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# High-sensitivity methane detection based on QEPAS and H-QEPAS technologies combined with a self-designed 8.7 kHz quartz tuning fork

Tiantian Liang <sup>a, b</sup>, Shunda Qiao <sup>a, b</sup>, Yanjun Chen <sup>a, b</sup>, Ying He <sup>a, b</sup>, Yufei Ma <sup>a, b, \*</sup>

<sup>a</sup> National Key Laboratory of Science and Technology on Tunable Laser, Harbin Institute of Technology, Harbin 150001, China
 <sup>b</sup> Zhengzhou Research Institute, Harbin Institute of Technology, Zhengzhou 450000, China

#### ARTICLE INFO

## ABSTRACT

Keywords: CH<sub>4</sub> detection Self-designed quartz tuning fork Quartz-enhanced photoacoustic spectroscopy Heterodyne quartz-enhanced photoacoustic spectroscopy Methane (CH<sub>4</sub>) is a greenhouse gas as well as being flammable and explosive. In this manuscript, quartzenhanced photoacoustic spectroscopy (QEPAS) and heterodyne QEPAS (H-QEPAS) exploring a self-designed quartz tuning fork (QTF) with resonance frequency ( $f_0$ ) of ~8.7 kHz was utilized to achieve sensitive CH<sub>4</sub> detection. Compared with the standard commercial 32.768 kHz QTF, this self-designed QTF with a low  $f_0$  and large prong gap has the merits of long energy accumulation time and low optical noise. The strongest line located at 6057.08 cm<sup>-1</sup> in the 2v<sub>3</sub> overtone band of CH<sub>4</sub> was chosen as the target absorption line. A diode laser with a high output power of > 30 mW was utilized as the excitation source. Acoustic micro-resonators (AmRs) were added to the sensor architecture to amplify the intensity of acoustic waves. Compared to the bare QTF, after the addition of AmRs, a signal enhancement of 149-fold and 165-fold were 711 ppb and 1.06 ppm for QEPAS and H-QEPAS sensors. Furthermore, based on Allan variance analysis the MDLs can be improved to 19 ppb and 27 ppb correspondingly. Compared to the QEPAS sensor, the H-QEPAS sensor shows significantly shorter measurement timeframes, allowing for measuring the gas concentration quickly while simultaneously obtaining  $f_0$  of QTF.

## 1. Introduction

Methane (CH<sub>4</sub>), as an important carbon-based resource, is closely related to human life and widely utilized in various fields, such as energy [1,2], medical diagnosis [3,4], chemical industry [5,6], and so on. However, CH<sub>4</sub> deserves special attention during extraction, transportation and use. On the one hand, CH<sub>4</sub> is not only the second most significant greenhouse gas in the world next to carbon dioxide  $(CO_2)$  [7], but its Global Warming Potential (GWP) is 34 times greater than CO<sub>2</sub> on a century timescale [8]. According to statistics, CH<sub>4</sub> concentration in the atmosphere reached 1922.26 ppb by May 2023, which is more than 260% of pre-industrial levels [9]. Unnoticed methane leaks from both biosphere and anthropogenic activity further contribute to the global greenhouse effect [10]. Additionally, CH<sub>4</sub> is a colorless, odorless, non-toxic, but flammable gas, with lower and upper explosive limit concentrations of 4.95% and 14.95% respectively [11]. When the concentration of CH<sub>4</sub> falls within this range and encounters an open flame, it can cause an explosion. Therefore, real-time detection of CH<sub>4</sub> content is crucial.

