

# Synthesis and Photoluminescence Properties of MoS<sub>2</sub>/Graphene Heterostructure by Liquid-Phase Exfoliation

Durai Murugan Kandhasamy,\* Paulpandian Muthu Mareeswaran, Selvaraju Chellappan, Dhenadhayalan Namasivayam, Afaf Aldahish, and Kumarappan Chidambaram



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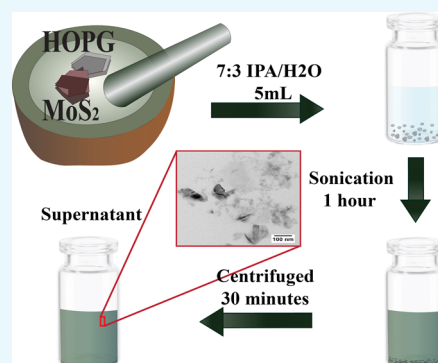
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**ABSTRACT:** Here, we report the synthesis of MoS<sub>2</sub>/graphene heterostructure in single-stage, liquid-phase exfoliation using a 7:3 isopropyl alcohol/water mixture. Further, the synthesized heterostructure was characterized using UV–visible and micro-Raman spectroscopies, transmission electron microscopy (TEM), and dynamic light scattering (DLS) analysis. UV–visible and micro-Raman analyses confirmed that the synthesized heterostructure had mostly few-layered (two-to-four sheets) MoS<sub>2</sub>. The photophysical properties of the heterostructure were analyzed using steady-state and time-resolved luminescence techniques. Enhanced photoluminescence was observed in the case of the heterostructure probably due to an increase in the defect sites or reduction in the rate of nonradiative decay upon formation of the sandwiched heterostructure. Applications of this heterostructure for fluorescence live-cell imaging were carried out, and the heterostructure demonstrated a better luminescence contrast compared to its individual counterpart MoS<sub>2</sub> in phosphate-buffered saline (PBS).



## INTRODUCTION

Electronic and optical properties of atomically thin, two-dimensional (2D) layered graphene and transition metal dichalcogenides (TMDCs) have attracted significant interest due to their dependence on layer thickness and crystallite size.<sup>1–6</sup> The presence of unsaturated d-orbitals and chemically active edge sites makes them act as semiconductors, semimetals, ferromagnets, and even superconductors.<sup>2</sup> Specifically, the inherent properties of two-dimensional (2D) materials are altered with a decrease in layer numbers to a single layer or a few layers, making them promising candidates in various fields of science.<sup>5,7–9</sup> Moreover, these 2D materials show enhanced biosensing and catalytic applications owing to their large surface area.<sup>6,10</sup> In particular, graphene is a widely known 2D material, which has become the research focus due to its excellent mechanical, thermal, electrical, and optical properties and its high specific surface area since its discovery in 2004.<sup>11–13</sup> It has shown potential applications in biosensors, batteries, photovoltaics, and supercapacitor applications.<sup>14</sup> Unlike graphene, MoS<sub>2</sub> is a direct-band-gap semiconductor with tunable electronic properties as a function of layer thickness and size.<sup>8,15,16</sup> MoS<sub>2</sub> shows potential applications in photocatalysis, energy storage, bioimaging, and biosensor applications.<sup>4,10,17–19</sup> Lateral and vertical heterostructures of graphene and TMDCs have attracted much interest due to their improved performance and stability in various fields ranging from energy conversion to biomedical applications.<sup>20–24</sup> The characteristic, high-quality heterointerface

strongly affects the built-in electric field and hence the charge transfer and dissociation of excitons. This results in the manipulation of the optical and electronic properties of 2D materials with improved performance toward specific applications in catalysis, sensing, and optoelectronic devices.<sup>25–30</sup>

Various methods have been proposed for the synthesis of high-quality heterostructures such as direct transfer of mechanically exfoliated materials, chemical vapor deposition, van der Waals epitaxy, and in situ catalytic growth using presynthesized precursor materials.<sup>31–34</sup> The large-scale synthesis of MoS<sub>2</sub>/graphene composite by a hydrothermal method using Mo and S precursors in the presence of exfoliated graphene for electrocatalytic water splitting was reported by Pandey et al.<sup>35</sup> Also, there are several reports on the large-scale synthesis of composite-like materials of graphene with other 2D materials.<sup>2,36,37</sup> However, synthesizing vertical/lateral stacks of these 2D heterostructures with controlled size and purity on a large scale is challenging and very rarely reported in the literature.<sup>34,38,39</sup>

Liquid-phase exfoliation (LPE) using *N*-methylpyrrolidone (NMP) and other organic solvents is shown to produce various

