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Small-scale displacement fluctuations of vesicles in fibroblasts

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The intracellular environment is a dynamic space filled with various organelles moving in all directions. Included in this diverse group of organelles are vesicles, which are involved in transport of molecular cargo throughout the cell. Vesicles move in either a directed or non-directed fashion, often depending on interactions with cytoskeletal proteins such as microtubules, actin filaments, and molecular motors. How these proteins affect the local fluctuations of vesicles in the cytoplasm is not clear since they have the potential to both facilitate and impede movement. Here we show that vesicle mobility is significantly affected by myosin-II, even though it is not a cargo transport motor. We find that myosin-II activity increases the effective diffusivity of vesicles and its inhibition facilitates longer states of non-directed motion. Our study suggests that altering myosin-II activity in the cytoplasm of cells can modulate the mobility of vesicles, providing a possible mechanism for cells to dynamically tune the cytoplasmic environment in space and time.

The cell interior is a dynamic environment where vesicles and other organelles traverse the cytoplasm to localize for specific biological processes or to evenly disperse throughout the cell¹. Vesicle motion can be categorized into two main regimes: directed and non-directed^{2–5}. Directed intracellular transport is critical to maintain normal cell function and it is typically driven by molecular motors that consume energy and transport cargo along cytoskeletal filaments^{6–8}. Non-directed vesicle motion is largely driven by thermal fluctuations and results in random diffusion in the complex cytoplasmic environment; although, recent studies show that it can be driven by molecular motors^{9–12} and possibly other athermal processes^{13,14} and is called active diffusion. Observations of vesicle movement in cells reveal vesicles stochastically alternating between directed and non-directed motion, mixing the dynamics and leading to anomalous motion^{9,11,15–23}. Sub-diffusive motion of vesicles in cells is widely observed and often attributed to the soft-glassy nature of the cytoskeleton²⁴.

Several techniques have been developed to analyze this stochastic motion^{3–5,25,26}. Most studies of vesicle movement focus on timescales of seconds, minutes, or even hours where long range transport occurs^{8,27,28}. In this study, we focus on the small-scale displacement fluctuations of vesicle motion on timescales from milliseconds to seconds to provide insight on the local mobility of the cytoplasmic microenvironment. We note that we have not specifically labeled the vesicles tracked in this study and do not have information regarding their molecular identity. We simply employ these vesicles as generic tracers that report on the local mobility in the cytoplasm.

Vesicle movement is an inherently mechanical process, where motion in the cytoplasm requires a force and the vesicle experiences a corresponding resistance (viscoelastic drag force) from the surrounding environment²⁹. In this picture, the cytoskeleton can simultaneously serve to facilitate mobility via molecular motor driven transport, and to hinder mobility since its dense meshwork could act as a barrier. In either case, the cytoskeleton is expected to affect vesicle motion in the cytoplasm. Further, cells exhibit different cytoskeletal structures and higher bulk stiffness on increasing substrate stiffness³⁰. To this end, we investigate two stimuli known to affect cytoskeletal mechanical properties: substrate stiffness and pharmacological drugs targeted to the cytoskeleton. We used glass and polyacrylamide gels ($E = 40$ kPa and 10 kPa) to study the effect of substrate stiffness. To discover the roles of the major cytoskeletal components in vesicle motion we introduced the pharmacological agents cytochalasin-D (cyto-D), colchicine, and blebbistatin, to target and depolymerize actin filaments, microtubule filaments, and inactivate myosin-II, respectively.

We used high-speed differential interference contrast microscopy to record vesicle motion in cells and analyze the effect of substrate stiffness and pharmacological drugs. Vesicle dynamics were quantified via mean squared displacement (MSD), velocity autocorrelation function (VACF), and van Hove correlation (VHC) analyses. We find that vesicle mobility in normal cells is independent of substrate stiffness, however, when pharmacological

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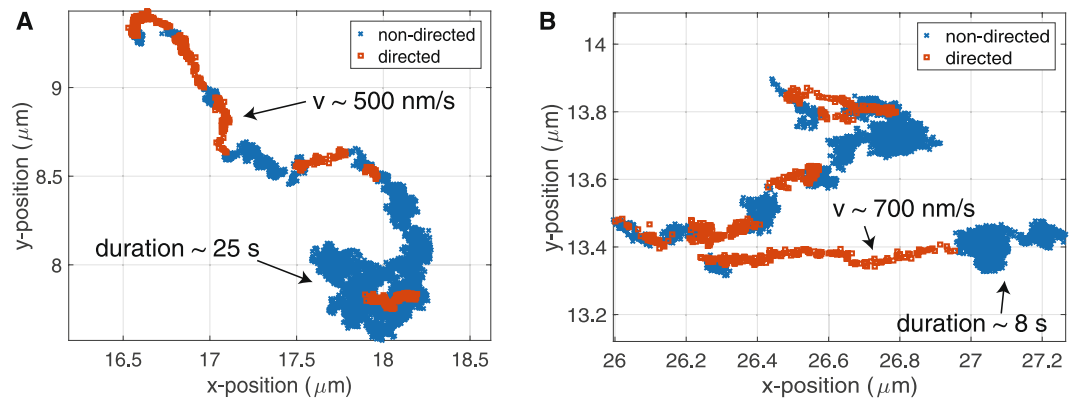


Figure 1. Trajectories of vesicle motion in fibroblasts. Two examples of raw data for the vesicle positions are shown. The temporal mean squared displacement ($tMSD$) analysis was used to identify directed and non-directed motion. **(A)** A trajectory exhibiting short bursts of directed motion ($v \sim 500$ nm/s) followed by a long period of non-directed motion (~ 25 s). **(B)** Another trajectory showing directed motion ($v \sim 700$ nm/s) and non-directed motion (~ 8 s). The maximum possible trajectory duration in this study is 30 s due to acquisition length.

treatments are used their effect increases on softer substrates. Furthermore, our results suggest that small-scale directed vesicle motion is robust, and the actin-myosin-II network plays a role in modulating local vesicle diffusion in the cytoplasm.

