoScan Software, USA) was used to perform electrochemical cyclic voltammetry experiments on the HDLC thin-film surface, which was the working electrode. The working electrode was connected through an L-shaped pogo electrode. The pogo contacts the sample working electrode through a separate access hole outside the liquid cell chamber. Moreover, it is not in contact with the electrolyte and does not require any special cleaning. A platinum–iridium counter electrode (CE) encompasses as much of the inner rim of the liquid cell as possible. The reference electrode (RE) is located within the electrolyte but is not in contact with the working electrode. Using cyclic voltammetry between the potentials of  $-1.0$  and  $+1.0 \text{ V}$ , the experiment was run in EC-AFM<sup>14,15</sup> with different scan rates ( $\text{mV s}^{-1}$ ) using the  $[Fe(CN)_6]^{4-}/[Fe(CN)_6]^{3-}$  redox couple in an unperturbed solution of 1 mM ferricyanide dissolved in 0.1 M  $H_2SO_4$  solution on the HDLC film. The surface morphology during electrochemical change was investigated using AFM *in situ* in the non-contact mode. The Fourier transform infrared (FTIR) spectrum of the sample was obtained on an infrared Fourier vacuum prism spectroscope Vertex 70v (Bruker Optik GmbH Germany) with ATR crystal (Zn–Se) for the solid HDLC samples. A KBr beam splitter, DLaTGS detector, MIR source, Rock Solid measuring device system, spectral resolution of  $> 0.4 \text{ cm}^{-1}$ , and contamination-free vacuum ( $< 0.2 \text{ mbar}$ ) are the key technical features of a Vertex 70v. To distinguish the transformational change and calculate the bonded hydrogen content due to the annealing of the HDLC structure, an FTIR spectrophotometer was employed.

## 3. Results and discussion

The FTIR spectra of the annealed HDLC films are shown in Fig. 1. A couple of characteristic absorption peaks of the HDLC films could be observed at 2850 and 2921  $\text{cm}^{-1}$  owing to the hybridized  $sp^3$  C–H arrangement.<sup>19</sup> Because these peaks are attributed to the attachment of H to carbon preferring the  $sp^3$  structure more than  $sp^2$ , no other peak is recorded. The intensity of percentage transmittance changes reduces as the annealing temperature rises. This observation is attributed to the change in the hydrogen content in the films due to the tempering or diminishing  $sp^3$  content in the HDLC films.

It is well known that igniting non-crystalline carbon films leads to a preferential release of hydrogen, which is the foundation of structural modification. Some authors report a similar decrease in the intensities of the aforementioned peaks.<sup>20–23</sup> Significantly, after annealing, there is no marked shift in the peak position. This implies that the interatomic space between C and H is unchanged after annealing. The net H content

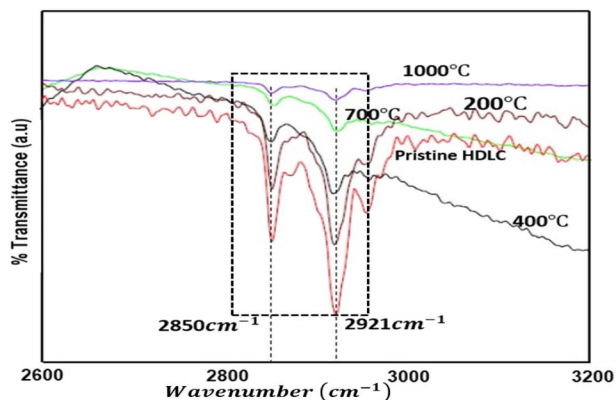


Fig. 1 FTIR spectra of pristine HDLC and annealed HDLC thin-films at different temperatures.

present in the films accounts for the cumulatively bonded hydrogens, those directly adhered to the carbon atoms, and those indirectly bound to the carbon atoms.<sup>7</sup> Some direct methods have been used to measure the H-content in diamond films, such as ERDA and NRA,<sup>24</sup> and other methods, such as NMR and IR absorption, can be used to monitor the H-content in diamond-like materials indirectly compared with the NRA method.<sup>24</sup> The effect of annealing was demonstrated using FTIR analysis herein. The total hydrogen content attached to C as C–H was estimated from the sum of the indicated areas of the infrared absorption peaks in the FTIR spectra using the following eqn (1):<sup>25</sup>

$$\begin{aligned} \text{C–H bond content} &\propto \int \frac{\alpha(\nu)}{\nu} d\nu \text{ and C–H bond content} \\ &= A \int \frac{\alpha(\nu)}{\nu} d\nu \end{aligned} \quad (1)$$

where  $A$  is the proportionality factor proportional to the inverse of the oscillator length; its value changes based on the position of the absorption peak. The values of  $A$  are  $1.085 \times 10^{19} \text{ cm}^{-2}$  and  $6.899 \times 10^{19} \text{ cm}^{-2}$  for the absorption peaks at  $2850 \text{ cm}^{-1}$  and  $2921 \text{ cm}^{-1}$ , respectively,<sup>25</sup> and  $\alpha(\nu)$  is the absorption coefficient at frequency  $\nu$ . Fig. 2 shows that thermal annealing leads to an indirect quantitative measurement of the evolution of

