RSC Advances



PAPER



Cite this: RSC Adv., 2022, 12, 11190

Revealing the different performance of Li₄SiO₄ and Ca₂SiO₄ for CO₂ adsorption by density functional theory†

Wenjing Yu, D^a Qian Xu, *^a Shenggang Li, D *^b Xiaolu Xiong, ^a Hongwei Cheng, D^a Xingli Zou D^a and Xionggang Lu^a

To reveal the difference between Li_4SiO_4 and Ca_2SiO_4 in CO_2 adsorption performance, the CO_2 adsorption on Li_4SiO_4 (010) and Ca_2SiO_4 (100) surfaces was investigated using density functional theory (DFT) calculations. The results indicate that the bent configuration of the adsorbed CO_2 molecule parallel to the surface is the most thermodynamically favorable for both Li_4SiO_4 and Ca_2SiO_4 surfaces. The Li_4SiO_4 (010) surface has greater CO_2 adsorption energy ($E_{\text{ads}} = -2.97 \, \text{eV}$) than the Ca_2SiO_4 (100) surface ($E_{\text{ads}} = -0.31 \, \text{eV}$). A stronger covalent bond between the C atom of adsorbed CO_2 and an O_5 atom on the Li_4SiO_4 (010) surface is formed, accompanied by more charge transfer from the surface to CO_2 . Moreover, the Mulliken charge of O_5 atoms on the Li_4SiO_4 (010) surface is more negative, and its p-band center is closer to the E_{f} , indicating O_5 atoms on Li_4SiO_4 (010) are more active and prone to suffering electrophilic attack compared with the Ca_2SiO_4 (100) surface.

Received 16th February 2022 Accepted 4th April 2022

DOI: 10.1039/d2ra01021f

rsc li/rsc-advances

1. Introduction

CO2 capture, storage and utilization (CCSU) is considered as one of the most promising technologies for reducing anthropogenic CO₂ emission, which can lead to global warming. Solid inorganic sorbents have been proven to efficiently remove CO₂ at high temperatures, and are more economical and effective than low-temperature amine-based materials in CO₂ capture from high temperature exhaust gas.1 Lithium orthosilicate (Li₄SiO₄) is one of the best CO₂ capture sorbents due to its significant advantages, such as large adsorption capacity, low regeneration temperature, and good adsorption and desorption cycle stability.2-4 There have been a lot of experimental studies on Li₄SiO₄ as an adsorbent to capture CO₂, including the synthesis method,5,6 kinetic behavior7-9 and modification of Li₄SiO₄.¹⁰⁻¹² However, lithium is relatively expensive and not very abundant in the earth's crust. In particular, lithium batteries have been widely used as a source of power or energy for a lot of things from portable electronics to electric vehicles. As a result, the demand for lithium is increased which leads directly to an increase of its price. Accordingly, it is very difficult to apply lithium-based ceramics on a huge scale to capture CO₂ economically and sustainably. Meanwhile there are abundant resources basic silicates all over the world, especially calcium silicates (Ca₂SiO₄) are often found in the industrial by-products named as slags generated during iron and steel production. Furthermore, Ca₂SiO₄, similar to Li₄SiO₄, is thermodynamically favorable for CO₂ capture from room temperature to 572 °C at ambient pressure, and Gibbs free energy changes for the carbonation of Ca2SiO4 and Li4SiO4 were calculated by HSC Chemistry 6.0 and shown in Fig. 1. However, the slow diffusion and reaction of carbonation between CO₂ and calcium silicates is a common issue even at high temperatures in the case of no participation of water.¹³ It was found that the amount of CO₂ captured with Ca2SiO4 is little at the temperature range from room temperature to 572 °C in our previous study. Zhao et al. 14,15 applied Ca2SiO4 as the inert material to enhance the sintering resistance and cyclic stability of CaO during multiple sorption/ desorption of CO₂. It could be deduced that Ca₂SiO₄ is much inerter for carbonation compared with Li₄SiO₄ at the high temperatures.

The investigation on the effect of the electronic structure of the silicates on their carbonation reactivity should be very important for understanding deeply the different carbonation behaviors for Li_4SiO_4 and Ca_2SiO_4 , and developing new approaches to improve the carbonation activity of silicates.

There are some investigations about CO₂ adsorption on the surface of the oxides and silicates with the first-principles calculations. Kim *et al.*¹⁶ made an assessment of Li₂O and Na₂O surfaces for CO₂ adsorption based on DFT calculations. They found that the introduction of dopant atoms larger than host metal atoms of the surfaces can negatively increase CO₂ adsorption energies. Kumar *et al.*¹⁷ studied the CO₂ adsorption on different terminations of Cr₂O₃ surfaces with DFT calculations and found that carboxylate species are formed on O layer

^aShanghai University, China

^bShanghai Advanced Research Institute, China

[†] Electronic supplementary information (ESI) available. See https://doi.org/10.1039/d2ra01021f

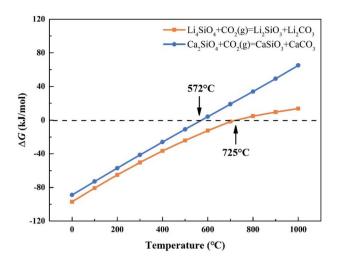


Fig. 1 Relationship between ΔG and temperature of carbonation reaction between CO_2 and Li_4SiO_4 and Ca_2SiO_4 .

terminated-(0001), and carbonate species are formed on O layer terminated-(10 $\bar{1}2$) and Cr layer terminated-(0112), indicating that the formation of physisorbed and chemisorbed species depends on different surface terminations. Kang *et al.*¹⁸ thermodynamically evaluated the CO₂ capture potential of Mg₂MO₄ (M = Si, V, and Ge). Their results indicated that the critical temperature at which CO₂ can be absorbed, increased with decreasing Pauling electronegativity of the M site.

