


Article

High Thermoelectric Performance Achieved in Sb-Doped GeTe by Manipulating Carrier Concentration and Nanoscale Twin Grains

Chao Li ^{1,2,3,*}, Haili Song ^{4,*}, Zongbei Dai ³, Zhenbo Zhao ³, Chengyan Liu ⁵, Hengquan Yang ^{6,*}, Chengqiang Cui ^{1,2,*} and Lei Miao ^{5,7,*} 

- ¹ School of Electromechanical Engineering, Guangdong University of Technology, Guangzhou 510006, China
² Ji Hua Laboratory, Foshan 528299, China
³ The Fifth Electronics Research Institute of Ministry of Industry and Information Technology, Guangzhou 510006, China; dai_zongbei@126.com (Z.D.); zhenbozhao83@163.com (Z.Z.)
⁴ Key Laboratory of Bioinorganic and Synthetic Chemistry of Ministry of Education, School of Chemistry, Sun Yat-Sen University, Guangzhou 510275, China
⁵ Guangxi Key Laboratory of Information Material, Guangxi Collaborative Innovation Center of Structure and Property for New Energy and Materials, School of Material Science and Engineering, Guilin University of Electronic Technology, Guilin 541004, China; chengyanliu@guet.edu.cn
⁶ Jiangsu Key Laboratory of Modern Measurement Technology and Intelligent Systems, School of Physics and Electronic & Electrical Engineering, Huaiyin Normal University, Huai'an 223300, China
⁷ SIT Research Laboratories, Innovative Global Program, Department of Materials Science and Engineering, Faculty of Engineering, Shibaura Institute of Technology, Tokyo 135-8548, Japan
* Correspondence: chaoli@fzxsmc.com (C.L.); songhli@mail.sysu.edu.cn (H.S.); yhq@hytc.edu.cn (H.Y.); cqccui@gdut.edu.cn (C.C.); miaolei@guet.edu.cn (L.M.)
† These authors contributed equally to this work.



Citation: Li, C.; Song, H.; Dai, Z.; Zhao, Z.; Liu, C.; Yang, H.; Cui, C.; Miao, L. High Thermoelectric Performance Achieved in Sb-Doped GeTe by Manipulating Carrier Concentration and Nanoscale Twin Grains. *Materials* **2022**, *15*, 406. <https://doi.org/10.3390/ma15020406>

Academic Editor: Je-Hyeong Bahk

Received: 23 October 2021

Accepted: 20 December 2021

Published: 6 January 2022

Publisher's Note: MDPI stays neutral with regard to jurisdictional claims in published maps and institutional affiliations.



Copyright: © 2022 by the authors. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (<https://creativecommons.org/licenses/by/4.0/>).

Abstract: Lead-free and eco-friendly GeTe shows promising mid-temperature thermoelectric applications. However, a low Seebeck coefficient due to its intrinsically high hole concentration induced by Ge vacancies, and a relatively high thermal conductivity result in inferior thermoelectric performance in pristine GeTe. Extrinsic dopants such as Sb, Bi, and Y could play a crucial role in regulating the hole concentration of GeTe because of their different valence states as cations and high solubility in GeTe. Here we investigate the thermoelectric performance of GeTe upon Sb doping, and demonstrate a high maximum zT value up to 1.88 in $\text{Ge}_{0.90}\text{Sb}_{0.10}\text{Te}$ as a result of the significant suppression in thermal conductivity while maintaining a high power factor. The maintained high power factor is due to the markable enhancement in the Seebeck coefficient, which could be attributed to the significant suppression of hole concentration and the valence band convergence upon Sb doping, while the low thermal conductivity stems from the suppression of electronic thermal conductivity due to the increase in electrical resistivity and the lowering of lattice thermal conductivity through strengthening the phonon scattering by lattice distortion, dislocations, and twin boundaries. The excellent thermoelectric performance of $\text{Ge}_{0.90}\text{Sb}_{0.10}\text{Te}$ shows good reproducibility and thermal stability. This work confirms that $\text{Ge}_{0.90}\text{Sb}_{0.10}\text{Te}$ is a superior thermoelectric material for practical application.

Keywords: GeTe; Sb-doping; optimizing carrier concentration; nanoscale twin grains

1. Introduction

Thermoelectric (TE) materials can convert heat directly into available electricity based on the Seebeck effect, providing an alternative way to utilize fossil energy more efficiently [1–6]. The thermoelectric material performance can be estimated by the dimensionless figure of merit $zT = S^2T\sigma/\kappa$, where S is the Seebeck coefficient, σ is the electrical conductivity, T is the absolute temperature, and κ is the thermal conductivity, which usually stems from the crystal lattice vibrations (κ_l) and the charge carriers drift (κ_e). Hence, a good thermoelectric material should have a large σ and S along with a low κ [7].

In recent decades, IV–VI semiconductors with a narrow band gap have made a great progress in terms of theory and TE performance, showing a bright prospect in TE application [7–11]. These compounds exhibit a unique TE-favorable property portfolio such as high valley degeneracy, large dielectric function, and strong lattice anharmonicity empowered by metavalent bonding [12]. For example, the maximum zT exceeds 2.5 and it has been attained in PbTe and SnSe TE materials [13–17]. However, the toxicity of Pb and the easy cleavage of SnSe limit their commercialization [18]. Alternatively, GeTe is an eco-friendly counterpart in the family of IV–VI semiconductors, which also exhibits an excellent TE performance [18–20]. GeTe undergoes ferroelectric structural transition, transforming from low-temperature rhombohedral phase ($R\bar{3}m$, the primitive cell could be regarded as a slightly distorted cubic along $\langle 111 \rangle$ direction) into high-temperature cubic one ($Fm\bar{3}m$) with Curie temperature between 600 and 700 K [21]. It is believed that such a slight symmetry breaking is conducive to converging the splitting valence bands and reducing the κ_l [22,23]. Additionally, some elements such as Sb, Bi, and Y generally have a high doping limit in GeTe, for instance, the solubility of Sb exceeds 10% in GeTe. Undoubtedly, such features are beneficial to adequately optimize the electronic transport properties, and, in turn, yield a high zT , e.g., the maximum zT over 2 is attainable in (Bi, Cu), (Bi, Pb), (Sb, Pb), and (Sb, In) codoping systems [24–28].

