



## Article

# Nitrogen-Containing Gas Sensing Properties of 2-D Ti<sub>2</sub>N and Its Derivative Nanosheets: Electronic Structures Insight

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**Abstract:** In this work, the potentials of two-dimensional Ti<sub>2</sub>N and its derivative nanosheets Ti<sub>2</sub>NT<sub>2</sub>(T=O, F, OH) for some harmful nitrogen-containing gas (NCG) adsorption and sensing applications have been unveiled based on the quantum-mechanical Density Functional Theory calculations. It is found that the interactions between pure Ti<sub>2</sub>N and NCGs (including NO, NO<sub>2</sub>, and NH<sub>3</sub> in this study) are very strong, in which NO and NO<sub>2</sub> can even be dissociated, and this would poison the substrate of Ti<sub>2</sub>N monolayer and affect the stability of the sensing material. For the monolayer of Ti<sub>2</sub>NT<sub>2</sub>(T=O, F, OH) that is terminated by functional groups on surface, the adsorption energies of NCGs are greatly reduced, and a large amount of charges are transferred to the functional group, which is beneficial to the reversibility of the sensing material. The significant changes in work function imply the good sensitivity of the above mentioned materials. In addition, the fast response time further consolidates the prospect of two-dimensional Ti<sub>2</sub>NT<sub>2</sub> as efficient NCGs' sensing materials. This theoretical study would supply physical insight into the NCGs' sensing mechanism of Ti<sub>2</sub>N based nanosheets and help experimentalists to design better 2-D materials for gas adsorption or sensing applications.

**Keywords:** atomic structure; electronic structure; Density Functional Theory (DFT); MXenes; two-dimensional materials



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## 1. Introduction

Industrial pollutant emissions account for a large share of environmental pollution, and one of the most harmful is nitrogen-containing gas (NCGs), such as NO, NO<sub>2</sub>, and NH<sub>3</sub>. At the same time, the nitrogen-containing gases themselves are generally toxic and can damage the respiratory cells of humans and animals [1–4]. Therefore, there is an urgent need to explore efficient and alternative materials for the sensing or capturing of these gases. Although some current solid-state sensors can also detect toxic gases, there is still much room for improvement due to the increasingly high demand for sensitivity, selectivity, stability, and other factors [5,6]. Two-dimensional (2-D) materials have large contact areas with gas and special electronic properties; thus, their interactions with gas molecules have been widely studied, especially in gas sensing and capturing [7]. For example, Zhang et al. studied the adsorption behavior between four modified graphene substrates and small gas molecules (CO, NO, NO<sub>2</sub>, and NH<sub>3</sub>) and significantly improved the performance of the gas sensor by introducing doping or defects [8]. Liu et al. found that the GeSe monolayer had potential for NH<sub>3</sub> gas capture and storage, as well as for the detection and catalysis of SO<sub>2</sub> and NO<sub>2</sub>. The optical properties of this GeSe monolayer can also be tuned as a candidate

material for optical sensors by adsorbing different small molecules [9]. Li et al. compared the structures and electronic properties of pure and two vacancy-defected MoS<sub>2</sub> nanosheets with NO adsorption. Defects can effectively improve the NO-sensing performance [10].

Transition metal carbide or nitride MXenes are also one family of promising gas-sensing materials [11,12], where M is a transition metal and X is carbon or nitrogen. The general formula of MXenes is written as M<sub>n</sub>X<sub>n-1</sub>. During the preparation process in an aqueous solution containing fluoride ions, there may be some O, F, and OH functional groups formed [13–15]. This type of MXene with functional groups is written as M<sub>n</sub>X<sub>n-1</sub>T<sub>n</sub> (T=O, F, OH). MXenes are also widely reported in energy storage, electrode materials, ion batteries, and the spintronic and optical devices, due to their unique structures and properties [16–24].

After our review of previous literature, we found that some carbide-based MXenes have the potential to be used as sensors for molecules, such as amino acids and nucleobase [25,26]. With less exfoliation energy compared with other types, the synthesizability of the nitride MXenes is good [27], while there are fewer studies on their use in sensing NCGs [28–32]. In order to obtain a theoretical understanding of the nitride MXenes in terms of their sensing or capturing performance towards nitrogen-containing gas (NCGs), in this work, we conducted a detailed computational study on the adsorption behaviors of NCGs (including NO, NO<sub>2</sub>, and NH<sub>3</sub>) of Ti<sub>2</sub>N monolayer with bare surface and terminated with O, F, and OH functional groups. Density Functional Theory–based first-principles calculations were performed to unveil the adsorption structures, charge transfer, band structures, density of states, and work functions to stimulate and inspire the deep mechanism investigations.

## 2. Methods

DFT calculations were employed by using the projector augmented wave (PAW) [33] pseudo-potentials in a Vienna *Ab initio* Simulation Package (VASP) [34]. For the exchange-correlation functional, we used the Perdew–Burke–Ernzerhof (PBE) [35] form of the generalized gradient approximation (GGA). Grimme’s DFT-D2 method was used to describe the van der Waals interactions [36]. This method adds semi-empirical dispersion on the basis of conventional Kohn–Sham DFT energy. In calculations, Ti<sub>2</sub>N and Ti<sub>2</sub>NT<sub>2</sub> (T=O, F, OH) were constructed as 4 × 4 × 1 supercells, and a 20 Å vacuum layer was added in Z-axis direction to eliminate the interlayer interaction. The energy cutoff of 520 eV was used for wave functions’ expansion. The Brillouin zone integration was sampled with 5 × 5 **k**-grid mesh for geometry optimizations and 15 × 15 **k**-grid mesh for electronic properties calculations to achieve high accuracy. Geometry optimizations were performed by using the conjugated gradient method, and the convergence criterion was set to be 0.02 eV/Å on force and 1 × 10<sup>-5</sup> eV/atom on energy. The adsorption energy (E<sub>ads</sub>) of nitrogen-containing gas (NCGs) on Ti<sub>2</sub>N or Ti<sub>2</sub>NT<sub>2</sub> (T=O, F, OH) is defined as follows:

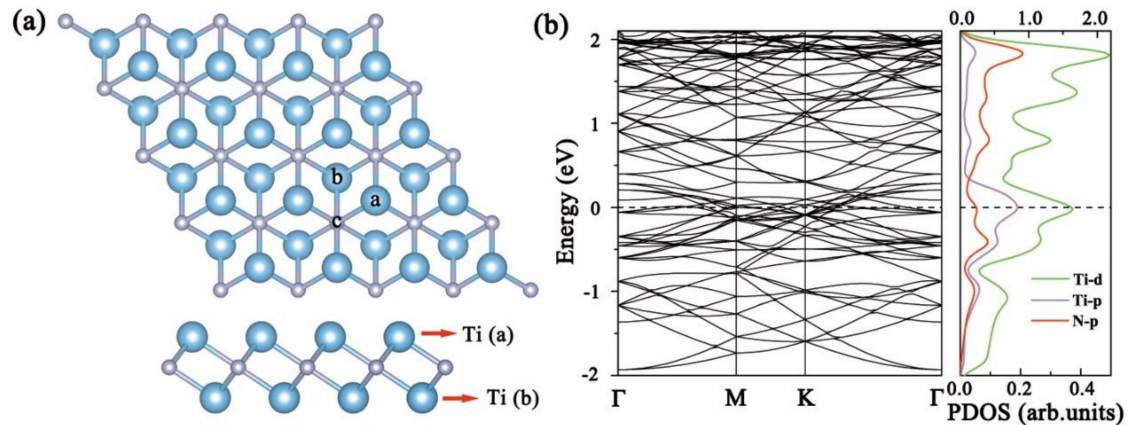
$$E_{\text{ads}} = E_{\text{MXene} + \text{NCGs}} - E_{\text{NCGs}} - E_{\text{MXene}} \quad (1)$$

where E<sub>ads</sub>, E<sub>MXene + NCGs</sub>, E<sub>MXene</sub>, and E<sub>NCGs</sub> stand for the adsorption energy, total energies of optimized Ti<sub>2</sub>N or Ti<sub>2</sub>NT<sub>2</sub> (T=O, F, OH) MXene with adsorbed NCGs molecules, energies of pure MXene, and isolated NCGs (NO, NO<sub>2</sub>, and NH<sub>3</sub>) molecules respectively.

## 3. Results and Discussion

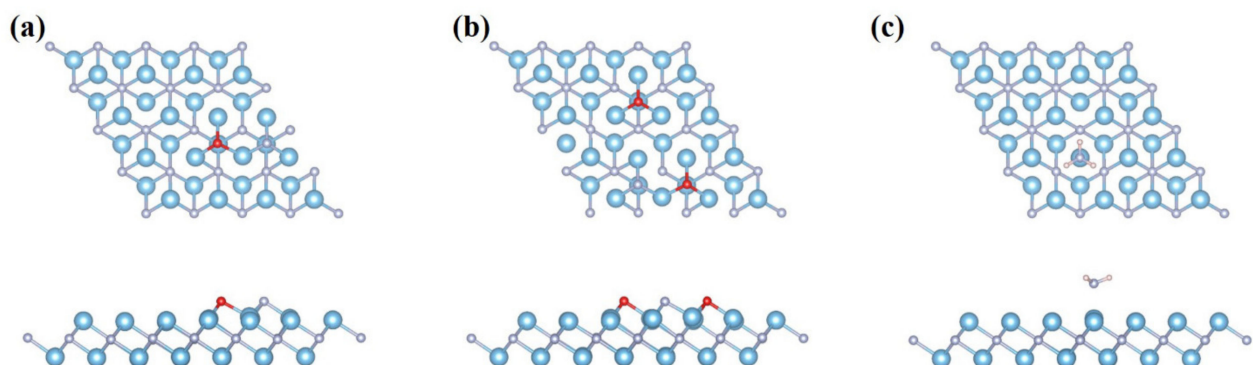
First of all, the Ti<sub>2</sub>N MXene has a hexagonal structure with three layers of atoms: the middle is N-atoms layer; and the upper (a) and lower (b) layers are Ti atoms, as is shown in Figure 1a. After structural optimizations, the Ti–N bond length is 2.07 Å, the bond angle of the same layer Ti and N (Ti–N–Ti) and N–Ti–N are both 92.4°, and the bond angle of different-layer Ti and N (Ti–N–Ti) is 87.6°. The lattice parameter of the 4 × 4 supercell is a = b = 11.98 Å. These parameters are very consistent with the results of Ti<sub>2</sub>N studied in other works in the literature [37,38], indicating the reliability of the calculation results. The band structure and PDOS of the material are shown in Figure 1b; it can be seen that there

are many electronic states available for occupation at the Fermi level, and  $\text{Ti}_2\text{N}$  exhibits metallic properties. The states near the Fermi level are occupied by Ti electrons, where N's contribution is relatively small, meaning that low-energy carriers mainly originate from Ti.