There are several investigations for Li₄SiO₄ on its structural, electronic, lattice dynamical and thermodynamic properties. Duan¹⁹ and Tang *et al.*²⁰ found the covalency properties of Li₄SiO₄ mainly resulting from the overlap of O 2p and Si 3p orbitals. Kong *et al.*²¹ studied the adsorption mechanism of H₂O on the Li₄SiO₄ (010) surface. It was suggested an interaction between adsorbed H₂O and Li₄SiO₄ (010) surface, including an electrophilic interaction of hydrogen atom in water with oxygen

atoms on the surface and a nucleophilic interaction of oxygen atoms in water with Li atoms on the surface.

Ca₂SiO₄, as the industrial cement clinkers, has been investigated by DFT calculations extensively, and many studies focused on the hydration of Ca₂SiO₄ phases. Qi et al.²² investigated H₂O adsorption on low-index surfaces of Ca₂SiO₄, indicating that electron are mainly transferred from surface atoms to H₂O molecule. Wang et al.²³ evaluated H₂O adsorption on β-Ca₂SiO₄ surfaces and found a dual interaction between H₂O and β-Ca₂SiO₄ (100) surface. Wang et al.²⁴ also studied the relationship between reactivity and electronic structure of α'_{L} -, β and γ-Ca₂SiO₄ for hydration process. They found that the higher hydration reactivity of α'_{L} - and β -Ca₂SiO₄ compared with γ -Ca2SiO4 are attributed to the higher charge density and larger local state density of the active oxygen atoms in α'_{L} - and β -Ca₂SiO₄. However, there are few investigations about explanation of the different behaviors of carbonation of Ca₂SiO₄ and Li₄SiO₄ on the base of their structural and electronic properties.

Herein, we have systematically investigated the structural and electronic properties of Li_4SiO_4 and Ca_2SiO_4 , and the adsorption of CO_2 on the most stable surfaces of Li_4SiO_4 and Ca_2SiO_4 on the base of density functional theory calculations. We tried to reveal the relationship between the electronic structures of Li_4SiO_4 and Ca_2SiO_4 and their reactivity for CO_2 adsorption on the molecular scale. The results of this investigation could converge to a proposed mechanism of CO_2 capture with the orthosilicates, on which the more reactive silicates for CO_2 capture can be screened out as the candidates for CO_2 capture.

2. Computational details

The calculations were performed based on density functional theory (DFT), using Cambridge Series Total Energy Package (CASTEP) code.²⁵ The exchange-correlation potential was approximated within the generalized gradient approximation

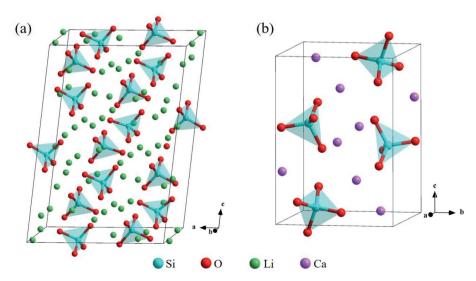


Fig. 2 Crystal structures of (a) Li₄SiO₄ and (b) Ca₂SiO₄. The Si, O, Li and Ca atoms are shown by blue, red, green and purple spheres, respectively.

Table 1 Comparison of calculated lattice constants of Li_4SiO_4 and Ca_2SiO_4 with experimental lattice constants

	Li ₄ SiO ₄			Ca ₂ SiO ₄			
	Cal	Expt ¹⁹	Δ (%)	Cal	Expt ⁴⁰	Δ (%)	
a (Å) b (Å)	11.511 6.080	11.532 6.075	0.18 0.08	5.571 6.800	5.502 6.745	1.25 0.82	
c (Å)	16.708	16.678	0.18	9.354	9.297	0.61	
β (°)	99.15	99.04	0.11	94.295	94.590	0.31	

(GGA)²⁶ using the Perdew–Burke–Ernzerhof (PBE) functional.²⁷ Dispersion-corrected calculations²⁸ were performed with Grimme's DFT-D3 methodology.²⁹ To model Li_4SiO_4 and Ca_2SiO_4 , the unit cell $(1 \times 1 \times 1)$ was applied for the calculation.

In order to optimize the crystal structures, the plane wave truncation energy and k-points were tested. A cutoff-energy of 650 eV was used for plane wave expansions. The k-points meshes within Monkhorst-Pack³⁰ framework were set as $3\times6\times2$ and $4\times3\times2$ for Li₄SiO₄ and Ca₂SiO₄ respectively. The Broyden–Fletcher–Goldfarb–Shenno (BFGS)³¹ minimization algorithm was used to optimize the primitive unit cell. The surfaces of Li₄SiO₄ and Ca₂SiO₄ were cleaved from the optimized bulk structure. All surfaces were kept stoichiometric and neutral to avoid the polarizing electric field. The thicknesses of vacuum layer were set as 15 Å to avoid the interaction between

slabs. The convergence criteria were fixed, specifically: the energy change within 1×10^{-5} eV per atom, the force on the atoms within 0.03 eV Å⁻¹, the stress on the atoms within 0.05 GPa, and the displacement of atoms within 1×10^{-3} Å. All the initial crystal structures and date were obtained from the Inorganic Crystal Structure Database (ICSD).³²

Results and discussion

3.1. Structural and electronic properties of bulks

The bulk structure of Li_4SiO_4 with the monoclinic phase, which space group is P21/m (no. 11), $^{33-35}$ was optimized. The unit cell of Li_4SiO_4 contains 126 atoms, including 14 $\left[\text{SiO}_4\right]^{4-}$ tetrahedra and 56 Li atoms, as shown in Fig. 2a, which are centrally symmetrical. Meanwhile, the bulk structure of Ca_2SiO_4 with the monoclinic phase, which space group is P21/n (no. 14), $^{36-38}$ was optimized as well. The unit cell of Ca_2SiO_4 consists of 28 atoms, including 4 $\left[\text{SiO}_4\right]^{4-}$ tetrahedra and 8 Ca atoms, as shown in Fig. 2b. The calculated and experimentally measured lattice parameters of Li_4SiO_4 and Ca_2SiO_4 are presented in Table 1, and they are in good agreement, implying the simulation settings are reliable and give reasonable results.

