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Aluminium incorporation in polar, semi- and non-polar AlGa_xN layers: a comparative study of x-ray diffraction and optical properties

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Growth of Al_xGa_{1-x}N layers ($0 \leq x \leq 1$) simultaneously on polar (0001), semipolar (10 $\bar{1}$ 3) and (11 $\bar{2}$ 2), as well as nonpolar (10 $\bar{1}$ 0) and (11 $\bar{2}$ 0) AlN templates, which were grown on planar sapphire substrates, has been investigated by metal-organic vapour phase epitaxy. By taking into account anisotropic in-plane strain of semi- and non-polar layers, their aluminium incorporation has been determined by x-ray diffraction analysis. Optical emission energy of the layers was obtained from room-temperature photoluminescence spectra, and their effective bandgap energy was estimated from room-temperature pseudo-dielectric functions. Both x-ray diffraction and optical data consistently show that aluminium incorporation is comparable on the polar, semi- and non-polar planes.

AlGa_xN-based light-emitting diodes (LEDs) operating in the deep ultraviolet (UV) spectral region ($\lambda < 360$ nm) have a wide range of potential applications such as water purification¹, disinfection of medical tools², as well as photo-therapy and medical diagnostics³. However, polar (0001) *c*-plane UV LEDs operating at below 240 nm still have very low external efficiency of less than 0.2%⁴. It is well-known that LEDs epitaxially grown along the [0001] *c*-direction have strong polarization fields across quantum-well (QW) structures⁵. These fields reduce the electron-hole wave-function overlap resulting in a reduction of the radiative recombination rate. Additionally, polar LEDs operating at $\lambda < 250$ nm, the light emission mode changes from transverse electric polarization ($E \perp [0001]$) to transverse magnetic polarization ($E \parallel [0001]$) resulting in a reduced light extraction efficiency⁶⁻⁸.

Growth on semi- and non-polar planes results in a reduction of built-in fields⁵, which should increase the radiative recombination efficiency of QW active regions. Additionally, the optical polarization of nonpolar (10 $\bar{1}$ 0) *m*-plane Al_xGa_{1-x}N QWs ($0 \leq x \leq 1$) grown on bulk AlN substrates has been found to be dominant along the *c*-direction over the entire range of composition⁹. However, despite these advantages, there are only a few reports about semi- and non-polar UV LEDs, e.g., *a*-plane AlN/non-UV-transparent-SiC LEDs operating at 210 nm¹⁰ and semipolar (11 $\bar{2}$ 2) AlGa_xN/UV-transparent-sapphire LEDs operating at 307 nm¹¹. One of the reasons is lack of high-quality high-UV-transparent AlGa_xN and AlN templates prepared on sapphire substrates. Recently, high-temperature thermal annealing has been used to improve the material and optical properties of semi- and non-polar AlN/sapphire templates¹²⁻¹⁶.

In order to realize UV (visible) emitters, growth of AlGa_xN (InGa_xN) QWs with a desired Al (In) composition needs to be well-controlled. In contrast to widely compositional studies for up to twenty various InGa_xN surface orientations grown on GaN substrates using metal-organic vapour phase epitaxy (MOVPE)¹⁷⁻²² and ammonia molecular beam epitaxy (MBE)²³, very limited study has been performed for semi- and non-polar AlGa_xN grown on AlN substrates, e.g., *m*-plane²⁴, semipolar (10 $\bar{1}$ 2)²⁵ and (20 $\bar{2}$ 1)²⁶. This is mainly due to limit of available non-*c*-plane AlN substrates. Additionally, it should be noted that these substrates are very small and expensive, and UV-light transparency still remains a challenge^{27,28}. Semi- and non-polar AlGa_xN layers have also been heteroepitaxially grown on sapphire substrates, e.g., *a*-plane layers on (10 $\bar{1}$ 2) *r*-plane sapphire²⁹, *m*-plane³⁰ and (11 $\bar{2}$ 2) layers^{31,32} on *m*-plane sapphire.

