

# Alumina Graphene Catalytic Condenser for Programmable Solid Acids

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and selectivity to products. Here, an ultrathin catalytic film of amorphous alumina (4 nm) was integrated into a catalytic condenser device that enabled tunable electron depletion from the alumina active layer and correspondingly stronger Lewis acidity. The catalytic condenser had the following structure: amorphous alumina/graphene/HfO<sub>2</sub> dielectric (70 nm)/p-type Si. Application of positive voltages up to +3 V between graphene and the p-type Si resulted in electrons flowing out of the alumina; positive charge accumulated in the catalyst. Temperatureprogrammed surface reaction of thermocatalytic isopropanol (IPA) dehydration to propene on the charged alumina surface revealed a shift in the propene formation peak temperature of up to



 $\Delta T_{\text{peak}} \sim 50$  °C relative to the uncharged film, consistent with a 16 kJ mol<sup>-1</sup> (0.17 eV) reduction in the apparent activation energy. Electrical characterization of the thin amorphous alumina film by ultraviolet photoelectron spectroscopy and scanning tunneling microscopy indicates that the film is a defective semiconductor with an appreciable density of in-gap electronic states. Density functional theory calculations of IPA binding on the pentacoordinate aluminum active sites indicate significant binding energy changes ( $\Delta BE$ ) up to 60 kJ mol<sup>-1</sup> (0.62 eV) for 0.125 e<sup>-</sup> depletion per active site, supporting the experimental findings. Overall, the results indicate that continuous and fast electronic control of thermocatalysis can be achieved with the catalytic condenser device. **KEYWORDS:** catalysis, dehydration, alumina, catalytic condenser, graphene

# INTRODUCTION

Manipulation of charge within catalytic materials permits tunable surface chemistry and electronic optimization of chemical reaction rates.<sup>1</sup> Here, we put forward the concept of "catalytic condensers" that allow for reversible tuning of charge density in a catalyst thin film by means of an applied voltage. The catalyst film, in this case, an amorphous 4 nm thick layer of alumina, is deposited on top of a graphene/ insulator/conductor stack, and application of a voltage between graphene and the bottom electrode induces charge in the catalyst, thereby tuning the surface catalytic chemistry. The extent of hole/electron density accumulated in the catalyst layer depends on the applied potential, composition, and electronic characteristics of the dielectric and catalyst films, as well as quality of film deposition. We propose that this will be a generally applicable strategy for manipulating the reactivity of catalyst thin films that can be tuned for different heterogeneous chemistries under static or dynamic modulation.

Precedents for this work include the traditional metal-oxidesemiconductor field effect transistor that alters charge carrier density in semiconductor layers with high-*k* dielectric oxide layers; gate potential variation manipulates the electronic bands and the charge density of the semiconductor to achieve high and low conductivity.<sup>2</sup> This device structure has been used recently to manipulate the conduction band edge position and electron accumulation in ultrathin ZnO layers in electrochemical devices. In that case, a 5 nm ZnO film on a 300 nm SiO<sub>2</sub> dielectric layer allowed for electron accumulation throughout the entire ZnO layer, such that the exposed top surface of ZnO, which was in contact with the electrolyte, exhibited tunable rates of electron transfer with a soluble redox agent, tetrabromo-1,4-benzoquinone.<sup>3</sup> This transistor method-

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ology was also applied to two-dimensional (2D) semiconductor crystals such as single layer (SL)  $MoS_2$  incorporated in a  $MoS_2/SiO_2/p$ -Si stack, wherein charge modulation of active sites (S vacancies) was shown to increase the activity of electrocatalytic hydrogen evolution by a factor of 4.<sup>4-6</sup> However, to our knowledge, the use of capacitive charging to enhance thermocatalytic activity has not been reported.

Here, we consider a catalytic condenser comprised of an ultrathin film of alumina, a metal oxide catalyst commonly used to carry out acid-catalyzed chemical transformations, such as alcohol dehydration<sup>7-9</sup> and as a support for metal catalysts.<sup>10,11</sup> As depicted in Scheme 1, the amorphous

Scheme 1. Charge Condensation in the Catalyst Alumina/ Graphene Active Surface Layer of a Thin Film Catalytic Condenser



alumina film (4 nm) supported on graphene is separated from a conductive p-type silicon wafer by 70 nm of  $HfO_2$  dielectric. When a voltage  $V_{CAT}$  is applied, charge accumulated in the bottom silicon electrode is balanced by a countercharge in the thin catalytic alumina/graphene film (as in a typical capacitor or condenser). Because of the extreme thinness of the alumina layer, its accumulated charge is accessible to reactive fluids (adsorbates in plasmas, vapors, gases, and liquids) on the top surface. The tunable activity of the alumina catalyst derives from the accumulated charge, which is controlled by  $V_{CAT}$ .

Alumina as the active layer within the catalytic condenser is desirable because it is a commonly used solid Lewis acid catalyst.<sup>12,13</sup> However, *bulk* alumina is an electronic insulator with a band gap of 8.7 eV,<sup>14–16</sup> with valence and conduction bands far from the Fermi level. Thus, it is not obvious that charge can be accumulated in an alumina film in a condenser stack. However, the insulating nature of alumina is a function of its structure. For example, the (100) and (110) surfaces of crystalline  $\gamma$ -Al<sub>2</sub>O<sub>3</sub> are predicted to have valence bands near the Fermi level and band gaps of only ~3 eV.<sup>17</sup> Similarly, the band gap for ultrathin (7–10 nm) films of *amorphous* alumina has been measured to be 2.5 eV, which is in the semiconductor range.<sup>15</sup> We will demonstrate below that the alumina we deposit can be viewed as a defective semiconductor, and so, charge accumulation is possible.