Results

All experiments were conducted using differential interference contrast microscopy to visualize vesicle motion in fibroblasts. Vesicles were tracked using an algorithm based on polynomial fitting with Gaussian weight, specifically developed for tracking endogenous vesicles in the cell with nanometer precision using widefield microscopy³¹. To quantify small-scale vesicle displacements, the temporal MSD ($tMSD$) analysis is used as described previously⁴. In brief, the local MSD for each vesicle is calculated and fit to a power-law to determine the state of motion, where a power-law of 1 indicates non-directed motion, and a power-law of 2 indicates ballistic directed motion (see Materials and Methods). The $tMSD$ analysis allows determination of directed and non-directed motion in a single vesicle trajectory (Fig. 1). Subsequently this allows the extraction of several parameters such as velocity and duration of each state of vesicle motion. Throughout the study, directed and non-directed motion was determined via the $tMSD$ analysis.

Vesicle motion is independent of substrate stiffness. Control cells were cultured on substrates of varying rigidity ($E = 10$ kPa, 40 kPa, and glass). Small-scale displacement fluctuations of vesicles in untreated fibroblasts were tracked and analyzed to characterize their overall motion on the three substrates. The MSD was calculated for each substrate and showed that the small-scale fluctuations of vesicles were not strongly affected by substrate stiffness (Fig. 2A). On all three surfaces, vesicle motion is sub-diffusive at short timescales and transitions to diffusive-like behavior at longer timescales. Using the $tMSD$ analysis, vesicle trajectories were split into directed and non-directed states to examine the full displacement distributions, also known as VHC functions³². Figure 2B,C indicates that the displacement distribution of both directed and non-directed vesicle motions are not strongly affected by substrate stiffness. Furthermore, the velocity distribution of vesicles undergoing directed motion also does not depend on the substrate (Fig. 3A). This result is unexpected, since substrate stiffness is known to affect cytoplasmic mechanics^{30,33}, however, overall vesicle mobility within the cytoplasm remains unperturbed.

Myosin-II inhibition decreases effective diffusivity of vesicles. To quantify the diffusivity of vesicles, the VACF was calculated and the effective diffusivity (D_{eff}) was extracted via the Green-kubo relation (see Materials and Methods). When cells are treated with blebbistatin, the D_{eff} of vesicles decreases by nearly 40-fold for cells on 10 kPa gels as evident by the blebbistatin data in Fig. 3B and Table 1. On stiffer surfaces (40 kPa and glass) the same effect is observed, however D_{eff} is decreased by 10-fold (Table 1). Thus, in the absence of myosin-II activity, small-scale vesicle motion is greatly decreased. The effect of perturbing the cytoskeleton itself seems to have less of an effect; depolymerizing actin (via cyto-D) or microtubules (via colchicine) does not significantly affect vesicle diffusivity, as D_{eff} remains within a factor of 2. Overall these results (Table 1) suggest that myosin-II plays a significant role in maintaining vesicle diffusivity for small-scale motion.

Directed motion of vesicles is robust. To take a closer look at the local dynamics of vesicles in the cytoplasm, we used a $tMSD$ analysis to analyze the statistics (e.g. velocity and run-time) of directed and non-directed vesicle motion. The run-time of a particular state (directed or non-directed) is the amount of time a vesicle remains in that state of motion as determined by the $tMSD$ analysis⁴. We found that vesicles typically undergo small-scale directed motions with average velocities of ~ 550 nm/s and run-times of less than 500 ms (100 video

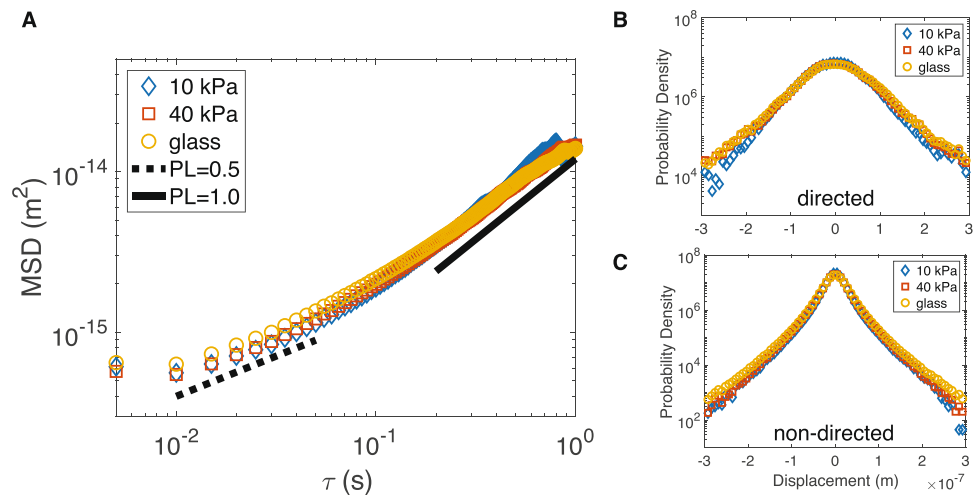


Figure 2. Vesicle dynamics does not depend on substrate stiffness. **(A)** Mean squared displacement (MSD) of vesicle motion in untreated cells indicates no strong dependence on substrate stiffness. Vesicles in all cases show sub-diffusive behavior at shorter timescales transitioning to diffusive behavior at longer timescales. van Hove Correlation (VHC) functions for vesicles undergoing directed motion **(B)** and non-directed motion **(C)** show the full distribution of vesicle displacements is not sensitive to substrate stiffness. VHCs calculated for a timescale of 150 ms. Notice that VHCs of both directed and non-directed motion are strongly non-Gaussian, indicating cytoplasmic activity³². (Number of experiments/cells/trajectories - 10 kPa: 7/28/1211, 40 kPa: 7/24/3021, glass: 10/31/5605; S.E.M. smaller than symbol size).