However, the pristine GeTe commonly exhibits an inferior TE performance due to its low S and high κ , i.e., the large number of Ge vacancies causes a high hole concentration ($\sim 10^{21} \text{ cm}^{-3}$), a low S ($\sim 34 \mu\text{V K}^{-1}$), and a high κ ($\sim 8 \text{ W m}^{-1} \text{ K}^{-1}$) at room temperature [29]. To suppress Ge vacancies and hole concentration, Cu interstitials and Sb/Bi/Y substitutions were introduced into GeTe matrix to compensate for the excess holes or increase the formation energy of Ge vacancy. For example, the hole concentration approaching $\sim 10^{20} \text{ cm}^{-3}$ could be obtained by Sb/Bi/Y doping, and the available maximum zT is up to 1.7–1.85 [30–33]. Aside from holes concentration optimization, the resonant doping and valence band alignment could also increase the S . It has been demonstrated that In/Ga/Ti doping could introduce the resonant level around the Fermi level [34–37], and Mn/Cd/Zn doping could decrease the split energy between light and heavy valence bands for GeTe [22,23,25]. In practice, the strategy that combines carrier concentration optimization with band alignment or resonant doping has been extensively used to optimize the electronic transport properties [34,35,38,39].

In addition to the optimization of electrical transport properties, doping or alloying can introduce significant lattice disorder, resulting in atomic mass and lattice strain fluctuations that endow a low κ_l by the reinforcement of phonon scattering [40]. For example, the room temperature κ_l of pristine GeTe ($2.7 \text{ W m}^{-1} \text{ K}^{-1}$) can be decreased within the range of $1.0\text{--}1.5 \text{ W m}^{-1} \text{ K}^{-1}$ by doping 10 at% Sb/Bi/Pb [41,42]. This value can even be further suppressed to $\sim 0.5 \text{ W m}^{-1} \text{ K}^{-1}$ in codoping systems, for instance, $\text{Ge}_{0.88-y}\text{Sb}_{0.08}\text{Pb}_y\text{Te}$ and $\text{Ge}_{0.85}\text{Bi}_{0.05}\text{Sb}_{0.10}\text{Te}$ alloys [43]. This prominent reduction in κ_l was mainly attributed to the atomic mass and lattice strain fluctuations induced by point defects. Yet, with the substitution of extrinsic atoms on Ge sites, the evolution of microstructures is complicated, often involving precipitates [27,31,44,45], stacking faults [42,46,47], inversion domain and twin [21,31,48–50], as well as herringbone domain structures [48,51]. These structural defects should also play an important role in strengthening the phonon scattering. Overall, doping/alloying is an imperative strategy for the enhancement of TE performance of GeTe through regulating the electron and phonon transport properties via energy band and defects engineering.

Here, we systematically investigated the TE performance of $\text{Ge}_{1-x}\text{Sb}_x\text{Te}$ ($x = 0, 0.05, 0.10, 0.15, \text{ and } 0.20$), which were prepared by the solid-state reaction and spark plasma sintering (SPS) route. We proved that a maximum zT of 1.88 at 773 K could be attained in $\text{Ge}_{0.90}\text{Sb}_{0.10}\text{Te}$ with good reproducibility and thermal stability. With Sb substitution on Ge sites, the Seebeck is greatly improved due to the reduction in hole concentration and valence band convergence. Moreover, a significantly reduced κ is achieved, which originates from both the decrease in κ_e and κ_l . Transmission electron microscopy (TEM)

observations verify the existence of nanoscale twins and dislocations in the matrix, which together with point defects (Sb substitutions and residual Ge vacancies) result in a low κ_l , i.e., $\sim 0.89 \text{ W m}^{-1} \text{ K}^{-1}$ in $\text{Ge}_{0.90}\text{Sb}_{0.10}\text{Te}$ at 773K. This work confirms $\text{Ge}_{0.90}\text{Sb}_{0.10}\text{Te}$ is a reliable candidate for TE application.

2. Materials and Methods

Polycrystalline samples $\text{Ge}_{1-x}\text{Sb}_x\text{Te}$ ($x = 0, 0.05, 0.10, 0.15, 0.20$) were prepared through a solid-state reaction followed by spark plasma sintering (SPS) procedure. High-purity ($>4 \text{ N}$) elements of Ge, Sb, and Te were weighted in the corresponding mole ratio and flame-sealed in quartz tubes under a vacuum lower than $5 \times 10^{-4} \text{ Pa}$. The tubes were heated to 1173 K and soaked for 6 h, both of heating and cooling rate at 1 K min^{-1} . These as-synthesized ingots were hand-ground into fine powders and then consolidated into cylindrical bulks by SPS at 823 K for 3 min under a uniaxial pressure of $\sim 50 \text{ MPa}$ in a vacuum ($<10 \text{ Pa}$). The relative density for these sintered samples is no less than 95%. Rectangular specimens ($\sim 2 \text{ mm} \times 2 \text{ mm} \times 10 \text{ mm}$) and square-shaped pellets ($\sim 6 \text{ mm} \times 6 \text{ mm} \times 1.5 \text{ mm}$) were cut off for electrical and thermal transport properties measurements, respectively.

The crystal structure of sintered samples was analyzed using a PANalytical X-ray diffractometer ($\text{Cu K}\alpha$, $\lambda = 0.154 \text{ nm}$) operated at 45 kV and 40 mA with a step size of 0.01313° . The high-angle annular dark field scanning transmission electron microscopy (HAADF-STEM) and the energy-dispersive X-ray (EDS) mapping were performed on a probe-corrected transmission electron microscopy (TEM; FEI Titan G2) equipped with a super EDS system at 300 kV. The backscattered electron image (BSE) and EDS mapping scanning were employed with electron microscopy (SEM, FEI Helios G4) equipped with a super EDS system.

Room temperature carrier concentration (p) and mobility (μ_H) were examined using a direct-current (dc) Hall effect measurement system (model 8404; Lakeshore Cryotronics, USA) at the excitation current up to 100 mA and magnetic field of 1.5 T. Electrical resistivity (ρ) and Seebeck coefficient (S) were measured by a static dc method (LSR-3, Germany) under a helium (99.999%) atmosphere. The thermal conductivity (κ) is calculated by $\kappa = D d C_p$, where D is the thermal diffusivity, C_p is the specific heat, and d is the volume density. The thermal diffusivity (D) was measured by a laser flash method (LFA 457; Netzsch, Germany); the specific heat (C_p) was estimated based on Dulong–Petit law, which is approximate to the measurement result; and the volume density (d) was determined by the Archimedes method.

3. Results and Discussion

Figure 1 shows the XRD patterns of the as-pulverized sintered $\text{Ge}_{1-x}\text{Sb}_x\text{Te}$ samples. The major diffraction peaks of each sample could be indexed to the rhombohedral structure of GeTe (JCPDS 047–1079). Particularly, the presence of double peaks in the 2θ range of $23\text{--}27^\circ$ and $41\text{--}45^\circ$ further confirms the rhombohedral phase of pristine GeTe. The peaks located at $\sim 27.3^\circ$ and $\sim 45.3^\circ$ correspond to cubic Ge (JCPDS 065–0333) secondary phase. The obviously enlarged Ge peaks with increasing the Sb content indicates more Ge precipitates in the matrix, which is probably due to that Sb atoms occupy on Ge sublattice sites. To examine the impact of Sb doping on the crystal structure of GeTe, we calculated lattice parameters of the sintered $\text{Ge}_{1-x}\text{Sb}_x\text{Te}$ samples and the results are shown in Figure 1b. We observed the increase in lattice parameter a and decrease in c with increasing the Sb content, suggesting a smaller c/a ratio. This phenomenon indicates that Sb doping could lower the structural transition temperature of GeTe and promote the formation of cubic phase, which is consistent with the convergence of (024) and (220) diffraction peaks (Figure 1a), and the observations in the Bi- and V-doped cases [32,47].