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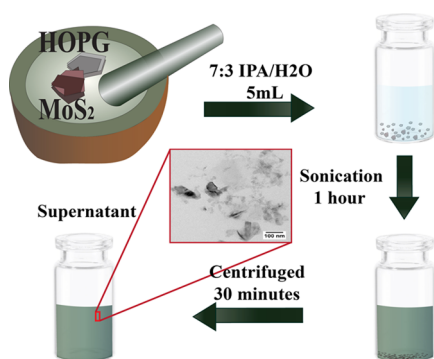


monolayer materials including graphene, hBN, MoS<sub>2</sub>, WS, etc.<sup>40–43</sup> Moreover, LPE using water and other low-boiling organic solvents can efficiently produce defect-free nanosheets of 2D materials including monolayers in suspension by a wet chemical method, followed by ultrasonication and centrifugation. Recently, it has been reported that liquid-phase exfoliation (LPE) effectively produces dielectric nanosheets by probing and matching the surface tension components of comparatively low-boiling solvents.<sup>44,45</sup> Monolayer enriched, luminescent nanosheets of MoS<sub>2</sub> and WS<sub>2</sub> have been synthesized in single-stage liquid-phase exfoliation with varying percentages of water/isopropanol mixture, and a 3:7 combination exhibits an enhanced exfoliation for these dichalcogenides.<sup>46,47</sup>

The large-scale synthesis of quality heterostructures is challenging and in some cases, solvents used for exfoliation may also chemically modify these materials.<sup>28,34,38</sup> But, in this work, for the first time we have successfully synthesized the MoS<sub>2</sub>/graphene heterostructure in a water/isopropyl alcohol solvent mixture via a scalable, single stage. Further, the method is benign to the environment and cost-effective compared to other methods reported. Thus, the synthesized material was characterized by spectroscopic techniques such as UV–visible and micro-Raman spectroscopies, and the structural aspects were probed using transmission electron microscopy (TEM) with small-angle electron diffraction (SAED) patterns. The luminescence properties were investigated by steady-state and time-resolved photoluminescence measurements. The applicability of this material toward live-cell imaging was studied, and better performance was observed for the heterostructure.

## RESULTS AND DISCUSSION

**UV–Visible and Micro-Raman Spectroscopic Analysis.** The heterostructure of MoS<sub>2</sub>/graphene and its individual counterparts MoS<sub>2</sub> and graphene suspensions was synthesized using MoS<sub>2</sub> and highly oriented pyrolytic graphite (HOPG) crystals as detailed in the **Materials and Methods** section (Figure 1). The UV–visible absorption spectra recorded for



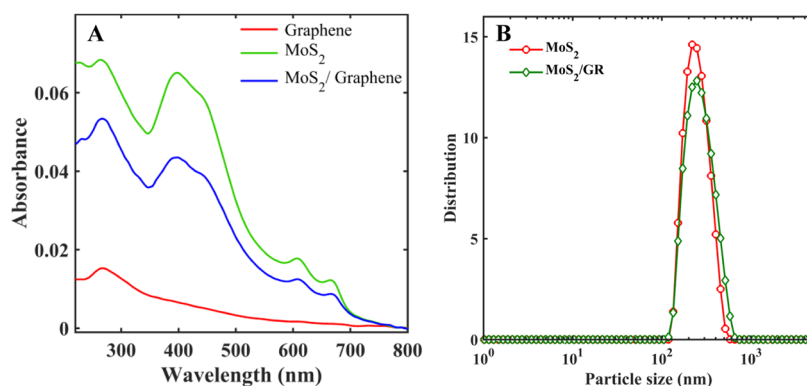
**Figure 1.** Synthesis of the MoS<sub>2</sub>/graphene heterostructure via the liquid-phase exfoliation method.

the MoS<sub>2</sub>/graphene heterostructure and its individual components exhibit the characteristic absorption bands of these materials. The absorption spectrum (Figure 2A, red) of exfoliated graphene in the IPA/H<sub>2</sub>O mixture shows a peak at around 270 nm, tailing until the end of the measured spectral range. This reveals that the graphene sample has both thin-layered graphene and thick-multilayered ones. The signature absorption maxima due to the exciton states of MoS<sub>2</sub> was

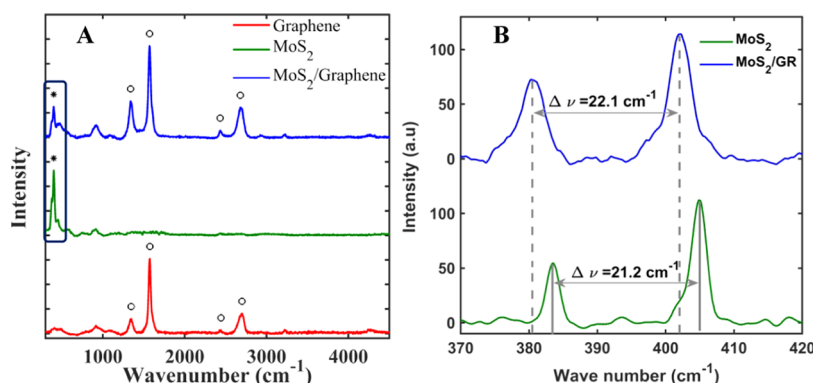
observed in the liquid-phase exfoliated MoS<sub>2</sub> and the MoS<sub>2</sub>/graphene heterostructure at 430, 607, and 665 nm, and they are assigned to C, B, and A exciton states of molybdenum disulfide, respectively.<sup>48</sup> A detailed investigation on the effect of exfoliated sheet thickness and size on the absorption maximum of MoS<sub>2</sub> has been documented before.<sup>5</sup> The observed “A” exciton maximum at 665 nm in both MoS<sub>2</sub> and its heterostructure clearly indicates that the obtained suspension has mostly thin layers of one-to-four sheet thickness. The ratio between the normalized absorption low at ~340 and B exciton maximum at 607 nm ( $OD_B/OD_{340}$ ) gives us the lateral size and the concentration of the particles in the suspension.<sup>5</sup> Estimated ratios of 0.38 and 0.34 indicate a ~200 nm lateral size for MoS<sub>2</sub> and MoS<sub>2</sub>/graphene dispersions. These observations were further confirmed by particle size measurement using dynamic light scattering.