Fig. 3 shows the total density of states (TDOS) and partial density of states (PDOS) for Li_4SiO_4 and Ca_2SiO_4 . Electrons occupying the orbitals below and near the Fermi level (E_f) is of great significance to the activity of the crystal materials for chemical reactions, ³⁹ so we focused on the electrons on the

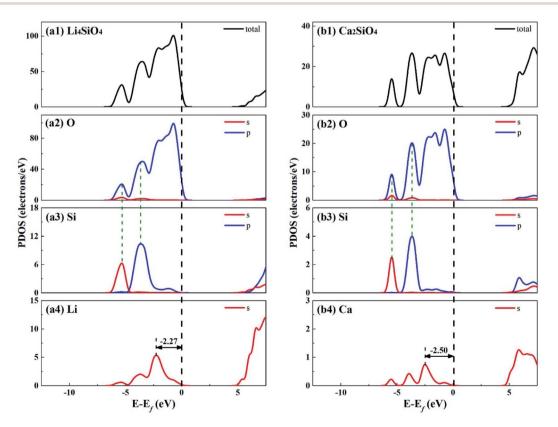


Fig. 3 DOS analysis for (a) Li_4SiO_4 and (b) Ca_2SiO_4 . The black dashed line shows the Fermi level.

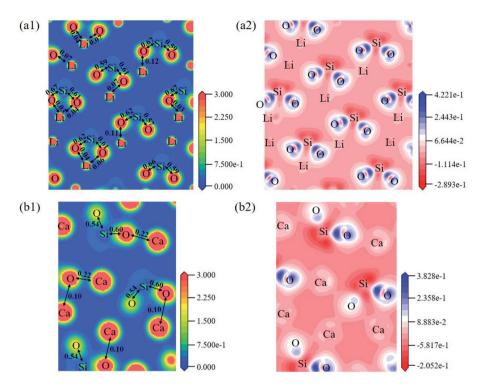


Fig. 4 Contour maps of electron density distributions and differential charge density of (a) Li_4SiO_4 in the plane (010) and (b) Ca_2SiO_4 in the plane (100).

Table 2 The surface energies (E_{surf} in J m $^{-2}$) of low Miller index surfaces of Li₄SiO₄ and Ca₂SiO₄

Surface	(100)	(010)	(001)	(110)	(101)	(011)	(111)
Li ₄ SiO ₄		0.78	1.34	0.80	1.28	0.87	0.84
Ca ₂ SiO ₄		1.19	0.80	0.86	0.66	0.75	0.82

orbitals below and near the $E_{\rm f}$. For Li₄SiO₄ and Ca₂SiO₄, their TDOS near the $E_{\rm f}$ is mainly contributed by the p orbitals of O atoms, suggesting that O atoms are more active and more likely serve as the electron donors. Their PDOS in the region between $-6.55\sim-2.92$ eV and $-6.25\sim-3.02$ eV are overlapped with the

Si s and p bands, implying orbital hybridization and Si–O binding in the bulks of Li_4SiO_4 and Ca_2SiO_4 . However, the states in the region between -3 and 0 eV, near to E_f is about 73% of the total PDOS for O p orbitals for Li_4SiO_4 , while that for Ca_2SiO_4 is about 70%. It can be deduced that there are more electronic states for O p orbitals below and near to E_f in Li_4SiO_4 than those in Ca_2SiO_4 , Then the electron transfer from O atoms occurs more easier in Li_4SiO_4 than in Ca_2SiO_4 .

Furthermore, the first high peak position⁴¹ in the PDOS for Li s orbital of Li_4SiO_4 , is closer to the E_f compared with that for Ca s orbital of Ca_2SiO_4 , implying that the outer electron of Li

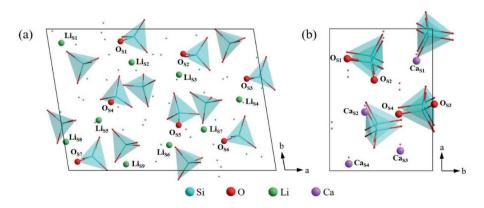


Fig. 5 Atomic arrangement of (a) $\text{Li}_4 \text{SiO}_4$ (010) surface and (b) $\text{Ca}_2 \text{SiO}_4$ (100) surface from the top view. The lines and dots represent the underlying atoms and tetrahedra respectively.

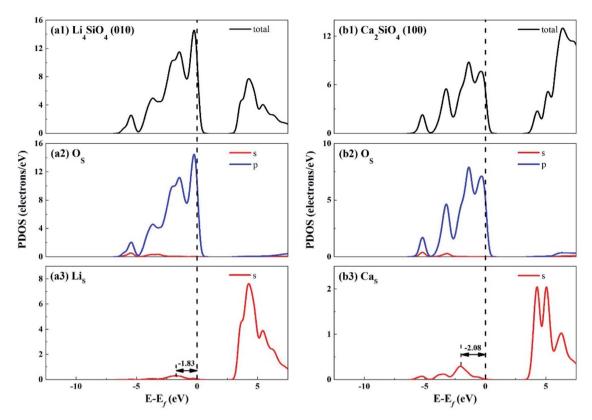


Fig. 6 DOS analysis of the topmost surface layers of (a) Li_4SiO_4 (010) surface and (b) Ca_2SiO_4 (100) surface without CO_2 adsorption. The black dashed line shows the Fermi level.