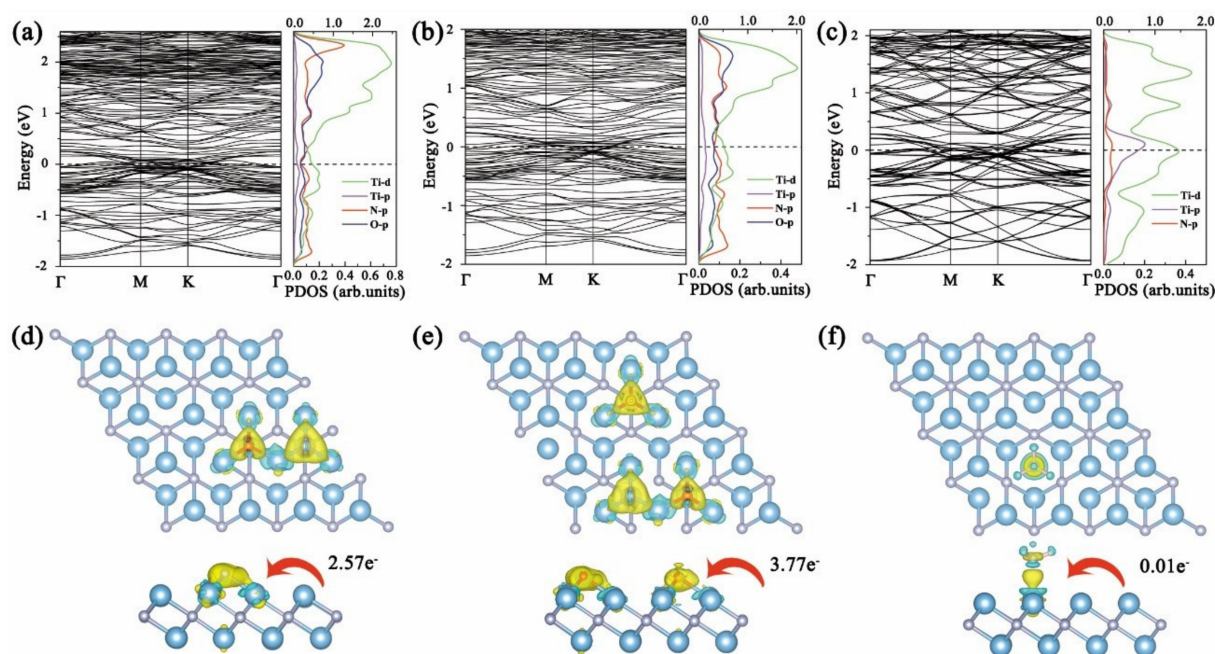


**Figure 1.** (a) Top and side view of the optimized structures for  $\text{Ti}_2\text{N}$  monolayer (blue and silver balls represent Ti and N atoms, respectively). (b) Band structure and projected density of states of  $\text{Ti}_2\text{N}$  monolayer (Ti *d*-orbital uses top abscissa and other orbitals use bottom coordinate).

In order to study the most stable adsorption configurations of NCGs on  $\text{Ti}_2\text{N}$  MXene, the NO,  $\text{NO}_2$ , and  $\text{NH}_3$  molecules were respectively placed at three different adsorption sites, namely on top of the Ti(a) atom shown in Figure 1a (site “a”), on top of the hollow Ti(b) atom (site “b”), or on top of the N atom (site “c”). The screened most stable configurations after adsorbing NCGs are shown in Figure 2. Both NO and  $\text{NO}_2$  show dissociative adsorption. The adsorption energy is far greater than those of other two-dimensional materials [39–44]. When NO is adsorbed on  $\text{Ti}_2\text{N}$ , N and O atoms are adsorbed at the adjacent site “b” respectively. The distance between N and O is 3.07 Å, which is much larger than 1.15 Å of free NO gas molecule, indicating that the N–O bond is broken and NO dissociates. The dissociated N and O atoms form new bonds with Ti atoms in the upper layer of  $\text{Ti}_2\text{N}$ , with bond lengths of 1.93 and 1.97 Å, respectively. Ti atoms protrude slightly upwards from its binding, and the Ti–N bond length on the surrounding matrix increases by about 0.14 Å. It can be seen from Figure 3a that, after dissociation, the N 2*p* and O 2*p* in NO and the nearby Ti 3*d*-orbitals have similar peak positions and obvious *p*–*d* hybridization in PDOS diagram. There is a strong bonding between NO and  $\text{Ti}_2\text{N}$ , and the adsorption energy is as high as −9.809 eV. At the same time, the band structure shows that the dissociative adsorption of NO on  $\text{Ti}_2\text{N}$  does not change the original metallic properties.



**Figure 2.** Top and side view of the optimized atomic structures for  $\text{Ti}_2\text{N}$  monolayer after adsorbing (a) NO, (b)  $\text{NO}_2$ , and (c)  $\text{NH}_3$  (blue, silver, red, and pink balls represent Ti, N, O, and H atoms, respectively).



**Figure 3.** Band structures, PDOS and differential charge densities of  $\text{Ti}_2\text{N}$  after adsorbing (a,d) NO, (b,e)  $\text{NO}_2$ , and (c,f)  $\text{NH}_3$  (Ti-d-orbital uses top abscissa and other orbitals use bottom coordinate; isovalue of NO and  $\text{NO}_2$  is set to  $0.01 \text{ e}/\text{\AA}^3$ ,  $\text{NH}_3$  is set to  $0.005 \text{ e}/\text{\AA}^3$ ).