Even though semipolar (10 $\bar{1}$ 3) AlN templates can be grown on *m*-plane sapphire¹⁶, crystal twinning has been observed for the layers that might lead to difficulties for AlGa_xN growth and composition determination. Recently,

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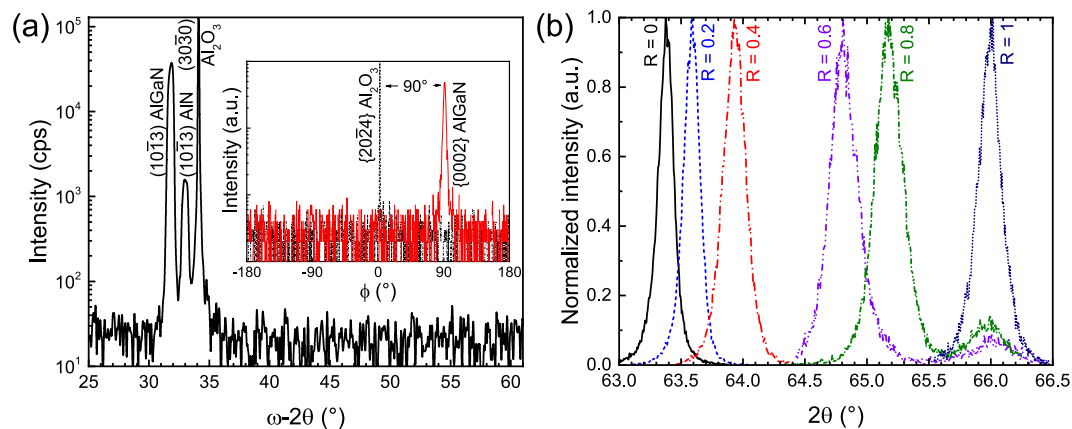


Figure 1. (a) Symmetric ω - 2θ XRD scan of an AlGaN layer grown on a $(10\bar{1}3)$ AlN/ m -plane sapphire template with $R_{\text{AlGaN}} = 0.2$. The inset shows azimuthal scans of the $\{20\bar{2}4\}$ sapphire and $\{0002\}$ Al_xGa_{1-x}N diffraction peaks performed in skew symmetry. (b) 2θ scans of the $(10\bar{1}3)$ Al_xGa_{1-x}N diffraction peaks of the layers grown on $(10\bar{1}3)$ AlN templates with different R_{AlGaN} .

we have successfully produced untwinned semipolar $(10\bar{1}3)$ AlN templates on m -plane sapphire using directional sputtering³³. Therefore, to extend compositional study of AlGaN with different surface orientations, in this paper, we report on MOVPE-growth of Al_xGa_{1-x}N layers simultaneously on polar (0001), semipolar $(10\bar{1}3)$ and $(11\bar{2}2)$, as well as nonpolar $(10\bar{1}0)$ and $(11\bar{2}0)$ AlN/sapphire templates over the entire range of composition. Compositional study of these layers has been investigated by x-ray diffraction (XRD), room-temperature photoluminescence (RT-PL) and pseudo-dielectric functions (DFs) measurements.

Results

Determination of Al incorporation by XRD. To vary the aluminium mole fraction (x_{AlN}) of the grown AlGaN layers, different $R_{\text{AlGaN}} = 2 \cdot [\text{TMAI}] / (2 \cdot [\text{TMAI}] + [\text{TMGa}])$ gas phase ratios ($0 \leq R_{\text{AlGaN}} \leq 1$) were employed while keeping NH₃ flow rate constantly. Figure 1(a) shows a symmetric ω - 2θ XRD scan of an AlGaN layer grown on a $(10\bar{1}3)$ AlN/sapphire template with $R_{\text{AlGaN}} = 0.2$. Besides the $(30\bar{3}0)$ diffraction peak of sapphire at 34.1° , there are only the $(10\bar{1}3)$ AlGaN and AlN peaks, indicating that this layer is indeed single phase.