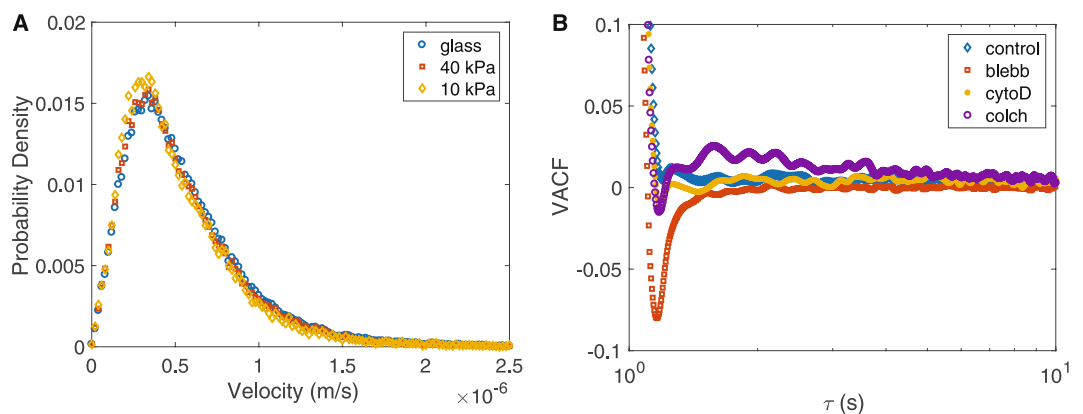


Figure 3. Velocity statistics indicate robustness to substrate stiffness and sensitivity to cytoskeletal perturbation. **(A)** The velocity distribution of directed motion is independent of substrate stiffness. This provides further support, in addition to MSDs and VHCs, that substrate stiffness does not affect vesicle motion in fibroblasts. **(B)** The velocity autocorrelation function (VACF) quantifies persistence and indicates that vesicle motion is strongly perturbed by blebbistatin, which inactivates myosin-II. The strong negative peak (red squares) indicates vesicle motion is decreased when myosin-II is inactivated. Data shown for cells cultured on 10 kPa substrate.

Substrate Stiffness	Control ($\mu\text{m}^2/\text{s}$)	Blebbistatin ($\mu\text{m}^2/\text{s}$)	Cyto-D ($\mu\text{m}^2/\text{s}$)	Colchicine ($\mu\text{m}^2/\text{s}$)
10 kPa	1.80E-3	4.34E-5	8.12E-4	1.49E-3
40 kPa	1.90E-3	2.10E-4	8.69E-4	1.44E-3
Glass	1.54E-3	2.02E-4	7.37E-4	1.58E-3

Table 1. Effective diffusion coefficients (D_{eff} in $\mu\text{m}^2/\text{s}$) Diffusion of vesicles is strongly affected by blebbistatin, which decreases D_{eff} by a factor of ~ 10 – 40 depending on substrate stiffness. In control cells, as well as cyto-D and colchicine treated, diffusion is not strongly affected.

frames) (Tables 2 and 3). The histograms of directed run-times were normalized to create a probability distribution function (Fig. 4A). Thus, integrating over specific time intervals provides the probability of that event occurring.

Substrate Stiffness	Control (nm/s)	Cyto-D (nm/s)	Colchicine (nm/s)
10 kPa	525 ± 385	595 ± 674	456 ± 288
40 kPa	539 ± 370	585 ± 507	562 ± 385
Glass	555 ± 371	589 ± 478	641 ± 446

Table 2. Average velocities of vesicles undergoing directed motion. (mean ± stdev). Velocity distributions of directed motion are broad (see Fig. 3A) and are not strongly affected by cytoskeletal perturbation. Note that blebbistatin data is not available, due to the low occurrences of directed motion.

Time (ms)	Control	Cyto-D	Colchicine	Blebbistatin
500	94.3%	93.9%	93.9%	100.0%
250	54.8%	57.2%	61.7%	71.4%

Table 3. Statistics of directed run-times within a certain interval time. Time intervals correspond to areas of integration shown in Fig. 4A. Most directed run-times were of duration 500 ms or less. (data shown for 10 kPa substrate).

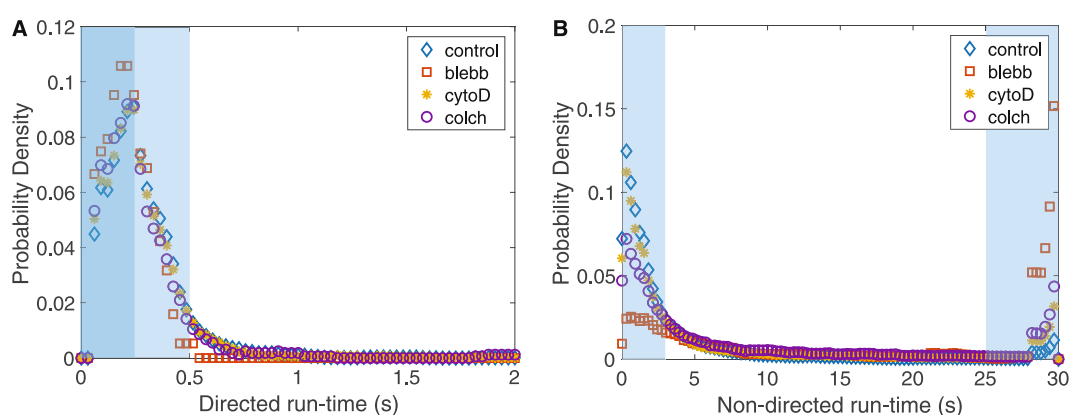


Figure 4. Directed run-times are robust, but non-directed run-times are sensitive to cytoskeletal perturbation. (A) The directed run-time distribution indicates vesicles spend a similar amount of time undergoing directed motion independent of cytoskeletal perturbation. These durations are relatively short, typically less than 500 ms. (B) Non-directed run-times indicate two timescales of random motion, short (<3 s) and long (>25 s) for all conditions. When cytoskeletal perturbations are introduced, there is a shift in non-directed motion from short to long run-times. Inactivation of myosin-II has the strongest effect, shifting nearly half of vesicles to experience mainly non-directed motion.