The particles dried and resuspended in MilliQ water were used for dynamic light scattering measurements, and the profile of the lateral size is shown in Figure 2B. The mean lateral size of the particles was found to be  $243 \pm 72$  and  $261 \pm 93$  nm for MoS<sub>2</sub> and the MoS<sub>2</sub>/graphene heterostructure, respectively. This is again consistent with the estimated lateral size using UV–visible spectra of the sample as mentioned in the previous section. The lateral size of the particles depends on the nature of the exfoliating medium and ultrasonication time.<sup>40</sup> The lateral size observed is comparable to the previous report with a larger polydispersity index for MoS<sub>2</sub> alone under identical conditions.<sup>46</sup> Moreover, the synthesized heterostructure reveals an increase in lateral size and polydispersity index. This may be attributed to the formation of a heterostructure with varying dimensions and thicknesses. MoS<sub>2</sub> and the MoS<sub>2</sub>/graphene heterostructure show a  $\zeta$ -potential value of  $-59$  and  $-62$  mV, respectively. This indicates that the surfaces of the exfoliated materials are negatively charged.  $\zeta$ -Potential of the exfoliated nanosheets of MoS<sub>2</sub> and graphene is known to be dependent on the pH of the solution.<sup>49,50</sup> Both show a wide range of  $\zeta$ -potential values between  $+10$  and  $-60$  mV as a function of the pH of the solution. A value of around  $-50$  mV was reported under neutral pH. In the present case, a slightly higher value of  $-63$  mV was observed for the heterostructure. This shows that the nature of charge density does not change upon the formation of the heterostructure.

The Raman spectrum of the two-dimensional materials provides a wealth of information regarding the thickness, defects, and more specifically, the nature of the interaction between layers in the case of heterostructures. The Raman spectrum of liquid-phase exfoliated graphene exhibits peaks at 1347, 1575, 2444, and 2701 cm<sup>-1</sup> assignable to D, G, D\*, and 2D bands, respectively (Figure 3A). An intense D band observed in the Raman signal indicates an increase in defect sites leading to more number of sp<sup>3</sup>-hybridized carbon atoms.<sup>14</sup> The heterostructure exhibits Raman features at 1340, 1571, 2434, and 2671 cm<sup>-1</sup>. The peak positions are due to the same vibronic bands of graphene mentioned above with the shift in the peak positions toward lower wavenumbers. This shift may be due to the formation of the heterostructure, which in turn affects the out-of-plane vibrations. The integrated intensity ratio  $I_D/I_G$  for the D band and G band is widely used for characterizing the defect quantity in graphitic materials. The estimated values of the  $I_D/I_G$  ratio were 0.1970 and 0.3946 for exfoliated graphene and the MoS<sub>2</sub>/graphene heterostructure, respectively.



**Figure 2.** (A) UV–visible spectra of liquid-phase exfoliated graphene (red line), MoS<sub>2</sub> (green line), and the MoS<sub>2</sub>/graphene (blue line) heterostructure in a 3:7 water/isopropanol mixture. (B) Lateral size histogram measured by dynamic light scattering of liquid-phase exfoliated MoS<sub>2</sub> (red line) and the MoS<sub>2</sub>/graphene heterostructure (green line) in a 3:7 water/isopropanol mixture.



**Figure 3.** Micro-Raman spectra of liquid-phase exfoliated graphene (red line), MoS<sub>2</sub> (green line), and the MoS<sub>2</sub>/graphene heterostructure in a 3:7 water/isopropanol mixture in a (A) wider scan range and a (B) high-resolution scan.

The quantitative amount of defects in the exfoliated graphene was estimated using Tuinstra–Koenig relation as follows<sup>51,52</sup>

$$L_a = (2.4 \times 10^{-10})\lambda^4(I_D/I_G)^{-1} \quad (1)$$

where  $L_a$  is the in-plane crystallite size in nanometer and  $\lambda$  is the wavelength of the excitation source (532 nm in the present measurement). The estimated in-plane crystallite sizes were 97 and 48 nm for the exfoliated graphene and the MoS<sub>2</sub>/graphene heterostructure, respectively. The translational symmetry of the graphene sheets was directly related to the presence of defects and structural disorders.<sup>53</sup> An increase in defects and/or disorders in the structure leads to a reduced in-plane crystallite size. The density of defects is calculated from the in-plane crystallite size using the following relation