In this work, the device shown in Scheme 1 was developed and evaluated experimentally and theoretically to understand the degree of electronic control that could be achieved for a model catalytic reaction, isopropanol (IPA) dehydration. On alumina surfaces, IPA is thought to adsorb and react on pentacoordinate Lewis acid sites  $(Al_V^{3+})$  with activity that correlates with acid strength, dictated by the partial positive charge  $(\delta^+)$ on the surface aluminum and partial negative charge on the bonded oxygen  $(\delta^{-})$ .<sup>18</sup> IPA dehydrates on alumina predominately unimolecularly to propene and water with minimal formation of either acetone via dehydrogenation or diisopropyl ether via bimolecular dehydration. Here, electronic modulation of the active site chemistry was interrogated by experimental and theoretical calculations, evaluating the charge accumulated in the catalytic layer and its effect on IPA dehydration kinetics as a function of device design and applied potential bias. Application of potential bias of +3 V is shown here to dramatically enhance the surface acidity of the catalytic layer and lower the reaction temperature. These results are presented following a thorough description of the device design and characterization.

## RESULTS AND DISCUSSION

The alumina-graphene catalytic condenser device depicted in Scheme 1 and Figure 1a was fabricated with full synthetic details available in the methods section and the Supporting Information (Figures S1 and S2). On top of the p-type Si substrate, a  $\sim$ 70 nm HfO<sub>2</sub> dielectric layer was grown by atomic layer deposition (ALD) at 100 °C using a flow-type ALD reactor. The growth rate of the HfO2 film, as measured by spectroscopic ellipsometry (Figure S3), was consistent with literature values.<sup>19,20</sup> Next, graphene was transferred onto the  $HfO_2$  layer using a previously reported polymethyl meth-acrylate (PMMA)-based method, <sup>21,22</sup> leaving a ca. 1.0 cm<sup>2</sup> conductive sheet across the device. To make good electrical contact to graphene, a small area (0.01 cm<sup>2</sup>) contact (Au/Ti, 45 nm/5 nm) was deposited on the graphene by e-beam evaporation through a shadow mask. Finally, the  $\sim$ 4 nm catalytic alumina layer was grown via ALD at 100 °C using the same flow-type ALD reactor as mentioned above.

Cross-sectional transmission electron microscopy (TEM) and high-angle annular dark-field (HAADF) scanning TEM (STEM) images, combined with compositional analysis using energy-dispersive X-ray (EDX) spectroscopy, confirmed the layered structure of the catalytic condenser device, as shown in Figures 1b-d and S4. From the cross section, as expected, a multilayer (ML) of graphene was observed between the  $HfO_2$ and the alumina layer, which was 0.4 nm thick. The aluminum oxide layer appears to be predominantly amorphous, as anticipated for films made by ALD at low deposition temperatures.<sup>19,23,24</sup> There were some regions, as identified in Figure S5, where surface HfO<sub>2</sub> formed crystalline grains  $\sim 10$ nm in diameter at the interface, possibly during the device fabrication process. Scanning electron microscopy (SEM) analysis in Figures 1c and Figure S6 confirms that the alumina film above the graphene is continuous across the device. The successful growth of the aluminum oxide layer on the graphene could be attributed to the use of H<sub>2</sub>O in ALD, which created OH sites on the graphene surface at low deposition temperatures.<sup>25-27</sup> The presence of a continuous alumina layer over large lateral dimensions (~0.1  $\mu$ m) with thicknesses of  $\sim$ 4 nm was further confirmed by atomic force microscopy (AFM). It should be noted that at large lateral scales of order 1 cm, graphene is expected to have wrinkles and folds, which were observed in our SEM-EDX analysis (Figures 1c and S6), resulting from its chemical vapor deposition (CVD) growth on the initial Cu substrate<sup>28</sup> and its transfer. Hexagonal domains,



**Figure 1.** Alumina catalytic condenser design and composition. (a) Schematic illustration of the device showing a 70 nm thick HfO<sub>2</sub> layer sandwiched between the p-type silicon substrate and a thin (4 nm) graphene/alumina bilayer. (b) Conventional bright-field TEM image of the device cross section showing all the functional layers in the stack: HfO<sub>2</sub>, graphene, and am-Al<sub>2</sub>O<sub>3</sub>. The scale bar is 5 nm. The red dashed line shows the boundary between the am-Al<sub>2</sub>O<sub>3</sub> and deposited protective am-C layer. The yellow arrows and dashed line highlight the location of the graphene layer. (c) Top-view SEM image of the device. The cracks in the graphene are visible as dark contrast. Yellow arrows show the hexagon-shaped second layers of graphene. The scale bar is 25  $\mu$ m. (d) HAADF-STEM image and the complementary EDX elemental maps from layer stack: Hf (blue), Al (green), and C (red). The C signal from graphene can be seen under the am-Al<sub>2</sub>O<sub>3</sub> layer. The C signal above am-Al<sub>2</sub>O<sub>3</sub> is from the deposited am-C protection layer. The scale bar is 5 nm. (e) AFM image and (f) height profile of the condenser device. The scale bar is 1  $\mu$ m. (g) Raman spectra of graphene collected at different sites on the device. Spot 1 is a representative spectrum of ML graphene, whereas Spot 2 is a representative spectrum of SL graphene.

also visible on these SEM images (Figures 1c and S6), represent additional graphene layers, indicating that these transferred graphenes are composed of regions of monolayers and bilayers/trilayers. Raman spectra measured over 30 spot sites on the device (Figure 1g) validated the existence of monolayer graphene and bilayers/trilayers based on the intensity ratio of the G (1587 cm<sup>-1</sup>) and 2D (2658 cm<sup>-1</sup>) band and showed minimal defects due to the weak D band intensity (1350 cm<sup>-1</sup>). Spot 1 is a representation of bilayer graphene with the ratio of intensity  $I_{2D}/I_G$  of ~1, whereas spot 2 represents a monolayer of graphene with the  $I_{2D}/I_G$  of ~2.