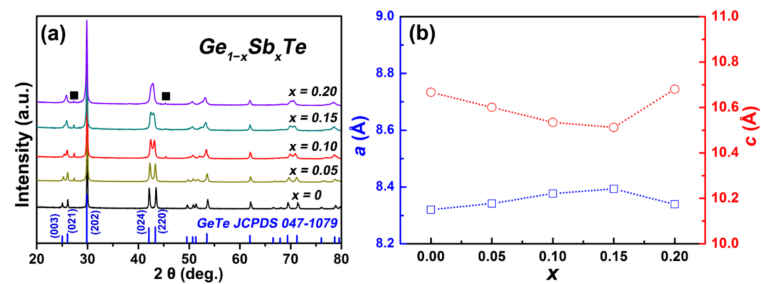


Figure 1. (a,b) Powder XRD patterns and lattice parameters a and c of $\text{Ge}_{1-x}\text{Sb}_x\text{Te}$. The black square (■) represents Ge precipitates.

We performed SEM and EDS to illustrate the microstructure evolution of GeTe upon Sb doping, as shown in Figure 2. Figure 2a shows a BSE image of the pristine GeTe, it can be seen that a small amount of nanoscale precipitates (dark contrast) is embedded into the host matrix (gray contrast). A BSE image of $\text{Ge}_{0.90}\text{Sb}_{0.10}\text{Te}$ (Figure 2b) shows that there are many micron-sized precipitates (dark contrast) in the matrix. The EDS mapping images for $\text{Ge}_{0.90}\text{Sb}_{0.10}\text{Te}$ shown in Figure 2d unveil that the precipitates are Ge particles, which agree with the XRD characterizations. Moreover, the electron channeling contrast image (ECCI) is taken on the region marked with the red dotted rectangle to characterize the microstructure of the matrix [52], as shown in Figure 2c, which suggests that the matrix is comprised by 2–10 micron-sized grains, and each grain includes several lamellar configurations of the size of 20–500 nm.

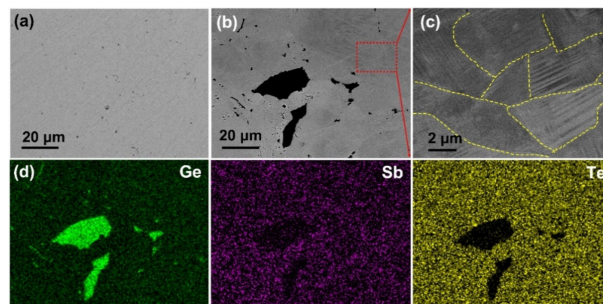


Figure 2. BSE images of the pristine GeTe (a) and $\text{Ge}_{0.90}\text{Sb}_{0.10}\text{Te}$ (b); (c) magnified ECC image of $\text{Ge}_{0.90}\text{Sb}_{0.10}\text{Te}$ taken on the region marked by the red dotted rectangle in Figure 2b; (d) elemental mapping images taken on the region that displayed in Figure 2b.

The electrical transport properties of the samples $\text{Ge}_{1-x}\text{Sb}_x\text{Te}$ are shown in Figure 3. From Figure 3a, we can see that the resistivity ρ of all the samples increases with rising temperature, showing a degenerate semiconductor behavior, coincides with its incipient metal nature of GeTe empowered by metavalent bonding [53]. Typically, for the pristine GeTe, the ρ is $\sim 1.26 \times 10^{-6} \Omega \text{ m}$ at 323 K, which increases to $\sim 5.11 \times 10^{-6} \Omega \text{ m}$ at 773 K. As Sb substitution on Ge sites could generate a donor type point defect and increase the formation energy of Ge vacancy, the ρ increases with increasing the Sb content. For example, the ρ at 323 K increases from $\sim 1.26 \times 10^{-6} \Omega \text{ m}$ for the pristine GeTe to $5.01 \times 10^{-5} \Omega \text{ m}$ for the sample $\text{Ge}_{0.80}\text{Sb}_{0.20}\text{Te}$. The abnormal change in $\rho \sim T$ curves shows a signature of structural phase transition. The transition temperature could be shifted to a lower temperature with increasing the Sb content, which is also observed in the Bi-doped GeTe system. This makes it possible to enhance the TE performance by controlling the distortion degree of the crystal structure from cubic to rhombohedral [39].

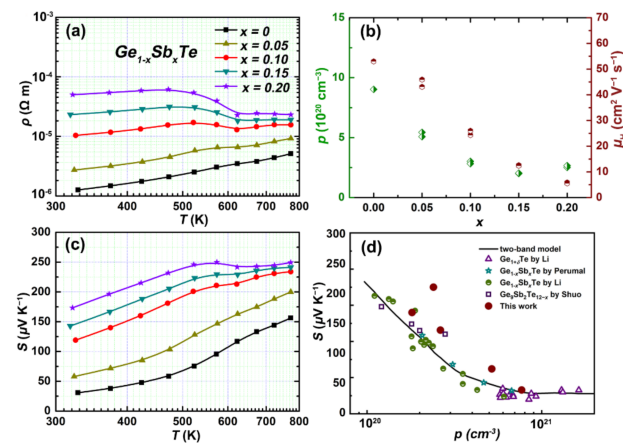


Figure 3. Electrical properties of $\text{Ge}_{1-x}\text{Sb}_x\text{Te}$: (a) electrical conductivity; (b) carrier concentration and mobility as a function of the Sb content x at room temperature; (c) Seebeck coefficient; (d) Pisarenko plot.

The Hall measurement was carried out on the samples $\text{Ge}_{1-x}\text{Sb}_x\text{Te}$, as shown in Figure 3b. It clearly shows that the carrier concentration and mobility are significantly suppressed upon Sb doping, mainly due to the generation of the charged substitutional point defects. For example, the room temperature carrier concentration gradually declines from $\sim 9.02 \times 10^{20}$ to $\sim 2.01 \times 10^{20} \text{ cm}^{-3}$ as the Sb content increases from $x = 0$ to $x = 0.15$, then rises to $\sim 2.55 \times 10^{20} \text{ cm}^{-3}$ when $x = 0.20$, indicating that the doping limit of Sb in GeTe is about 15 at%, and the doping efficiency is greatly reduced when the Sb content exceeds 10 at% [43,50]. The room-temperature carrier mobility monotonically drops from $53 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ for the pristine GeTe to $5.9 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ for the sample $\text{Ge}_{0.80}\text{Sb}_{0.20}\text{Te}$. Both the declined carrier concentration and mobility lead to the increase in ρ for the Sb-doped GeTe system.