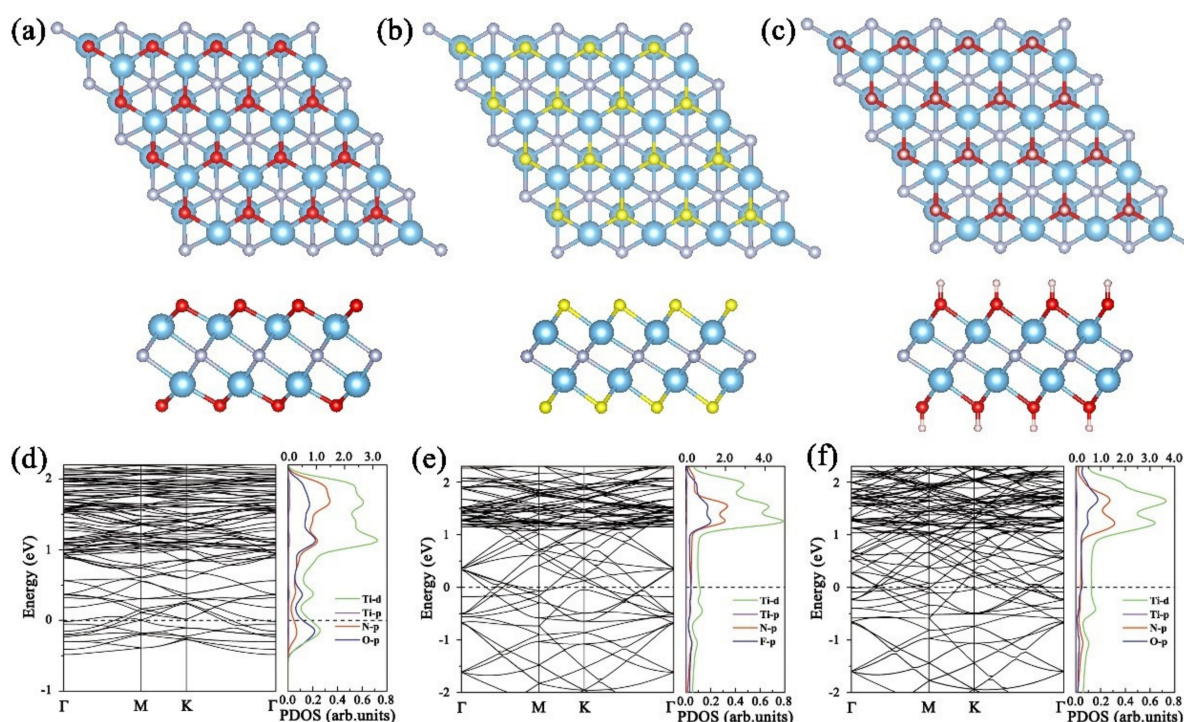
When  $\text{NO}_2$  interacts with  $\text{Ti}_2\text{N}$ , it is interesting that  $\text{NO}_2$  will also dissociate similar to NO adsorption, and all the N and O atoms are adsorbed at the site “b”. One O-atom keeps the same shape adsorbed at adjacent position, while the another O-atom is adsorbed away from it. The N-O bond length in  $\text{NO}_2$  is increased to  $3.07 \text{ \AA}$ . The dissociative adsorption energy is up to  $-14.336 \text{ eV}$ . As can be seen from the electronic structure in Figure 3b, the *p-d* hybridization of N  $2p$ -, O  $2p$ -, and Ti  $3d$ -orbitals corresponds with the strong bonding of Ti-N and Ti-O, respectively. The metallic behavior of  $\text{Ti}_2\text{N}$  remains after  $\text{NO}_2$  dissociative adsorption. From the differential charge density in Figure 3d–f, it can be seen that the dissociation of NO and  $\text{NO}_2$  is caused by the excessive charge transfer from  $\text{Ti}_2\text{N}$  MXene to the gas molecule. The Bader charge analysis shows that  $2.57 \text{ e}^-$  and  $3.77 \text{ e}^-$  charges were transferred from  $\text{Ti}_2\text{N}$  to NO and  $\text{NO}_2$ , respectively. The bulk of these transferred charges comes from the surrounding Ti atoms, and this is also the fundamental reason why the N-O bond is broken.

The adsorption of  $\text{NH}_3$  is completely different from those of NO and  $\text{NO}_2$ . It is bound by associative adsorption with the binding distance of  $2.2 \text{ \AA}$  and the  $E_{\text{ads}}$  value of  $-1.286 \text{ eV}$ . As shown in Figure 3c,  $\text{NH}_3$  tends to be adsorbed at site “a” with N-atom below and three H-atoms evenly distributed above. The distance between the Ti atom at site “a” and the surrounding N atom is elongated to  $2.14 \text{ \AA}$ , and the N-Ti-N bond angle is reduced to  $88.3^\circ$ . To figure out the binding and charge transfer mechanism between  $\text{NH}_3$  molecule and  $\text{Ti}_2\text{N}$  monolayer, the band structure and PDOS were analyzed. It is difficult to find a pronounced overlap between N *p* and Ti *d* close to the Fermi level in Figure 3c. The amount of charge transfer is only  $0.01 \text{ e}^-$ , which is far less than those of NO and  $\text{NO}_2$ . The bonding between  $\text{NH}_3$  and  $\text{Ti}_2\text{N}$  is very weak, and the band structure that is almost the same as the bare  $\text{Ti}_2\text{N}$  monolayer also confirms this.

For ideal gas sensing application, excessive adsorption energy is obviously not feasible. The dissociative adsorption of NO and  $\text{NO}_2$  on  $\text{Ti}_2\text{N}$  monolayer will obviously poison the substrate and make it fail to work cyclically. The dissociation of gas molecules results from the strong interaction between the bare  $\text{Ti}_2\text{N}$  surface and O atoms. Thus, the pure  $\text{Ti}_2\text{N}$  monolayer can be used as a kind of gas-trapping material instead of highly efficient sensor for NO and  $\text{NO}_2$ . Compared with those two gases, the adsorption effect of  $\text{NH}_3$

is satisfactory. The associative adsorption with lower energy will not cause chemical poisoning, which is beneficial to  $\text{Ti}_2\text{N}$  as a promising candidate material for reversible  $\text{NH}_3$  sensors. In actual preparation process, some functional groups are usually introduced onto the  $\text{Ti}_2\text{N}$  monolayer, and these functional groups may have significant impacts on adsorption properties of the material. Considering this reality, in the following part, we recount how we studied the adsorption behaviors of the  $\text{Ti}_2\text{N}$  monolayer terminated by O, F, and OH towards NCGs.