A high quality  $HfO_2$  dielectric layer is important for voltagetunable thermocatalysis to ensure that charges are distributed uniformly along the graphene–catalyst layer. A particular consideration is a possible change in the  $HfO_2$  film structure during programmed temperature ramps, which were used in our reaction analysis described later. Consequently, the thermal stability of the  $HfO_2$  film in the catalytic capacitor device was assessed by measuring the X-ray diffraction (XRD) pattern of the  $HfO_2$ -Si structure as a function of temperature, as shown in Figure 2a. These  $HfO_2$ -Si samples were heated in a muffle furnace under ambient conditions for 1 h at different temperatures prior to XRD data collection. The 70 nm  $HfO_2$ XRD-amorphous ALD films were heated from 100 to 350 °C; the films started to crystallize *only* after heating above 400 °C. The detected peaks were associated with the monoclinic polymorphs of HfO<sub>2</sub>, with the crystallites aligning preferentially along certain Miller indices. This is expected because selection of substrate composition, substrate orientation, deposition temperature, and choice of precursors are known to affect the crystallization process and structural orientation of ALD films.<sup>29</sup> Crystallite formation above 400 °C is known to impact the HfO<sub>2</sub> structural integrity, increase surface roughness, introduce pinholes, and affect the electronic properties.<sup>30–34</sup> So, practically speaking, 400 °C is an upper bound for examining catalytic behavior in these particular devices.

To further establish a working temperature range for the catalytic condenser, the dielectric constant and capacitance of HfO<sub>2</sub> were determined as a function of processing temperature (Figure 2b). For the electrical measurements, a metal-HfO<sub>2</sub>-Si capacitor was fabricated by depositing a 0.01 cm<sup>2</sup> Au/Ti metal contact (45 nm/5 nm) via shadow-masked e-beam evaporation on HfO2-Si after the samples were heat-treated at selected temperatures. At low treatment temperatures (100–300  $^\circ$ C), a specific capacitance of  $\sim$ 330 nF cm<sup>-2</sup> was calculated from experimental displacement current-voltage curves (Figures 2b, S8 and S9) at varying sweep rates (0.25-1.25 V/s); this specific capacitance decreased to ~200 nF cm<sup>-2</sup> for treatment at 400 °C. This corresponds to dielectric constants of ~26 and ~16 (Figure 2b) as the HfO<sub>2</sub> layer converted from amorphous to crystalline states, in agreement with the range of reported dielectric constants of bulk HfO2. 23,24,35 It must be noted that



**Figure 2.** Characterization of the alumina–graphene catalytic condenser device. (a) XRD patterns of ALD-HfO<sub>2</sub> on Si after being heated in air. Peaks associated with crystallized monoclinic HfO<sub>2</sub> are shown as the reference. (b) Average dielectric constant of HfO<sub>2</sub> as a function of annealing pretreatment temperature in air. The capacitive current density for each point can be found in Figure S8. A change in capacitive current density occurs at ~350 °C before the ALD HfO<sub>2</sub> film begins to crystallize. (c) Capacitive displacement current density of graphene/HfO<sub>2</sub>/Si as a function of voltage bias at different voltage sweep rates. (d) Capacitive current density of alumina/graphene/HfO<sub>2</sub>/Si as a function of voltage and sweep rates. (e) Average number of induced holes on graphene (blue) and alumina–graphene (red) as a function of  $V_{CAT}$ . Error bars are within the size of the data point. (f) Averaged capacitance of the catalytic condenser with varying active layers above HfO<sub>2</sub>: Au\* (Au/Ti contact, 45 nm/5 nm), graphene (blue), and alumina on graphene (Al<sub>2</sub>O<sub>3</sub>/G, red). (g) Number of holes on the catalytic condenser at +3 V with varying active layers above the HfO<sub>2</sub> layer; the estimated catalytic active site density of alumina is provided for comparison (purple, sites/cm<sup>2</sup>). Au\*/HfO<sub>2</sub> denotes an Au/Ti contact (45/5 nm) on HfO<sub>2</sub>.

there was significant current leakage *only* after heating above 400  $^{\circ}$ C, again confirming that 400  $^{\circ}$ C is an upper limit for reaction studies.

To assess charge accumulation in our catalytic condensers, we performed additional displacement current measurements on two different condenser stacks, one without aluminum oxide (graphene/HfO<sub>2</sub>/p-Si) and the other being the complete catalytic condenser (alumina/graphene/HfO<sub>2</sub>/p-Si). Displacement currents as a function of voltage and sweep rate for both stacks are shown in Figure 2c,d, respectively, and analysis of these curves leads to the electron density versus voltage plot shown in Figure 2e. As expected, the stored electron density in the devices increases linearly with voltage,  $V_{CAT}$ . Note that for the alumina/graphene device shown in red that the maximum electron density is nearly  $5.3 \times 10^{12} \text{ e}^{-}/\text{cm}^{2}$ . Note also that the slope observed in Figure 2e is the specific capacitance. It is evident that the device with the alumina/graphene top contact has an eightfold larger capacitance than the device with just

graphene as the top contact layer (289 vs 36 nF cm<sup>-2</sup>). In other words, the charge accumulated in the condenser is substantially larger when alumina is present, which is not unexpected as the electronic density of states in the ultrathin graphene is likely insufficient to fully screen the applied electric field.<sup>36</sup> The alumina layer provides additional density of states and thus larger charge accumulation.

