



# From Euler disk to phonon bottleneck effect: Excited state physics

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The Euler disk, known since 1765,<sup>1</sup> describes a solid disk that rotates on a smooth, firm surface, similar to a coin spun by a finger on a table, before its energy dissipates and it falls on the table.<sup>2</sup> It represents a typical nonequilibrium excited state in classical physics. The most striking characteristic of an Euler disk is that it exhibits a sharply enhanced sound frequency when approaching the critical time,  $t_c$ , after which the disk instantly ceases rotation and precession. At the excited-state-to-ground-state transition, its frequency abruptly changes from infinite to zero. This quasi-divergent and suddenly disappearing behavior is intriguing, revealing that this is a shared physical property in the microscopic quantum regime.

The phonon bottleneck effect (PBE) is a well-known phenomenon in quantum materials, which is caused by the closing or sharp shrinkage of an electronic energy gap<sup>3</sup> and has various effects. In ultrafast spectroscopy, the PBE plays the same role as the "U-shaped" double peaks in scanning tunneling spectroscopy when detecting a superconducting energy gap. When the PBE occurs at a critical temperature or pressure, the slow component lifetime of the photoexcited carriers sharply increases, manifesting quasi-divergent behavior (defined as an increasing rate of exhibiting an even higher value when approaching the critical point). Simultaneously, the amplitude, proportional to the density of the photocarriers, decreases noticeably, reaching a minimum value at identical critical state variables  $T_c$  or  $P_c$ . Similar to the Euler disk frequency, the carrier lifetime undergoes a sharp change immediately after the phase transition, reaching a constant value irrespective of the state before the phase transition. The amplitude reaches a constant value immediately after the phase transition.

In this work, the peculiar quasi-divergent behaviors of the classical regime Euler disk and quantum phenomenon PBE are compared, revealing a shared physical mechanism that bridges the classical Euler disk and quantum PBE, shedding light on the excited state physics.

The rolling Euler disk in the experiment is illustrated schematically in Figure 1A. The disk is made of aluminum alloy with a radius  $r = 33.5$  mm, depth  $d = 7$  mm, and mass  $m = 37.4$  g. The Euler disk rotates fast when triggered by a finger, simultaneously precessing around the  $z$  axis with a speed of  $\Omega$  for 12 s and then suddenly decreasing in frequency to zero. A microphone was used to record the sound signal caused by the collision with the glass surface. The signal is analyzed using Fourier transformation. A series of peak curves is identified, which becomes steeper near  $t_c$  (Figure 1B). One curve is selected for quantitative analysis (Figure 1D). The energy of the Euler disk steadily decreases, down to 0 at  $t_c$  (Figure 1C). The energy becomes 0 for  $t > t_c$ , and the definition of the frequency is no longer applicable. If one considers the period to be infinite, the frequency is 0. Thus, there is no connection between the  $t < t_c$  and  $t > t_c$  regimes. The physics and model of excited state relaxation only apply to  $t < t_c$  for an Euler disk.

When approaching  $t_c$ , the Euler disk mechanical energy is given by

$$E = E_0(1 - t/t_c)^s, \quad (\text{Equation 1})$$

where  $E_0$  is the energy at  $t = 0$ , and  $s$  is a coefficient jointly determined by energy dissipation channels, including both dry and viscous contour friction.<sup>4</sup> The precession frequency is given by

$$f = f_0(1 - t/t_c)^{-s/2}, \quad (\text{Equation 2})$$

where  $f_0$  is the precession frequency at  $t = 0$ . Fitting the experimental data (red curves in Figures 1C and 1D) yields  $s = 0.564$  and  $s = 0.552$ , respectively.

A similar peculiar behavior occurs in quantum materials. Ultrafast spectroscopy has revealed that the PBE in superconductors exhibits similar behavior.