Various technologies for gas detection such as the thermal, mechanical, electrochemical, and optical have been developed in recent years [12–23]. Among these, laser absorption spectroscopy (LAS) stands out as an optical technology that offers numerous benefits compared to others, such as fast responses, long-term stability, low baseline drift, and high sensitivity and selectivity [24-32]. Quartz-enhanced photoacoustic spectroscopy (QEPAS) is a commonly utilized LAS based gas sensing technology, which was first proposed in 2002 [33]. In QEPAS method, after absorbing the specific wavelength of the laser, the target gas molecules will release heat to the surrounding environment through the non-radiative transition. Based on the photoacoustic effect [34], the continuous emission of acoustic waves will occur at the position of contact between the laser and the gas molecules due to the periodic modulated light beam [35-39]. These acoustic waves will induce vibrations in the detector of quartz tuning fork (QTF) within QEPAS sensing system, thereby exploiting its piezoelectric effect to generate the piezoelectric signal on its surface that can be utilized for gas concentration inversion. Compared to traditional acoustic detection devices, such as microphone [40-43], QTF possesses the advantages of high

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<sup>\*</sup> Corresponding author at: National Key Laboratory of Science and Technology on Tunable Laser, Harbin Institute of Technology, Harbin 150001, China. *E-mail address:* mayufei@hit.edu.cn (Y. Ma).

Q-factor, narrow bandwidth, strong anti-noise capability, affordability, and compact size [44–48]. Consequently, QEPAS technology gained significant attention in the field of gas sensing [49–51].

The resonance frequency ( $f_0$ ) of QTF is a very important parameter for QEPAS based gas sensing. It determines the frequency for laser modulation and signal demodulation, and therefore it should be calibrated frequently, especially when the environment such as pressure and temperature changes. However, the frequency calibration in traditional QEPAS not only interrupts the monitoring process but also consumes a significant amount of time. The heterodyne quartz-enhanced photoacoustic spectroscopy (H-QEPAS), also known as beat frequency quartz-enhanced photoacoustic spectroscopy (BF-QEPAS), was initially put forward in 2017 by Dong et al. [49], offering the capability to simultaneously measure both gas concentration and  $f_0$  of QTF in just a few seconds [52,53]. By using H-QEPAS technology to measure QTF frequency, the monitoring process remains uninterrupted and the time needed for frequency calibration is greatly reduced.

In this manuscript, a highly sensitive CH<sub>4</sub> sensor based on QEPAS and H-QEPAS technologies was demonstrated respectively. A selfdesigned QTF with a low resonance frequency of ~8.7 kHz and a large prong gap possessing the merits of long energy accumulation time and low optical noise was employed as the acoustic wave detector. A continuous wave, distributed feedback (CW-DFB) diode laser operating at a center emission wavelength of 1650.96 nm with a high output power of > 30 mW was utilized as the excitation source. A set of acoustic micro-resonators (AmRs) were incorporated on both sides of the QTF to amplify the intensity of acoustic waves. The enhancement effect of different sizes of AmRs was measured separately in the two sensors. Detailed research for CH<sub>4</sub> sensing was performed.

## 2. Experimental setup

It is crucial for QEPAS system to carefully choose the gas absorption line taking into consideration compatibility with the existing laser source, minimal overlap with other gas absorption lines, and maximizing line strength. Given the well-developed manufacturing technology for DFB laser diodes in the near-infrared band [54,55], the  $2v_3$ overtone band of CH<sub>4</sub> around 1660 nm is noticed [56]. According to the HITRAN 2020 database [57], simulations were conducted for CH<sub>4</sub>'s absorption coefficient in this band at 296 K temperature, 1 atm pressure, and a 1 cm optical path length. As shown in Fig. 1(a), the absorption line located at 6057.08 cm<sup>-1</sup> (1650.96 nm) with the strongest line strength within this band was selected as the target absorption line. To avoid



Fig. 1. Simulation for absorption coefficient based on the HITRAN 2020 database at 296 K temperature, 1 atm pressure, and 1 cm optical path length. (a) The  $2\nu_3$  overtone band of CH<sub>4</sub> around 6024 cm<sup>-1</sup>. (b) The comparison of absorption coefficient of CH<sub>4</sub>, CO<sub>2</sub> and H<sub>2</sub>O in the range from 6048 to 6066 cm<sup>-1</sup>.

interference from common atmospheric gases, simulations for absorption coefficien were performed for  $CO_2$  and  $H_2O$ . As depicted in Fig. 1 (b), there is negligible overlap among  $CH_4$ ,  $CO_2$  and  $H_2O$ .