A direct comparison of capacitances and maximum charge density at  $V_{CAT}$  of +3 V for the three types of stack devices discussed above is shown in Figure 2f,g. It is noteworthy that the condensers with Au versus alumina/graphene top contacts are comparable in capacitance and charge accumulation. Because the capacitances of the graphene-only versus the alumina/graphene devices are so different, we infer that the majority of the induced charge in the alumina/graphene device is actually in the alumina and therefore available to impact surface thermochemistry.



**Figure 3.** Electronic Characterization of the Alumina/Graphene Layer by UPS and STM. (a) UPS with energy cutoff for work function and valence band onset (2.3 eV) highlighted. (b) Magnified view of the UPS data near the Fermi level shown in the inset. The slope change near 0.25 eV indicates the presence of in-gap states above the valence band edge. (c) STM constant-current distance–voltage curves for the alumina/graphene film with valence band edge ( $E_V = -2.3 \text{ eV}$ ), conduction band edge ( $E_C = 2.2 \text{ eV}$ ), and work function (WF = 4.4 eV) identified. (d) Differential conductance near  $E_F$  revealing a discrete gap state (~300 meV). (e) AFM and (f) STM image (tip bias = -710 mV, I = 622 pA) of the alumina surface showing similar morphologies.

# **Alumina Film Electronic Characteristics**

The electronic states of the top layer of the alumina–graphene catalytic condenser were characterized by ultraviolet photoelectron spectroscopy (UPS) and scanning tunneling microscopy (STM), as depicted in Figure 3. Based on the UPS spectra, the ALD-deposited alumina thin film is electronically very different from insulating bulk  $Al_2O_3$ . The UPS spectrum in Figure 3a indicates the valence band edge ( $E_V$ ) is at 2.3 eV below the Fermi level. However, closer inspection between -0.5 and +0.8 eV in Figure 3b indicates the presence of in-gap states near the Fermi level at ~250–350 meV. These gap states are present on the surface of the alumina even for devices without graphene, as shown in Figure S7.

The electronic states of the alumina-graphene catalytic condenser active layer were further probed by STM in Figure 3c-f. While bulk crystalline alumina is an insulator, the  $\sim$ 4 nm alumina film in the condenser was sufficiently conductive for mapping via STM (Figure 3f); the resulting image was consistent with the morphology of the surface observed by AFM with the same resolution (Figure 3e). The electronic characteristics of the surface at large biases were evaluated as constant-current distance-voltage curves (Figure 3c), identifying the valence band edge ( $E_V$  of -2.3 eV), the conduction band edge ( $E_{\rm C}$  of 2.2 eV), and the work function (WF of 4.4 eV). It is clear from Figure 3c that the amorphous alumina film has a band gap of ~5 eV, much smaller than crystalline  $\gamma$ -Al<sub>2</sub>O<sub>3</sub>, which has a band gap over 8 eV.<sup>14-16</sup> Furthermore, STM differential conductance measurements under varying bias (Figure 3d) also indicated a gap state about 300 meV

below the Fermi level, in general accordance with the UPS spectrum of Figure 3b.

The combined spectroscopy data of Figure 3 point to the existence of in-gap electronic states in the alumina film that can be occupied by (or depleted of) electrons. These in-gap electronic states exist with sufficient density to permit STM imaging via tunneling current between the STM tip and the aluminum oxide film. The ALD-deposited alumina film acts much like a "poor insulator" or defective semiconductor, allowing charge to pass through the layer to the conductive graphene layer below. Graphene can move charge laterally across the surface of the device, and the aluminum oxide in turn has sufficient defect state densities near the Fermi level to conduct in the normal direction.

# Catalytic Condenser Alumina Surface Chemistry

The unique chemical properties of the catalytic condenser's alumina film were then characterized by thermocatalytic dehydration of IPA. Catalytic properties of the alumina–graphene condenser were probed through temperature-programmed surface reaction (TPSR) under varying voltage bias; details of the TPSR experimental design and method are provided in the Supporting Information (Figures S11–S17). Briefly, the catalytic condenser was housed within a chamber capable of evacuation to ultrahigh vacuum (UHV, ~10<sup>-9</sup> Torr, Figure S13), where IPA was dosed into the chamber to adsorb on the alumina film; 24 h of pumping removed physisorbed IPA from the chamber, leaving chemisorbed IPA on the aluminum oxide active sites. As the alumina/graphene condenser surface temperature was increased, aluminum



**Figure 4.** TPSR of IPA dehydration. (a) Temperature-programmed IPA dehydration on the alumina catalytic condenser device (4 nm alumina/ graphene/70 nm HfO<sub>2</sub>/p-type Si) to form propene at varying applied voltages,  $V_{CAT}$ . (b) Temperature-programmed IPA dehydration to propene on the catalytic condenser at  $V_{CAT}$  of +3 V and the powders of  $\gamma$ -alumina, titania, and zirconia with activation energy,  $E_a$  (kJ mol<sup>-1</sup>), to form propene determined by the Redhead equation. (c) Activation energy,  $E_a$ , of propene formation from IPA on the alumina catalytic condenser as determined by the Redhead equation. Error bars represent a 90% confidence interval.

oxide catalyzed the dehydration of chemisorbed IPA, which was followed by mass spectrometry. The propene signal was tracked at a mass-to-charge (m/z) ratio of 41, correcting for the influence of simultaneously desorbing unreacted IPA (m/z) = 45, Figure S19) and subtraction of a sigmoidal baseline derived from the chamber pressure; full details are provided in the Supporting Information. Catalyst surface temperature was measured using a type-K thermocouple placed near the location of the catalytic condenser within the UHV chamber (Figures S16 and S17).