To quantify the statistics of motion, we integrate the probability density functions. Integrating the probability density function of directed run-times shows 94% of vesicles in control, cyto-D, and colchicine treated cells experience directed motion of duration less than 500 ms, and in the case of blebbistatin-treated cells the number was slightly larger, 100% (Fig. 4A, light blue shaded region and Table 3). This result suggests the possibility that local directed transport of vesicles is not directly dependent on myosin-II, actin filaments, or microtubules. Alternatively, it is possible that a subpopulation of stable actin and microtubules still remain after drug treatment and facilitate local directed motion. It is important to note that this result applies to small-scale directed motion only, which may be strongly dependent on local cytoplasmic geometry (e.g. mesh size). Long-range directed transport must undoubtedly depend on the presence of cytoskeletal filaments.

Myosin-II inhibition increases the duration of non-directed motion of vesicles. In contrast to small-scale directed motion, non-directed motion of vesicles in the cytoplasm is not robust and does appear to be affected by drug treatments. We identified two distinct groups for the run-time duration of non-directed motion: short duration (~3 seconds) and long duration (~25 seconds) (Fig. 4B). A short duration non-directed run-time indicates a vesicle spends only a short amount of time moving randomly before transitioning to directed movement. A long duration non-directed run time indicates a vesicle undergoes random motion for the majority of the measurement period. To quantify the statistics, we integrate the probability density functions. In control cells, 73% of vesicles exhibit short run-times of non-directed motion (blue diamonds) and only 4% exhibit long run-times. However, when myosin-II is inhibited (Fig. 4B, red squares), there is a significant shift towards more vesicles undergoing non-directed motion for longer run-times: only 24% display short run-times while 47% display long run-times. Thus, we see a shift from 4% (control) to 47% (blebbistatin) of vesicles exhibiting

Time (s)	Control	Cyto-D	Colchicine	Blebbistatin
<3	73.4%	65.5%	50.2%	24.4%
>25	4.0%	11.1%	15.1%	47.0%

Table 4. Statistics of non-directed run-times within a certain interval time. Time intervals correspond to areas of integration shown in Fig. 4B. In control cells, the majority of non-directed motions lasted less than 3 s before switching to directed motion. When cytoskeletal perturbations are introduced, the distribution shifts indicating vesicles spend longer durations in non-directed states. This effect is particularly strong for blebbistatin.

non-directed run-times when myosin-II activity is inhibited. Treatments affecting the cytoskeleton have an intermediate effect. In cyto-D and colchicine treated cells 65% and 50% of vesicles, respectively, experience short run-times and 11% and 15% experience long run-times of random small-scale motion (Table 4). These results suggest that in the unperturbed state (control), the majority of vesicles spend only a short amount of time undergoing non-directed motion before transitioning to directed motion. When myosin-II is inhibited, this behavior shifts so that a large proportion of vesicles undergo non-directed motion for the majority of the measurement time.

Small-scale vesicle motion in cells on glass substrates is not strongly affected by drug treatments.

For cells cultured on a glass substrate, the vesicles in treated and untreated cells exhibit similar small-scale motion as shown by the MSD (Fig. 5A). This result is interesting because it suggests that perturbation of actin, myosin-II, or microtubules does not induce a strong effect on small-scale fluctuating vesicle motion for cells cultured on glass surfaces. Any change in vesicle motion is significantly greater on softer substrates.

On softer substrates, depolymerization of F-actin increases vesicle motion and depolymerization of microtubules decreases vesicle motion.

In contrast to cells on glass, vesicle motion in cells on softer substrates is affected by cytoskeletal perturbations. This effect increases with decreasing stiffness of the substrate (Fig. 5B,C). On softer substrates, depolymerizing microtubules leads to a decrease in vesicle mobility (yellow circles) and depolymerizing actin leads to an increase in vesicle mobility (purple stars). The observed effect is stronger on 10 kPa gels (Fig. 5C) relative to 40 kPa (Fig. 5B) and is most evident at longer timescales. This result aligns with the idea that microtubules function as highways for vesicle transport and that actin may serve to cage vesicles^{34,35}.

Discussion

Substrate stiffness is known to affect cell mechanics, thus it was expected to influence vesicle motion. However, the MSD analysis (Fig. 2) and effective diffusion coefficients (Table 1) show substrate stiffness does not strongly affect vesicle motion. A recent study of size and speed-dependent vesicle motion in cells³⁶ suggests that in this regime of vesicle velocity ($v/r \sim 1$) vesicle motion due to an applied force is independent of vesicle size. Their finding implies that in this regime, the mechanical properties of the cytoplasm may not have a strong effect on vesicle motion. Our results provide further evidence supporting the study by Hu *et al.*³⁶, since we do not observe a change in vesicle motion with substrate stiffness (which is expected to change cell mechanical properties). While vesicle dynamics in healthy cells seems unaffected by substrate stiffness, the same cannot be said for cells treated with cytoskeletal perturbing agents.

For instance, when myosin-II activity is inhibited the vesicle diffusivity decreases and this effect is stronger on softer substrates. The VACF is a convenient tool to quantify the persistent dynamics of a system, for example, whether it behaves more like a solid, liquid, or gas^{32,37}. When blebbistatin is present, the VACF exhibits the strongest initial anti-correlation of all treatments (Fig. 3). One interpretation is that the cytoplasm exhibits more solid-like behavior when myosin-II activity is inhibited. Or in other words, myosin-II activity helps maintain a soft cytoplasmic environment. Myosin-II is known to form myosin heavy chain isoforms which attach to actin filaments and pull them toward each other³⁸. Myosin-II activity may fluidize the cytoplasm (softening the environment) allowing vesicles to diffuse more freely, as supported by experimental and theoretical studies^{39,40}. Since these motors can be activated locally, this is a possible mechanism for cells to tune their local cytoplasmic environment to facilitate or impede vesicle motion. A similar effect of myosin-II inhibition has been observed in amoeba⁴¹.