A CW-DFB diode laser with a center emission wavenumber of 6057.08 cm<sup>-1</sup> was employed as an excitation source to excite the target line. The output characteristic curve of the used laser is presented in Fig. 2. As the injection current increases, the output power exhibits an upward trend whereas the emission wavenumber of the laser shows a downward trend. When it is operated at a temperature of 31 °C and an injection current of 162.9 mA, the laser's emission wavenumber co-incides with the target CH<sub>4</sub> absorption line.

The schematic diagram of both QEPAS and H-QEPAS sensing systems is illustrated in Fig. 3. The laser emitted from the pigtail is collimated first by a fiber collimator. The beam enters the gas cell and passes through the two AmRs and QTF. The gap between the two prongs of QTF is 1.73 mm, which is larger than the usually used QTF and is beneficial to reducing optical loss and noise [58,59]. Ultimately the laser beam is detected by a power meter to monitor the optical alignment. The absorption of the modulated laser by the gas molecules in the gas cell generates acoustic waves that induce vibration in the OTF. The presence of AmRs enhances the vibration amplitude on the OTF, resulting in a larger piezoelectric signal and better detection performance [60–62]. This signal is subsequently demodulated with a lock-in amplifier. In the QEPAS system, wavelength modulation and second harmonic (2 f) demodulation technology were employed to suppress effectively background noise. A high-frequency sine wave and a low-frequency ramp wave were combined and then sent into the laser controller for wavelength modulation of the diode laser. To obtain the steady-state response of QTF, a long period ( $\sim 1 \text{ min}$ ) ramp wave is often used to sweep the entire gas absorption line. The frequency of the sine wave is set to half of the resonance frequency  $(f_0/2)$  of the used QTF. By employing lock-in amplifier demodulation, the peak value of the 2f signal can be utilized for gas concentration retrieval. In the H-QEPAS system, wavelength modulation and first harmonic (1 f) demodulation technology were adopted. The laser wavelength is concurrently modulated by a high-frequency sine wave and a low-frequency ramp wave, with the latter being divided into fast-scanning stage and stable stage as depicted in Fig. 3. The frequency of the sine wave is denoted as  $f_m = f_0 \pm \Delta f$  [49]. When the ramp wave, adding a sine wave with frequency of  $f_m$ , rapidly scans across the gas absorption line, it allows for the acquisition of transient response from the QTF itself. The frequency of the transient response signal is equal to the resonance frequency  $f_0$  of the OTF, and its amplitude decays exponentially with time. Subsequently, this signal is demodulated at  $f_m$  by a lock-in amplifier, resulting in a heterodyne signal with  $\Delta f$  frequency. Gas concentration inversion can be determined based on the peak value of the heterodyne signal while simultaneously enabling measurement of QTF resonance frequency based on the equation  $f_0 = f_m \pm \Delta f$ .

## 3. Results and discussion

The frequency response curve of the self-designed QTF used in this manuscript is illustrated in Fig. 4(a). After square normalization of the measured data, Lorentz fitting was performed. The  $f_0$  was determined to be 8685.44 Hz with a bandwidth w = 0.85 Hz for the bare QTF. QTF's quality factor (Q value) could be calculated as  $Q = f_0/w \approx 10218$ . After adding AmR with length (*L*) of 18 mm and inner diameter (*ID*) of 0.9 mm the  $f_0$  and Q of the QTF are remeasured as 8685.22 Hz and 5949, respectively, which is shown in Fig. 4(b). The changes of  $f_0$  and Q resulted from the acoustic coupling between AmR and QTF. In the QEPAS system, the modulation frequency  $f_m$  of sine wave was set to  $f_0/2$ . In the H-QEPAS system, to make sure that the peak value of the heterodyne signal is large enough and the number of positive peaks is sufficient to calculate the QTF frequency,  $f_m = f_0 - \Delta f$  was chosen. The ramp wave utilized in the H-QEPAS system had been optimized. The optimized modulation frequency was 0.25 Hz, and the fast-scanning



Fig. 2. The output characteristic curve of CW-DFB diode laser for CH<sub>4</sub> sensing. (a) The emission wavenumber varies with the injection current and working temperature.



Fig. 3. (a) The schematic diagram of both QEPAS and H-QEPAS systems for CH<sub>4</sub> sensing. (b) The actual image of self-designed QTF.