Under UHV-TPSR conditions, unimolecular IPA dehydration to propene over the alumina catalytic condenser is the dominant pathway, with negligible formation of acetone and diisopropyl ether through dehydrogenation and bimolecular dehydration, respectively. The dominance of unimolecular dehydration was consistent with ambient pressure TPSR over bulk alumina powders, where 99% selectivity to propene as the product of unimolecular dehydration was observed (Table S1). Additionally, IPA dehydration kinetics on the 1 cm<sup>2</sup> alumina catalytic condenser, in the absence of applied potential bias, was found to be in agreement with  $\gamma$ -alumina powder catalysts. Without applied potential bias, the rate of propene formation on the alumina condenser peaked at ~130 °C ( $T_{\text{peak}}$ , Figures 4a and S18). Applying a first-order Polanyi-Wigner equation to the zero-bias experimental TPSR profile,<sup>37</sup> an activation energy ( $E_a$ ) of 115 ± 1 kJ mol<sup>-1</sup> (1.19 eV) for IPA unimolecular dehydration over the alumina condenser was measured in the absence of a potential bias ( $V_{CAT} = 0$  V, Figure 4c). Similarly, we measured an activation energy of  $117 \pm 1$  kJ mol<sup>-1</sup> (1.21 eV) for unimolecular IPA dehydration over bulk  $\gamma$ alumina powder via ambient pressure TPSR (Figure 4b), consistent with prior measurements by Gorte and coworkers.<sup>13</sup>

Condensation of holes (positive charge) for catalytic modulation of alumina was then evaluated by IPA dehydration TPSR under  $V_{CAT}$  biases of 0, +1, +2, and +3 V. Our central result is that application of positive voltage bias  $(V_{CAT} > 0 \text{ V})$ significantly reduced the peak temperature of the rate of propene formation (Figure 4a). While propene formation peaked at ~130 °C under zero bias ( $V_{CAT}$  of 0 V), propene formation peak temperatures reduced by as much  $\sim 50$   $^\circ C$  $(\Delta T_{\text{peak}})$  at  $V_{\text{CAT}}$  of +3 V; the process was reversible when the voltage was reverted to  $V_{CAT}$  of 0 V. The change in peak temperature is commensurate with a change in the kinetics of IPA unimolecular dehydration; positive potential bias reduced the activation energy to  $109 \pm 2$ ,  $104 \pm 3$ , and  $99 \pm 4$  kJ mol<sup>-1</sup> for  $V_{\text{CAT}}$  of +1, +2, and +3 V, respectively, constituting an overall reduction in the activation energy ( $\Delta E_a$ ) from  $V_{CAT}$  of 0 V of ~16 kJ mol<sup>-1</sup> (0.16 eV, Figure 4c). It is worth noting that in the absence of an aluminum oxide layer, no propene formation was observed in the course of a TPSR (Figures S20 and S21).

One potential explanation for the experimentally observed higher activity for unimolecular IPA dehydration over the graphene-alumina catalytic condenser is ohmic heating of the active surface. Charge separation across the 70 nm HfO<sub>2</sub> insulating layer introduces the possibility for "leakage current" through the dielectric, potentially leading to resistive heating of the catalytic layer. We have demonstrated earlier that leakage current is negligible when HfO<sub>2</sub> is maintained below 400 °C. The power supply (Keithley 2450) applying the potential bias to the active device surface during the experiment was limited to 105  $\mu$ A to eliminate thermal artifacts due to ohmic heating. The possibility of ohmic heating of the catalytic condenser was also experimentally evaluated by measuring the surface temperature of the catalytic condenser with an infrared camera (FLIR A655) under varying potential bias (Figure S22); pubs.acs.org/jacsau



**Figure 5.** DFT calculations of IPA adsorption and dehydration on the  $Al_V^{3^+}$  site of  $\gamma$ -alumina (100) as a function of charge. (a) 0.125 holes (h<sup>+</sup>) per site with the  $Al_V^{+3}$ - $O^{IPA}$  bond distance of 1.963 Å, (b) neutral with the  $Al_V^{+3}$ - $O^{IPA}$  bond distance of 2.016 Å, (c) IPA binding energy versus hole density, and (d) associated activation energy to form propene and water from IPA. Color scheme: O—red, C—black, Al—grey, and H—white.

surface temperature did not vary more than one degree over 10 min for all considered applied potentials ( $0 \le V_{\text{CAT}} \le +3$  V), thereby ruling out potential thermal artifacts.

Alternatively, the enhanced rate of IPA unimolecular dehydration under positive bias can be rationalized based on increased electron-hole density in the thin alumina-graphene layer, as indicated by capacitance measurements (Figure 2) and spectroscopy of the alumina film (Figure 3). Reduced electron density at the Lewis acid catalytic site of the aluminum oxide surface  $(Al_V^{3+})$  increases site acidity, binding IPA stronger and reducing the catalytic activation barrier for its unimolecular dehydration to propene and water. The extent of positive charge condensation in the graphene-alumina active layer depends on the extent of applied potential bias as well as the design of the catalytic condenser stack. As summarized in Figure 2g, the number of electrons depleted in the aluminum oxide film at  $V_{\rm CAT}$  of +3 V was measured to be 5.3  $\times$  10<sup>12</sup> e<sup>-</sup> cm<sup>-2</sup>; this is within an order of magnitude of pentacoordinate aluminum active site density available for acid catalysis reported on  $\gamma$ -Al<sub>2</sub>O<sub>3</sub> (10<sup>13</sup>-10<sup>14</sup> sites cm<sup>-2</sup>).<sup>13,38-40</sup> It is therefore possible to deplete 0.05-0.50 electrons per active site at a potential bias  $V_{CAT}$  of +3 V. It should be noted that the amorphous 4 nm film atop the condenser may have an active site density distinct from  $\gamma$ -Al<sub>2</sub>O<sub>3</sub>, but it is unlikely to vary by more than an order of magnitude. The effect of this charge depletion at the active site can be further evaluated by computation, as described next.