Figure 3c shows the temperature-dependent Seebeck coefficient (S) of the samples $\text{Ge}_{1-x}\text{Sb}_x\text{Te}$. The positive S suggests that all the samples have a p-type conductive behavior and the inability of single Sb doping in inverting such a conductive behavior. As expected, the S is greatly improved upon Sb doping, increasing from $\sim 31 \mu\text{V K}^{-1}$ for the pristine GeTe to $\sim 119 \mu\text{V K}^{-1}$ for $\text{Ge}_{0.90}\text{Sb}_{0.10}\text{Te}$, and further to $\sim 173 \mu\text{V K}^{-1}$ for $\text{Ge}_{0.80}\text{Sb}_{0.20}\text{Te}$ at 323 K. A maximum S up to $249 \mu\text{V K}^{-1}$ at 574 K is attained in $\text{Ge}_{0.80}\text{Sb}_{0.20}\text{Te}$.

To find out the underlying reason for the enhanced S of GeTe upon Sb doping, the Pisarenko curve based on a two-valence-band model is plotted in Figure 3d, and the representative data from the literature are given for comparison. It can be seen that the data for the pristine GeTe fall exactly on the Pisarenko curve, while the data for the Sb-doped samples lie above the curve, indicating that the valence band convergence could be induced by Sb doping. According to the results of band structure calculations, the maximum of the heavy valence band is higher than that of the light one for the rhombohedral GeTe, which inverses for the cubic GeTe [31]. As mentioned above, Sb doping promotes the structural transition of GeTe from rhombohedral to cubic, leading to band convergence by decreasing the split energy between the heavy and light valence bands. Overall, the greatly enhanced S in the Sb-doped GeTe samples could be attributed to the valence band convergence and the hole concentration suppression, as the S is inversely coupled with the carrier concentration.

Figure S1 shows the temperature-dependent power factor (PF) of the samples $\text{Ge}_{1-x}\text{Sb}_x\text{Te}$. The maximum PF as high as $\sim 4.74 \times 10^{-3} \text{ W m}^{-1} \text{ K}^{-2}$ at 773 K is attained in the pristine GeTe. The high PF is maintained when the Sb content does not exceed 10 at% due to the greatly enhanced S , despite the evident increase in ρ . The maximum PF up to $\sim 3.23 \times 10^{-3} \text{ W m}^{-1} \text{ K}^{-2}$ achieved in $\text{Ge}_{0.90}\text{Sb}_{0.10}\text{Te}$ is comparable to V-, Y-, Bi-, and Sb-doped counterparts [30,32,33,47].

Figure 4a shows the temperature-dependent thermal conductivity (κ) of the samples $\text{Ge}_{1-x}\text{Sb}_x\text{Te}$. The κ is greatly reduced over the whole temperature range (323–773 K) upon

Sb doping, which decreases from $\sim 7.40 \text{ W m}^{-1} \text{ K}^{-1}$ for the pristine GeTe to $1.61 \text{ W m}^{-1} \text{ K}^{-1}$ for the sample $\text{Ge}_{0.90}\text{Sb}_{0.10}\text{Te}$ at 323 K. To understand the details regarding the great reduction of κ , the κ is generally divided into two parts of electronic thermal conductivity (κ_e) and lattice thermal conductivity (κ_l). The κ_e is calculated according to the Wiedemann–Franz law, $\kappa_e = L \rho^{-1} T$, where L is the Lorentz number estimated based on a two-valence-band model (Figure S2). As shown in Figure 4b, it is observed that the κ_e for GeTe could be remarkably suppressed upon Sb doping due to the significant increase in ρ . For example, the κ_e as high as $5.07 \text{ W m}^{-1} \text{ K}^{-1}$ at 323 K is attained in the pristine GeTe, which dramatically drops to $0.43 \text{ W m}^{-1} \text{ K}^{-1}$ for $\text{Ge}_{0.90}\text{Sb}_{0.10}\text{Te}$. The κ_l is obtained by subtracting κ_e from κ , i.e., $\kappa_l = \kappa - \kappa_e$. The temperature-dependent κ_l of the $\text{Ge}_{1-x}\text{Sb}_x\text{Te}$ is plotted in Figure 4c. It indicates that the κ_l of GeTe could be substantially suppressed upon Sb doping, e.g., the highest κ_l up to $\sim 2.33 \text{ W m}^{-1} \text{ K}^{-1}$ at 323 K is attained in the pristine GeTe, which decreases to $\sim 1.19 \text{ W m}^{-1} \text{ K}^{-1}$ for the sample $\text{Ge}_{0.90}\text{Sb}_{0.10}\text{Te}$.

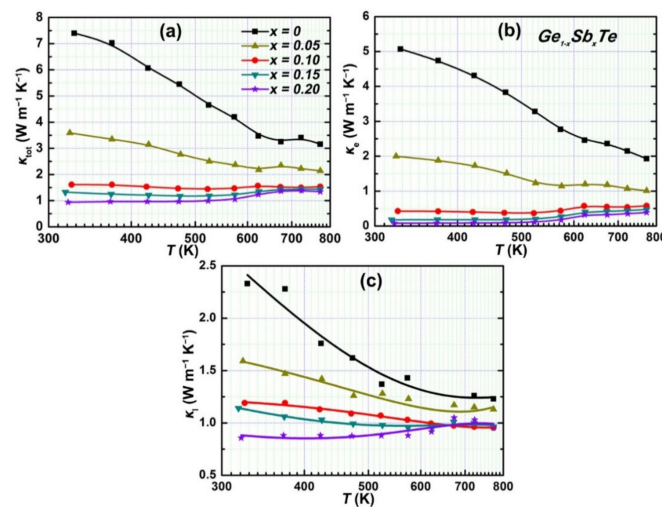


Figure 4. Temperature-dependent thermal properties of $\text{Ge}_{1-x}\text{Sb}_x\text{Te}$: (a) thermal conductivity; (b) electronic thermal conductivity; (c) lattice thermal conductivity.