Fig. 4. Frequency response characteristics of the self-designed QTF. (a) Bare QTF. (b) QTF with an L of 18 mm and an ID of 0.9 mm AmRs.

stage lasted 1.25 s. The filter slope and time constant of the filter in the lock-in amplifier were set to 30 dB/Oct and 20 ms respectively in both systems. The corresponding bandwidth was ~3 Hz which ensures that heterodyne signals with frequency  $\Delta f$  were not filtered out by the filter in lock-in amplifier.

In addition, the modulation depth of sine wave will influence the

detection performance [34]. When the operating temperature of the laser is fixed, the laser's emission wavenumber is determined by its injection current. Therefore, the modulation depth of sine wave was replaced by modulation current. The signal amplitudes of both QEPAS and H-QEPAS systems were measured respectively when they changed with the modulation current. As shown in Fig. 5, the signal amplitudes of



Fig. 5. The curve that depicts the variation of peak signal with the modulation current. (a) QEPAS system. (b) H-QEPAS system.

both systems initially rose and then declined as the modulation current increased. The optimal modulation current of the QEPAS system was determined as 4.34 mA. While for the H-QEPAS system, it was 5.90 mA. The poor high-frequency modulation performance of the CW-DFB laser results in a larger optimal modulation current of the latter [63].

The enhancement effect of QTF vibration is jointly influenced by the size of AmR, including its length (L) and inner diameter (ID). The optimal length ( $L_{opt}$ ) of AmR depends on  $f_0$  of the used QTF, and it should satisfy  $\lambda_s/4 < L < \lambda_s/2$  [60], where  $\lambda_s$  could be calculated as  $\lambda_s = c_s/f_0$  and  $c_s$  represents the sound velocity in the medium. When assigning a value of 340 m/s to  $c_s$ , the range for  $L_{opt}$  spans from 9.8 mm to 19.6 mm. To achieve optimal enhancement effect, AmRs with different lengths (L =17 mm, 18 mm, 19 mm, and 20 mm), different inner diameters (ID =0.8 mm, 0.9 mm, 1.0 mm, 1.1 mm, 1.3 mm, 1.5 mm, 1.6 mm, and 1.7 mm) and same outer diameters (OD = 2.5 mm) were designed. All kinds of AmRs were made of stainless steel. The coupling degree between QTF and AmRs varies depending on the sizes of the latter. In Fig. 4 (b), taking an example with L = 18 mm and ID = 0.9 mm, the QTF coupled with this set of AmRs exhibited a wider bandwidth of 1.46 Hz, a lower Q value of 5949, and a drift in frequency. Therefore, after each replacement of the set of AmRs, it is necessary to re-calibrate  $f_0$  for the OTF.

To assess the improvement of AmR on the signal level of the QEPAS sensor, measurements were conducted with varying *ID* and *L*, as depicted in Fig. 6. When keeping the *L* constant for AmR, the peak signal value of the QEPAS system initially increased and then decreased as the *ID* decreased. The optimal *ID* was determined to be 0.9 mm. Conversely, when maintaining a constant *ID* for AmR, the signal peak value of the



Fig. 6. QEPAS peak signal as a function of L and ID of the AmRs.

QEPAS system also exhibited an initial increase followed by a decrease with an increase in *L*. With the exception of the 1.7 mm *ID*, an optimal *L* of 19 mm was observed in other *IDs* of AmRs. However, for an *ID* of 1.7 mm, an optimal *L* of 18 mm was found. This is because the *L* of the optimal AmR gradually increases as the *ID* decreases [64].