In the excited state above the Fermi surface, a microscopic quasi-equilibrium balance forms between photocarriers and phonons (Figure 1E). Photocarriers release high-frequency phonons (HFPs; defined as  $\hbar\omega_{ph} > 2\Delta(T)$ ) to combine into Cooper pairs. Conversely, HFPs can break Cooper pairs to excite photocarriers. The closed-cycle process proceeds until the HFP propagates away or decays.<sup>3</sup> Rothwarf-Taylor equations<sup>5</sup> describe these behaviors well: (1)  $dn/dt = \eta N - \beta n^2$  and (2)  $dN/dt = -\eta N/2 + \beta n^2/2 - \Lambda(N - N_T)$ , where  $\Lambda$  is the HFP decay rate.<sup>3</sup>

As the temperature approaches  $T_c$ ,  $\Delta(T)$  gradually becomes 0, strikingly enhancing the HFP density. Thus, the quasi-equilibrium is sustained, exhibiting a dramatically elongated lifetime (Figure 1G). This is the *slow* component lifetime  $\tau_{slow}$ ,<sup>3</sup> which can only be detected by ultrafast experimental means. In parallel, the quasi particle density decreases, reaching a minimum at  $T_c$ . Simultaneously observing the behaviors in  $A_{slow}$  and  $\tau_{slow}$  near  $T_c$  is necessary for identifying the transition (Figures 1F and 1G).

At temperatures near  $T_c$ ,  $A_{slow}$  and  $\tau_{slow}$  are given by<sup>3</sup>

$$A_{slow} \propto \left[ \sqrt{\Delta(T)} \exp\left(-\Delta(T)/k_B T + C\right) \right]^{-1}, \quad (\text{Equation 3})$$

and

$$\tau_{slow} \propto \left\{ \left[ \delta + 2n_T(T) \right] \left[ \Delta(T) + \alpha T \Delta(T)^4 \right] \right\}^{-1}, \quad (\text{Equation 4})$$

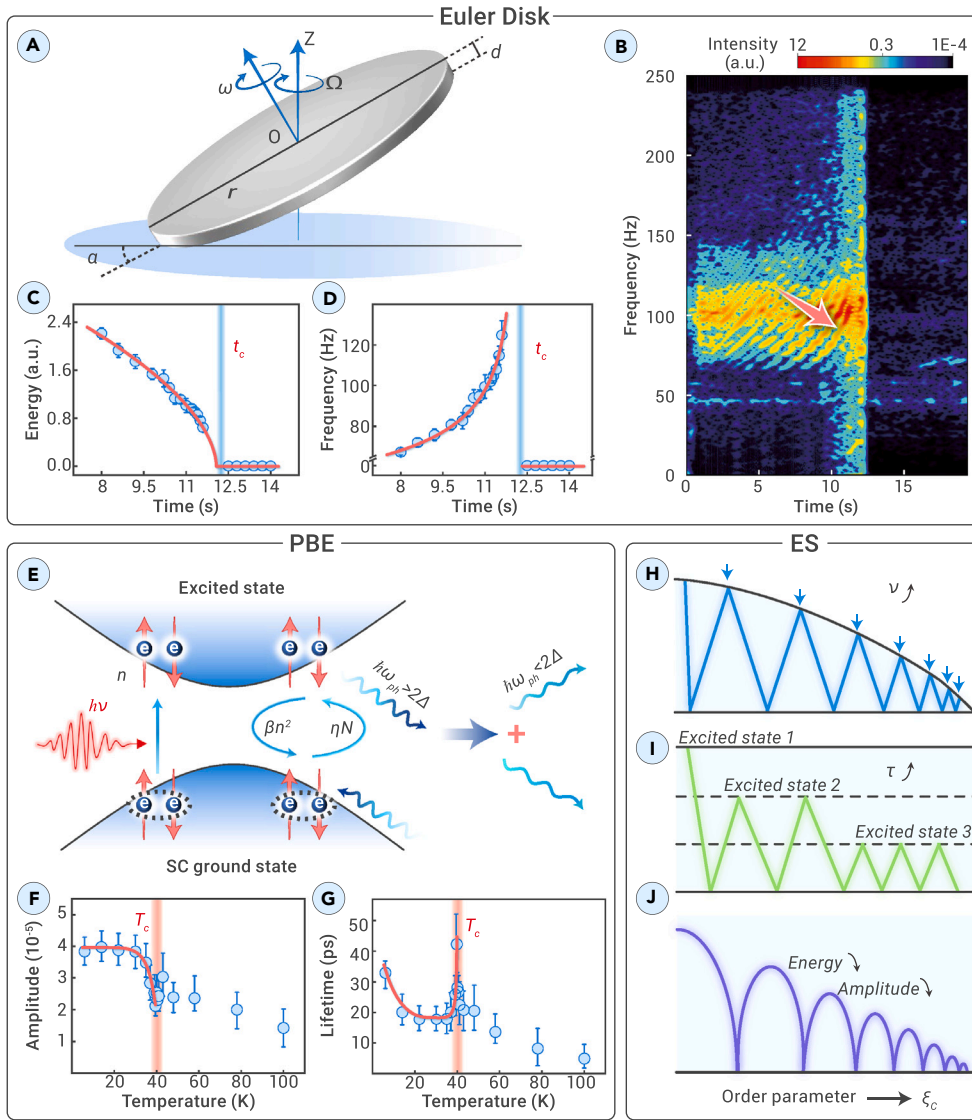
where  $C$  is related to  $\eta$  and  $\beta$ ;  $n_T$  is the density of thermal carriers,  $\delta$  and  $\alpha$ ; and  $\Delta(T) = \Delta_0 \tanh(\Theta \sqrt{T_c/T - 1})$  is the empirical SC gap,<sup>5</sup> with  $\Theta$  being related to the coupling strength. The SC  $\Delta_0$  can be obtained precisely by fitting a large number of data points. No coefficient is assumed between  $\Delta_0$  and  $T_c$ , enabling ultrafast spectroscopy to detect the strong coupling in superconductors.