$$n_D (\text{cm}^{-2}) = 10^{14}/\pi L_a^2 \quad (2)$$

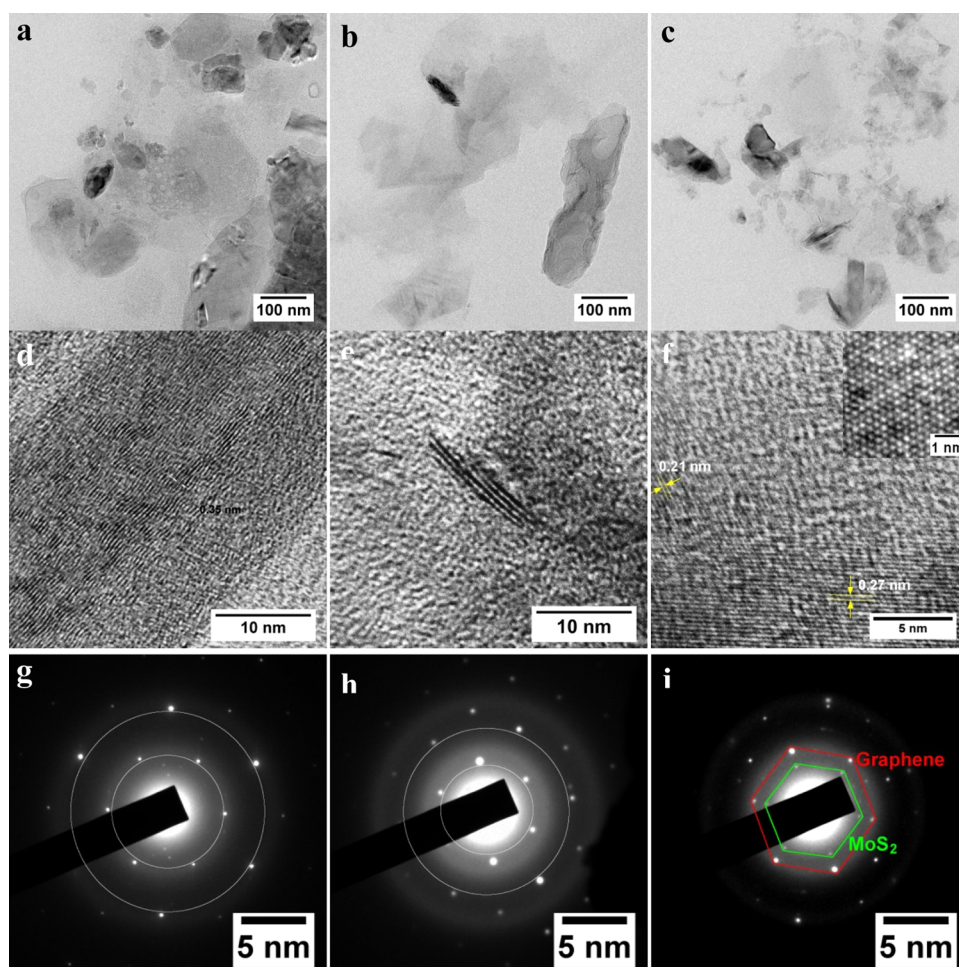
where  $n_D$  is the defect density and estimated to be  $0.33 \times 10^{10}$  and  $1.34 \times 10^{10} \text{ cm}^{-2}$  for the exfoliated graphene and the MoS<sub>2</sub>/graphene heterostructure, respectively. The decrease in the size of the in-plane crystallites and the higher defect density upon the formation of the heterostructure may be attributed to the effective exfoliation under the chosen solvent mixture and/or partial oxidation of graphene during ultrasonic treatment. The width of the “G” peak also implies the nature of defects and the presence of other impurities. In the present case, the observed width of  $\sim 15 \text{ cm}^{-1}$  for both exfoliated graphene and the MoS<sub>2</sub>/graphene heterostructure implies that the material is almost equal to mechanically exfoliated graphene in purity.

Such quality production of few-layered heterostructures is attributed to the usage of low-boiling and environmentally benign solvent mixtures during exfoliation.

Raman spectra of MoS<sub>2</sub> and the MoS<sub>2</sub>/graphene heterostructure, shown in Figure 3A, have a peak at  $\sim 400 \text{ cm}^{-1}$  accompanied by a shoulder. To resolve the spectral signs of MoS<sub>2</sub>, a narrow scan was carried out (Figure 3B), which clearly reveals peaks at 383.6 and 404.8  $\text{cm}^{-1}$  with a peak separation of 21.2  $\text{cm}^{-1}$ . Lattice vibrations and the intralayer bonding in stacked few-layered crystallites were perturbed by weak van der Waals interlayer interactions. The E<sub>2g</sub> mode assigned to the opposite vibration of S atoms with respect to the Mo atom and the out-of-plane vibration of only S atoms in opposite directions yields the A<sub>1g</sub> mode.<sup>54</sup> The layer thickness-dependent frequency shift of both modes has been examined, and the frequency of the E<sub>2g</sub> peak decreases while that of the A<sub>1g</sub> peak increases with increasing layer thickness.