To investigate crystal twinning and the epitaxial in-plane relationship of the grown $(10\bar{1}3)$ Al_xGa_{1-x}N layers and m -plane sapphire, XRD off-axis ϕ -scans were measured. The skew-symmetric $\{20\bar{2}4\}$ sapphire diffraction peak of m -plane sapphire substrate was measured with a tilt angle of: 32.4° , which indicates $[0001]_{\text{sapphire}}$. To indicate $[0001]_{\text{AlGaN/AlN}}$ of the $(10\bar{1}3)$ layers, the skew-symmetric $\{0002\}$ AlGaN peak was measured with a tilt angle of: 31.6° . The inset of Fig. 1(a) shows ϕ -scans of the $\{20\bar{2}4\}$ sapphire and $\{0002\}$ AlGaN diffraction peaks of a $(10\bar{1}3)$ layer. Only one peak of $\{0002\}$ AlGaN is found, which tilts exactly 90° with respect to the $\{20\bar{2}4\}$ sapphire peak, indicating that this layer is untwinned. The relationship is found to be $[30\bar{3}\bar{2}]_{\text{AlGaN/AlN}} \parallel [11\bar{2}0]_{\text{sapphire}}$ and $[11\bar{2}0]_{\text{AlGaN/AlN}} \parallel [0001]_{\text{sapphire}}$.

Figure 1(b) shows 2θ scans of the $(10\bar{1}3)$ AlGaN layers grown with different R_{AlGaN} . Various diffraction peak positions indicate different x_{AlN} of these layers. A similar result has been found for the other layers with different surface orientations.

Semi- and non-polar AlGaN layers hetero-epitaxially grown on sapphire substrates generally have triclinic and orthorhombic distortions of their wurtzite unit cells, respectively. Anisotropy in the lattice and thermal expansion mismatches along two in-plane directions results in anisotropic in-plane strain causing these distortions. This makes lattice parameter measurements, and thus x_{AlN} determination, difficult. By taking into account these distortions, XRD methods have been developed to determine x_{AlN} of nonpolar³⁴ and semipolar AlGaN layers³⁵.

For the differently oriented AlGaN layers studied here, their a and c lattice constants have been calculated by measuring different symmetric, skew-symmetric and asymmetric 2θ diffraction peaks, as shown in Table 1. An example of lattice measurements for m -plane AlGaN can be seen in ref. 30. Figure 2(a) shows the measured lattice constants of the layers as a function of R_{AlGaN} . The lattice constants of all the layers show a linear behaviour with R_{AlGaN} . Additionally, all layers in this study have an expected ratio of a to c lattice constant with a corresponding composition, indicating that they are fully relaxed.

Based on these measured lattice constants, x_{AlN} of all the layers with different surface orientations has been estimated, as shown in Fig. 2(b). At each growth condition (R_{AlGaN}), x_{AlN} values of these layers are slightly different. For example, maximum differences (Δx) of 0.02/0.08/0.03 are estimated for the layers grown with $R_{\text{AlGaN}} = 0.1/0.4/0.8$, respectively. Given these scattered data points, x_{AlN} values of all the layers can be considered to be comparable. In contrast to a linear behaviour of $R_{\text{AlGaN}}-x_{\text{AlN}}$ observed for c - and m -plane layers grown at 1050°C reported in ref. 30, for the samples studied here grown at 1150°C , a non-linear behaviour has been observed. This is attributed to TMAI:NH₃ pre-reactions and gallium desorption^{36,37}.