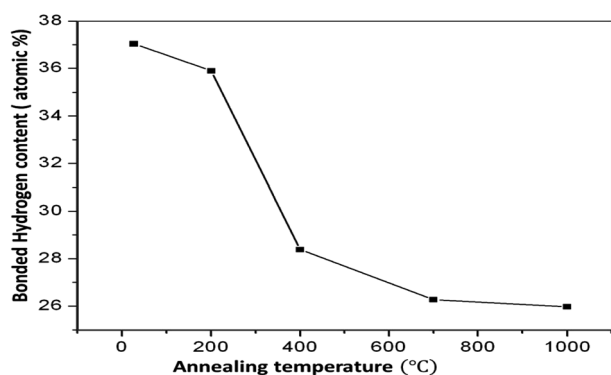


Fig. 2 Bonded hydrogen vs. annealing temperature.

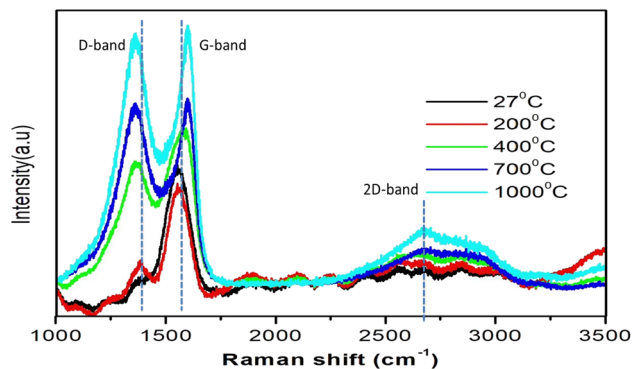


Fig. 3 Typical Raman spectra of HDLC with different annealing temperatures.

Table 1 Variation of the  $sp^2/sp^3$  ratio and force of attraction with annealing temperature

Sample	Annealing temperature (°C)	Force of attraction (nN)	$sp^2$	$sp^3$	$sp^2/sp^3$
HDLC	27	50.50	61.33	38.67	1.58
	200	48.80	63.77	36.22	1.76
	400	28.70	77.96	22.04	3.54
	700	21.20	83.82	16.18	5.18
	1000	21.00	84.80	15.20	5.58

hydrogen from both groups *via* the breaking of C–H bonds in the films.

One of the commonly used non-destructive techniques to characterize graphitic materials is Raman spectroscopy, particularly to specify the defects and disorders as well as the ordered structure of graphene and graphite-like structures. Fig. 3 shows two Raman bands, one at  $1369 \text{ cm}^{-1}$  (D band) and another at  $1551 \text{ cm}^{-1}$  (G band), which correspond to the vibrational modes of the CH and CH<sub>2</sub> groups at 27 °C on the HDLC surface. The G band is associated with the symmetric  $E_{2g}$  C=C stretching mode of  $sp^2$  bonding in carbonaceous materials,<sup>26</sup> whereas the D band can be ascribed to the angle

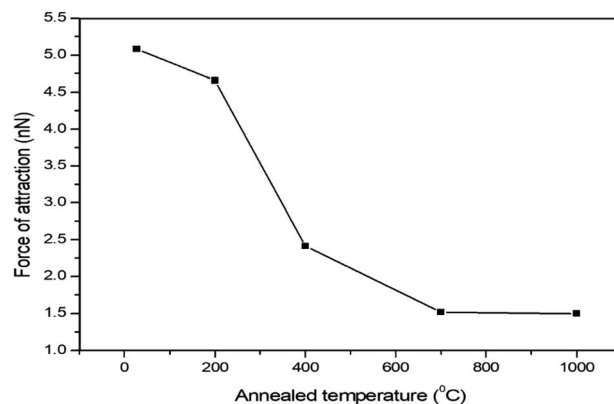


Fig. 4 Force of attraction vs. annealed temperature.