With the optimized size configuration for AmR, signals from both QEPAS and H-QEPAS systems were measured at a CH<sub>4</sub> concentration level of 1%. The results are shown in Fig. 7(a) and (b), respectively. The noise levels were separately measured in pure nitrogen (N<sub>2</sub>) environments for both QEPAS and H-QEPAS systems, as shown in Fig. 8(a) and (b). The peak value obtained from the bare QTF-based QEPAS signal was recorded as 7.68 µV. However, when utilizing a QTF coupled with optimized AmRs, it increased significantly to 1155.66 µV, resulting in approximately a 149-fold enhancement. The system noise level was calculated to be 82.17 nV while yielding a signal-to-noise ratio (SNR) of 14064 and a minimum detection limit (MDL) of 711 ppb. In the H-QEPAS sensor system, the peak value of the signal obtained using a bare QTF was 4.61 µV. After adding the AmRs with optimum size, the signal peak value increased to 765.20 µV, leading to a 165-fold improvement. The calculated system noise, SNR, and MDL were 81.04 nV, 9442, and 1.06 ppm respectively. The QEPAS system can achieve a lower detection limit, but it requires approximately 1 min to obtain a complete 2fwaveform. Whereas the H-OEPAS system can capture the entire heterodyne signal waveform within just 4 s and simultaneously measure  $f_0$ for the QTF. In Fig. 7(b), the heterodyne signal frequency could be calculated as  $\Delta f = 1/\Delta t = 1.44$  Hz, and then the  $f_0$  of QTF can be retrieved according to the formula:  $f_0 = f_m + \Delta f = 8685.18$  Hz. Compared with the optical excitation method shown in Fig. 4, the measurement error is less than 0.16 Hz.

Finally, the long-term stability for both QEPAS and H-QEPAS sensor systems was evaluated when continuous monitoring of pure N2 was performed. The Allan deviation plot illustrates the impact of noise on system performance, and it was adopted for both systems. The obtained results are depicted in Fig. 9(a) and (b), respectively. It could be observed that white noise dominates and the Allan deviation decreased with increasing integration time at small integration time. However, as the integration time increased, random-walk drift started to dominate [65]. At this time, the Allan deviation instead increased, indicating an optimal integration time exists. Increasing the integration time of the QEPAS system to 100 s and 4560 s resulted in the MDLs of 103 ppb and 19 ppb, respectively. In contrast, when extending the H-QEPAS system's integration time to 100 s and 3200 s, the corresponding MDLs were 141 ppb and 27 ppb, respectively. In Table 1, we compared the performance obtained through QEPAS and H-QEPAS technologies in other references with the results in this manuscript. The selected CH<sub>4</sub> absorption line reported in this paper is located in the near-infrared band, and the obtained MDL of CH<sub>4</sub>-QEPAS sensor is superior to that of the other studies



Fig. 7. Comparison between the bare QTF and QTF with AmRs. (a) 2 f signal in the QEPAS system. (b) Heterodyne signal in the H-QEPAS system.



**Fig. 8.** The measured noise level for the two systems at pure N<sub>2</sub> environments. (a) QEPAS system. (b) H-QEPAS system.

in this band, but not as good as the results in the mid-infrared band because of the stronger absorption of mid-infrared band. The MDL of  $CH_4$  detection using H-QEPAS in this study is better than the literature when H-QEPAS technology is employed.

#### 4. Conclusion

In this paper, highly sensitive detection of CH<sub>4</sub> was achieved based on QEPAS and H-QEPAS technologies, respectively, by utilizing a selfdesigned QTF with  $f_0$  of ~8.7 kHz. This used QTF has such a low  $f_0$ and big prong gap of 1.73 mm possessing the merits of long energy accumulation time and low optical noise. A CW-DFB diode laser with high output power and a central emission wavelength of 1650.96 nm was selected as the excitation source to match the strongest absorption line of  $CH_4$  in the  $2\nu_3$  overtone band. To further enhance the detection performance of CH<sub>4</sub>, AmRs were equipped on both sides of the QTF to amplify its vibration amplitude. Through experimental investigation, the optimal size of AmRs was determined, which resulted in an enhancement of signal peak values by a factor of 149-fold and 165-fold compared to the bare OTF for QEPAS and H-QEPAS systems, respectively. The corresponding MDLs for CH<sub>4</sub> detection were 711 ppb and 1.06 ppm for these two sensors. Furthermore, the long-term stability was analyzed using Allan deviation analysis. Increasing the integration time of the QEPAS system to 4560 s led to an MDL improvement down to 19 ppb. In contrast, when the integration time of the H-QEPAS sensor was increased to 3200 s, its MDL improved to 27 ppb. Although slightly inferior in terms of detection performance compared to the QEPAS system, the H-QEPAS system offers advantages such as fast measurement and obtaining  $f_0$  simultaneously along with gas concentration. In future research, we will continue to optimize the shape and size of the tuning fork, while exploring alternative materials with higher piezoelectric coefficients to substitute quartz.