Surface orientations	Diffraction peaks
(0001)	(0002), (10 $\bar{1}$ 5), (20 $\bar{2}$ 5)
(10 $\bar{1}$ 0)	(10 $\bar{1}$ 0), (10 $\bar{1}$ 1), (11 $\bar{2}$ 0), (11 $\bar{2}$ 2), (12 $\bar{3}$ 0), (20 $\bar{2}$ 1), (20 $\bar{2}$ 1), (21 $\bar{3}$ 1), (21 $\bar{3}$ 3)
(11 $\bar{2}$ 0)	(11 $\bar{2}$ 0), (10 $\bar{1}$ 0), (10 $\bar{1}$ 1), (10 $\bar{1}$ 2), (11 $\bar{2}$ 2), (12 $\bar{3}$ 0), (21 $\bar{1}$ 0), (21 $\bar{3}$ 1), (21 $\bar{3}$ 2)
(10 $\bar{1}$ 3)	(10 $\bar{1}$ 3), (0002), (10 $\bar{1}$ 0), (10 $\bar{1}$ 1), (11 $\bar{2}$ 0), (11 $\bar{2}$ 2), (11 $\bar{2}$ 2)
(11 $\bar{2}$ 2)	(11 $\bar{2}$ 2), (0002), (0004), (10 $\bar{1}$ 0), (11 $\bar{2}$ 0), (10 $\bar{1}$ 3), (21 $\bar{1}$ 2)

Table 1. 2θ XRD peaks used to measure lattice constants of the AlGa $_x$ N layers with different surface orientations.

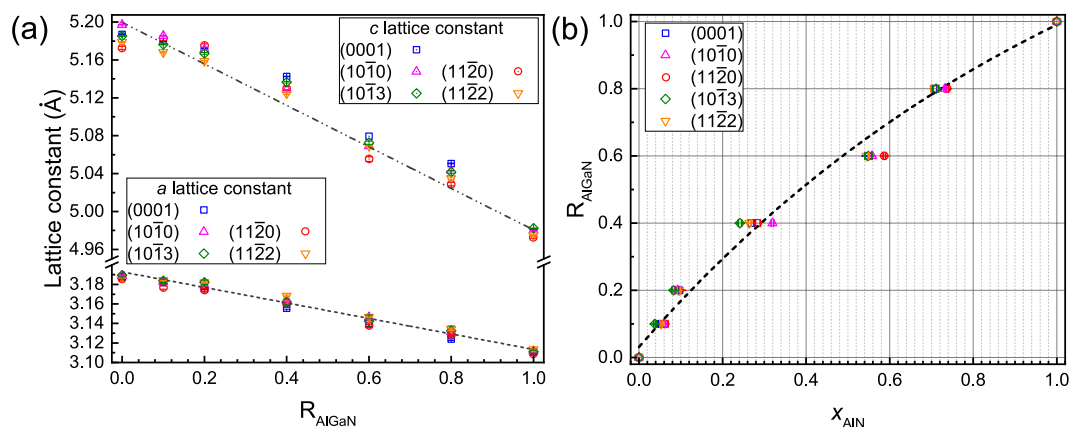


Figure 2. (a) Measured lattice constants of the differently oriented Al $_x$ Ga $_{1-x}$ N layers grown with different R_{AlGaN} . (b) Calculated x_{AlN} from XRD data of the differently oriented Al $_x$ Ga $_{1-x}$ N layers plotted as a function of R_{AlGaN} . Error bars in (a) and (b) are standard errors estimated from the calculations.