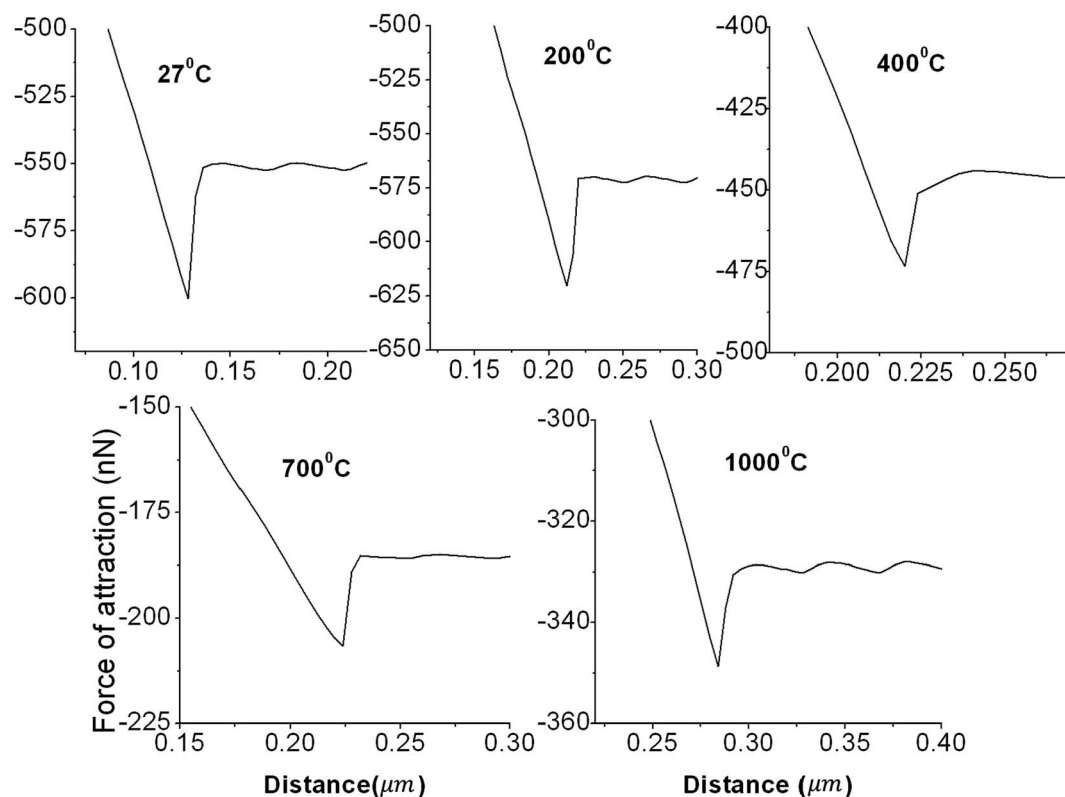


Fig. 5 Typical force vs. distance curves at different annealing temperatures of the HDLC surface using a spring constant of  $0.765 \text{ N m}^{-1}$  in air for the  $\text{Si}_3\text{N}_4$  tip.

discrepancy in the graphite-like microdomains affected by the  $\text{sp}^3$  bonds or  $j$ -point phonons of  $A_{1g}$  symmetry.<sup>27</sup> Due to the increase in  $\text{sp}^2$  carbon after annealing, the second-order Raman (2D) spectrum increases. The intensity ratio of the D-band and G-band ( $I_D/I_G$ ) is calculated to ascertain the microstructure of the graphite-like sheets and specify the lattice disorder. The

shifting of the G and D bands with increasing  $I_D/I_G$  ratio due to annealing clearly indicates dehydrogenation (*i.e.*, oxidation) of the HDLC film surface. Using the G peak position ( $\omega_G$ ) of the Raman spectra, the  $\text{sp}^2/\text{sp}^3$  ratio is estimated by the following empirical equation,<sup>27</sup> as shown in Table 1.

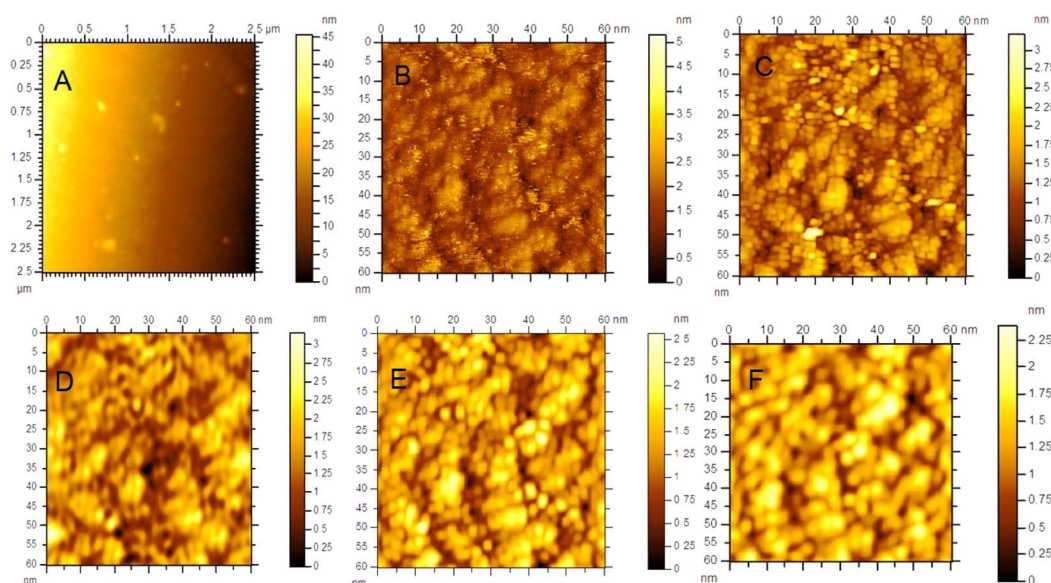


Fig. 6 (A) Non-contact-mode AFM topography image of the HDLC surface; scan size:  $2.5 \mu\text{m} \times 2.5 \mu\text{m}$ . STM images of the annealed HDLC film at (B)  $200^\circ\text{C}$ , (C)  $300^\circ\text{C}$ , (D)  $400^\circ\text{C}$ , (E)  $700^\circ\text{C}$  (F) and  $1000^\circ\text{C}$ ; scan size:  $60 \text{ nm} \times 60 \text{ nm}$ .