To illustrate the low κ_l of $\text{Ge}_{0.90}\text{Sb}_{0.10}\text{Te}$, its microstructural details were investigated by TEM. A typical bright-field TEM image, as shown in Figure 5a, reveals the stacking lamellae in alternate bright and dark contrast, and there exists fluctuation in contrast at the edges of the lamellae (marked with the yellow arrows). However, the elemental Ge, Sb, and Te are uniformly distributed in the host matrix according to the EDS mapping results (Figure 5b), suggesting that the change in bright and dark contrast is unrelated to chemical compositions. A typical selected area electron diffraction (SAED) pattern of the lamellar structures can only be indexed to the $[1\bar{1}2]$ zone axis with rhombohedral symmetry (see the inset in Figure 5a). A HRTEM image is obtained in the same zone axis, which matches with the lattice fringes of the (111) planes of rhombohedral GeTe (Figure 5c). Moreover, the split diffraction spots far away from the transmitted beam are observed in the SAED pattern, together with the fast Fourier transform (FFT) image (the inset in Figure 5c), demonstrating that the lamellar structure is caused by the twin grains, which is similar to the herringbone structure observed in Bi-/Sb-doped GeTe samples [30,32]. It has been demonstrated that twin boundaries can effectively scatter phonons while maintaining a high charge carrier mobility [54]. Furthermore, the inverse FFT (IFFT) image discloses that there are dislocations at the twin boundaries, resulting from the release of strain. Additionally, geometric phase analysis (GPA) is executed to calculate the distribution of strain in the matrix based on the HRTEM image (Figure 5c), as shown in Figure 5d, confirming that an obvious strain fluctuation distributes around the dislocations. It is known that dislocations are very conducive to reducing the κ_l by scattering the mid-frequency phonons [52]. Overall, the low κ_l of $\text{Ge}_{0.90}\text{Sb}_{0.10}\text{Te}$ could

be attributed to its complex microstructures such as point defects (Sb substitutions and residual Ge vacancies), dislocations, and twin boundaries.

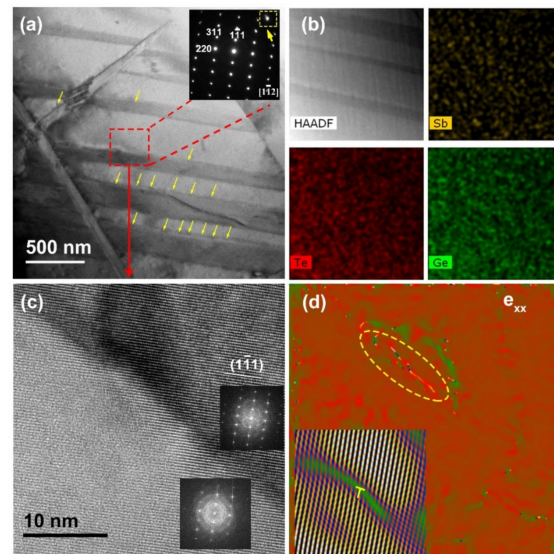


Figure 5. The microstructure characterization of $\text{Ge}_{0.90}\text{Sb}_{0.10}\text{Te}$: (a) a low-magnification bright-field image, the inset is the SAED pattern taken on the marked region; (b) a HAADF image and the corresponding elemental mappings; (c) a HRTEM image from the marked region in Figure 5a, the insets are the corresponding FFT images; (d) the distribution of strain calculated by GPA, the inset is the corresponding IFFT image.

The TE performance of GeTe is greatly enhanced upon Sb doping as a result of holes concentration optimization, band convergence, and the significant decrease in thermal conductivity. A highest zT up to 1.88 is achieved in the sample $\text{Ge}_{0.90}\text{Sb}_{0.10}\text{Te}$ (Figure 6a), which is higher than Bi-/In-/Ti-doped counterparts (single extrinsic elemental doping systems) [32,37,55]. It demonstrates that elemental Sb is an outstanding dopant for enhancing the TE performance of GeTe.

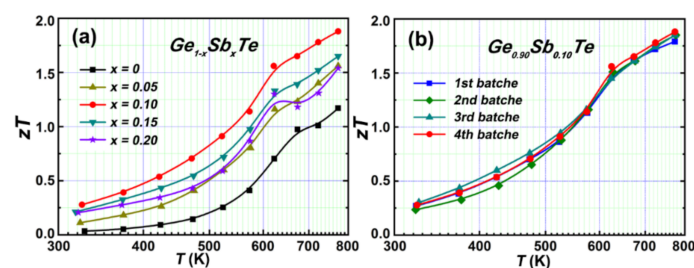


Figure 6. (a) zT value of $\text{Ge}_{1-x}\text{Sb}_x\text{Te}$ ($x = 0, 0.05, 0.10, 0.15, 0.20$); (b) zT value of the different batch samples $\text{Ge}_{0.90}\text{Sb}_{0.10}\text{Te}$.

More importantly, the high TE performance of $\text{Ge}_{0.90}\text{Sb}_{0.10}\text{Te}$ is reproducible, which is verified by the repeatable zT of the different batches samples, as shown in Figure 6b. Additionally, the thermal cycling test was performed on a sample $\text{Ge}_{0.90}\text{Sb}_{0.10}\text{Te}$, the electrical properties remain unchanged during the 2 times heating and cooling cycles (Figure S3), suggesting excellent stability of $\text{Ge}_{0.90}\text{Sb}_{0.10}\text{Te}$. These results indicate $\text{Ge}_{0.90}\text{Sb}_{0.10}\text{Te}$ has promising application in commercial TE devices.

4. Conclusions

The $\text{Ge}_{1-x}\text{Sb}_x\text{Te}$ samples were fabricated by the solid-state reaction followed by spark plasma sintering procedure. Sb doping not only leads to an optimized hole concentra-

tion approaching 10^{20} cm^{-3} but also results in valence band convergence for the GeTe system. Moreover, the thermal conductivity is greatly reduced due to the suppression of both electronic and lattice thermal conductivity. While the reduction in electronic thermal conductivity is due to the increase of electrical resistivity, the decline in lattice thermal conductivity is caused by the modification in microstructures, that is, the phonon scattering is remarkably strengthened by substitutional point defects, dislocations, and twin boundaries. As a result, a maximum zT up to 1.88 is achieved in $\text{Ge}_{0.9}\text{Sb}_{0.1}\text{Te}$. This work demonstrates that $\text{Ge}_{0.90}\text{Sb}_{0.10}\text{Te}$ is a promising candidate for TE commercial application.

Supplementary Materials: The following are available online at <https://www.mdpi.com/article/10.3390/ma15020406/s1>, this article at Figure S1: Temperature-dependent power factor of $\text{Ge}_{1-x}\text{Sb}_x\text{Te}$ ($x = 0, 0.05, 0.10, 0.15, 0.20$), Figure S2: Temperature-dependent Lorentz number of $\text{Ge}_{1-x}\text{Sb}_x\text{Te}$ ($x = 0, 0.05, 0.10, 0.15, 0.20$), and Figure S3: Temperature-dependent electrical properties of the $\text{Ge}_{0.90}\text{Sb}_{0.10}\text{Te}$ under thermal cycles: (a) electrical conductivity; (c) Seebeck coefficient.

Author Contributions: Conceptualization, C.L. (Chao Li); Data curation, H.S.; Supervision, H.Y., C.C. and L.M.; Writing—review & editing, Z.D., Z.Z. and C.L. (Chengyan Liu). All authors have read and agreed to the published version of the manuscript.