While the VACF was used in this study to quantify the persistent dynamics and calculate D_{eff} , plots of MSD quantify the overall vesicle motion and suggest that actin filaments may impede local motion. MSD analysis suggests that on softer substrates, vesicle motion increases in the absence of actin filaments as is visible by the increased amplitude of motion for cyto-D treated cells (Fig. 5B,C purple stars). Note that while this result is clear in the MSD, it is not reflected in the D_{eff} obtained from integrating the VACF (Table 1) because this calculation is dominated by the initial anticorrelation (Fig. 3B). The notion of actin filaments acting as a barrier impeding motion is consistent with previous results where F-actin was found to trap non-functionalized quantum dots in the actin meshwork of COS-7 cells³⁵.

In addition to the VACF and MSD, the run-times of directed and non-directed motion shed light on the cytoskeleton's influence on vesicle motion. It was observed that small-scale directed motion for vesicles in all treated and untreated cells lasted only about 500 ms (Fig. 4A), and velocities of these vesicles were not significantly changed by any drug treatment (Fig. 3A and Table 2). Since these directed motions may not be driven by cargo transport motors (as they are not influenced by depolymerization of actin or microtubules), this robustness may reflect an underlying structure of the cytoskeletal network. For instance, the mean directed run-time (~200

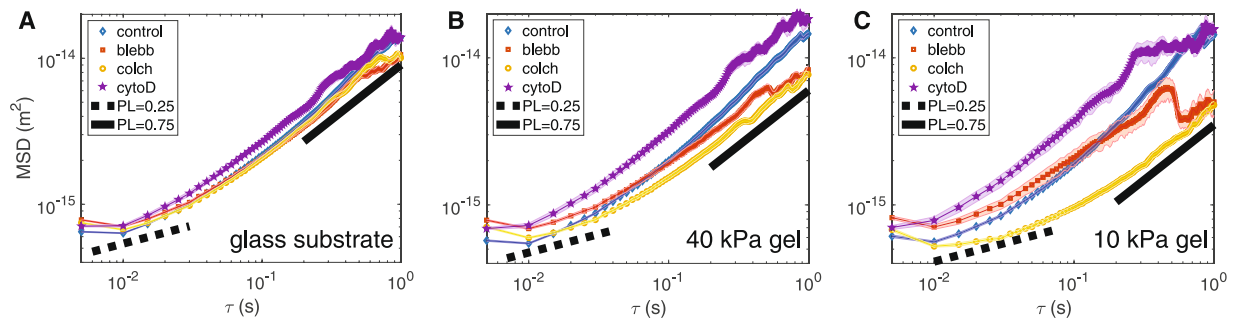


Figure 5. Vesicle motion is more sensitive to cytoskeletal perturbation on softer substrates. (A) Vesicle dynamics on glass is not strongly affected by cytoskeletal perturbation. However, the effect becomes more pronounced on (B) 40 kPa gels and (C) 10 kPa gels. Relative to the control vesicle motion, MSD plots suggest colchicine treatment decreases motion and cyto-D increases motion. (Number of experiments/cells/trajectories - glass: control 13/31/6738, blebb 8/19/4498, colch 8/16/2111, cytoD 8/18/3984; 40 kPa: control 10/30/3484, blebb 8/16/2412, colch 8/16/1324, cytoD 8/16/2139; 10 kPa: control 10/27/1855, blebb 8/15/773, colch 8/10/532, cytoD 8/12/797; S.E.M. shown as shaded background).

ms) and the mean directed velocity (~ 500 nm/s) provides a characteristic length of ~ 100 nm of the directed motion. That value is on-par with the expected mesh size of the cytoskeleton in cells and in this case may be related to the density of intermediate filaments^{42,43}, because we observe the same behavior in the absence of actin filaments or microtubules.

As aforementioned, a recent study demonstrated that quantum dots ($r_{QD} = 15$ nm) in the cytoplasm of fibroblast-like cells (COS-7) can undergo either fast or slow diffusion³⁵, and fast diffusion occurs when the actin is destabilized. Our recorded measurements of fluctuating vesicles ($r_{ves} \sim 600$ nm) are in agreement with the slow diffusing quantum dots trapped in the actin meshwork. In our case, the diffusion coefficient of vesicles was $D_{ves} = 1.5 \times 10^{-3} \mu\text{m}^2/\text{s}$ whereas for quantum dots, $D_{QD} = 0.6 \times 10^{-1} \mu\text{m}^2/\text{s}$ ³⁵. This difference, $D_{QD}/D_{ves} \sim 40$, can be accounted for by the difference in radii of our objects, $r_{ves}/r_{QD} \sim 40$. Similarly, a separate study measured the diffusion of peroxisomes ($r_{PO} \sim 200$ nm) in COS-7 cells to be $D_{PO} = 6 \times 10^{-3} \mu\text{m}^2/\text{s}$ ⁴⁴, where that difference can also be accounted for by the radii, $r_{ves}/r_{PO} \sim 3$. The results we found are also remarkably similar to measurements of carbon nanotube diffusion in COS-7 cells ($l_{CNT} \sim 100\text{--}300$ nm), including the MSD and the velocity distribution⁹. Together, our measurements of small fluctuating motion of vesicles, combined with studies of quantum dots³⁵, peroxisomes⁴⁴, and carbon nanotubes⁹, suggest the local diffusion in the fibroblast cytoplasm may be similar for objects varying in size from tens of nanometers to micron-sized. While this result may be expected in a simple viscous fluid, how this occurs in a heterogeneous complex viscoelastic medium is not clear. Further studies are necessary to probe why seemingly diffusive-like behavior is repeatedly observed in these complex nonequilibrium systems. Future studies would also benefit from specific labeling of vesicles since analyzing a non-homogeneous, undefined set of vesicular cargoes limits interpretation of data. The vesicles tracked in this study are of mixed identity (e.g. peroxisomes, early endosomes, late endosomes, lysosomes, autophagosomes) and each could have distinct motility behavior.

Conclusion

In this study, we found that small-scale displacement fluctuations of vesicles in fibroblasts do not depend on their substrate stiffness. However, drugs perturbing the cytoskeleton (actin, microtubules, and myosin-II) do affect vesicle dynamics and this affect is increased on softer substrates. It was found that myosin-II has the largest effect on vesicle motion, as it decreases the diffusion coefficient of the local cytoplasm. Actin and microtubule filaments also modulate the mobility of vesicles but to a lesser extent. Our results support the idea that actin may trap vesicles and impede local movement³⁵. This increased understanding of the role of the cytoskeleton in local vesicle motion may help inform developments for targeted drug therapies in the future.