## IPA Binding and Dehydration Computed Energy

To simulate direct electron transfer to the alumina layer, we carried out periodic density functional theory (DFT) calculations whereby 0.25, 0.50, and 1.0 electrons were removed from a thin (6 Å) model Al<sub>2</sub>O<sub>3</sub> slab. More specifically, we examined the dehydrated  $\gamma$ -Al<sub>2</sub>O<sub>3</sub> (100) surface as a model alumina surface, as it is the lowest energy surface at the reaction conditions used in the experiments carried out herein with temperatures greater than 200 °C and partial pressures of water ( $P_{\rm H2O}$ ) > 10<sup>-3</sup> bar.<sup>13,41</sup>

The removal of electrons from the thin  $Al_2O_3$  slab increases the positive charge on the active  $Al_V^{3+}$  sites and increases the Lewis acidity. IPA then binds to the surface stronger. The adsorption energy decreases from  $-1.16 \text{ eV} (112 \text{ kJ mol}^{-1}, \text{ at}$ zero charge) to  $-1.78 \text{ eV} (-172 \text{ kJ mole}^{-1})$  as positive charge increases from 0 to  $0.125 \text{ h}^+$  site<sup>-1</sup>, as shown in Figure 5a–c. The IPA binding energy further strengthens as additional electron density is removed (Figure 5c). Naturally, the increased charge also results in shorter IPA  $O^{IPA}$ - $Al_V^{3+}$  bond distances; the  $O^{IPA}$ - $Al_V^{3+}$  bond decreases from 2.016 Å under neutral conditions to 1.963 Å for 0.125 h<sup>+</sup>/site, as shown in Figure 5a,b.

Alcohol dehydration over alumina has been proposed to proceed over the Al<sub>Va</sub><sup>3+</sup>-O<sub>3b</sub><sup>2-</sup> site pairs in a concerted E2 mechanism in which the C-OH and C-H bonds are simultaneously activated by the Alv3+-O2- acid-base site pairs to directly eliminate propene, which then desorbs.<sup>13,38</sup>  $Al-(OH)^{\delta-}$  and  $AlO-(H)^{\delta+}$  subsequently react to form and eliminate water. The calculated reactant, transition, and product states on the dehydrated  $\gamma$ -Al<sub>2</sub>O<sub>3</sub> (100) surface are shown in Figure S29. The energetic pathways to form propene and water from IPA are summarized in Figure 5d, depicting the activation energy as a function of the IPA binding energy. The neutral alumina site dehydrates IPA with  $E_a$  of 143 kJ mol<sup>-1</sup>, while the depleted alumina site (0.125  $h^+/site$ ), which has a higher Lewis acidity, has a lower barrier  $E_a$  of 119 kJ mol<sup>-1</sup>. The barriers decrease monotonically with the positive charge (i.e., acidity) of the  $Al_V^{3+}$  sites and show the same trends as those found experimentally and shown in Figure 3c. The calculated barriers also decrease monotonically with the heat of IPA adsorption.

# **Catalytic Condenser Performance**

The combined device characterization, spectroscopy, reaction experiments, and computation suggest that the active site of alumina has tunable acidity with the applied voltage,  $V_{CAT}$ . As suggested in Scheme 1, positive  $V_{CAT}$  depletes electrons in the catalytic alumina layer, most likely from the occupied in-gap states. The resulting electron-deficient aluminum site then binds IPA stronger and dehydrates it to propene with a lower activation barrier as larger V<sub>CAT</sub> is applied. Calculated activation energies (Figure 5d) indicated that the IPA-topropene activation barrier should reduce about 24 kJ mol<sup>-1</sup> when 12.5% of an electron is depleted from an  $Al_V^{3+}$  active site. Comparison with experimental activation energies (Figure 3c) exhibiting  $\Delta E_a$  of 16 kJ mol<sup>-1</sup> for differences in applied biases of 0 and +3 V indicates that about 5-10% of an electron has been depleted per active site at an applied  $V_{CAT}$  of +3 V; this is consistent with the ratio of electron depletion (~5  $\times$  10<sup>12</sup> h<sup>+</sup> cm<sup>-2</sup>, Figure 2) and alumina active site density (~5  $\times$  10<sup>13</sup> sites  $cm^{-2}$ ).

Enhancing electron depletion at the active site for more acidic solid materials beyond the current catalytic condenser requires application of higher voltages, higher quality insulating

 $(e.g., HfO_2)$  films, or alternative device design. The considered alumina/graphene catalytic condenser (70 nm  $HfO_2$ ) stabilizes charge in a 4 nm alumina film. Thicker alumina films reduce the capability for electronically altering the surface acid site strength, as shown by experiments with a  $\sim$ 50 nm alumina film on the alumina/graphene/HfO2/p-Si catalytic condenser (Figure S20); alternatively, thinner alumina films can potentially stabilize more charge per active site. Additionally, application of a higher voltage  $V_{\text{CAT}}$  to the considered catalytic condenser (4 nm alumina film, graphene, 70 nm HfO<sub>2</sub>) is problematic because leakage current through the device becomes substantial; this concept can be extended to higher voltages by further improving the quality of the HfO<sub>2</sub> film by reduction of pinholes. As a third design option, the condenser design can utilize different insulating materials such as higher-k dielectric materials (e.g., ferroelectrics) in thinner (<70 nm) insulating layers, both of which yield higher device capacitance and charge accumulation at the active site.