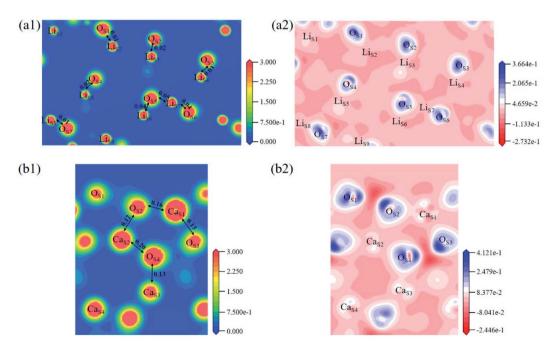


Fig. 7 Contour maps of electron density distributions and differential charge density in the cross sections perpendicular to the (001) plane in (a) Li_4SiO_4 (010) surface and (b) Ca_2SiO_4 (100) surface.

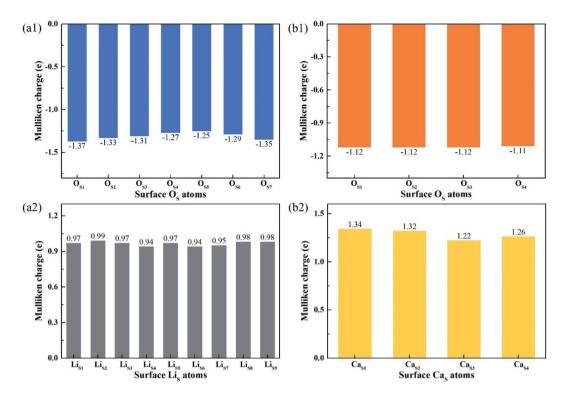


Fig. 8 Mulliken charge analysis of atoms on the topmost surface layers of (a) Li_4SiO_4 (010) surface and (b) Ca_2SiO_4 (100) surface.

atoms in Li_4SiO_4 can be transferred away easier than Ca atoms in Ca_2SiO_4 .

The electron density distribution can show the bonding between atoms and differential charge density can show the accumulation and depletion of electrons. In Fig. 4a, electron density between Si and O atoms is higher than surrounding and electrons accumulate in the middle of Si and O atoms, demonstrating that a covalent interaction of Si–O. And according to the charge population marked in the figure, the Si–O covalent interaction is stronger in Li_4SiO_4 . Li/Ca–O have a certain covalent interaction and the strength of interaction of Li–O is weaker than Ca–O.

3.2. Structural and electronic properties of surfaces

In order to find out the most stable surface, the surface energies of seven low Miller index surfaces were calculated. The surface energy ($E_{\rm surf}$) can be calculated according to the eqn (1):^{42,43}

$$E_{\text{surf}} = (E_{\text{slab}} - nE_{\text{bulk}})/2A \tag{1}$$

where $E_{\rm slab}$ and $E_{\rm bulk}$ are the total energy of relaxed slab model and unit cell, respectively. n is the number of formula units contained in the slab. A is the area of the slab. According to the calculation results listed in Table 2, the (010) surface and (100) surface were the most stable surface of ${\rm Li_4SiO_4}$ and ${\rm Ca_2SiO_4}$ respectively, due to their lowest values of surface energy, which is consistent with the previous calculations. 21,22,42 Fig. 5 shows the atomic arrangement of these two surfaces from the top view. The topmost surface layer consists of ${\rm Li}$ and ${\rm O}$ atoms for

 $\text{Li}_4 \text{SiO}_4$, and Ca and O atoms for $\text{Ca}_2 \text{SiO}_4$. The atoms of Li, Ca, and O on the topmost surface layer are referred to as Li_S, Ca_S and O_S hereafter, respectively.

The electronic properties of the surfaces should differ from those of the bulks due to the dangling bonds or surface reconstruction. The electronic properties of Li₄SiO₄ (010) surface and Ca₂SiO₄ (100) surface were calculated as well. Fig. 6 shows the TDOS and PDOS of the topmost surface layers of Li₄SiO₄ (010) surface and Ca₂SiO₄ (100) surface. Although the surface TDOS resemble the bulk TDOS shown in Fig. 3 for both Li₄SiO₄ and Ca₂SiO₄, the states of O_S p orbitals are shifted up in energy, and the states in the region between -3 and 0 eV are about 79% of the total PDOS for O_S p orbitals for Li₄SiO₄ (010) surface, while that for Ca₂SiO₄ (100) surface is about 77%. Furthermore, p-band center of O_S atoms increased to −1.725 eV from -1.936 eV of O atoms in bulk Li₄SiO₄, and -1.939 eV from -2.103 eV in bulk Ca₂SiO₄. It can be deduced that the reactivity of the surface O_S atoms is enhanced compared with the O atoms in the bulks. Considering the states near to E_f and the p-band center levels, the O_S atoms in Li₄SiO₄ (010) surface are more prone to suffer the electrophilic attacks with respect to Ca2SiO4 (100) surface.

The electron density and differential charge density of the topmost surface layer atoms are shown in Fig. 7. The covalent interaction between Li_{S} and O_{S} atoms is weaker than that in bulk, which can be seen from the charge population. While the covalent interaction between Ca_{S} and O_{S} atoms is stronger than that in bulk.