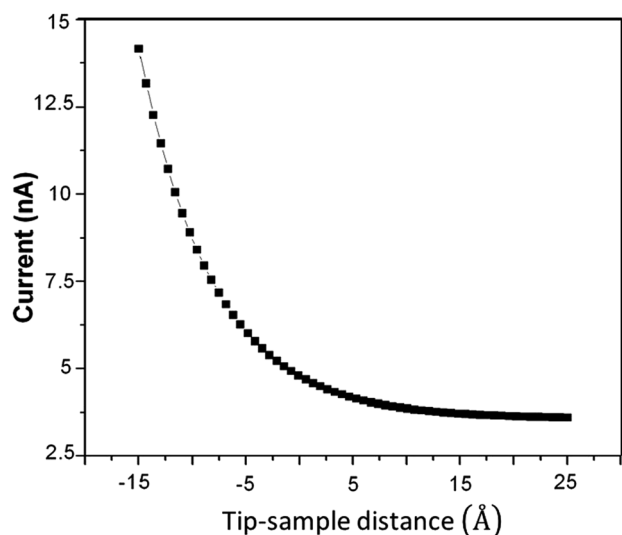


Fig. 7 Typical tunneling current vs. distance curve for the annealed (400 °C) sample.

$$sp^3 \text{ content} = 0.24 - 48.9 (\omega_G - 0.1580) \quad (2)$$

Fig. 4 presents the force of attraction vs. annealing temperature of the sample using a  $\text{Si}_3\text{N}_4$  cantilever tip with a spring constant of  $\sim 0.765 \text{ N m}^{-1}$  obtained by AFM to probe the distant assigned forces between the  $\text{Si}_3\text{N}_4$  and HDLC surfaces separated by air. Fig. 5 presents the distance-dependent force curves along the  $X$ -axis cantilever tip to the sample distance and along the  $Y$ -

axis attraction force between the tip and sample in nano-newtons (nN). From Fig. 5, it is seen that the values of the force of attraction between HDLC and the annealed<sup>28</sup> HDLC surfaces and silicon nitride probe decrease gradually as the annealing temperature goes from 27 °C to 1000 °C, as shown in Table 1. This gives the impression that with the decrease of bonded hydrogen, as we increase the sample temperature from 27 °C to 1000 °C, a decrease in the force of attraction between the sample and the  $\text{Si}_3\text{N}_4$  tip is observed. This is caused by a decrease in the density after annealing or an increase in the electron density on the surface of the annealed sample due to graphitization, as shown in Fig. 4. This is due to the influence of the  $sp^2/sp^3$  ratio<sup>29</sup> ( $I_D/I_G$ ), which increases with temperature, as supported by the Raman spectra.

It is impossible to measure the fixed-current STM image for the as-prepared HDLC films due to their non-conductive nature, but it is possible with the annealed one due to its conductive nature, as shown in Fig. 6. Fig. 6(A) shows a non-contact-mode AFM topography image of the as-prepared HDLC surface (scan size  $2.5 \mu\text{m} \times 2.5 \mu\text{m}$ ), which significantly differs from the images in Fig. 6(B)–(F), which are nanoscale STM images of the annealed HDLC films with a scan size of  $60 \text{ nm} \times 60 \text{ nm}$ . Fig. 7 presents a typical tunneling current vs. tip-sample distance plot and Fig. 8 presents a  $\ln(\text{current})$  vs. tip-sample distance plot for the annealed HDLC film surface. This outcome points to the usual current vs. height relation.<sup>30</sup> The calculated work function of the annealed (400 °C) film obtained from the slope of the curve<sup>31</sup> (Fig. 8) is  $\sim 0.021 \text{ eV}$ , which is about 214 times lower than the work function of a pristine graphene film ( $\sim 4.5 \text{ eV}$ ).<sup>32</sup> Another reason for this significant change in work function can