Fig. 9. Allan deviation analysis. (a) QEPAS system. (b) H-QEPAS system.

#### Table 1

Comparison of CH<sub>4</sub> detection limits in QEPAS and H-QEPAS methods.

Technology	Wave band	Wavenumber ( $cm^{-1}$ )	MDL (ppm) (@Integration time)	Reference
QEPAS	Near- infrared	6057.08	3.2 (@1 s)	[66]
QEPAS	Near- infrared	4290.22	1.2 (@100 s)	[67]
QEPAS	Near- infrared	6046.95	0.288 (@125 s)	[68]
QEPAS	Mid-infrared	2788.64	55.57 (@1 s)	[49]
QEPAS	Mid-infrared	1275.04	0.028 (@0.1 s)	[69]
QEPAS	Mid-infrared	1295.5	0.048 (@0.1 s)	[70]
QEPAS	Mid-infrared	2989.5	0.09 (@1 s)	[71]
QEPAS	Mid-infrared	3038.5	0.05 (@1 s)	[72]
QEPAS	Mid-infrared	1275.04	0.013 (@1 s)	[73]
QEPAS	Near- infrared	6057.08	0.103 (@100 s)	This study
H-QEPAS	Near- infrared	6046.95	28.35 (@114 s)	[74]
H-QEPAS	Mid-infrared	2788.64	40.75 (@ 1 s)	[49]
H-QEPAS	Near- infrared	6057.08	0.141 (@100 s)	This study

#### CRediT authorship contribution statement

Ma Yufei: Writing – review & editing, Supervision, Conceptualization. Chen Yanjun: Validation. He Ying: Resources. Liang Tiantian: Writing – original draft, Investigation. Qiao Shunda: Formal analysis.

## **Declaration of Competing Interest**

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

## Data availability

Data will be made available on request.

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**Tiantian Liang** received his B.S. degree in optoelectronic information science and engineering from Yanshan University, China, in 2021. Then, he received his M.S. degree in physical electronics from Harbin Institute of Technology in 2023. He is now pursuing a PhD degree in physical electronics from Harbin Institute of Technology. His research interest is focused on photoacoustic spectroscopy and the design of quartz tuning fork.



Shunda Qiao received his B.S. degree in electronic science and technology from Yanshan University, China, in 2018. In 2020, he received his M.S. degree and began to pursue a PhD degree of physical electronics from Harbin Institute of Technology. His research interests include photoacoustic spectroscopy and its applications.



Yanjun Chen received his bachelor's degree in electronic science and technology from Harbin Institute of Technology in 2023. In the same year he pursued a master's degree in physical electronics from Harbin Institute of Technology. His main research area is quartz-enhanced photoacoustic spectroscopy

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Ying He obtained his master's and doctorate degrees in physical electronics from Harbin Institute of Technology in 2017 and 2022, respectively. Currently, he is an associate professor at the Harbin Institute of Technology, China. His research interests including trace gas sensing based on QEPAS, LITES and other laser spectroscopy.



Yufei Ma received his PhD degree in physical electronics from Harbin Institute of Technology, China, in 2013. From September 2010 to September 2011, he spent as a visiting scholar at Rice University, USA. Currently, he is a professor at Harbin Institute of Technology, China. He is the winner of National Outstanding Youth Science Fund. His research interests include optical sensors, trace gas detection, laser spectroscopy, solid-state laser and optoelectronics. He has published more than 100 publications and given more than 20 invited presentations at international conferences. He serves as associate editor for Optica Optica Express, SPIE Optical Engineering, Wiley Microwave and Optical Technology Letters and Frontiers in Physics. He also serves as topical editor for CLP

Chinese Optics Letters and editorial board member for Elsevier Photoacoustics MDPI Sensors and Applied Sciences.