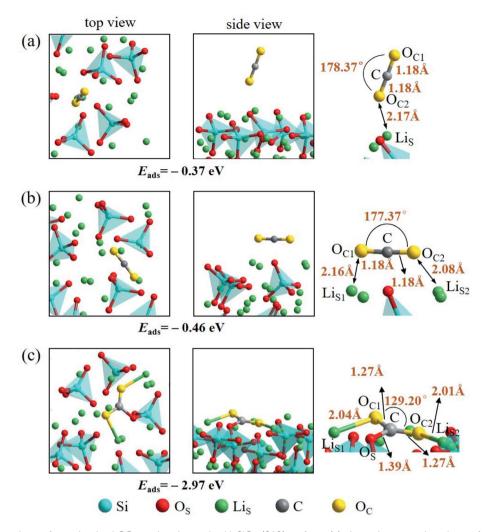


Fig. 9 Three configurations of an adsorbed CO_2 molecule on the Li_4SiO_4 (010) surface: (a) along the normal to the surface, (b) parallel to the surface, (c) bent configuration.

Fig. 8 shows the Mulliken charge of atoms on the topmost surface layer. There is a considerable difference of the charge between bulk atoms (listed in ESI†) and atoms on the topmost surface layer. The positive charge of surface Li_{S} and Ca_{S} atoms increases, and the negative charge of the O_{S} atoms increases compared with bulk atoms. And the deviation of Mulliken charge of surface Li_{S} atoms from the bulk atoms is relatively large, whereas the Mulliken charge of surface Ca_{S} and O_{S} atoms differ from their bulk atoms slightly. Furthermore, the O_{S} of $\text{Li}_{4}\text{SiO}_{4}$ (010) surface carry more negative charge than $\text{Ca}_{2}\text{SiO}_{4}$ (100) surface. According Lewis acid/base theory, O_{S} atoms of $\text{Li}_{4}\text{SiO}_{4}$ (010) surface are more basic and easier to lose electrons.

3.3. CO₂ adsorption on the surfaces

The adsorption energy ($E_{\rm ads}$) of a CO₂ molecule on Li₄SiO₄ (010) surface and Ca₂SiO₄ (100) surface is calculated according to the eqn (2):

$$E_{\text{ads}} = E_{\text{slab+CO}_2} - (E_{\text{slab}} + E_{\text{CO}_2}) \tag{2}$$

where $E_{\rm slab+CO_2}$ is the total energy of the surface with CO₂ adsorption, $E_{\rm CO_2}$ is the total energy of an isolated CO₂ molecule. The lower adsorption energy describes the stronger binding between the adsorbed CO₂ molecule and the surface, which reflects the stability of adsorption.

An isolated CO_2 molecule has linear configuration with the length of C–O bond of 1.18 Å. Three adsorption configurations are presented in this study when a CO_2 molecule is adsorbed on the Li_4SiO_4 (010) surface as shown in Fig. 9. In the first configuration, the adsorbed CO_2 molecule almost remains linear configuration along the normal to the Li_4SiO_4 (010) surface. The distance between O atom in CO_2 and the nearest Li_S atom is 2.17 Å, and adsorption energy is -0.37 eV. In the second configuration shown in Fig. 9b, the adsorbed CO_2 molecule has a linear configuration parallel to the Li_4SiO_4 (010) surface, and the adsorption energy is -0.46 eV. In the third configuration shown in Fig. 9c, the adsorbed CO_2 molecule is lying flat on the surface with the bent configuration. The C atom in CO_2 forms a bond with a surface O_S atom with the O_S –C distance of 1.39 Å, and its two oxygen atoms (O_C) are

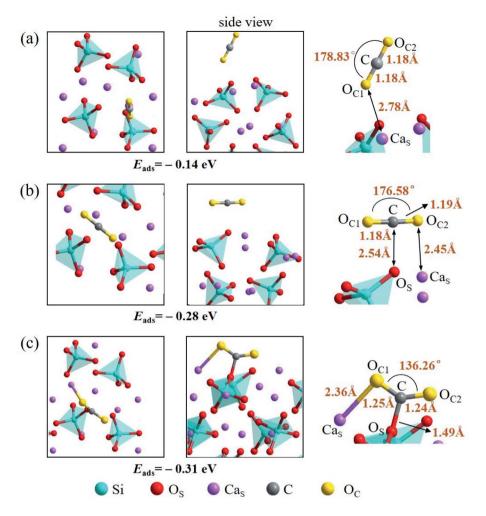


Fig. 10 Three configurations of an adsorbed CO_2 molecule on the Ca_2SiO_4 (100) surface: (a) along the normal to the surface, (b) parallel to the surface, (c) bent configuration.

coordinated to the two surface Li $_{\rm S}$ atoms, with the O $_{\rm C_1}$ –Li $_{\rm S_1}$ and O $_{\rm C_2}$ –Li $_{\rm S_2}$ distances are 2.04 and 2.01 Å, respectively. The adsorbed CO $_{\rm 2}$ molecule is bent with an O $_{\rm C_1}$ –C–O $_{\rm C_2}$ angle of 129.20° and the C–O $_{\rm C}$ length of 1.27 Å. The adsorption energy is –2.97 eV. The third bent configuration of the adsorbed CO $_{\rm 2}$ on the Li $_{\rm 4}$ SiO $_{\rm 4}$ (010) surface is energetically favorable over the other two configurations.

There are also three adsorption configurations considered for an adsorbed $\rm CO_2$ molecule on the $\rm Ca_2SiO_4$ (100) surface, as shown in Fig. 10. The first configuration, where an almost linearly $\rm CO_2$ is adsorbed, is in vertical orientation and tilted slightly to the $\rm Ca_2SiO_4$ (100) surface, and the adsorption energy is -0.14 eV. The absorbed $\rm CO_2$ molecule is bent a little bit in the second configuration, which is in parallel orientation, and the adsorption energy is -0.28 eV. When the absorbed $\rm CO_2$ molecule lying on the $\rm Ca_2SiO_4$ (100) surface in a bent configuration shown in Fig. 10c, the distance of $\rm C-O_S$ is 1.49 Å, and $\rm O_{C_1}$ –C-O $_{C_2}$ angle is 136.26°. The bent configuration has the lowest value of adsorption energy, -0.31 eV in the third configurations.