Funding: This work was supported by the National Key Research and Development Program of China (Grant No. 2017YFE0198000), National Natural Science Foundation of China (Grant No. 61974042, and 51772056), Guangdong Basic and Applied Basic Research Foundation (Grant No. 2020A1515110178), Independent research Foundation of Guangxi Key Laboratory of Information Materials (Grant No.201007-K), Natural Science Research Project of Colleges and Universities in Jiangsu Province (Grant No. 20KJB430039), Jiangsu Shuang Chuang program for Doctor (Grant No. 202031025).

Institutional Review Board Statement: Not applicable.

Informed Consent Statement: Not applicable.

Data Availability Statement: The data presented in this study are available on request from the corresponding author.

Conflicts of Interest: The authors declare no conflict of interest.

References

1. Snyder, G.J.; Toberer, E.S. Complex thermoelectric materials. *Nat. Mater.* **2008**, *7*, 105–114. [[CrossRef](#)] [[PubMed](#)]
2. Shi, X.L.; Zou, J.; Chen, Z.G. Advanced Thermoelectric Design: From Materials and Structures to Devices. *Chem. Rev.* **2020**, *120*, 7399–7515. [[CrossRef](#)]
3. Xiang, H.; Xing, Y.; Dai, F.Z.; Wang, H.; Su, L.; Miao, L.; Zhang, G.; Wang, Y.; Qi, X.; Yao, L.; et al. High-entropy ceramics: Present status, challenges, and a look forward. *J. Adv. Ceram.* **2021**, *10*, 385–441. [[CrossRef](#)]
4. Liu, Z.-Y.; Zhu, J.-L.; Tong, X.; Niu, S.; Zhao, W.-Y. A review of CoSb_3 -based skutterudite thermoelectric materials. *J. Adv. Ceram.* **2020**, *9*, 647–673. [[CrossRef](#)]
5. An, D.; Wang, J.; Zhang, J.; Zhai, X.; Kang, Z.; Fan, W.; Yan, J.; Liu, Y.; Lu, L.; Jia, C.-L.; et al. Retarding Ostwald ripening through Gibbs adsorption and interfacial complexions leads to high-performance SnTe thermoelectrics. *Energy Environ. Sci.* **2021**, *14*, 5469–5479. [[CrossRef](#)]
6. Smorodin, B.L.; Velarde, M.G. Convective instability of an Ohmic liquid layer in an unsteady thermal field. *Phys. Fluids* **2008**, *20*, 44101. [[CrossRef](#)]
7. Yu, Y.; Cagnoni, M.; Cojocar-Mirédin, O.; Wuttig, M. Chalcogenide Thermoelectrics Empowered by an Unconventional Bonding Mechanism. *Adv. Funct. Mater.* **2019**, *30*, 1904862. [[CrossRef](#)]
8. Xu, W.; Zhang, Z.; Liu, C.; Gao, J.; Ye, Z.; Chen, C.; Peng, Y.; Bai, X.; Miao, L. Substantial thermoelectric enhancement achieved by manipulating the band structure and dislocations in Ag and La co-doped SnTe. *J. Adv. Ceram.* **2021**, *10*, 860–870. [[CrossRef](#)]
9. Tang, J.; Yao, Z.; Chen, Z.; Lin, S.; Zhang, X.; Xiong, F.; Li, W.; Chen, Y.; Pei, Y. Maximization of transporting bands for high-performance SnTe alloy thermoelectrics. *Mater. Today Phys.* **2019**, *9*, 100091. [[CrossRef](#)]
10. Zhou, C.; Yu, Y.; Lee, Y.L.; Ge, B.; Lu, W.; Cojocar-Mirédin, O.; Im, J.; Cho, S.P.; Wuttig, M.; Shi, Z.; et al. Exceptionally High Average Power Factor and Thermoelectric Figure of Merit in n-type PbSe by the Dual Incorporation of Cu and Te. *J. Am. Chem. Soc.* **2020**, *142*, 15172–15186. [[CrossRef](#)]
11. Pei, Y.; Shi, X.; LaLonde, A.; Wang, H.; Chen, L.; Snyder, G.J. Convergence of electronic bands for high performance bulk thermoelectrics. *Nature* **2011**, *473*, 66–69. [[CrossRef](#)] [[PubMed](#)]