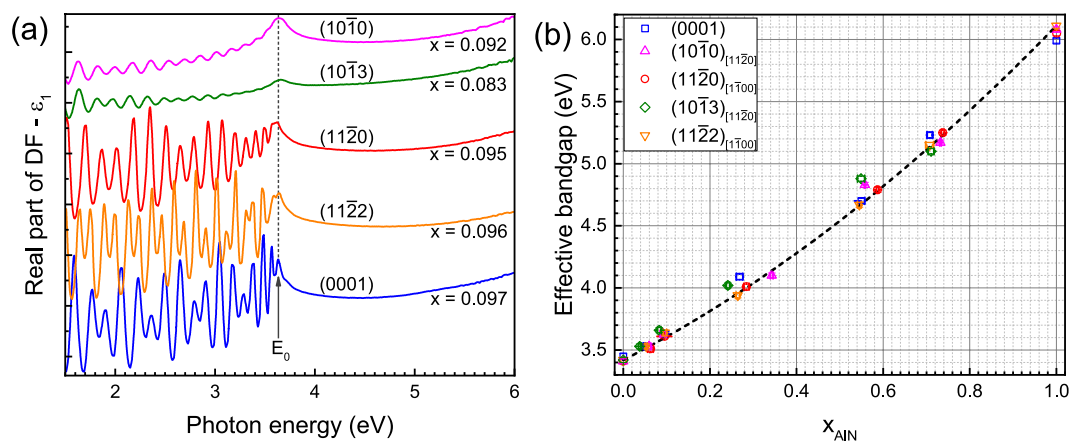


Figure 3. (a) Real part ($\langle \epsilon_1 \rangle$) of the pseudo-dielectric functions of semi- and non-polar AlGa $_x$ N layers ($E \perp [0001]$), and of a c -plane co-loaded layer ($E \parallel [0001]$) grown with $R_{\text{AlGaN}} = 0.2$. Effective bandgap (E_0) of the band structure is indicated by an arrow. (b) Bandgap of these layers plotted as a function of x_{AlN} . The dashed line is a bandgap-bowing fitting of the experimental data with a bowing parameter of 0.9 eV.

Optical properties. *Optical bandgap energy.* For wurtzite nitrides, the valence band maximum is split both by spin-orbit interaction and non-cubic crystal field, resulting in three valence-band states (i.e., Γ_{7-}^v , Γ_{7+}^v and Γ_9^v) at the Brillouin zone centre³⁸. For m -plane AlN³⁹ and (11 $\bar{2}$ 2) AlGa $_x$ N⁴⁰, the absorption origin for $E \parallel [0001]$ indicates transitions from Γ_{7+}^v , while the one for $E \perp [0001]$ indicates transitions mainly from Γ_{7-}^v and/or Γ_9^v . Figure 3(a) exemplifies real parts ($\langle \epsilon_1 \rangle$) of the DFs measured on the differently oriented AlGa $_x$ N layers grown with $R_{\text{AlGaN}} = 0.2$ ($x_{\text{AlN}} \approx 0.1$). The $\langle \epsilon_1 \rangle$ parts of m -plane and (10 $\bar{1}$ 3) layers were measured along [11 $\bar{2}$ 0]_{AlGa $_x$ N}, while they were measured along [1 $\bar{1}$ 00]_{AlGa $_x$ N} for the a -plane and (11 $\bar{2}$ 2) layers. Compared to the c -plane, a -plane and (11 $\bar{2}$ 2) layers, the interference fringes of the $\langle \epsilon_1 \rangle$ parts of the m -plane and (10 $\bar{1}$ 3) layers have weaker ampli-