It is found that the bent configuration consisting of a CO_2 molecule absorbed parallel along to the surface is the most thermodynamically stable among the three configurations considered here for both Li_4SiO_4 and Ca_2SiO_4 surfaces. Furthermore, the Li_4SiO_4 (010) surface has greater adsorption to CO_2 than the Ca_2SiO_4 (100) surface due to the stronger bond between C atom in CO_2 and the surface O_S atom for Li_4SiO_4 .

3.4. Partial density of state analysis

Considering that the adsorption energy of adsorbed CO_2 in bent configuration is the lowest, the PDOS calculations were only performed for this configuration, as shown in Fig. 11 and 12. It can be seen that the s and p orbitals of C and O_{C} of adsorbed CO_2 molecule both move towards lower energy level and broaden compared with the isolated CO_2 molecule, indicating that CO_2 molecule becomes more stable after adsorption. C s and p orbitals of adsorbed CO_2 are hybridized with O_{S} p orbitals, having bonding character between C and O_{S} atoms. The states in the region between -3 and 0 eV are about 72% for O_{S} p orbital for $\mathrm{Li}_4\mathrm{SiO}_4$ (010) surface after adsorption, which is

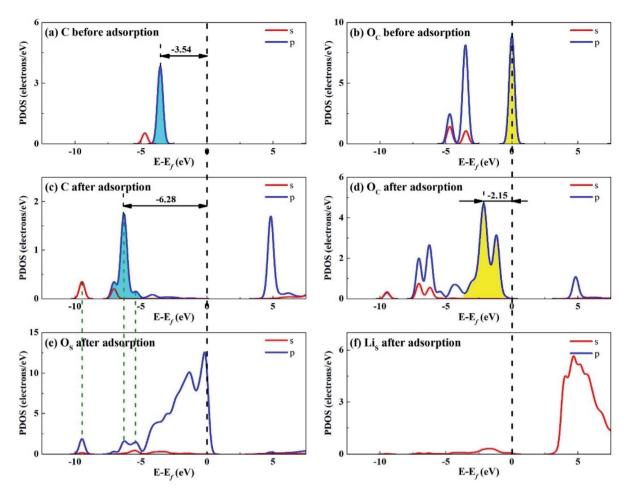


Fig. 11 PDOS analysis of (a) C and (b) O_C atoms of CO_2 before adsorption, (c) C and (d) O_C atoms of CO_2 and (e) O_S and (f) Li_S atoms of Li_4SiO_4 (010) surface after adsorption.

decreased from 79% before adsorption, demonstrating that PDOS of O_S p orbitals is moved to lower energy level after CO_2 adsorption. Furthermore, the PDOS peak of Li_S s orbital become weaker and broader after adsorption.

Fig. 12 shows the PDOS for CO_2 and Ca_2SiO_4 (100) surface after adsorption. Similarly, the s and p orbitals of C and O_C of adsorbed CO_2 move towards lower energy and broaden, but not as far as Li_4SiO_4 (010) surface. It can be deduced that CO_2 adsorption on Li_4SiO_4 (010) surface is more stable. The states in the region between -3 and 0 eV are about 73% for O_S p orbitals for Ca_2SiO_4 (100) surface. The PDOS peaks for Ca_S s orbital become weaker and broader after adsorption, similar to Li_S s orbital.

To better elucidate the different CO_2 absorption behaviors of Li_4SiO_4 (010) and Ca_2SiO_4 (100) surfaces, the p-band centers of C and O_C in CO_2 and O_S on the surfaces with and without CO_2 absorption were calculated, and the results were shown in Fig. 13. Comparing the p-band centers of O_S atoms on two clean surfaces, it can be found that the p-band center of O_S atoms on the Li_4SiO_4 (010) surface is closer to the E_f than that of Ca_2SiO_4 (100) surface, which means the O_S atoms of Li_4SiO_4 (010) surface are more active and easier to transfer electrons to CO_2

absorbed. When CO_2 is adsorbed on the Li_4SiO_4 (010) surface, it is obvious that p-band centers of C, O_C and O_S atoms are farther away from the E_f , indicating that CO_2 absorbed and O_S become stable. On the other hand, energy-down shift of the p-band centers of C, O_C and O_S atoms of Ca_2SiO_4 (100) surface due to CO_2 adsorption is much smaller compared with Li_4SiO_4 (010) surface, probably leading to its higher absorption energy and less CO_2 absorption.

3.5. Mulliken charge analysis

To understand the interactions and charge distributions associated with $\rm CO_2$ adsorbed with the most energetically favorable configurations on $\rm Li_4SiO_4$ (010) and $\rm Ca_2SiO_4$ (100) surfaces, a Mulliken charge anlysis was performed. The detail date of Mulliken charges for $\rm CO_2$ and $\rm Li_4SiO_4$ (010) and $\rm Ca_2SiO_4$ (100) surfaces before and after adsorption were shown in Fig. 14. For $\rm CO_2$ adsorbed on $\rm Li_4SiO_4$ (010) surface, it is found that the Mulliken charge on a surface $\rm O_S$ atom changes from -1.33e to -0.86e, while the charges on the surface $\rm Li_{S_1}$ and $\rm Li_{S_2}$ atoms increase from 0.93e to 0.93e to 0.93e, and from 0.93e to 0.95e, respectively. It can be deduced that $\rm CO_2$ adsorption induces the net