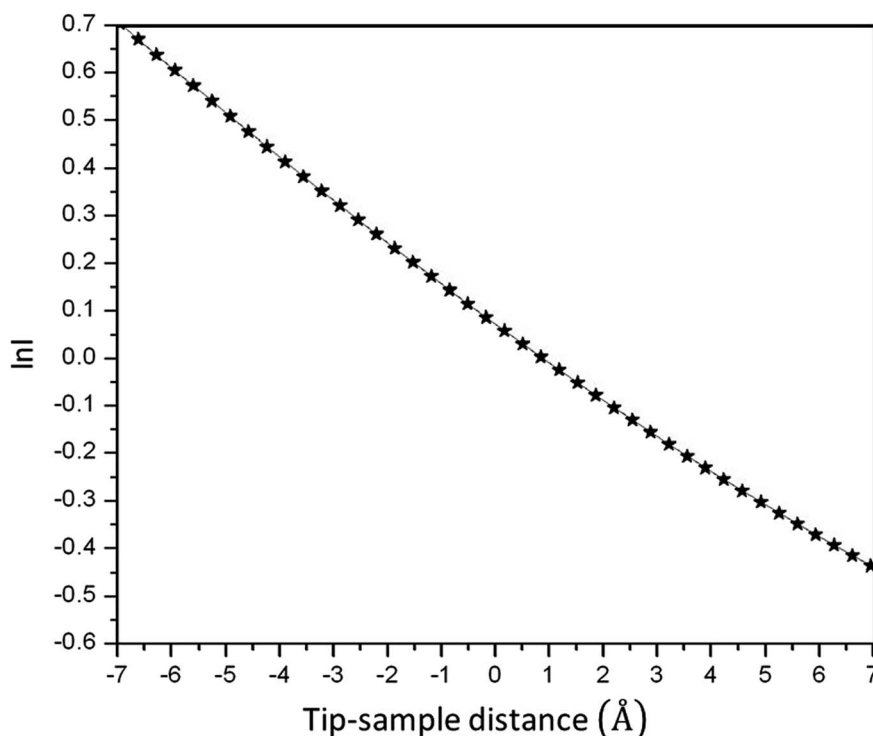


Fig. 8 Semilogarithmic plot of the tunneling current vs. distance for the annealed (400 °C) HDLC sheet.

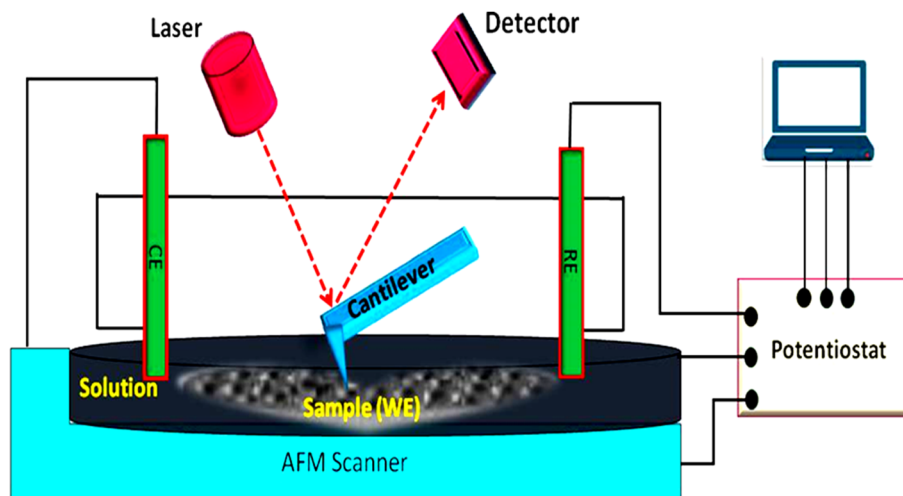


Fig. 9 Schematic drawing of the electrochemical AFM (EC-AFM) with a potentiostat.