12. Kooi, B.J.; Wuttig, M. Chalcogenides by Design: Functionality through Metavalent Bonding and Confinement. *Adv. Mater.* **2020**, *32*, e1908302. [[CrossRef](#)]
13. Korkosz, R.J.; Chasapis, T.C.; Lo, S.H.; Doak, J.W.; Kim, Y.J.; Wu, C.I.; Hatzikraniotis, E.; Hogan, T.P.; Seidman, D.N.; Wolverton, C.; et al. High ZT in p-Type $(\text{PbTe})_{(1-2x)}(\text{PbSe})_x(\text{PbS})_x$ Thermoelectric Materials. *J. Am. Chem. Soc.* **2014**, *136*, 3225–3237. [[CrossRef](#)] [[PubMed](#)]
14. Pei, Y.; Tan, G.; Feng, D.; Zheng, L.; Tan, Q.; Xie, X.; Gong, S.; Chen, Y.; Li, J.-F.; He, J.; et al. Integrating Band Structure Engineering with All-Scale Hierarchical Structuring for High Thermoelectric Performance in PbTe System. *Adv. Energy Mater.* **2017**, *7*, 1601450. [[CrossRef](#)]
15. Sarkar, S.; Zhang, X.; Hao, S.; Hua, X.; Bailey, T.P.; Uher, C.; Wolverton, C.; Dravid, V.P.; Kanatzidis, M.G. Dual Alloying Strategy to Achieve a High Thermoelectric Figure of Merit and Lattice Hardening in p-Type Nanostructured PbTe. *ACS Energy Lett.* **2018**, *3*, 2593–2601. [[CrossRef](#)]
16. Xiao, Y.; Wu, H.; Cui, J.; Wang, D.; Fu, L.; Zhang, Y.; Chen, Y.; He, J.; Pennycook, S.J.; Zhao, L.-D. Realizing high performance n-type PbTe by synergistically optimizing effective mass and carrier mobility and suppressing bipolar thermal conductivity. *Energy Environ. Sci.* **2018**, *11*, 2486–2495. [[CrossRef](#)]
17. Zhou, C.; Lee, Y.K.; Yu, Y.; Byun, S.; Luo, Z.Z.; Lee, H.; Ge, B.; Lee, Y.L.; Chen, X.; Lee, J.Y.; et al. Polycrystalline SnSe with a thermoelectric figure of merit greater than the single crystal. *Nat. Mater.* **2021**, *20*, 1378–1384. [[CrossRef](#)]
18. Zhang, X.; Bu, Z.; Lin, S.; Chen, Z.; Li, W.; Pei, Y. GeTe Thermoelectrics. *Joule* **2020**, *4*, 986–1003. [[CrossRef](#)]
19. Hong, M.; Zou, J.; Chen, Z.G. Thermoelectric GeTe with Diverse Degrees of Freedom Having Secured Superhigh Performance. *Adv. Mater.* **2019**, *31*, e1807071. [[CrossRef](#)]
20. Bai, G.; Yu, Y.; Wu, X.; Li, J.; Xie, Y.; Hu, L.; Liu, F.; Wuttig, M.; Cojocaru-Mirédin, O.; Zhang, C. Boron Strengthened GeTe-Based Alloys for Robust Thermoelectric Devices with High Output Power Density. *Adv. Energy Mater.* **2021**, *11*, 2102012. [[CrossRef](#)]
21. Tran, X.Q.; Hong, M.; Maeno, H.; Kawami, Y.; Toriyama, T.; Jack, K.; Chen, Z.-G.; Zou, J.; Matsumura, S.; Dargusch, M.S. Real-time observation of the thermally-induced phase transformation in GeTe and its thermal expansion properties. *Acta Mater.* **2019**, *165*, 327–335. [[CrossRef](#)]
22. Hong, M.; Wang, Y.; Feng, T.; Sun, Q.; Xu, S.; Matsumura, S.; Pantelides, S.T.; Zou, J.; Chen, Z.G. Strong Phonon-Phonon Interactions Securing Extraordinary Thermoelectric $\text{Ge}_{1-x}\text{Sb}_x\text{Te}$ with Zn-Alloying-Induced Band Alignment. *J. Am. Chem. Soc.* **2019**, *141*, 1742–1748. [[CrossRef](#)]
23. Zheng, Z.; Su, X.; Deng, R.; Stoumpos, C.; Xie, H.; Liu, W.; Yan, Y.; Hao, S.; Uher, C.; Wolverton, C.; et al. Rhombohedral to Cubic Conversion of GeTe via MnTe Alloying Leads to Ultralow Thermal Conductivity, Electronic Band Convergence, and High Thermoelectric Performance. *J. Am. Chem. Soc.* **2018**, *140*, 2673–2686. [[CrossRef](#)]
24. Bu, Z.; Li, W.; Li, J.; Zhang, X.; Mao, J.; Chen, Y.; Pei, Y. Dilute Cu_2Te -alloying enables extraordinary performance of r-GeTe thermoelectrics. *Mater. Today Phys.* **2019**, *9*, 100096. [[CrossRef](#)]
25. Hong, M.; Wang, Y.; Liu, W.; Matsumura, S.; Wang, H.; Zou, J.; Chen, Z.-G. Arrays of Planar Vacancies in Superior Thermoelectric $\text{Ge}_{1-x-y}\text{Cd}_x\text{Bi}_y\text{Te}$ with Band Convergence. *Adv. Energy Mater.* **2018**, *8*, 1801837. [[CrossRef](#)]
26. Yue, L.; Fang, T.; Zheng, S.; Cui, W.; Wu, Y.; Chang, S.; Wang, L.; Bai, P.; Zhao, H. Cu/Sb Codoping for Tuning Carrier Concentration and Thermoelectric Performance of GeTe-Based Alloys with Ultralow Lattice Thermal Conductivity. *ACS Appl. Energy Mater.* **2019**, *2*, 2596–2603. [[CrossRef](#)]
27. Li, P.; Ding, T.; Li, J.; Zhang, C.; Dou, Y.; Li, Y.; Hu, L.; Liu, F.; Zhang, C. Positive Effect of Ge Vacancies on Facilitating Band Convergence and Suppressing Bipolar Transport in GeTe-Based Alloys for High Thermoelectric Performance. *Adv. Funct. Mater.* **2020**, *30*, 1910059. [[CrossRef](#)]
28. Zhang, X.; Li, J.; Wang, X.; Chen, Z.; Mao, J.; Chen, Y.; Pei, Y. Vacancy Manipulation for Thermoelectric Enhancements in GeTe Alloys. *J. Am. Chem. Soc.* **2018**, *140*, 15883–15888. [[CrossRef](#)]
29. Levin, E.M.; Besser, M.F.; Hanus, R. Electronic and thermal transport in GeTe: A versatile base for thermoelectric materials. *J. Appl. Phys.* **2013**, *114*, 083713. [[CrossRef](#)]
30. Perumal, S.; Roychowdhury, S.; Negi, D.S.; Datta, R.; Biswas, K. High Thermoelectric Performance and Enhanced Mechanical Stability of p-type $\text{Ge}_{1-x}\text{Sb}_x\text{Te}$. *Chem. Mater.* **2015**, *27*, 7171–7178. [[CrossRef](#)]
31. Wu, D.; Zhao, L.D.; Hao, S.; Jiang, Q.; Zheng, F.; Doak, J.W.; Wu, H.; Chi, H.; Gelbstein, Y.; Uher, C.; et al. Origin of the high performance in GeTe-based thermoelectric materials upon Bi_2Te_3 doping. *J. Am. Chem. Soc.* **2014**, *136*, 11412–11419. [[CrossRef](#)] [[PubMed](#)]
32. Perumal, S.; Roychowdhury, S.; Biswas, K. Reduction of thermal conductivity through nanostructuring enhances the thermoelectric figure of merit in $\text{Ge}_{1-x}\text{Bi}_x\text{Te}$. *Inorg. Chem. Front.* **2016**, *3*, 125–132. [[CrossRef](#)]
33. Gao, W.; Liu, Z.; Zhang, W.; Sato, N.; Guo, Q.; Mori, T. Improved thermoelectric performance of GeTe via efficient yttrium doping. *Appl. Phys. Lett.* **2021**, *118*, 33901. [[CrossRef](#)]
34. Perumal, S.; Samanta, M.; Ghosh, T.; Shenoy, U.S.; Bohra, A.K.; Bhattacharya, S.; Singh, A.; Waghmare, U.V.; Biswas, K. Realization of High Thermoelectric Figure of Merit in GeTe by Complementary Co-doping of Bi and In. *Joule* **2019**, *3*, 2565–2580. [[CrossRef](#)]
35. Srinivasan, B.; Gellé, A.; Gucci, F.; Boussard-Pledel, C.; Fontaine, B.; Gautier, R.; Halet, J.-F.; Reece, M.J.; Bureau, B. Realizing a stable high thermoelectric $zT \sim 2$ over a broad temperature range in $\text{Ge}_{1-x-y}\text{Ga}_x\text{Sb}_y\text{Te}$ via band engineering and hybrid flash-SPS processing. *Inorg. Chem. Front.* **2019**, *6*, 63–73. [[CrossRef](#)]