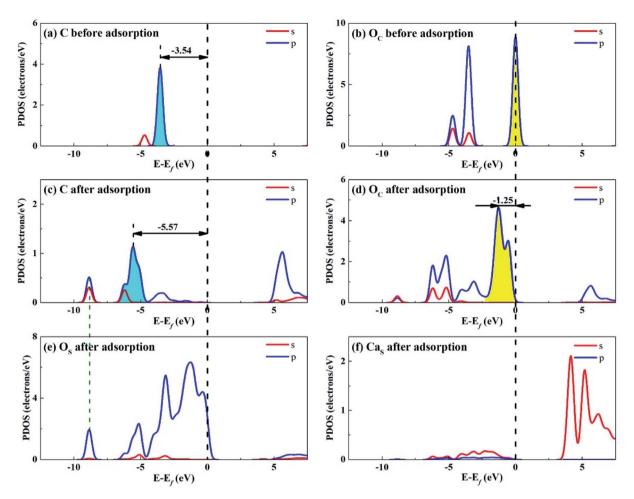


Fig. 12 PDOS analysis of (a) C and (b) O_C atoms of CO_2 before adsorption, (c) C and (d) O_C atoms of CO_2 and (e) O_S and (f) Ca_S atoms of Ca_2SiO_4 (100) surface after adsorption.

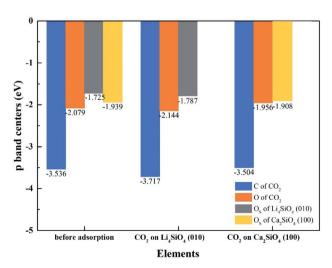


Fig. 13 The p-band centers of C 2p and O 2p orbitals of CO_2 , Li_4SiO_4 (010) surface and Ca_2SiO_4 (100) surface before and after adsorption.

electron loss of Li_4SiO_4 (010) surface. On the other side, the Mulliken charge of C atom in CO_2 adsorbed decreases from 0.98e to 0.76e, and the charges of O_{C_1} and O_{C_2} from -0.49e to -0.78e and -0.76e, respectively, which means that CO_2 adsorbed gains the charges from Li_4SiO_4 (010) surface.

In the case of CO_2 adsorption on Ca_2SiO_4 (100) surface, it is found that charge gain for C atom in CO_2 adsorbed is similar to that on Li_4SiO_4 (010) surface, but charges gained by O_{C1} and O_{C2} atoms are fewer. Furthermore, the Mulliken charge on an O_S atom on Ca_2SiO_4 (100) surface changes from -1.12e to -0.88e, and the charge on the surface Ca_S from 1.34e to 1.40e. The net charge transfer from Ca_2SiO_4 (100) surface to CO_2 adsorbed is much less compared to that Li_4SiO_4 (010) surface.

The charge population between C in CO_2 adsorbed and a surface O_S atom was calculated to be 0.62 and 0.45 for Li_4SiO_4 (010) and Ca_2SiO_4 (100) surfaces, respectively. It can be deduced that the $C-O_S$ covalent interaction on Li_4SiO_4 (010) surface is even stronger, which leads to the stronger adsorption of CO_2 .

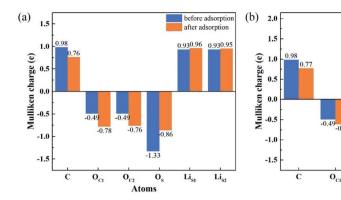


Fig. 14 Mulliken charge analysis of adsorbed CO_2 molecule and (a) Li_4SiO_4 (010) surface, and (b) Ca_2SiO_4 (100) surface before and after adsorption.

4. Conclusions

A density functional theory calculation was conducted to research the $\rm CO_2$ adsorption on the $\rm Li_4SiO_4$ (010) and $\rm Ca_2SiO_4$ (100) surfaces. The bent configuration consisting of a $\rm CO_2$ molecule adsorbed parallel along to the surface is the most thermodynamically favorable for $\rm Li_4SiO_4$ and $\rm Ca_2SiO_4$ surfaces. And the adsorption energy of $\rm Li_4SiO_4$ (010) surface is -2.97 eV, more negative than $\rm Ca_2SiO_4$ (100) surface, -0.31 eV. $\rm Li_4SiO_4$ (010) surface is more favorable for forming a stronger covalent bond between a surface $\rm O_S$ atom to the C atom of $\rm CO_2$ adsorbed and transferring more charges to adsorbed $\rm CO_2$. In addition, it was found that the Mulliken charge of $\rm O_S$ atoms on the $\rm Li_4SiO_4$ (010) is more negative, and its p-band center is closer to the $\rm E_f$, which implies $\rm O_S$ atoms of $\rm Li_4SiO_4$ (010) are more active and more likely serve as the electron donors with respect to $\rm Ca_2SiO_4$ (100) surface.

Conflicts of interest

There are no conflicts to declare.

Acknowledgements

This work was financially supported by the Natural Science Foundation of China (No. 51574163).

References

- 1 R. Ben-Mansour, M. A. Habib, O. E. Bamidele, M. Basha, N. A. A. Qasem, A. Peedikakkal, T. Laoui and M. Ali, *Appl. Energy*, 2016, **161**, 225.
- 2 M. Seggiani, M. Puccini and S. Vitolo, *Int. J. Greenhouse Gas Control*, 2013, 17, 25.
- 3 M. Kato, S. Yoshikawa and K. Nakagawa, *J. Mater. Sci. Lett.*, 2002, 21, 485.
- 4 M. Kato, K. Nakagawa, K. Essaki, Y. Maezawa, S. Takeda, R. Kogo and Y. Hagiwara, *Int. J. Appl. Ceram. Technol.*, 2005, 2, 467.
- 5 K. Wang, Z. Yin and P. Zhao, Ceram. Int., 2016, 42, 2990.