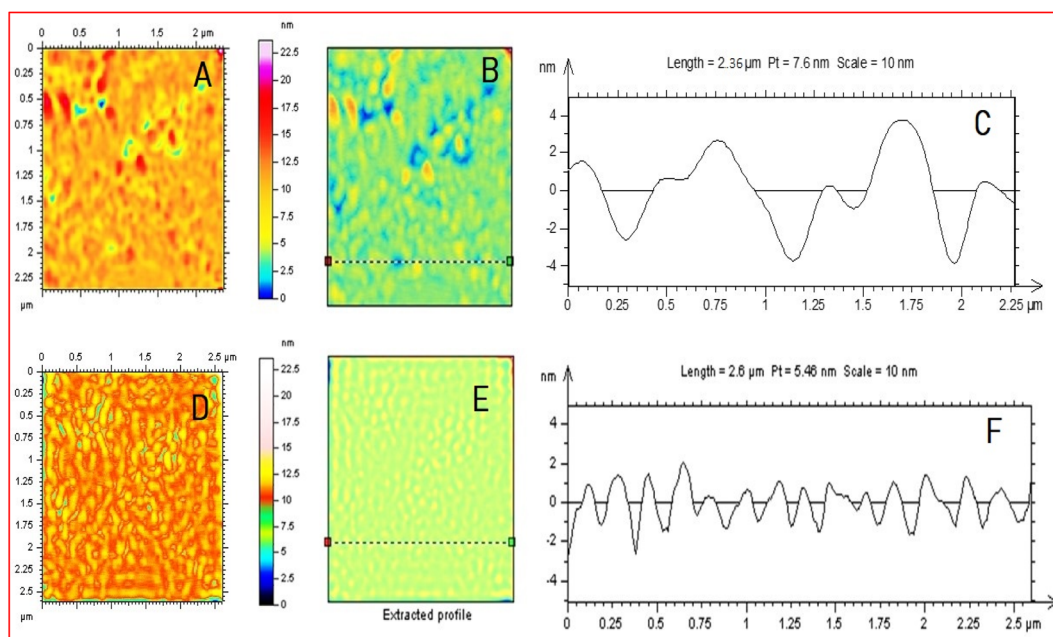


Fig. 10 (A) Non-contact-mode AFM image in the range of  $-1.0\text{ V}$ – $0\text{ V}$  with the cathodic peak potential (B) indicated by a line in the HDLC image along the surface and (C) the corresponding height profile. (D) Non-contact mode AFM image in the range of  $0\text{ V}$ – $-1.0\text{ V}$  with the anodic peak potential (E) indicated by a line in the HDLC image along the surface and (F) the corresponding height profile. Scan size:  $2.5\text{ }\mu\text{m} \times 2.5\text{ }\mu\text{m}$ .

be ascribed to the preferential contact of graphene with an insulating or metallic substrate, which can alter the Fermi energy level.<sup>32,33</sup> The conductivity and mechanical properties of annealed HDLC increase due to the increase in the number of graphene layers with the graphitization of HDLC at high temperatures.

Fig. 9 presents a schematic diagram of the EC-AFM with a potentiostat. Fig. 10(A)–(C) present a non-contact-mode AFM image in the range of  $-1.0$  to  $0\text{ V}$  in the reduction process, a line of the image of HDLC along with the surface measure roughness (RMS  $\sim 0.721\text{ nm}$ ), and the corresponding height

profile, while Fig. 10(D)–(F) present a non-contact-mode AFM image in the range of  $0.0$  to  $1.0\text{ V}$  in the oxidation process, a line of the image of HDLC along with the surface measure roughness (RMS  $\sim 0.571\text{ nm}$ ) and corresponding height profile; both have a scan size of  $2.5\text{ }\mu\text{m} \times 2.5\text{ }\mu\text{m}$ . Fig. 10(A) and (D) significantly differ from the non-contact-mode image of HDLC in Fig. 6(A) due to the change in the surfaces by the cathodic peak potential and anodic peak potential.

Fig. 11 presents the cyclic voltammograms with different scan rates. The peak potential of the working electrode is  $+0.216\text{ V}$  during a forward-biased scan from negative to positive

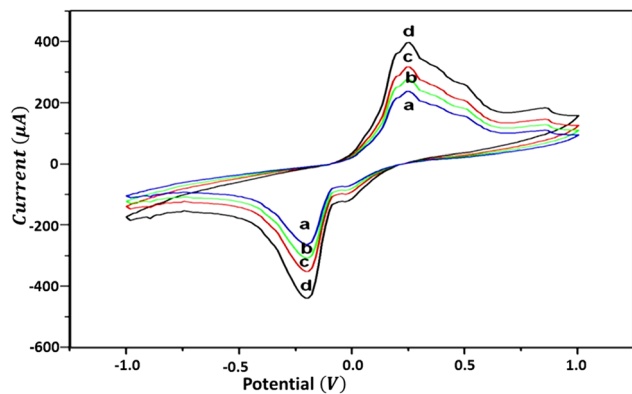


Fig. 11 Cyclic voltammograms of the HDLC electrode at various scan rates: (a)  $50 \text{ mV s}^{-1}$ , (b)  $100 \text{ mV s}^{-1}$ , (c)  $150 \text{ mV s}^{-1}$  and (d)  $200 \text{ mV s}^{-1}$ . The electrolyte is  $1 \text{ mM K}_3\text{Fe}(\text{CN})_6/0.1 \text{ M H}_2\text{SO}_4$ .

( $-1.0$  to  $+1.0 \text{ V}$ ), which is referred to as the anodic peak potential  $E_{\text{pa}}$ . Similarly, the peak potential of the working electrode is  $-0.210 \text{ V}$  during a reverse-biased scan from positive to negative ( $+1.0$  to  $-1.0 \text{ V}$ ), which is referred to as the cathodic peak potential  $E_{\text{pc}}$ . The separation of peak potentials ( $\Delta E = E_{\text{pa}} - E_{\text{pc}}$ ) is  $426 \text{ mV}$ , which is considerably large, indicating a quasi-reversible biased behavior of the working electrode (HDLC surface) at a scan rate range of  $50 \text{ mV s}^{-1}$  to  $200 \text{ mV s}^{-1}$ .<sup>34</sup> This indicates that the HDLC film surface after irreversibly undergoing a surface redox couple reaction can donate and receive electrons owing to the presence of  $\text{sp}^2 \text{ C}=\text{C}$  (donor) and  $\text{sp}^3 \text{ C-H}$  (acceptor) carbons present homogeneously in the HDLC film with a thickness of  $\sim 168 \text{ nm}$ .<sup>35</sup> The quasi-reversible process on the HDLC film suggests that during the reduction of HDLC, the effect is finite unless we provide an external potential to remedy oxidation, which means that more hydrogen can be stored in this process.