36. Li, M.; Hong, M.; Tang, X.; Sun, Q.; Lyu, W.-Y.; Xu, S.-D.; Kou, L.-Z.; Dargusch, M.; Zou, J.; Chen, Z.-G. Crystal symmetry induced structure and bonding manipulation boosting thermoelectric performance of GeTe. *Nano Energy* **2020**, *73*, 104740. [[CrossRef](#)]
37. Wu, L.; Li, X.; Wang, S.; Zhang, T.; Yang, J.; Zhang, W.; Chen, L.; Yang, J. Resonant level-induced high thermoelectric response in indium-doped GeTe. *NPG Asia Mater.* **2017**, *9*, e343. [[CrossRef](#)]
38. Shuai, J.; Sun, Y.; Tan, X.; Mori, T. Manipulating the Ge Vacancies and Ge Precipitates through Cr Doping for Realizing the High-Performance GeTe Thermoelectric Material. *Small* **2020**, *16*, e1906921. [[CrossRef](#)]
39. Li, J.; Zhang, X.; Chen, Z.; Lin, S.; Li, W.; Shen, J.; Witting, I.T.; Faghaninia, A.; Chen, Y.; Jain, A.; et al. Low-Symmetry Rhombohedral GeTe Thermoelectrics. *Joule* **2018**, *2*, 976–987. [[CrossRef](#)]
40. Carruthers, P. Scattering of Phonons by Elastic Strain Fields and the Thermal Resistance of Dislocations. *Phys. Rev.* **1959**, *114*, 995–1001. [[CrossRef](#)]
41. Li, J.; Zhang, X.; Lin, S.; Chen, Z.; Pei, Y. Realizing the High Thermoelectric Performance of GeTe by Sb-Doping and Se-Alloying. *Chem. Mater.* **2016**, *29*, 605–611. [[CrossRef](#)]
42. Xu, X.; Xie, L.; Lou, Q.; Wu, D.; He, J. Boosting the Thermoelectric Performance of Pseudo-Layered Sb₂Te₃(GeTe)_n via Vacancy Engineering. *Adv. Sci.* **2018**, *5*, 1801514. [[CrossRef](#)]
43. Li, J.; Zhang, X.; Wang, X.; Bu, Z.; Zheng, L.; Zhou, B.; Xiong, F.; Chen, Y.; Pei, Y. High-Performance GeTe Thermoelectrics in Both Rhombohedral and Cubic Phases. *J. Am. Chem. Soc.* **2018**, *140*, 16190–16197. [[CrossRef](#)] [[PubMed](#)]
44. Wu, D.; Feng, D.; Xu, X.; He, M.; Xu, J.; He, J. Realizing high figure of merit plateau in Ge Bi Te via enhanced Bi solution and Ge precipitation. *J. Alloy Compd.* **2019**, *805*, 831–839. [[CrossRef](#)]
45. Li, J.; Chen, Z.; Zhang, X.; Yu, H.; Wu, Z.; Xie, H.; Chen, Y.; Pei, Y. Simultaneous Optimization of Carrier Concentration and Alloy Scattering for Ultrahigh Performance GeTe Thermoelectrics. *Adv. Sci.* **2017**, *4*, 1700341. [[CrossRef](#)]
46. Chen, S.; Bai, H.; Li, J.; Pan, W.; Jiang, X.; Li, Z.; Chen, Z.; Yan, Y.; Su, X.; Wu, J.; et al. Vacancy-Based Defect Regulation for High Thermoelectric Performance in Ge₉Sb₂Te_{12-x} Compounds. *ACS Appl Mater. Interfaces* **2020**, *12*, 19664–19673. [[CrossRef](#)] [[PubMed](#)]
47. Sun, Q.; Li, M.; Shi, X.L.; Xu, S.D.; Liu, W.D.; Hong, M.; Lyu Wy Yin, Y.; Dargusch, M.; Zou, J.; Chen, Z.G. Versatile Vanadium Doping Induces High Thermoelectric Performance in GeTe via Band Alignment and Structural Modulation. *Adv. Energy Mater.* **2021**, *11*, 2100544. [[CrossRef](#)]
48. Lee, H.S.; Kim, B.-S.; Cho, C.-W.; Oh, M.-W.; Min, B.-K.; Park, S.-D.; Lee, H.-W. Herringbone structure in GeTe-based thermoelectric materials. *Acta Mater.* **2015**, *91*, 83–90. [[CrossRef](#)]
49. Bayikadi, K.S.; Sankar, R.; Wu, C.T.; Xia, C.; Chen, Y.; Chen, L.-C.; Chen, K.-H.; Chou, F.-C. Enhanced thermoelectric performance of GeTe through in situ microdomain and Ge-vacancy control. *J. Mater. Chem. A* **2019**, *7*, 15181–15189. [[CrossRef](#)]
50. Bayikadi, K.S.; Wu, C.T.; Chen, L.-C.; Chen, K.-H.; Chou, F.-C.; Sankar, R. Synergistic optimization of thermoelectric performance of Sb doped GeTe with a strained domain and domain boundaries. *J. Mater. Chem. A* **2020**, *8*, 5332–5341. [[CrossRef](#)]
51. Nevskii, S.; Sarychev, V.; Konovalov, S.; Granovskii, A.; Gromov, V. Formation Mechanism of Micro- and Nanocrystalline Surface Layers in Titanium and Aluminum Alloys in Electron Beam Irradiation. *Metals* **2020**, *10*, 1399. [[CrossRef](#)]
52. Abdellaoui, L.; Chen, Z.; Yu, Y.; Luo, T.; Hanus, R.; Schwarz, T.; Bueno Villoro, R.; Cojocar-Mirédin, O.; Snyder, G.J.; Raabe, D.; et al. Parallel Dislocation Networks and Cottrell Atmospheres Reduce Thermal Conductivity of PbTe Thermoelectrics. *Adv. Funct. Mater.* **2021**, *31*, 2101214. [[CrossRef](#)]
53. Zhu, M.; Cojocar-Mirédin, O.; Mio, A.M.; Keutgen, J.; Kupers, M.; Yu, Y.; Cho, J.Y.; Dronskowski, R.; Wuttig, M. Unique Bond Breaking in Crystalline Phase Change Materials and the Quest for Metavalent Bonding. *Adv. Mater.* **2018**, *30*, e1706735. [[CrossRef](#)]
54. Yu, Y.; He, D.-S.; Zhang, S.; Cojocar-Mirédin, O.; Schwarz, T.; Stoffers, A.; Wang, X.-Y.; Zheng, S.; Zhu, B.; Scheu, C.; et al. Simultaneous optimization of electrical and thermal transport properties of Bi_{0.5}Sb_{1.5}Te₃ thermoelectric alloy by twin boundary engineering. *Nano Energy* **2017**, *37*, 203–213. [[CrossRef](#)]
55. Shuai, J.; Tan, X.J.; Guo, Q.; Xu, J.T.; Gellé, A.; Gautier, R.; Halet, J.F.; Failamani, F.; Jiang, J.; Mori, T. Enhanced thermoelectric performance through crystal field engineering in transition metal-doped GeTe. *Mater. Today Phys.* **2019**, *9*, 100094. [[CrossRef](#)]