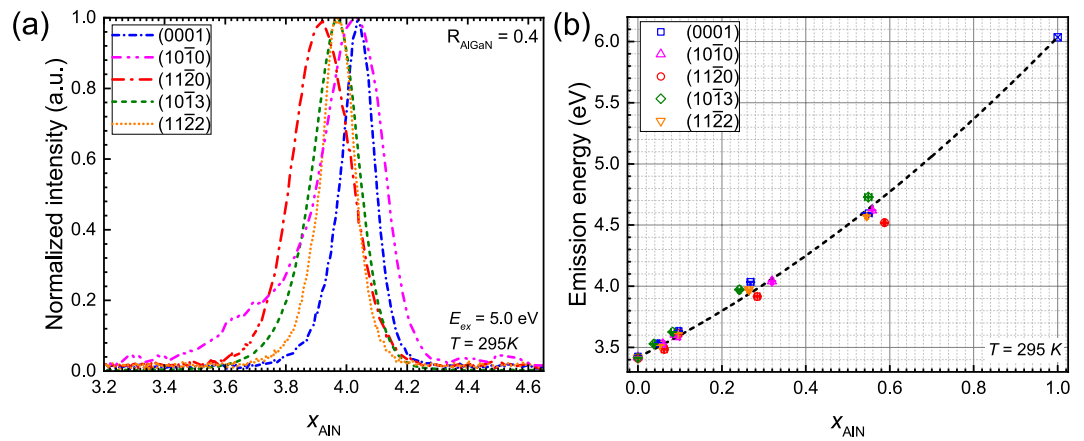


Figure 4. (a) RT-PL spectra of the differently oriented $\text{Al}_x\text{Ga}_{1-x}\text{N}$ layers grown with $R_{\text{AlGa}} = 0.4$. (b) PL peak energy of the layers plotted as a function of x_{AlN} . A near band-edge energy of AlN at 6.035 eV (\square) is taken from ref. 42. The dashed line is a bowing fitting of the experimental data with a bowing parameter of 0.9 eV.

tudes because of their rougher interfaces³⁰. From these $\langle \epsilon_i \rangle$ parts, the fundamental bandgap energy ($E_g^{\text{AlGa}}N$) of the grown layers is approximately estimated from a sharp excitonic E_0 peak^{39–41}.

$E_g^{\text{AlGa}}N$ of all the AlGa $_N$ layers is plotted as a function of x_{AlN} in Fig. 3(b). Their $E_g^{\text{AlGa}}N$ values are comparable over the entire range of composition. This indicates comparable x_{AlN} values, consistent with the values estimated by XRD (Fig. 2(b)). The dependence of $E_g^{\text{AlGa}}N$ on x_{AlN} can be described as:

$$E_g^{\text{Al}_x\text{Ga}_{1-x}\text{N}} = x \cdot E_g^{\text{AlN}} + (1 - x) \cdot E_g^{\text{GaN}} - b \cdot x \cdot (1 - x),$$

where b denotes the bandgap bowing parameter. To fit the experimental data, a measured E_g^{AlN} of 6.11 eV and a measured E_g^{GaN} of 3.42 eV on the (1122) AlN and GaN samples along $[1\bar{1}00]_{\text{AlGa}}N$ were used, respectively. The shift of $E_g^{\text{AlGa}}N$ with x_{AlN} is well reproduced with a bowing parameter of about 0.9 eV. This value is in good agreement with values reported for a -plane²⁹, m -plane³⁰, (1122)⁴⁰, and c -plane AlGa $_N$ layers^{30,41}.

Photoluminescence. A correlation between the bandgap energy with optical emission properties has also been investigated. Due to the excitation energy of laser ($E_{\text{ex}} = 5$ eV), only samples grown with $R_{\text{AlGa}} < 0.8$ (i.e., $x_{\text{AlN}} < 0.7$) can be measured. Figure 4(a) exemplifies RT-PL spectra measured on the differently oriented AlGa $_N$ layers grown with $R_{\text{AlGa}} = 0.4$. The near band-edge (NBE) emission energy of these samples, which was estimated from a Gaussian fit of the corresponding band, is following $(11\bar{2}0)_{3.92\text{ eV}} < (11\bar{2}2)_{3.98\text{ eV}} = (10\bar{1}3) < (0001) = (10\bar{1}0)_{4.04\text{ eV}}$. This order is slightly different from the order shown in XRD data (Fig. 2(b)). However, the maximum NBE difference is of about 120 meV, which is equal to about a difference of 0.06 in x_{AlN} . This composition difference is comparable with the maximum Δx of 0.08 estimated from the XRD data.

The PL emission energy vs x_{AlN} is also well reproduced with a bowing parameter of about 0.9 eV, as shown in Fig. 4(b). For the bowing fitting, an NBE of 3.42 eV obtained from the grown Ga $_N$ layers and an NBE of 6.035 eV of c -, a - and m -plane AlN homo-epilayers taken from ref. 42 were used. The PL data correlates very well with the optical bandgap data indicating a random alloy and almost negligible Ga clustering. This is different from the case of InGa $_N$ QWs, where a strong In clustering has often been reported, which results in a large Stokes shift between NBE and effective bandgap^{43–46}.