In addition, the peak separation and the intensity of the peaks E<sub>2g</sub> and A<sub>1g</sub> increase with the increase in the number of layers.<sup>15</sup> The liquid exfoliated MoS<sub>2</sub> sample prepared by the present mixed-solvent method has a sheet with a thickness of two to four layers as evident from the frequencies of the A<sub>1g</sub> and E<sub>2g</sub> modes and the separation frequency of the two peaks. However, a shift in the peak position compared to the mechanically exfoliated sample may be attributed to the nature of the solvent/material interaction or partial oxidation of MoS<sub>2</sub> nanosheets. The peak frequencies of both E<sub>2g</sub> and A<sub>1g</sub> show a shift toward lower wavenumbers in the presence of graphene, which is ascribed to the constrained out-of-plane vibrations





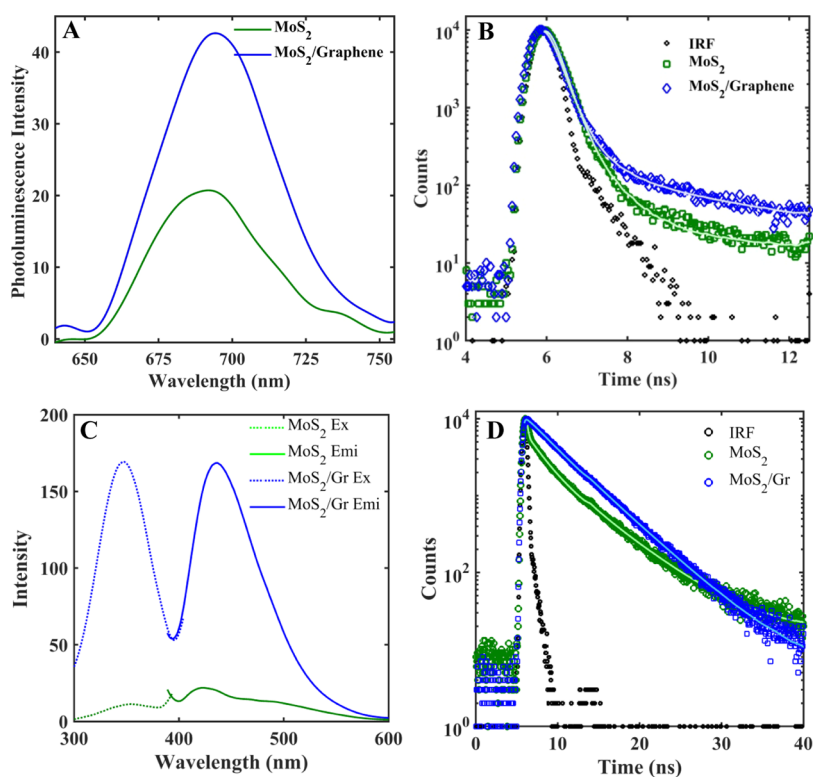
**Figure 4.** Transmission electron microscopy (TEM) images: (a, d) graphene, (b, e) MoS<sub>2</sub>, and (c, f) the MoS<sub>2</sub>/graphene heterostructure. Top panel: large-area scan and middle panel: high-resolution scan. SAED pattern of (g) graphene, (h) MoS<sub>2</sub>, and (i) the MoS<sub>2</sub>/graphene heterostructure.

due to the effective formation of the heterostructure. Also, every particle or position monitored during Raman measurement shows the features of both MoS<sub>2</sub> and graphene. This is a clear indication of the absence of free particles of both MoS<sub>2</sub> and graphene in the synthesized heterostructure. At present, the nature of interaction dominating the formation of the heterostructure is not clear. However, the balance between weak hydrogen bonding and hydrophobic and hydrophilic interaction might be the driving force for the effective formation of the heterostructure during liquid-phase exfoliation.

**Structural Analysis of the Heterostructure.** The size and lattice alignment of the MoS<sub>2</sub>/graphene heterostructure were analyzed using transmission electron microscopy and small-angle electron diffraction as shown in Figure 4. A typical hexagonal diffraction pattern was seen in both MoS<sub>2</sub> and graphene SAED patterns with a lattice spacing of 0.67 and 0.35 nm, respectively.<sup>55</sup> The observed lattice spacing confirmed that the sonication-assisted liquid-phase exfoliation did not alter the interlayer distance as seen in some of the previous studies. The vertical stacking of the synthesized heterostructure is also evident from the high-resolution image shown in Figure 4f. Vertical heterostructures of TMDCs with graphene have been achieved with controlled chemical vapor deposition and mechanical exfoliation followed by transfer.<sup>56,57</sup> However,

this is the first time vertical heterostructures of MoS<sub>2</sub>/graphene were synthesized by single-stage, scalable liquid-phase exfoliation. Moreover, the position of the SAED pattern reveals that the angle of lattice orientation between graphene and MoS<sub>2</sub> might be either 0 or 60° in many cases, as we did not notice any other orientation angle in the SAED pattern. The characterizations carried out above unambiguously confirmed that the mixed-solvent liquid-phase exfoliation method adopted to synthesize heterostructures of MoS<sub>2</sub>/graphene was effective, and the strategy is scalable for large-scale preparation of heterostructures with the advantage of the combined/modulated optical and electrical properties of both.