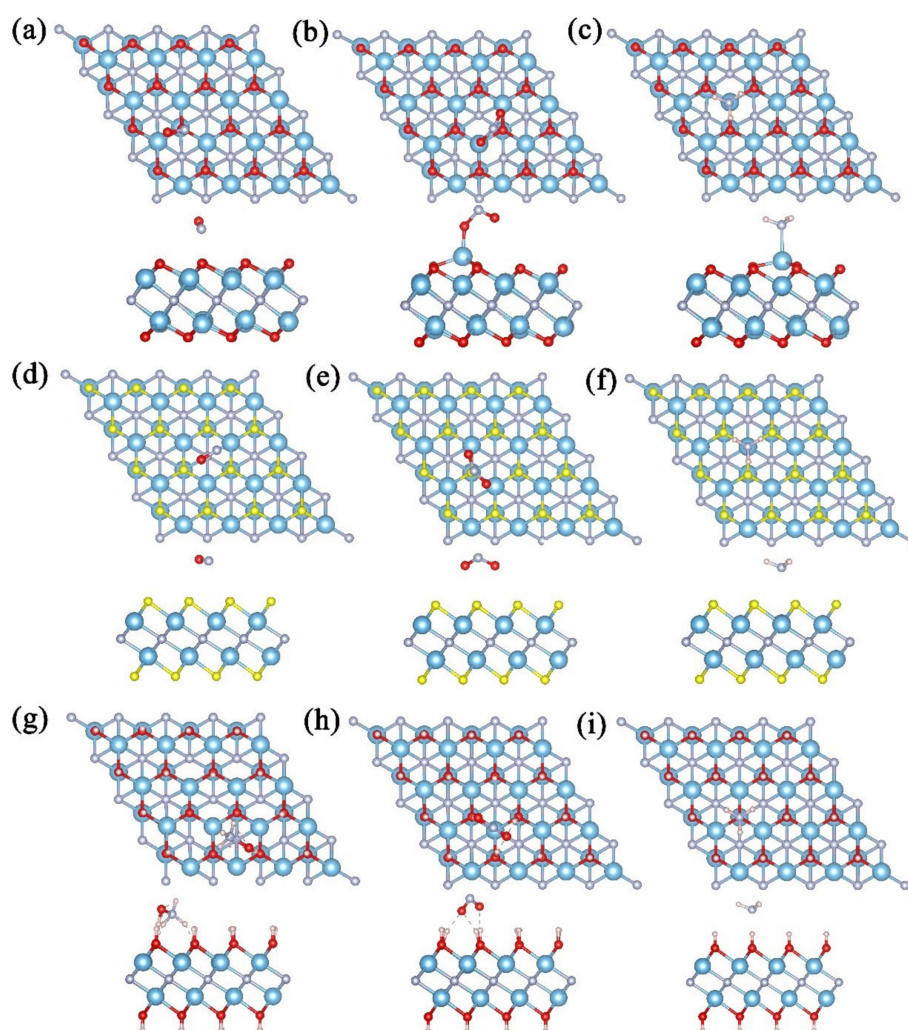
We introduce O, F, and OH functional groups at three adsorption positions (site “a”, “b”, and “c”) of  $\text{Ti}_2\text{N}$ , of which OH only considers the vertical direction (O atom is close to the matrix). All the functional groups tend to adsorb at site “b” with the lowest energy, as shown in Figure 4a–c. After adsorbing different functional groups, the Ti–N bonds show different degrees of extension: 2.164, 2.052, and 2.065 Å for  $\text{Ti}_2\text{NO}_2$ ,  $\text{Ti}_2\text{NF}_2$ , and  $\text{Ti}_2\text{N}(\text{OH})_2$ , respectively, which is due to the bonding between Ti and functional groups. Gouveia et al. reported possible atomic-layer-stacking-phase transitions of some MXenes [41]; thus, we also carefully checked this point in our research. It is found that the material system maintains phase stability and stacking way after introduction of those functional groups. The Bader charge analysis shows that 1.1, 0.8, and 0.8  $e^-$  are respectively transferred to O, F, and OH from  $\text{Ti}_2\text{N}$ . It can also be observed in Figure 4d–f that there is a strong overlap between the  $\text{Ti}d$ -orbital and the T (O, F, or OH) $d$ -orbitals.



**Figure 4.** Top and side view of the optimized structures, band structures, and PDOS of (a,d)  $\text{Ti}_2\text{NO}_2$ , (b,e)  $\text{Ti}_2\text{NF}_2$ , and (c,f)  $\text{Ti}_2\text{N}(\text{OH})_2$  monolayers (blue, silver, red, pink, and yellow balls represent Ti, N, O, H, and F atoms, respectively.  $\text{Ti}d$ -orbital uses top abscissa and other orbitals use bottom coordinate).

After obtaining the stable configurations of  $\text{Ti}_2\text{NT}_2$  (T=O, F, OH), we placed various NCGs molecules at different adsorption sites (site “a”, “b”, and “c”), such as those in pure  $\text{Ti}_2\text{N}$ . It can be seen from the most stable adsorption structures after optimizations in Figure 5 that all NCGs have no dissociative adsorption on  $\text{Ti}_2\text{NT}_2$  (T=O, F, OH). The binding distance is defined as the shortest direct distance between the NCGs molecule and the matrix; hereby, the binding distances of NO,  $\text{NO}_2$ , and  $\text{NH}_3$  on  $\text{Ti}_2\text{NO}_2$  are 2.082, 1.824, and 2.172 Å, respectively. The bond lengths of NO and  $\text{NH}_3$  after adsorption are the same as those of free gas molecules, while one of the N–O bonds of  $\text{NO}_2$  is stretched to 1.530 Å.