Electrodes that tend to be hydrophobic, such as diamond, HDLC, and other carbon-based electrodes, can promote hydrophobic interactions with hydrophobic analytic species, such as proteins and aromatic compounds, and foul the electrode.<sup>36</sup> Due to the chemical inertness and wide potential window of our HDLC film, it shows reproducible electron transfer and behaves as a quasi-reversible working electrode in the presence of a hydrophilic electrolyte. The  $I$ - $V$  characteristic behavior presents a wide potential window, and HDLC exhibits ambipolar behavior; the top view from the schematic drawing of HDLC surface<sup>37</sup> shows the coherent domains of  $\text{sp}^3 \text{ C-H}$  (non-conducting) and  $\text{sp}^2 \text{ C}=\text{C}$  (conducting) carbons. This indicates the resistive nature of the HDLC surface, which has been previously discussed.<sup>37</sup> Hydrogenated diamond-like carbon (HDLC) is a promising electrode material owing to its unique electrochemical properties, such as chemical inertness, resulting in a wide potential window and low background current. The basic electrochemical properties and responses of HDLC to several redox systems have also been investigated and recently reviewed. One of the main factors that determines the properties of a given HDLC material is the ratio of  $\text{sp}^3$  to  $\text{sp}^2$  bonded carbon in the film.

## 4. Conclusions

We have synthesized an HDLC thin-film *via* chemical vapor deposition. As the materials are synthesized from a gaseous system using radio frequency plasma (RF plasma), they did not require high heat and pressure. Hence, this method is cost-effective. The thermal annealing effect on the properties of the obtained HDLC thin-films has been studied by preferentially synthesizing the films on silicon (100) substrates using the plasma-enhanced chemical vapor deposition technique. The structural, electrical, and hybridization ( $\text{sp}^2/\text{sp}^3$ ) calculations of the as-deposited and annealed films were determined to investigate these effects. An appreciable reduction in the hydrogen concentration upon annealing of the samples was predicted using FTIR spectroscopy, which is supported by Raman spectroscopy. AFM, STM, and EC-AFM studies were performed to measure the transformation in the HDLC thin-films by annealing. Using Raman, STM, and EC-AFM analyses, the nature of the conductivity and functionality of the surface were determined. In conclusion, the application of the scanning probe measurement technique is highly effective and may be successfully used to characterize HDLC, graphene or graphite-like carbon materials, as has been demonstrated above.

## Author contributions

Hari Shankar Biswas: resources, conceptualization, investigation, formal analysis, validation, writing – review & editing; Jagannath Datta: investigation, formal analysis, writing; Prasenjit Mandal: writing – review & editing; Sandeep Poddar: formal analysis, writing – review & editing, validation; Amit Kumar Kundu: writing – review & editing; Indranil Saha: validation; proof checking, grammar correction.

## Conflicts of interest

There are no conflicts to declare.

## Acknowledgements

One of the authors (Hari Shankar Biswas) thanks the UGC Minor Project, Reference No. F.PSW-140/15-16, for funding. The authors also thank Dr Susanta Debnath, Saha Institute of Nuclear Physics, Mr U.S. Sil of PPD, SINP, Dr Suchandra Chatterjee, and Dr Lalita Das, Department of Chemistry, Surendranath College, for the characterization of the samples and their technical help during the experiments in this work.

## References

- 1 S. Aisenberg and R. Chobat, *J. Appl. Phys.*, 1971, **42**, 2935.
- 2 H. S. Biswas, *J. Indian Chem. Soc.*, 2019, **96**, 413–418.
- 3 W. S. Choi, K. Kim, J. Yi and B. Hong, *Mater. Lett.*, 2008, **62**, 577.
- 4 Z. Panosyan, *Thin Solid Films*, 2009, **517**, 5404.
- 5 M. I. Landstrass and K. V. Ravi, *Appl. Phys. Lett.*, 1989, **55**(10), 975.