The photoluminescence behavior of MoS<sub>2</sub> in suspension and its heterostructure with graphene was evaluated by steady-state and time-resolved measurements. MoS<sub>2</sub> is a direct-band semiconductor, and the band-gap emission occurs at 1.9 eV (~650 nm) due to the interband transition of the K point.<sup>47</sup> This interband photoluminescence is observed at 688 nm for MoS<sub>2</sub> on excitation using 532 nm wavelength in suspension (Figure 5A). The mechanically exfoliated MoS<sub>2</sub> monolayer exhibits strong photoluminescence with a peak at 650–690 nm and a shoulder at around 610 nm due to A and B excitonic bands, respectively. The photoluminescence peak maximum of the monolayer and few layers shifts on different substrates such as sapphire, SiO<sub>2</sub>, Si, and fused silica.<sup>58</sup> Moreover, the nature of



**Figure 5.** (A) Photoluminescence excitation (dotted line) and emission spectra (solid line) of MoS<sub>2</sub> (green line) and the MoS<sub>2</sub>/graphene heterostructure in a 3:7 water/isopropanol mixture, (B, D) photoluminescence decay profile monitored at 450 and 680 nm on excitation at 375 and 470 nm wavelengths using picosecond laser diodes, respectively, and (C) photoluminescence emission spectra corresponding to the band-gap excitation measured by exciting with a 532 nm laser source.

impurities also plays a critical role in the photoluminescence intensity and peak position. Chemically exfoliated MoS<sub>2</sub> exhibits metallic behavior due to lithium intercalation under normal conditions and upon mild annealing results in the restoration of the intrinsic photoluminescence of the thin film.<sup>47</sup> In the present case, the observed photoluminescence slightly shifted to the red region, which may be due to solvent–solute interaction in the suspended MoS<sub>2</sub>. However, the MoS<sub>2</sub>/graphene heterostructure shows a peak maximum at 695 nm, which shows a further red shift compared to MoS<sub>2</sub> alone in the solvent system. Further, it is observed that the photoluminescence intensity increased to twice that of MoS<sub>2</sub> alone. This could be explained by the fact that interlayer coupling between MoS<sub>2</sub> and graphene results in an increase in layer thickness.

The photoluminescence intensity and the peak position depend on the layer thickness, and a red shift in the emission intensity is observed upon increasing the layer thickness.<sup>47</sup> On the other hand, the increase in the PL intensity may be attributed to the interlayer coupling of MoS<sub>2</sub> and graphene layers, which in turn reduces the rate of nonradiative decay.<sup>59</sup> Further, the photoluminescence behavior was analyzed by time-resolved PL decay as shown in Figure 5B and Table 1. The decay obeys a biexponential function with a prominent shorter decay time of 0.2 ns (95%) and a longer decay component of 4.0 ns (5%). The average decay time was estimated to be 0.4 ns. The MoS<sub>2</sub>/graphene heterostructure also follows biexponential decay with decay times of 4.10 ns (10%) and 0.25 ns (90%), and the average lifetime was found to be 0.64 ns in the suspension. An average lifetime of 410 ps was reported by Mohamed et al. for the monolayer MoS<sub>2</sub>, and

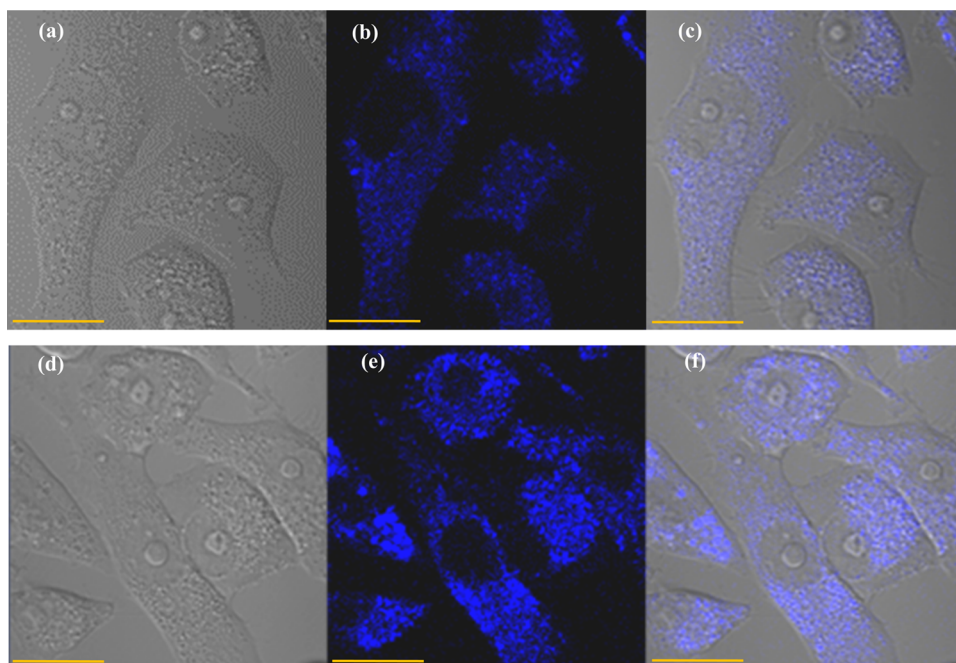
**Table 1. Photoluminescence Parameters of Liquid-Phase Exfoliated MoS<sub>2</sub> and the MoS<sub>2</sub>/Graphene Heterostructure<sup>a</sup>**

sample	$\lambda_{\text{max}}$ emission (nm)	decay time		
		$\tau_1$ (ns)	$\tau_2$ (ns)	$\langle\tau\rangle$ (ns)
MoS <sub>2</sub>	430 (360)	5.35 (67)	1.95 (33)	4.23
	688 (532)	4.08 (5)	0.23 (95)	0.42
MoS <sub>2</sub> /graphene	438 (355)	4.75 (71)	3.01 (29)	4.25
	695 (532)	4.10 (10)	0.25 (90)	0.64