The 2-D  $Ti_2NO_2$  does not change significantly under NO environment, while the Ti atoms exposed to the adsorption positions of  $NO_2$  and  $NH_3$  are obviously dragged out. For NCGs adsorption on the surface of  $Ti_2NF_2$ , the molecular bond lengths and the matrix are almost unchanged with the adsorption distances of 2.951, 2.800, and 2.947 Å, respectively. It is surprising when NO is adsorbed on  $Ti_2N(OH)_2$ . The H atoms on the substrate bind with NO, where three H atoms combine with the N atom and one H atom combines with the O atom to form a hydroxyl ammonium-like product. The adsorption distances of  $NO_2$  and  $NH_3$  are 1.625 and 1.702 Å, which are measured from the nearest H atom; thus, they are shorter than those of other  $Ti_2NT_2$  (T=O, F). While, the adsorption is not strong, especially the  $NH_3$  molecule has basically no change compared with the free molecule.



**Figure 5.** Top and side view of the optimized structures for  $Ti_2NO_2$  after adsorbing (a) NO, (b)  $NO_2$ , and (c)  $NH_3$ ;  $Ti_2NF_2$  after adsorbing (d) NO, (e)  $NO_2$ , and (f)  $NH_3$ ; and  $Ti_2N(OH)_2$  after adsorbing (g) NO, (h)  $NO_2$ , and (i)  $NH_3$  (blue, silver, red, pink, and yellow balls represent Ti, N, O, H, and F atoms, respectively).

We can clearly see that the absolute values of adsorption energy of NCGs on  $Ti_2NT_2$  (T=O, F, OH) decrease sharply compared with the pure substrate in Table 1, especially on the surface of  $Ti_2NF_2$ . Since the H atoms of the OH functional groups fall off the  $Ti_2N(OH)_2$  substrate and bind with NO, their decrease is not that large. On the pure  $Ti_2N$  monolayer, a large amount of charge is transferred from  $Ti_2N$  to NO and  $NO_2$ . After the introduction of the functional group O, F, or OH, most of the charge transfer from  $Ti_2N$  is captured by these functional groups, which leaves a very small charge amount for exchange with NCGs.

In addition, one can see that the difference in adsorption energy of each NCG on substrate is obvious, indicating that the selectivity is good. As mentioned above, NO and NO<sub>2</sub> are dissociatively adsorbed on pure Ti<sub>2</sub>N, which results in the poisoning of the Ti<sub>2</sub>N surface. Similarly, NO binds with H atoms of Ti<sub>2</sub>N(OH)<sub>2</sub> and destroys the surface integrity of the substrate, which is detrimental to the sensing material. Thus, Ti<sub>2</sub>N(OH)<sub>2</sub> is not suitable for NO sensing. From the perspective of adsorption energy, those of NO on Ti<sub>2</sub>NO<sub>2</sub>, NO<sub>2</sub> on Ti<sub>2</sub>NF<sub>2</sub>, and NH<sub>3</sub> on Ti<sub>2</sub>N(OH)<sub>2</sub> are between physisorption and chemisorption, which are important to make sensors work reversibly and increase their efficiency [45].

**Table 1.** Adsorption energy ( $E_{\text{ads}}$ ) and charge-transfer profiles of NCGs adsorption on 2-D Ti<sub>2</sub>N and Ti<sub>2</sub>NT<sub>2</sub>(T=O, F, OH).

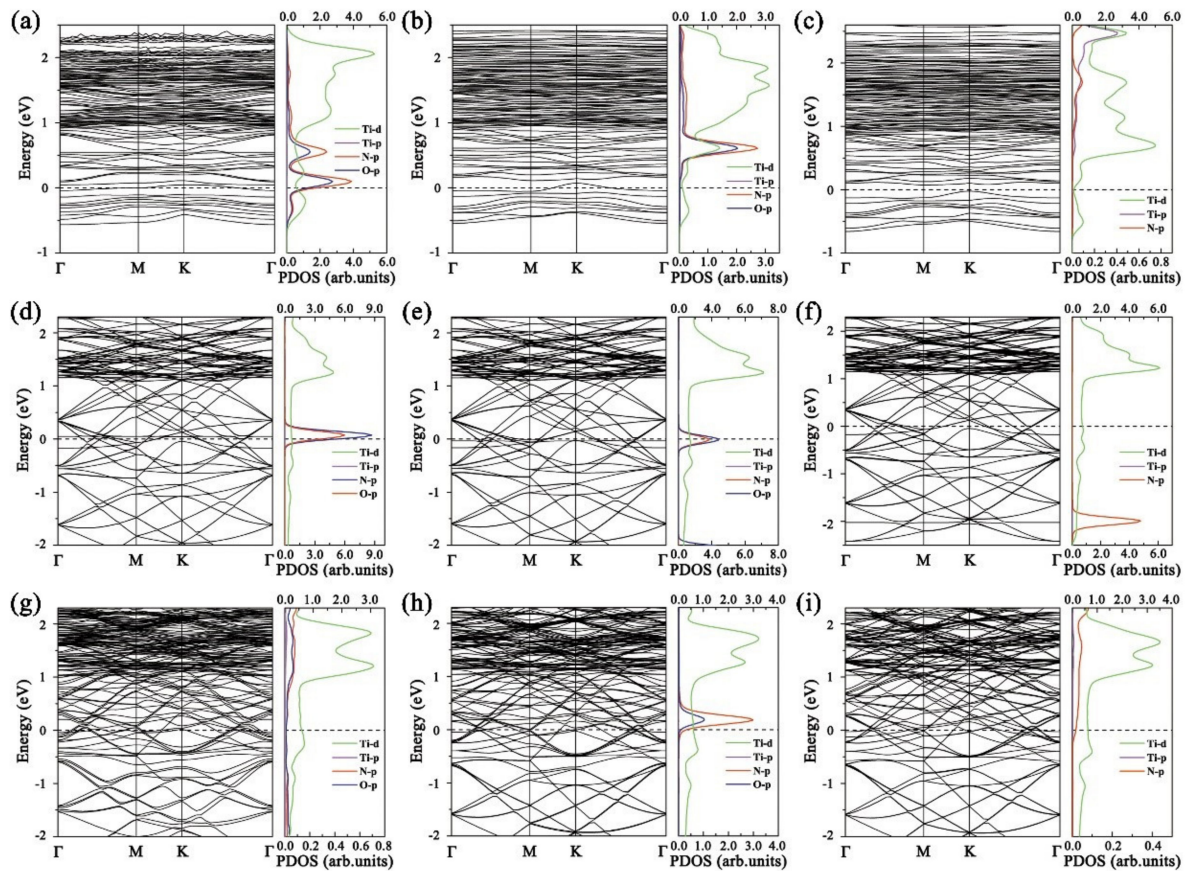
NCGs	Adsorption Energy (eV)				Charge Transfer (e <sup>-</sup> )			
	Ti <sub>2</sub> N	Ti <sub>2</sub> NO <sub>2</sub>	Ti <sub>2</sub> NF <sub>2</sub>	Ti <sub>2</sub> N(OH) <sub>2</sub>	Ti <sub>2</sub> N	Ti <sub>2</sub> NO <sub>2</sub>	Ti <sub>2</sub> NF <sub>2</sub>	Ti <sub>2</sub> N(OH) <sub>2</sub>
NO	-9.81	-0.68	-0.15	-5.55	2.57	-0.24	-0.07	0.50
NO <sub>2</sub>	-14.34	-2.38	-0.36	-3.78	3.77	0.49	0.29	1.78
NH <sub>3</sub>	-1.29	-1.16	-0.13	-0.57	0.01	-0.18	-0.03	0.15