Materials and Methods

Hydrogel fabrication and cell culture. Polyacrylamide (PA) gels of varying elastic moduli ($E = 10$ and 40 kPa) were fabricated as previously reported⁴⁵. Gel stiffness was modulated by mixing acrylamide and bisacrylamide according to specifications reported in a well-established protocol⁴⁶. Monkey kidney fibroblast cells (CRL-1651, ATCC, Manassas, VA) were plated sparsely on PA gels (2500 cells/ cm^2) and immersed in media comprised of Dulbecco's Modified Eagle's Medium (Corning, Corning, NY) (4.5 g/L glucose, 4 mM L-glutamine) supplemented with 10% fetal bovine serum (Sigma-Aldrich, St. Louis, MO) and 1% Penicillin Streptomycin (Corning, Corning, NY).

Cytoskeletal perturbation. Three agents were used to perturb the cytoskeleton: cytochalasin-D, colchicine, and blebbistatin. Cytochalasin-D ($5 \mu\text{M}$) was used to depolymerize F-actin filaments^{47,48}. Colchicine ($10 \mu\text{M}$) was used to depolymerize microtubules^{49,50}. Blebbistatin ($25 \mu\text{M}$) was used to inactivate myosin-II motors⁵¹. In all cases cells were incubated in the perturbing agent for at least one hour before, and throughout the imaging process. After imaging, the perturbing agent was washed out and replaced with fresh cell media to observe cell recovery and viability.

Microscopy. All experiments were performed using differential interference microscopy on an Olympus IX81 microscope with a 40X UApo N340 water immersion objective (NA 1.15) (Olympus America Inc., Center Valley, PA) mounted on a vibration isolation table (Newport Corporation, Irvine, CA). An environmental chamber enclosed the experimental platform, and maintained cell culture conditions throughout imaging (5% CO₂, 70% humidity and 37 °C). Images were acquired with a Neo sCMOS camera (active pixels 1392 × 1040, resolution of 165 nm per pixel) (Andor Technology, Belfast, Northern Ireland) at 200 frames per second for a total duration of 30 seconds.

Image analysis and processing. Image stack alignment to remove drift was completed with a subpixel registration algorithm⁵². Vesicle motion was tracked using an algorithm for precise particle tracking by polynomial fitting with Gaussian weight³¹ in MATLAB (The MathWorks, Natick, MA, USA). This algorithm was developed specifically for tracking endogenous vesicles in living cells with noisy backgrounds at nanometer precision. In this study, we applied the algorithm to differential interference contrast microscopy images (representative image shown in Fig. S2). To compute the precision of our vesicle tracking analysis, we calculate the error to be $\sqrt{\text{MSD}(\tau = 5 \text{ ms})} \sim 25 \text{ nm}$ ³¹. This tracking error is consistent with that measured in other comparable studies^{9,35,33}. Note that any trajectory of a vesicle that moved within 3 μm of the cell boundary or nucleus was discarded. Similarly, trajectories of neighboring vesicles that came within 500 nm were discarded. The radii of vesicles were calculated, as in Rogers *et al.*³¹, and found to be $r = 619 \pm 95 \text{ nm}$. Vesicle motion was not observed to depend on vesicle size, in the small range of vesicle sizes observed ($400 < r < 800 \text{ nm}$). All computations, statistical analysis, and plots were generated in MATLAB.

Temporal mean squared displacement analysis. To discriminate between directed and non-directed motion, vesicle dynamics were analyzed using a temporal Mean Squared Displacement (*t*MSD) algorithm to extract the local dynamics of the vesicle as a function of time as done previously⁴. Briefly, the algorithm estimates the local behavior for time, *t*, by calculating the *t*MSD of a rolling window centered about the time of interest,

$$t\text{MSD}(\tau) = \langle |\mathbf{r}(t' + \tau) - \mathbf{r}(t')|^2 \rangle_{(t-t_{\max}/2 < t' < t+t_{\max}/2)} \quad (1)$$

where $\langle \dots \rangle$ indicates an average over timescale τ , $\mathbf{r}(t)$ is the position vector as a function of time, *t'* is the time in the rolling window, τ is the timescale, and τ_{\max} is the width of the rolling window ($\tau_{\max} = 500 \text{ ms}$). The *t*MSD is then fitted to a power law of the form, $t\text{MSD}(\tau) = C\tau^\alpha$ on the interval 100–160 ms, and the power law scaling, $\alpha(t)$, can be extracted as a function of time to indicate the type of diffusion the vesicle is undergoing ($\alpha < 1$, sub-diffusive; $\alpha = 1$, Brownian; $\alpha > 1$, super-diffusive). Based on previous calibrations, $\alpha_{\text{directed}} \geq 1.4$ was chosen as a conservative threshold for directed motion and $\alpha_{\text{non-directed}} < 1.4$ was considered non-directed. Only trajectories of particles longer than 200 frames were kept for analysis. An example is shown in Fig. S1.

Correlation functions. In addition to the *t*MSD we used the velocity autocorrelation function (VACF) and the van Hove correlation function (VHC) to quantify dynamics. The VACF,

$$\psi(\tau) = \frac{\langle \mathbf{v}(t) \cdot \mathbf{v}(t + \tau) \rangle}{\langle [\mathbf{v}(t)]^2 \rangle} \quad (2)$$

quantifies persistence of directional motion by its characteristic decay with τ . The effective diffusion coefficient of the vesicle was calculated from the VACF using the Green-Kubo relation, $D_{\text{eff}} = v_0^2 \int_0^\infty \psi(\tau) d\tau$ ³⁷. We use this integral approach to calculate D_{eff} instead of fitting $\text{MSD} = 4D\tau + v^2\tau^2$ because we often observe our vesicle motion to be sub-diffusive. The VHC, quantifies the probability distribution of vesicle displacements, $P(\Delta x, \tau)$ where $\Delta x = x(t + \tau) - x(t)$ for a given timescale, τ . The overall shape of the VHCs provide valuable information on the vesicle fluctuations³².