<sup>a</sup>Excitation wavelength corresponding to the emission maximum is given in parentheses. Amplitudes for the respective decay times are given in parentheses.

the observed lifetime of 420 ps matches with their reports for MoS<sub>2</sub> suspension.<sup>60</sup> Moreover, the enhancement in decay time upon the formation of the heterostructure is consistent with the enhancement of luminescence intensity. This again validates that upon the formation of the heterostructure, the rate of nonradiative decay is decreased.

Earlier, it was observed that the suspension of MoS<sub>2</sub> in solution exhibits strong luminescence in the blue region due to defect sites or size quantization in addition to band-edge emission.<sup>61–63</sup> The effect of the heterostructure on the photoluminescence properties of MoS<sub>2</sub> due to defect sites has also been explored. Figure 5C,D shows the steady-state and time-resolved photoluminescence of MoS<sub>2</sub> and the MoS<sub>2</sub>/graphene heterostructure. Maximum photoluminescence intensity was observed upon excitation at 360 nm with a peak at 430 nm, and the respective excitation spectrum is also shown in Figure 5C. MoS<sub>2</sub> sheets exhibit excitation-dependent



**Figure 6.** (a, d) Bright-field images, (b, e) fluorescence images, and (c, f) merged images of HeLa cells incubated for 1 h in MoS<sub>2</sub> and the MoS<sub>2</sub>/graphene heterostructure in PBS buffer.

emission in the blue region attributed to inhomogeneous size distribution and the presence of multiple trap states.<sup>43,61</sup>

Interestingly, more than 20-fold enhancement in the emission intensity was observed for the heterostructure, as seen in Figure 5C. This may be due to the combined effect of an increase in the number of defect sites and the reduction in the nonradiative rate in the heterostructure. This is further ascertained by measuring the time-resolved decay profile as shown in Figure 5D. The decay profile displays a biexponential decay with decay times of 5.3 ns (67%) and 1.95 ns (33%) for MoS<sub>2</sub>. The average decay times were estimated to be 4.23 and 4.25 ns for MoS<sub>2</sub> and the MoS<sub>2</sub>/graphene heterostructure, respectively. The nearly identical decay time for both and an enhancement in the PL intensity may be correlated to the increase in defect sites upon the heterostructure formation.

The effect of MoS<sub>2</sub> on normal and cancer cell lines has been studied earlier, and the results indicate that MoS<sub>2</sub> adhesion leads to cell death in cancerous cell lines and has no effect on the normal cell lines after 24 h of incubation.<sup>18</sup> The uptake of MoS<sub>2</sub> and its localization inside the cell is probed using transient absorption and Raman scattering in a custom-built dual-mode microscope using HeLa cell lines.<sup>64</sup> However, live-cell imaging using MoS<sub>2</sub> as a luminescent probe is limited due to the poor photoluminescence of the material in biological medium.<sup>61,65</sup> The heterostructure obtained by the present strategy has more than 20-fold enhancement in luminescence intensity compared to MoS<sub>2</sub> nanosheets under identical conditions and might be a potential candidate for live-cell imaging using the photoluminescence of the material. HeLa cell lines incubated with 50 mL of the resuspended MoS<sub>2</sub>/graphene heterostructure were used for imaging as mentioned in the Materials and Methods section in a biological buffer medium (Figure 6). Bright-field images of the cells incubated with 50  $\mu$ L of the redispersed MoS<sub>2</sub> and its heterostructure are shown in Figure 6a,d, and it is noticed that the cells are alive up to 12 h in the incubated medium. Fluorescent and merged images of the cells reveal that both the materials internalized

into the cytoplasm of the cells. In addition, it is clearly visible that the heterostructure has an enhanced fluorescent contrast compared to MoS<sub>2</sub> alone in an identical environment.

## CONCLUSIONS

The heterostructure of MoS<sub>2</sub> with graphene was successfully fabricated by single-stage liquid-phase exfoliation using an IPA/water mixture (7:3 v/v). Thus, the obtained van der Waals heterostructure was characterized using Raman, UV-visible, TEM, and dynamic light scattering studies. The heterostructures have two to five layers of MoS<sub>2</sub> as evident from Raman and UV-visible analyses. DLS, UV-visible, and TEM analyses revealed that the lateral size of the material was about 200 nm and the surface of both MoS<sub>2</sub> and the heterostructure was negatively charged with  $\zeta$ -potential values of  $\sim -60$  mV. TEM with SAED confirm the formation of the vertical heterostructure of MoS<sub>2</sub>/graphene with the present strategy. About 20-fold enhancement in the photoluminescence intensity was attributed to the suppression of non-radiative decay pathways due to the formation of the heterostructure, which is also clear from higher decay times in time-resolved photoluminescence measurement. The material shows a potential application in live-cell imaging using photoluminescence compared to MoS<sub>2</sub> with the added advantage of graphene in the heterostructure.