The electronically tunable nature of the alumina site in the catalytic condenser offers new capability for enhancing catalysis. While catalytic materials are conventionally tuned by precise synthesis of catalyst active sites, the limitation of material composition and stable structures only allows discrete options in active site design. In contrast, the bias potentials  $(V_{\text{CAT}})$  to the active layer of a catalytic condenser can span a continuum of voltage bias values. Precise selection of the active site electronic state can potentially allow for optimization of selectivity to products or maximization of the catalytic rate.

Finally, the catalytic condenser design also allows for the capability to program catalysts that operate under continuous forced dynamic conditions.<sup>1,42</sup> Oscillation of the catalyst state yielding variation in the binding energy of the reactants and products has been predicted to increase the catalytic rate orders of magnitude at resonance conditions,<sup>43</sup> control the extent of reaction,<sup>44</sup> and allow for significant selectivity control to products.<sup>45</sup> Resonance conditions of the maximum catalytic rate are predicted to occur at varying frequencies depending on the extent of reactant binding energy variation ( $\Delta BE$ ), the composition and structure of the active site, and the nature of the chemistry of interest; but simulations have predicted catalytic resonance to occur for frequencies of ~100 Hz up to 1 MHz.<sup>42-47</sup> For the considered catalytic condenser device, this requires depleting and replacing the electrons in the active sites at a comparable frequency. The time constant associated with electrons moving across the bare graphene surface (length of 1 cm) on HfO<sub>2</sub> was calculated to be 625  $\mu$ s using the electron mobility (1600 cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup>), corresponding to a device operational frequency of 1600 Hz. Upon adding the alumina layer to the graphene device, electron mobility in the full device still permitted operation at ~1000 Hz (see Supporting Information for the full calculations). The capability to program a repeating electronic perturbation of the catalyst (e.g., sinusoidal waveform of voltage bias) surface at high frequency (>1000 Hz) and amplitude ( $\Delta BE > 0.4 \text{ eV}$ ) opens a new possibility for catalysts that change at the time scale of the catalytic turnover frequency.

# CONCLUSIONS

A 1.0 cm<sup>2</sup> ML catalytic condenser device was synthesized by sequential deposition of 70 nm of amorphous  $HfO_2$  on a conductive p-type Si wafer, followed by transfer of a graphene layer, on top of which was grown a 4 nm amorphous layer of the alumina catalyst. The  $HfO_2$  layer remained amorphous

below 400 °C with a dielectric constant of 26, providing a capacitance of  $\sim 300$  nF cm<sup>-2</sup> and capability for stabilizing charge up to  $\sim 5 \times 10^{12}$  h<sup>+</sup>/cm<sup>2</sup>. Comparison of capacitance between two devices, one with graphene only and the other with alumina on graphene, indicated that the majority of the induced positive charge in the alumina/graphene device was actually in the aluminum oxide film. The alumina/graphene catalytic condenser with the HfO<sub>2</sub> insulator was observed to be conductive by scanning-tunneling microscopy with gap states located ~0.3 eV below the Fermi level as determined by ultraviolet photoelectron spectroscopy. Dehydration of IPA on the alumina/graphene catalytic condenser top surface exhibited a variable peak temperature of propene formation in TPSR experiments. Propene peak temperatures,  $T_{peak}$ , shifted 50 °C with different potential biases ( $V_{CAT}$  of 0, +1, +2, +3 V), corresponding to a difference in activation energy of  $\Delta E_a = 16$ kJ mol<sup>-1</sup> (0.17 eV). Enhanced Lewis acidity with positive  $V_{CAT}$ bias was consistent with increased IPA binding energy and reactivity to propene due to lower activation barriers as electrons were depleted at the surface active site, as calculated via DFT. The performance of the alumina/graphene catalytic condenser under bias indicates substantial potential for improving catalytic activity and tuning thermocatalytic selectivity, while oscillation of the alumina/graphene catalytic condenser via applied potential with large variation in surface acidity at frequencies as high as 1000 Hz achieves the conditions required for catalytic resonance.

# METHODS

The alumina-graphene catalytic condenser was synthesized and then evaluated by electron microscopy, spectroscopy, experimental catalytic reaction, AFM, and computation. Full research methods are described in the Supporting Information.

# **Catalytic Condenser Fabrication**

Devices were fabricated on a p-type Si substrate (WaferPro) using sequential deposition of HfO2, graphene, an Au/Ti contact, and a thin amorphous alumina active layer. HfO2 and alumina layers were grown within a flow-type ALD reactor (Kurt J Lesker ALD-150 LE). The 70nm HfO<sub>2</sub> dielectric layer was grown via 500 HfO<sub>2</sub> ALD cycles at 100 °C. Each HfO<sub>2</sub> ALD cycle was composed of alternating exposures to Tetrakis(dimethylamido) hafnium (TDMAH, t = 13 ms) and H<sub>2</sub>O (t= 200 ms) with purges in between each step. Graphene, obtained commercially on a Cu foil (Graphene Supermarket), was transferred on top of the HfO<sub>2</sub> layer using an existing PMMA-based method<sup>21,22</sup> and retained electrical conductivity (Figure S10). A small area (0.01 cm<sup>2</sup>) contact (Au/Ti, 45 nm/5 nm) was then grown on the graphene by e-beam evaporation through a shadow mask. Finally, the catalytic alumina layer of  $\sim$ 4 nm, based on the Al<sub>2</sub>O<sub>3</sub> growth rate, was deposited via 50 Al<sub>2</sub>O<sub>3</sub> ALD cycles at 100 °C using the same flowtype ALD reactor as mentioned above. Each Al<sub>2</sub>O<sub>3</sub> ALD cycle was composed of alternating exposures to trimethylaluminum (TMA, t =12 ms) and  $H_2O$  (t = 10,000 ms) with purges in between each step. Devices retained their structure after experimental reaction trials (Figures S23 and S24).