Data Availability

The datasets generated during and/or analysed during the current study are available from the corresponding author on reasonable request.

References

- Ross, J. L., Ali, M. Y. & Warshaw, D. M. Cargo transport: molecular motors navigate a complex cytoskeleton. *Curr. opinion cell biology* **20**, 41–47 (2008).
- Caspi, A., Granek, R. & Elbaum, M. Diffusion and directed motion in cellular transport. *Phys Rev E* **66**, 011916, <https://doi.org/10.1103/PhysRevE.66.011916> (2002).
- Arcizet, D., Meier, B., Sackmann, E., Rädler, J. O. & Heinrich, D. Temporal analysis of active and passive transport in living cells. *Phys Rev Lett* **101**, 248103, <https://doi.org/10.1103/PhysRevLett.101.248103> (2008).
- Ahmed, W. W., Williams, B. J., Silver, A. M. & Saif, T. A. Measuring nonequilibrium vesicle dynamics in neurons under tension. *Lab on a Chip* **13**, 570–578, <https://doi.org/10.1039/C2LC41109A> (2013).
- Koslover, E. F., Chan, C. K. & Theriot, J. A. Cytoplasmic flow and mixing due to deformation of motile cells. *Biophys. J* **113**, 2077–2087 (2017).
- Hirokawa, N., Nitta, R. & Okada, Y. The mechanisms of kinesin motor motility: lessons from the monomeric motor KIF1A. *Nat Rev Mol Cell Bio* **10**, 877–884, <https://doi.org/10.1038/nrm2807> (2009).
- Roberts, A. J., Kon, T., Knight, P. J., Sutoh, K. & Burgess, S. A. Functions and mechanics of dynein motor proteins. *Nat Rev Mol Cell Bio* **14**, nrm3667, <https://doi.org/10.1038/nrm3667> (2013).
- Schuh, M. An actin-dependent mechanism for long-range vesicle transport. *Nat Cell Biol* **13**, <https://doi.org/10.1038/ncb2353> (2011).
- Fakhri, N. *et al.* High-resolution mapping of intracellular fluctuations using carbon nanotubes. *Science* **344**, 1031–1035, <https://doi.org/10.1126/science.1250170> (2014).