## MATERIALS AND METHODS

**Synthesis of the MoS<sub>2</sub>/Graphene Heterostructure.** HOPG and MoS<sub>2</sub> crystals used in this work were purchased from 2D Semiconductors. All other reagents were of analytical grade and used as received. Millipore water was used throughout the study. The individually exfoliated materials were prepared following the methods described earlier.<sup>46</sup> The heterostructure was synthesized by slightly modifying the above protocol (Figure 1). Briefly, 5 mg of graphene and MoS<sub>2</sub> were taken in a mortar and ground into a powder. Then, the



contents were transferred to a glass vial and sonicated for 1 h in a 3:7 v/v water/IPA mixture. The resultant suspension was centrifuged at 2000 rpm for 10 min, and the supernatant was used for further characterization and studies.

**Characterization of the MoS<sub>2</sub>/Graphene Heterostructure.** Absorption spectra for the suspensions at room temperature were recorded on a Shimadzu 2401 UV–visible spectrophotometer with a spectral range of 250–1100 nm. Micro-Raman analysis at room temperature was carried out on individual sheets drop-cast on a pre-cleaned borosilicate coverslip. The spectrographs were recorded using a BLZE-100 HR EMCCD camera (Princeton Instruments) equipped with an Acton SP2500 monochromator and a confocal microscope. A 532 nm continuous-wave, tunable power laser was used for excitation through a 60× objective. The lateral size and  $\zeta$ -potential of the exfoliated samples were recorded using a Horiba Nanopartica SZ-100 instrument by dynamic light scattering and  $\zeta$ -potential measurement, respectively. Transmission electron microscopy images with small-angle electron diffraction were recorded for the samples coated on a copper grid and dried in vacuum using a JEOL-2100 + high-resolution transmission electron microscope. Confocal fluorescence images were acquired with a Zeiss LSM 710 under the excitation wavelength of 405 nm.

**Photoluminescence Studies of the MoS<sub>2</sub>/Graphene Heterostructure.** Photoluminescence emission and excitation spectra were recorded on a Horiba–Jobin Yvon FluoroMax-4P spectrofluorometer equipped with a Xe lamp for excitation and an R928P photon counting module for the detection of light emitted from the sample. The emission from the direct-band-gap excitation was monitored using a 5 mW, 532 nm CW laser as the excitation source and a USB spectrometer (RI Instruments and Innovation, India). Photoluminescence decay was monitored using the TCSPC technique (Model 5000U, IBH, U.K.) and a 375 nm, picosecond pulse LED as the excitation source, and an MCP-PMT (Hamamatsu R3809U) was used for detection of photons emitted after passing through a monochromator. The decay was fitted with Horiba Scientific decay analysis software DAS6 to a bi/triexponential behavior.

**In Vitro Fluorescence Bioimaging.** HeLa cells were obtained from American Type Culture Collection and cultured in RPMI-1640 medium supplemented with FBS (10%), antibiotic-antimycotic solution (1%), L-glutamine (2 mM), and nonessential amino acids (1%). The live cells were maintained at 37 °C and 5% CO<sub>2</sub> in a humidified chamber. HeLa cells were incubated with MoS<sub>2</sub> (50  $\mu$ L) for 30 min to equilibrate the experimental conditions. After incubation, the cells were washed with PBS buffer to remove extracellular MoS<sub>2</sub> prior to recording images. The fluorescence images of MoS<sub>2</sub> and the MoS<sub>2</sub>/graphene heterostructure incubated cells were recorded at the excitation wavelength of 405 nm with confocal fluorescence microscopy.

## AUTHOR INFORMATION

### Corresponding Author

Durai Murugan Kandhasamy – Department of Bioelectronics and Biosensors, Alagappa University, Karaikudi 630003 Tamil Nadu, India; [orcid.org/0000-0003-2925-4728](https://orcid.org/0000-0003-2925-4728); Email: [kdmurugan@gmail.com](mailto:kdmurugan@gmail.com)

## Authors

Paulpandian Muthu Mareeswaran – Department of Industrial Chemistry, Alagappa University, Karaikudi 630003 Tamil Nadu, India; [orcid.org/0000-0002-2901-4110](https://orcid.org/0000-0002-2901-4110)

Selvaraju Chellappan – National Centre for Ultrafast Processes, University of Madras, Chennai 600113, India; [orcid.org/0000-0002-3625-4243](https://orcid.org/0000-0002-3625-4243)

Dhenadhayalan Namasivayam – Department of Chemistry, National Taiwan University and Institute of Atomic and Molecular Sciences, Academia Sinica, Taipei 10617, Taiwan; [orcid.org/0000-0003-4494-5746](https://orcid.org/0000-0003-4494-5746)

Afaf Aldahish – Department of Pharmacology, School of Pharmacy, King Khalid University, Abha 62529, Saudi Arabia

Kumarappan Chidambaram – Department of Pharmacology, School of Pharmacy, King Khalid University, Abha 62529, Saudi Arabia

Complete contact information is available at: <https://pubs.acs.org/10.1021/acsomega.1c05250>

## Notes

The authors declare no competing financial interest.

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