- 6 W. Zhang, J. Einstein and L. Ley, *Phys. Rev. E: Stat., Nonlinear, Soft Matter Phys.*, 2008, **78**(4), 041603.
- 7 W. Choi and B. Hong, *Renewable Energy*, 2008, **33**, 226.
- 8 H. Li, X. Tao, C. Wang, C. Jianmin, H. Zhou and H. Liu, *Thin Solid Films*, 2006, **515**, 2153–2160.
- 9 K. Taube, *Surf. Coat. Technol.*, 1998, **98**, 976.
- 10 J. L. Davidson, W. P. Kang, Y. Gurbuz, K. C. Holmes, L. G. Davis, A. Wisitor-at, D. V. Kerns, R. L. Eidson and T. Henerson, *Diamond Relat. Mater.*, 1999, **8**, 1741.
- 11 G. Dearnaley, *Clin. Mater.*, 1993, **12**, 237–244.
- 12 S. Ravi and P. Silva, *EMIS Datareviews series EMIS Processing Series*, ed. B. Weiss, INSPEC, The Institution of Electrical Engineers, London, United Kingdom, 2003.
- 13 F. Airoldi, A. Colombo, D. Tavano, G. Stankovic, S. Klugmann, V. Paolillo, *et al.*, *Am. J. Cardiol.*, 2004, **93**, 474–477.
- 14 H. S. Tran, M. M. Puc, C. W. Hewitt, D. B. Soll, S. W. Marra, V. A. Simonetti, *et al.*, *J. Invest. Surg.*, 1999, **12**, 133–140.
- 15 J. K. Luo, J. H. He, Y. Q. Fu, A. J. Flewitt, N. A. Fleck and W. I. Milne, *J. Micromech. Microeng.*, 2005, **15**, 1406.
- 16 J. Zhang, F. Zhang, H. Yang, X. Huang, H. Liu, J. Zhang and S. Guo, *Langmuir*, 2010, **26**(9), 6083–6085.
- 17 S. Fenghua, C. Guofu and J. Sun, *Tribol. Int.*, 2019, **130**, 1–8.
- 18 H. S. Biswas, J. Datta, D. P. Chowdhury, A. V. R. Reddy, U. C. Ghosh, A. K. Srivastava and N. R. Ray, *Langmuir*, 2010, **26**(22), 17413–17418.
- 19 M. R. Roy, E. Pascual and M. C. Polo, *Surf. Coat. Technol.*, 2008, **202**, 2354.
- 20 R. Paul, S. Dalui, S. N. Das, R. Bhar and A. K. Pal, *Appl. Surf. Sci.*, 2008, **255**, 1705.
- 21 R. Wächter and A. Cordery, *Diamond Relat. Mater.*, 1999, **8**, 504.
- 22 A. M. M. Omer, S. Adhikari, S. Adhikary, M. Rusop, H. Uchida, T. Soga and M. Umeno, *Diamond Relat. Mater.*, 2006, **15**, 645.
- 23 H. Weidong, D. Ding and Z. Rujuan, *Plasma Sci. Technol.*, 2004, **6**, 2255.
- 24 J. Datta, H. S. Biswas, P. Rao, G. L. N. Reddy, S. Kumar, N. R. Ray, D. P. Chowdhury and A. V. R. Reddy, *Nucl. Instrum. Methods Phys. Res., Sect. B*, 2014, **328**, 27–32.
- 25 E. Titusa, D. S. Misra and A. K. Sikder, *Diamond Relat. Mater.*, 2005, **14**, 476.
- 26 E. F. Sheka, Y. A. Golubev and N. A. Popova, *Nanomaterials*, 2020, **10**, 2021.
- 27 A. Singha, A. Ghosh, A. Roy and N. R. Ray, *J. Appl. Phys.*, 2006, **100**(8), 44910.
- 28 C. Casiraghi, A. C. Ferrari and J. Robertson, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2005, **72**(14), 85401.
- 29 Z. Lu, B. Zhang, N. Yao, X. Zhang, B. Ma and Z. Fan, *IEEE Cat.*, 2004, 04EX839.
- 30 N. Noei, A. Weismann and R. Berndt, *Beilstein J. Nanotechnol.*, 2018, **9**, 3048–3052.
- 31 <https://www.ntmdt.com/spm-basics/view/current-distance-characteristic>.
- 32 Y. J. Yu, Y. Zhao, S. Ryu, L. E. Brau, K. S. Kim and P. Kim, *Nano Lett.*, 2009, **9**(10), 3430–3434.
- 33 G. Giovannetti, P. A. Khomyakov, G. Brocks, V. M. Karpan, J. van den Brink and P. J. Kelly, *Phys. Rev. Lett.*, 2008, **101**, 026803.
- 34 A. J. Bard and L. R. Faulker, *Electrochemical methods: fundamentals and applications*, John Wiley & Sons Inc., New York, 2nd edn, 2001, vol. 239, ch. 6, ISBN: 978-0-471-04372-0.
- 35 J. Datta, N. R. Ray, P. Sen, H. S. Biswas and E. A. Vogler, *Mater. Lett.*, 2012, **71**, 131.
- 36 L. H. Benjamin, S. Shajahan and D. K. Y. Wong, *Rev. Anal. Chem.*, 2016, **35**(1), 1–28.
- 37 H. S. Biswas, *Int. J. Adv. Res. Sci. Eng. Technol.*, 2018, **9**(5), 87–94.