To better understand the effects of NCGs molecules on Ti<sub>2</sub>NT<sub>2</sub>(T=O, F, OH), the band structure and PDOS of all the adsorption systems are calculated and shown in Figure 6. Regardless of the type of NCGs adsorption, the system retains metallic properties. Only when Ti<sub>2</sub>NO<sub>2</sub> is exposed to NO<sub>2</sub>, the Ti *d*-, O *p*-, and N *p*-orbitals show a certain degree of d-p hybridization, while the *p*- and *d*-orbitals in other PDOS do not overlap. When Ti<sub>2</sub>NO<sub>2</sub> adsorbs NO, a flat band appears near the Fermi energy and indicates that the electrons are delocalized, which further proves that the interaction between NO and Ti<sub>2</sub>NO<sub>2</sub> is not strong. This phenomenon is more obvious when NCGs are adsorbed on Ti<sub>2</sub>NF<sub>2</sub> shown in Figure 6d–f. These results are consistent with the above discussion on energetics and charge-transfer profiles.

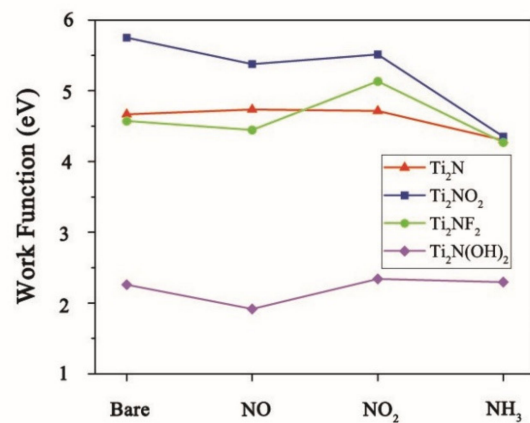
An ideal sensing material not only needs to possess excellent gas-binding capability but also has a requirement for good sensitivity. The conductivity of materials, especially two-dimensional materials, is closely related with work function. During adsorption, if the electron affinity of gas molecule surpasses the work function of substrate materials, the molecule will grab electrons from the surface of materials and the molecule would be charged negatively; in contrast, the adsorbed molecule would be charged positively. Both two cases would change the electronic structure, as well as conductivity of materials. The work function [46] is defined as the least energy required to free an electron from the surface of a system, which can be calculated by the following equation:

$$\varphi = V_{\infty} - E_f \quad (2)$$

where  $\varphi$ ,  $V_{\infty}$ , and  $E_f$  are the work function, electrostatic potential at the vacuum level, and the Fermi energy, respectively. Figure 7 shows the change of work function of Ti<sub>2</sub>N and Ti<sub>2</sub>NT<sub>2</sub>(T=O, F, OH) before and after adsorption of various NCGs. "Bare" represents the bare system (before adsorption), and the names of NCGs in abscissa represent various systems after adsorption of different gases. The work functions of Ti<sub>2</sub>NO<sub>2</sub> and Ti<sub>2</sub>N(OH)<sub>2</sub> decrease obviously after adsorbing NO, while that of Ti<sub>2</sub>NF<sub>2</sub> increases significantly after adsorbing NO<sub>2</sub>. After adsorbing NH<sub>3</sub>, the work functions of Ti<sub>2</sub>N, Ti<sub>2</sub>NF<sub>2</sub>, and Ti<sub>2</sub>NO<sub>2</sub> all decrease, and the change of Ti<sub>2</sub>NO<sub>2</sub> is the most obvious. Because all the systems are metallic, even small changes in work function can cause large variation in conductivity [43]. Comprehensively considering the above results on adsorption structure and energy, Ti<sub>2</sub>NO<sub>2</sub> can be used as a superior sensing candidate material for NO and NH<sub>3</sub>, and Ti<sub>2</sub>NF<sub>2</sub> can be utilized as that for NO<sub>2</sub>. In addition, the potential of sensing properties of Ti<sub>2</sub>N, Ti<sub>2</sub>NF<sub>2</sub>, and Ti<sub>2</sub>N(OH)<sub>2</sub> towards NH<sub>3</sub> is also good.