The critical order parameters, peculiar characteristic behaviors, and physical processes of the Euler disk and PBE exhibit similarities. For example, phase transitions occur in both, with well-defined order parameters. Both *extensive quantities* (Euler disk order parameter energy and PBE order parameter carrier density) decrease with increasing state variable and reach their minimum value at the critical state variable. Both *intensive quantities* (Euler disk order parameter frequency and PBE order parameter lifetime) exhibit quasi-divergent enhancement. Both the Euler disk and the quantum material enter a new regime that cannot be described by excited state physics when it exceeds the critical state variable.

These similarities reflect the shared physics: both are determined by how their excited states relax to ground states, where a microscopic dynamic quasi-equilibrium exists. The key enabling gradient is the existence of the excited state. Near the transition, the gap shrinks, leading to a higher frequency of the energy exchange (Figure 1H). The lifetime is also enhanced because more carriers exist in the excited state (Figure 1I)—these are two faces of one coin. Accordingly, a similar example is proposed: bouncing a ball on a surface. During the excited-state-to-ground-state relaxation, the bouncing frequency is enhanced, and the mechanical energy gradually dissipates (Figure 1J).

The formulae describing the excited states no longer apply to systems at  $T > T_c$ . All shared properties disappear with the disappearance of the microscopic quasi-equilibrium, which only exists because of the excited state. At  $T > T_c$ , all the physical quantities are determined by their new states. The excited state is spectacular, because when it is gone, a phase transition occurs.

In conclusion, a correlation was discovered between the classical Euler disk and quantum material dynamics PBE. The Euler disk frequency and PBE lifetime increase sharply when approaching the critical point, accompanied by a decrease



**Figure 1. Comparison of the Euler disk with the PBE** (A) Schematic of a rolling Euler disk. (B) Experimentally recorded sound signal of the Euler disk from 0 to 19.38 s in the frequency domain. The color (from black to red) represents the intensity of the sound signal. (C and D) Time-dependent Euler disk energy and frequency of a typical curve marked in (B). Blue dots, experimental data; red curves, fitting curves. The error bars in (C) and (D) represent the energy fluctuation and full width at half maxima of the frequency, respectively. (E) Illustration of the microscopic equilibrium between photocarriers and high-frequency phonons (HFPs). Here,  $N$  denotes the density of HFPs and  $n$  of photocarriers,  $\eta$  and  $\beta$  are parameters, and  $\Delta$  is the superconducting gap. (F and G) The amplitude and lifetime of the quasiparticle slow component, adapted from a previous work.<sup>3</sup> (H–J) Illustration of the increasing frequency  $\nu$  (H), lifetime  $\tau$  (I), and bouncing (relaxation) of a falling ball in its excited state (J).

in their energy and quasiparticle density. Critical extensive and intensive order parameters are identified. The findings reveal that excited state physics has shared properties in the classical and quantum regimes. The Euler disk is a classical visualization of the PBE, and the PBE is a quantum-regime Euler disk. The shared properties between the Euler disk and PBE exist ubiquitously in all excited states.

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## DECLARATION OF INTERESTS

The authors declare no competing interests.