#### STEM, TEM, and SEM Analyses

Focused Ion Beam (FEI Helios NanoLab G4 dual-beam) was employed to prepare an electron-transparent cross-sectional lamella for the TEM and STEM study. An am-C layer (~50 nm) was deposited on the device using the sputter coater as a protection layer. The additional protection layer consisted of am-C (2  $\mu$ m) and Pt (2  $\mu$ m) was deposited on the region of interest. FIB was operated at 30 kV with 5 kV electron beam. Conventional TEM, HAADF-STEM images, and STEM–EDX elemental maps were obtained using Thermo Fisher Talos F200X G2 S-TEM equipped with a Super-X energy-dispersive X-ray spectrometer. TEM and HAADF-STEM

images were acquired at 200 kV. Camera length was set to be 125 mm with the probe convergence angle of 10.5 mrad for HAADF-STEM imaging. For EDX-mapping, beam current ranged from 300 pA to 400 pA. SEM images and SEM–EDX maps were acquired using JEOL 6500 FEG-SEM. SEM images and SEM–EDX elemental maps were obtained under 5 kV. The EDX maps were quantified using the Aztec EDX analytical system (Oxford Instruments). For more details, refer to the Supporting Information.

#### Spectroscopy

Raman spectra were obtained from a Witec Alpha 300R confocal Raman microscope equipped with a UHTS300 spectrometer and a DV401 CCD detector. XRD patterns were recorded on a Bruker D8 Discover 2D diffractometer equipped with a Co K $\alpha$  source ( $\lambda$  = 0.179 nm). Full details are available in the Supporting Information. XPS measurements were performed on a PHI Versa Probe III XPS system (ULVAC-PHI) using a monochromated Al K $\alpha$  X-ray source (1486.6 eV). The base pressure was  $4.0 \times 10^{-8}$  Pa. During data collection, the pressure was ca.  $2.0 \times 10^{-6}$  Pa. All samples were mounted on the holder using a Cu pin. The measurements were conducted using an X-ray spot size of  $0.1 \times 0.1 \text{ mm}^2$  with a power of 25 W under 15 kV. The survey spectra were measured using 280 eV pass energy and 1.0 eV/step. The data were processed with Multipak software. UPS measurements were performed on the same XPS system using ultraviolet radiation from an ionized He source. The base pressure was  $8.0 \times 10^{-7}$  Pa during measurement. The measurement spot size was approximately  $6 \times 6 \text{ mm}^2$ . Survey spectra were measured using 1.3 eV pass energy and 0.025 eV per step.

#### **Temperature-Programmed Surface Reaction**

TPSR was carried out in a customized ultrahigh vacuum chamber reactor with IPA dehydration to propene as a reaction probe. The propene signal was monitored during the temperature ramp of the device, while voltage was applied. Full description of the equipment setup, device handling, temperature calibration, and the full TPSR process can be found in the Supporting Information.

#### Atomic Force Microscopy

AFM measurements were performed using a Bruker Dimension Icon closed-loop atomic force microscope using the force-modulation technique with a force feedback setpoint of 581 pN. Probes were standard Si cantilevers calibrated before measurement with a spring constant of 0.674 N m<sup>-1</sup>.

## Scanning Tunneling Microscopy

STM measurements were made using an Omicron low-temperature scanning tunneling microscope at room temperature under a vacuum of  $8.0 \times 10^{-10}$  mBar using a Pt/Ir probe. Samples were grounded through the alumina/graphene film. Differential spectra were collected from a Stanford Research Systems lock-in amplifier with a modulation signal of  $V_{\rm mod} = 10$  mV and  $f_{\rm mod} = 10$  kHz. Data were processed using the WSxM software.

#### Computation

First principles periodic DFT calculations were carried out to determine the binding energies and the activation barriers using the plane wave implementation of DFT in the Vienna ab initio simulation package (VASP).<sup>48-50</sup> The generalized gradient approximation form of the Perdew, Burke, and Ernzerhof functional was used to determine the energies, and dispersion corrections were included via the D3 corrections developed by Grimme.<sup>51,52</sup> An energy cutoff of 400 eV was used in the construction of plane waves, and projector augmented wave potentials were used to model interactions between core and valence electrons.<sup>53,54</sup> A periodic slab of  $2 \times 2$  times the primitive unit cell and consisting of seven layers of  $\gamma$ -Al<sub>2</sub>O<sub>3</sub> (100) was used for all calculations, with the bottom-most layer held fixed. All electronic energies were converged to within a tolerance of  $10^{-6}$  eV using a 2 ×  $1 \times 1$  gamma-centered k-point grid. The geometries of the structures were optimized until the maximum force on each atom was calculated to be less than 0.05 eV Å-1. Additional details are provided in the Supporting Information.

# ASSOCIATED CONTENT

# Supporting Information

The Supporting Information is available free of charge at https://pubs.acs.org/doi/10.1021/jacsau.2c00114.

Catalytic condenser fabrication methods, TEM and SEM characterization, reactor setup, reaction process, and control experiments, electronic characterization, XRD, Redhead analysis of TPSRs, and additional references 55 and 56 (PDF)

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#### Notes

The authors declare no competing financial interest.

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