**Figure 6.** Band structures and PDOS of  $\text{Ti}_2\text{NO}_2$  after adsorbing (a) NO, (b)  $\text{NO}_2$ , and (c)  $\text{NH}_3$ ;  $\text{Ti}_2\text{NF}_2$  after adsorbing (d) NO, (e)  $\text{NO}_2$ , and (f)  $\text{NH}_3$ ;  $\text{Ti}_2\text{N}(\text{OH})_2$  after adsorbing (g) NO, (h)  $\text{NO}_2$ , and (i)  $\text{NH}_3$  (*Ti**d*-orbital uses top abscissa and other orbitals use bottom coordinate).



**Figure 7.** Work function change of  $\text{Ti}_2\text{N}$  and  $\text{Ti}_2\text{NT}_2$  ( $\text{T}=\text{O}, \text{F}, \text{OH}$ ) before and after adsorption of various NCGs.

Finally, we discuss the recovery time ( $\tau$ ). An ideal gas-sensing material needs appropriate adsorption energy, and a shorter recovery time is better. To make a gas sensor work efficiently, a short recovery time is required. It can be calculated through transition state theory, namely by the following relationship [47,48]:

$$\tau = \nu^{-1} \exp(-E_{\text{ads}}/kT) \quad (3)$$



where  $\nu$ ,  $E_{\text{ads}}$ ,  $k$ , and  $T$  are the attempt frequency, adsorption energy, Boltzmann's constant ( $8.62 \times 10^{-5} \text{ eV K}^{-1}$ ), and operational temperature, respectively. Under the same conditions, the increase in adsorption energy will lead to an exponential increase in recovery time; then the higher operating temperature compensation is required, which is detrimental to the sensor performance. In order to get a relatively reasonable value, we set the attempt frequency here to  $10^{12} \text{ s}^{-1}$  based on traditional transition state theory [49–51]. The recovery time of NO on  $\text{Ti}_2\text{NO}_2$  is calculated to be  $2.16 \times 10^5 \text{ us}$ , and that of  $\text{NO}_2/\text{Ti}_2\text{NF}_2$  and  $\text{NH}_3/\text{Ti}_2\text{NF}_2$  is 1.29 and  $1.71 \times 10^{-4} \text{ us}$ , respectively. Here it can be seen in Figure 8 that the adsorption energy difference of less than 1 eV can be enlarged to several orders of magnitude in recovery time. Although NO on  $\text{Ti}_2\text{NO}_2$  has the longest recovery time, the value is only 0.216 s. For the system of  $\text{NH}_3/\text{Ti}_2\text{NF}_2$ , it has the shortest recovery time, which shows that the MXene system in this research can be promising. Therefore, when other conditions are met, the candidate materials with appropriate gas adsorption energy and short response time should be selected.

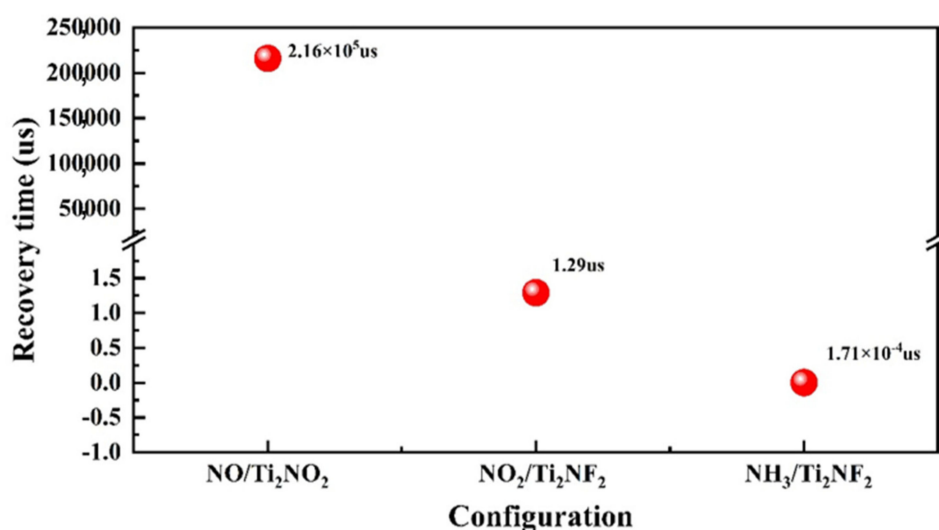


Figure 8. Recovery time of different sensing systems.

#### 4. Summary and Outlook

In this work, we used the first-principles DFT calculations to investigate the 2-D  $\text{Ti}_2\text{N}$  nanosheet and its  $\text{Ti}_2\text{NT}_2$  ( $T=\text{O}, \text{F}, \text{OH}$ ) derivative materials for harmful nitrogen-containing gas (NCGs) adsorption and sensing applications. When NCGs are exposed on the surface of  $\text{Ti}_2\text{N}$ , the dissociative adsorption of NO and  $\text{NO}_2$  occurs and destroys the stability of the nanosheet, which would make it hard to work reversibly, while  $\text{NH}_3$  does not have this effect on  $\text{Ti}_2\text{N}$ . When studying the interaction mechanisms between NCGs and  $\text{Ti}_2\text{NT}_2$  ( $T=\text{O}, \text{F}, \text{OH}$ ), it can be found that a large amount of charges accumulates at functional groups, thus greatly reducing the adsorption energy of NCGs in absolute value. The obvious change in work function, coupled with the metallic nature of the systems, would improve gas sensitivity. Appropriate gas adsorption energy can directly determine fast recovery time to ensure the efficiency and reversibility of sensors. This research gives the electronic structures origin and application prospect of  $\text{Ti}_2\text{NT}_2$  ( $T=\text{O}, \text{F}, \text{OH}$ ) as efficient NCGs sensing materials and would inspire experimentalists to explore better 2-D candidates in the field.

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