6 M. E. Bretado, V. G. Velderrain, D. L. Gutiérrez, V. Collins-Martínez and A. L. Ortiz, Catal. Today, 2005, 107–108, 863.

before adsorption

1.34 1.40

Cas

7 K. Essaki and M. Kato, J. Mater. Sci., 2005, 4, 1.

 O_{c_2}

- 8 R. Rodríguez-Mosqueda and H. Pfeiffer, J. Phys. Chem., 2010, 114, 4535.
- 9 M. J. Venegas, E. Fregoso-Israel, R. Escamilla and H. Pfeiffer, *Ind. Eng. Chem. Res.*, 2007, **46**, 2407.
- 10 S. Jeoung, J. H. Lee, H. Y. Kim and H. R. Moon, *Thermochim. Acta*, 2016, **637**, 31.
- 11 I. C. Romero-Ibarra, J. Ortiz-Landeros and H. Pfeiffer, *Thermochim. Acta*, 2013, 567, 118.
- 12 V. L. Mejía-Trejo, E. Fregoso-Israel and H. Pfeiffer, *Chem. Mater.*, 2008, **20**, 7171.
- 13 G. V. Tomarov, Y. V. Petrov, A. A. Shipkov, O. A. Dovgii, V. N. Semenov and A. V. Mikhailov, *Therm. Eng.*, 2008, 55,
- 14 M. Zhao, J. Shi, X. Zhong, S. Tian, J. Blamey, J. Jiang and P. S. Fennell, *Energy Environ. Sci.*, 2014, 7, 3291.
- 15 M. Zhao, Y. Song, G. Ji and X. Zhao, Energy Fuels, 2018, 32, 5443.
- 16 Y. S. Kim and S. G. Kang, Appl. Surf. Sci., 2019, 486, 571.
- 17 A. Kumar, F. Ropital, T. de Bruin and B. Diawara, *Appl. Surf. Sci.*, 2020, **529**, 147127.
- 18 S. G. Kang, J. CO₂ Util., 2020, 42, 101293.
- 19 Y. Duan and K. Parlinski, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2011, **84**, 104113.
- 20 T. Tang, P. Chen, W. Luo, D. Luo and Y. Wang, J. Nucl. Mater., 2012, 420, 31.
- 21 X. Kong, Y. Yu, S. Ma, T. Gao, C. Xiao and X. Chen, *Chem. Phys. Lett.*, 2018, **691**, 1.
- 22 C. Qi, D. Spagnoli and A. Fourie, *Appl. Surf. Sci.*, 2020, **518**, 146255.
- 23 Q. Wang, H. Manzano, I. López-Arbeloa and X. Shen, *Minerals*, 2018, **8**, 1.
- 24 Q. Wang, F. Li, X. Shen, W. Shi, X. Li, Y. Guo, S. Xiong and Q. Zhu, *Cem. Concr. Res.*, 2014, 57, 28.
- 25 M. C. Payne, M. P. Teter, D. C. Allan, T. A. Arias and J. D. Joannopoulos, *Rev. Mod. Phys.*, 1992, **64**, 1045.
- 26 J. P. Perdew, K. Burke and M. Ernzerhof, *Phys. Rev. Lett.*, 1996, 77, 3865.

27 B. Hammer, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1999, **59**, 7413.

- 28 J. P. P. Ramalho, J. R. B. Gomes and F. Illas, *RSC Adv.*, 2013, 3, 13085.
- 29 S. Grimme, J. Antony, S. Ehrlich and H. Krieg, *J. Chem. Phys.*, 2010, **132**, 154104.
- 30 D. J. Chadi, Phys. Rev. B, 1977, 16, 1746.
- 31 B. G. Pfrommer, M. Côté, S. G. Louie and M. L. Cohen, *J. Comput. Phys.*, 1997, **131**, 233.
- 32 A. Belsky, M. Hellenbrandt, V. L. Karena and P. Lukschb, *Acta Crystallogr.*, 2002, **58**, 364.
- 33 R. Zhang, S. Ma, Q. Wang, C. Xiao, C. Zhang and T. Gao, *Ceram. Int.*, 2020, **46**, 8192.
- 34 Q. Guan, T. Gao, Y. Shen, S. Ma, T. Lu, X. Chen, C. Xiao and X. Long, *Int. J. Mod. Phys. B*, 2015, **29**, 1550128.
- 35 D. Tranqui, R. D. Shannont and H. Y. Chen, *Acta Crystallogr. B*, 1979, 35, 2479.

- 36 C. Haitao, H. Xuefei and H. Weigang, *Rare Met. Mater. Eng.*, 2018, 47, 729.
- 37 C.-J. Chan, W. M. Kriven and J. F. Young, *J. Am. Ceram. Soc.*, 1988, **71**, 713.
- 38 Y. J. Kim, I. Nettleship and W. M. Kriven, *J. Am. Ceram. Soc.*, 1992, 75, 2407.
- 39 A. V. Marenich, S. V. Jerome, C. J. Cramer and D. G. Truhlar, *J. Chem. Theory Comput.*, 2012, **8**, 527.
- 40 K. H. Jost, B. Ziemer and R. Seydel, *Acta Crystallogr.*, 1977, 33, 1696.
- 41 W. Tao, C. Zhu, Q. Xu, S. Li, X. Xiong, H. Cheng, X. Zou and X. Lu, *ACS Omega*, 2020, **5**, 20090.
- 42 E. Durgun, H. Manzano, P. V. Kumar and J. C. Grossman, *J. Phys. Chem. C*, 2014, **118**, 15214.
- 43 W.-B. Zhang, C. Chen and S.-Y. Zhang, *J. Phys. Chem. C*, 2013, 117, 21274.