10. Almonacid, M. *et al.* Active diffusion positions the nucleus in mouse oocytes. *Nat. Cell Biol.* **17**, 470–479, <https://doi.org/10.1038/ncb3131> (2015).
11. Brangwynne, C. P., Koenderink, G. H., C, M. F. & Weitz, D. A. Intracellular transport by active diffusion. *Trends Cell Biol.* **19**, 423–427, <https://doi.org/10.1016/j.tcb.2009.04.004> (2009).
12. Ahmed, W. W. *et al.* Active mechanics reveal molecular-scale force kinetics in living oocytes. *Biophys. J.* **114**, 1667–1679 (2018).
13. Golestanian, R. Enhanced diffusion of enzymes that catalyze exothermic reactions. *Phys Rev Lett* **115**, 108102 (2015).
14. Sengupta, S. *et al.* Enzyme molecules as nanomotors. *J. Am. Chem. Soc.* **135**, 1406–1414 (2013).
15. Höfling, F. & Franosch, T. Anomalous transport in the crowded world of biological cells. *Rep Prog Phys* **76**, 046602, <https://doi.org/10.1088/0034-4885/76/4/046602> (2013).
16. Goychuk, I., Kharchenko, V. O. & Metzler, R. How molecular motors work in the crowded environment of living cells: Coexistence and efficiency of normal and anomalous transport. *Plos One* **9**, e91700, <https://doi.org/10.1371/journal.pone.0091700> (2014).
17. Gal, N. & Weihs, D. Experimental evidence of strong anomalous diffusion in living cells. *Phys Rev E* **81**, 020903, <https://doi.org/10.1103/PhysRevE.81.020903> (2010).
18. M, T. I., Munteanu, E. L., Thon, G., Oddershede, L. & Kirstine, B. Anomalous diffusion in living yeast cells. *Phys. Rev. Lett.* **93**, 078102, <https://doi.org/10.1103/PhysRevLett.93.078102> (2004).
19. Salman, H., Gil, Y., Granek, R. & Elbaum, M. Microtubules, motor proteins, and anomalous mean squared displacements. *Chem Phys* **284**, 389–397, [https://doi.org/10.1016/S0301-0104\(02\)00669-9](https://doi.org/10.1016/S0301-0104(02)00669-9) (2002).
20. Regner, B. M., Tartakovsky, D. M. & Sejnowski, T. J. Identifying transport behavior of Single-Molecule trajectories. *Biophys J* **107**, 2345–2351, <https://doi.org/10.1016/j.bpj.2014.10.005> (2014).
21. Saxton, M. J. A biological interpretation of transient anomalous subdiffusion. i. qualitative model. *Biophys J* **92**, 1178–1191, <https://doi.org/10.1529/biophysj.106.092619> (2007).
22. Weiss, M., Elsner, M., Kartberg, F. & Nilsson, T. Anomalous subdiffusion is a measure for cytoplasmic crowding in living cells. *Biophys J* **87**, 3518–3524, <https://doi.org/10.1529/biophysj.104.044263> (2004).
23. Dix, J. A. & Verkman, A. Crowding effects on diffusion in solutions and cells. *Annu. Rev Biophys* **37**, 247–263, <https://doi.org/10.1146/annurev.biophys.37.032807.125824> (2008).
24. Fabry, B. *et al.* Scaling the microrheology of living cells. *Phys Rev Lett* **87**, 148102 (2001).
25. Chen, K., Wang, B., Guan, J. & Granick, S. Diagnosing heterogeneous dynamics in single-molecule/particle trajectories with multiscale wavelets. *ACS Nano* **7**, 8634–44, <https://doi.org/10.1021/nn402787a> (2013).
26. Gal, N., Lechtman-Goldstein, D. & Weihs, D. Particle tracking in living cells: a review of the mean square displacement method and beyond. *Rheol. Acta* **52**, 425–443 (2013).
27. Hancock, W. O. Bidirectional cargo transport: moving beyond tug of war. *Nat. Rev. Mol. Cell. Bio.* **15**, 615 (2014).
28. Rogers, S. L. & Gelfand, V. I. Membrane trafficking, organelle transport, and the cytoskeleton. *Curr. opinion in cell biology* **12**, 57–62 (2000).
29. Campàs, O. A toolbox to explore the mechanics of living embryonic tissues. *Semin. Cell Dev Biol.* **55**, 119–130, <https://doi.org/10.1016/j.semcdb.2016.03.011> (2016).
30. Solon, J., Levental, I., Sengupta, K., Georges, P. C. & Janmey, P. A. Fibroblast adaptation and stiffness matching to soft elastic substrates. *Biophys. J.* **93**, 4453–4461 (2007).
31. Rogers, S. S., Waigh, T. A., Zhao, X. & Lu, J. R. Precise particle tracking against a complicated background: polynomial fitting with gaussian weight. *Phys. Biol.* **4**, 220 (2007).
32. Ahmed, W. W. & Saif, T. A. Active transport of vesicles in neurons is modulated by mechanical tension. *Sci Reports* **4**, srep04481, <https://doi.org/10.1038/srep04481> (2015).
33. Lo, C.-M., Wang, H.-B., Dembo, M. & Wang, Y.-I. Cell movement is guided by the rigidity of the substrate. *Biophys. J* **79**, 144–152 (2000).
34. Giner, D. *et al.* Vesicle movements are governed by the size and dynamics of f-actin cytoskeletal structures in bovine chromaffin cells. *Neurosci.* **146**, 659–669 (2007).
35. Katrukha, E. A. *et al.* Probing cytoskeletal modulation of passive and active intracellular dynamics using nanobody-functionalized quantum dots. *Nat Commun* **8**, ncomms14772, <https://doi.org/10.1038/ncomms14772> (2017).
36. Hu, J. *et al.* Size- and speed-dependent mechanical behavior in living mammalian cytoplasm. *Proc Natl. Acad. Sci.* **114**, 9529–9534, <https://doi.org/10.1073/pnas.1702488114> (2017).
37. Boon, J. P. & Yip, S. *Molecular hydrodynamics* (Courier Corporation, 1980).
38. Murrell, M., Oakes, P. W., Lenz, M. & Gardel, M. L. Forcing cells into shape: the mechanics of actomyosin contractility. *Nat. Rev. Mol. Cell Biol.* **16**, 486 (2015).
39. Humphrey, D., Duggan, C., Saha, D., Smith, D. & Käs, J. Active fluidization of polymer networks through molecular motors. *Nature* **416**, 413 (2002).
40. Oriola, D., Alert, R. & Casademunt, J. Fluidization and active thinning by molecular kinetics in active gels. *Phys Rev Lett* **118**, 088002 (2017).
41. Reverey, J. F. *et al.* Superdiffusion dominates intracellular particle motion in the supercrowded cytoplasm of pathogenic *acanthamoeba castellanii*. *Sci Reports* **5**, 11690, <https://doi.org/10.1038/srep11690> (2015).
42. Sivaramakrishnan, S., DeGiulio, J. V., Lorand, L., Goldman, R. D. & Ridge, K. M. Micromechanical properties of keratin intermediate filament networks. *Proc. Natl. Acad. Sci.* **105**, 889–894 (2008).
43. Schopferer, M. *et al.* Desmin and vimentin intermediate filament networks: their viscoelastic properties investigated by mechanical rheometry. *J. molecular biology* **388**, 133–143 (2009).
44. Lin, C. *et al.* Active diffusion and microtubule-based transport oppose myosin forces to position organelles in cells. *Nat Commun* **7**, 11814, <https://doi.org/10.1038/ncomms11814> (2016).
45. Knoll, S., Ahmed, W. & reports, S. T. Contractile dynamics change before morphological cues during fluorescence illumination. *Sci. Reports.* <https://doi.org/10.1038/srep18513> (2015).
46. Kadow, C. E., Georges, P. C., Janmey, P. A. & Beningo, K. A. *Methods in Cell Biology - Polyacrylamide Hydrogels for Cell Mechanics: Steps Toward Optimization and Alternative Uses*, vol. 83 (2007).
47. Goddette, D. & Frieden, C. Actin polymerization. the mechanism of action of cytochalasin d. *J. Biol. Chem.* **261**, 15974–15980 (1986).
48. Cooper, J. A. Effects of cytochalasin and phalloidin on actin. *The J.cell biology* **105**, 1473–1478 (1987).
49. Dalbeth, N., Lauterio, T. J. & Wolfe, H. R. Mechanism of action of colchicine in the treatment of gout. *Clin. therapeutics* **36**, 1465–1479 (2014).
50. Kemkemer, R., Teichgräber, V., Schrank-Kaufmann, S., Kaufmann, D. & Gruler, H. Nematic order-disorder state transition in a liquid crystal analogue formed by oriented and migrating amoeboid cells. *The Eur. Phys. J. E* **3**, 101–110 (2000).
51. Kovács, M., Tóth, J., Hetényi, C., András, M. & Sellers, J. R. Mechanism of blebbistatin inhibition of myosin II. *J Biol Chem* **279**, 35557–35563, <https://doi.org/10.1074/jbc.m405319200> (2004).
52. Thevenaz, P., Ruttimann, U. E. & Unser, M. A pyramid approach to subpixel registration based on intensity. *IEEE transactions on image processing* **7**, 27–41 (1998).
53. Guo, M. *et al.* Probing the stochastic, motor-driven properties of the cytoplasm using force spectrum microscopy. *Cell* **158**, 822–832 (2014).

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Author Contributions

D.P., P.B.-P., W.A. analyzed the data and wrote the manuscript. S.K. and W.A. designed and conducted all experiments. T.S. and W.A. supervised the project. All authors reviewed the manuscript.

Additional Information

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