Discussion

Of the twenty different semi- and non-polar InGa $_N$ QWs MOVPE-grown on bulk Ga $_N$ substrates investigated so far, compositional study shows different results^{17–22}. For example, it has been reported that In incorporation in (0001) $<$ (1122)¹⁹ or (0001) \approx (1122)^{18–22}, (1010) $<$ (1122)^{17–21} or (1010) \approx (1122)¹⁸, (1010) $<$ (0001)^{20,21} or (1010) \approx (0001)^{18,19}, and (1120) $<$ (0001)^{21,22} or (1120) \approx (0001)¹⁸. This is possible that these discrepancies are due to divergence in strain relaxation in the investigated samples and/or indium clustering. For InGa $_N$ layers grown by MBE²³, which has completely different growth kinetics from MOVPE, it has been reported that In incorporation in (1122) $<$ (0001) $<$ (1010). So far, there is only one compositional study for (1013) InGa $_N$ QWs by MOVPE²¹, indicating that this plane has the lowest In-content among all the aforementioned planes.

To study In incorporation in different InGa $_N$ surface orientations, a few theoretical calculations have also been performed by taking into account surface kinetics^{47,48} or strain energy dependent surface orientations^{49,50}. However, they also show contrary results, e.g., In incorporation in m -plane InGa $_N$ was found to be smaller⁴⁷ or higher than c -plane InGa $_N$ ^{49,50}. Additionally, it has been theoretically^{49,50} and experimentally^{19,23} reported that different growth conditions (e.g., pressure, temperature, and V/III) might result in different In incorporations on different surface orientations.

Of the relaxed AlGa $_N$ layers with five different surface orientations studied here, their x_{AlN} is comparable over the entire range of composition, as consistently confirmed by XRD and optical data. The comparable x_{AlN} of the

m-plane and *c*-plane layers is in good agreement with a previous report³⁰, even though the growth temperature used here is 100°C higher. Given the slightly scattered data points of the *a*-plane and *c*-plane layers, their comparable x_{AlN} also can be considered as a consistent result with a previous report²⁹, where only a slightly higher x_{AlN} of *c*-plane layers was found ($\Delta x_{\text{AlN}} \leq 0.05$).

For the *c*-plane and (11 $\bar{2}$ 2) AlGa_N layers studied here, their comparable x_{AlN} is contrary to previous results reported for (11 $\bar{2}$ 2) vs *c*-plane layers, where x_{AlN} of (11 $\bar{2}$ 2) layers was found to be lower ($\Delta x \leq 0.1$)³² or higher ($\Delta x \leq 0.2$)³¹ than that of *c*-plane layers. This might be due to different growth conditions and/or calculation methods used. So far no theoretical study about composition dependent surface orientations has been done for AlGa_N. In case of InGa_N, most experimental data seems to indicate a higher In incorporation for orientations with almost upright metal dangling bonds. This can indicate that the bonding and incorporation of In versus In desorption are the most important step. Since the AlGa_N layers studied here have a similar Al incorporation for all orientations, one may argue that the strong polarity of Al(Ga)N together with the lower total strain facilitates Ga incorporation and makes Ga desorption the less likely process. Further investigations and calculations need to be performed to clarify this.

Conclusions

Compositional study of relaxed co-loaded AlGa_N layers with polar (0001), semipolar (10 $\bar{1}$ 3) and (11 $\bar{2}$ 2), as well as nonpolar (10 $\bar{1}$ 0) and (11 $\bar{2}$ 0) surface orientations has been investigated. By taking into account the compositional effects of anisotropic in-plane strain, aluminium incorporation in semi- and non-polar layers was determined by x-ray diffraction analysis. The AlN mole fraction of all the co-loaded layers estimated by x-ray diffraction is comparable. This is consistent with their comparable optical bandgap energy and near band-edge emission energy, which were determined from room-temperature pseudo-dielectric functions and photoluminescence measurements, respectively. The dependence of the bandgap and emission energy on composition indicates a bowing parameter of 0.9 eV.

Experimental Methods

Growth was performed in an EpiQuest 3 × 2-inch close-coupled showerhead MOVPE reactor. Ammonia (NH₃), trimethylgallium (TMGa) and trimethylaluminium (TMAl) were used as precursors. Differently oriented AlN templates grown on sapphire substrates were used to grow AlGa_N layers, including (0001) AlN ($d \approx 800$ nm) on *c*-plane sapphire, (11 $\bar{2}$ 2) AlN ($d \approx 1000$ nm) on *m*-plane sapphire, (10 $\bar{1}$ 0) AlN ($d \approx 350$ nm) on *m*-plane sapphire, and (11 $\bar{2}$ 0) AlN ($d \approx 350$ nm) on *r*-plane sapphire. The (11 $\bar{2}$ 0) AlN template was grown simultaneously with the (10 $\bar{1}$ 0) AlN template. Growth parameters of these templates are reported elsewhere³⁰. To produce an Al-polar (10 $\bar{1}$ 3) AlN template, about 10-nm-thick (10 $\bar{1}$ 3) AlN layer was initially sputtered onto a 2-inch *m*-plane sapphire wafer using directional sputtering³³. Afterwards, this wafer was loaded into the reactor chamber to grow a 300-nm-thick AlN layer at a surface temperature of 1290°C.

All the 2-inch AlN/sapphire wafers were diced into 1 × 1 cm² pieces. These pieces were then co-loaded into the reactor chamber for AlGa_N epitaxy. Initially, about 100-nm-thick AlN layer was grown on these templates at 1290°C at a reactor pressure of 27 hPa. Afterwards, AlGa_N layers with a nominal thickness of 1.5 μm were grown on these templates at 1150°C at a reactor pressure of 100 hPa. To vary x_{AlN} , different R_{AlGa_N} ratios were employed ($0 \leq R_{\text{AlGa}_N} \leq 1$), while keeping NH₃ flow rate constantly. Growth parameters of these layers are reported in ref. 30; however, the AlGa_N growth temperature at 1050°C was used in that study.

The crystal orientation of the AlGa_N/AlN samples was characterized using a PANalytical X'pert triple-axis high-resolution X-ray diffraction (HR-XRD) system equipped with an asymmetric four-crystal monochromator (4 × Ge220) for CuK_{α1} source. On-axis ω -2 θ scans have been measured using an open detector without any receiving slit to distinguish between all possible orientations of the epilayers. For lattice calculations, different 2 θ diffraction peaks of the samples were measured using an HR analyzer detector, as shown in Table 1.

For room-temperature photoluminescence (RT-PL) measurements, the samples were excited by a Krypton Fluoride (KrF) excimer laser (ExciStar XS-200) with excitation wavelength of 248 nm ($E_{\text{ex}} = 5$ eV) and a spot size of 50 × 500 μm². During PL measurements, a pulse energy of 7 mJ and a repetition rate of 200 Hz were employed, giving a power density of 5.6 kW/cm². PL signals were recorded by a high-sensitivity Ocean Optics spectrometer (QE65 Pro).

The fundamental bandgap energy of the layers was estimated from real and imaginary parts of the pseudo-dielectric functions (DFs). DFs were recorded at RT using a Horiba UVISEL 2 spectroscopic ellipsometer at an incident angle of 70° and a spot size of 705 × 2030 μm². The photon energy was varied from 1.45 to 6.45 eV with the spectral resolution of 0.02 eV.

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Author contributions

D.V.D. conducted the design experiments, growth, performed and analysed all measurements, and wrote this manuscript. N.H. prepared the (10 $\bar{1}$ 3) AlN wafer by directional sputtering. Y.H., H.A. and M.P. organized the project. All authors discussed the results and reviewed the manuscript.

Competing interests

The authors declare no competing interests.

Additional information

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