


 Cite this: *RSC Adv.*, 2024, 14, 2226

Enhanced particle separation through ultrasonically-induced microbubble streaming for automated size-selective particle depletion†

 Amirabas Bakhtiari * and Christian J. Kähler

In this study, we present an automated method for achieving Size-Selective Particle Depletion in microchannels. This technique is notable for its label-free, sheath-free, and cost-effective attributes. It combines continuous Poiseuille flow with microbubble streaming to enable the manipulation of particles in an automatic or semi-automatic manner at periodic intervals. Larger particles are retained in proximity to the microbubble, with the option for subsequent eviction through a designated waste exit or their accumulation within a collection chamber for future analysis or manipulation. Unlike many conventional methods, our approach keeps the target particles in the vortices near the microbubble while the primary fluid flows continuously through the microchannel. Subsequently, these particles are ejected in just a few milliseconds, preserving the primary fluid and significantly reducing fluid wastage. We conducted an analysis covering multiple critical facets of the study. This included a rigorous statistical examination, flow characterization using volumetric micro PTV, high-frequency micro PTV for observing flow field transitions, evaluating the system's particle trapping capabilities across different sizes with a proprietary algorithm, and investigating the z-axis distribution of both incoming and escaped particles using volumetric micro PTV. The invaluable insights gleaned from this data played a pivotal role in refining the system and optimizing its operational parameters to achieve peak efficiency across various conditions, encompassing varying particle sizes, flow rates, and seeding densities.

 Received 23rd November 2023
 Accepted 2nd January 2024

DOI: 10.1039/d3ra08038b

rsc.li/rsc-advances

1 Introduction

1.1 Overview

Continuous and precise manipulation of particles and cells in microfluidic systems holds immense significance across a wide spectrum of scientific, medical, and industrial domains. This capability not only opens up new avenues for cutting-edge research but also addresses critical challenges in various practical applications. In this context, size-selective particle removal has emerged as a pivotal technique, enabling the separation of particles based on their size and contributing significantly to fields such as biology and medicine, and industry. Its applications extend to recycling, where it facilitates the recovery of valuable metals from electronic waste, thereby contributing to sustainable resource management.^{1–3} In the mining sector, the technique is employed for the enrichment of desired minerals, improving the efficiency of mineral extraction.^{4,5} Furthermore, in wastewater treatment, the separation of particles based on size is a key component of the purification process, helping to ensure the quality of treated water.⁵

Beyond the industrial arena, particle separation is fundamental to the progress of various biomedical applications. Researchers have recognized the importance of size-selective particle removal in the isolation and analysis of rare cells, which has profound implications in the diagnosis and treatment of diseases. For instance, the isolation of rare circulating tumor cells (CTCs) from blood samples has been instrumental in cancer research, providing insights into disease progression and treatment response. This technology has become a powerful tool for early cancer detection and monitoring, as evidenced by studies in the literature.^{6,7}

Moreover, in the realm of medical diagnostics, size-selective particle removal has been applied in the development of microfluidic devices for blood analysis. These devices offer rapid and accurate blood cell counting and differentiation, contributing to advancements in point-of-care testing.^{7,8} Additionally, in the context of microbiology, the isolation of bacteria and other microorganisms based on size is critical for various applications, such as pathogen detection and environmental monitoring.⁹

Microfluidic systems have emerged as highly sensitive and precise tools for particle and cell separation at the microscale, leveraging external fields such as acoustics, electricity, magnetism, and optics to achieve superior accuracy compared to their macroscopic counterparts.^{10–14} However, when dealing with

Institute for Fluid Mechanics and Aerodynamics, Bundeswehr University Munich, 85577 Neubiberg, Germany. E-mail: amirabas.bakhtiari@unibw.de

† Electronic supplementary information (ESI) available. See DOI: <https://doi.org/10.1039/d3ra08038b>



cells and tissues, specific circumstances may arise, such as the application of high-power lasers in optical techniques or intense electric fields in electrokinetic processes, which can result in severe damage to cell membranes or disrupt natural experimental conditions. This underscores the critical need for employing non-invasive and non-destructive force fields when aiming for non-contact particle separation at the cellular level.¹⁵

As a non-invasive and biocompatible technique, acoustically actuated microbubbles can be used as a tool to selectively capture microparticles, vesicles, and biomedical cells while maintaining cell viability.¹⁶ When microbubbles are subjected to acoustic excitation at or near their resonant frequency within microchannels, they induce a primary oscillatory flow. This primary flow, in turn, initiates a secondary flow characterized by the formation of counter-rotating vortices. Precise control over these vortices can be achieved by adjusting the frequency and amplitude of the excitation signal.^{17–20} In recent years, this streaming has been successfully used in many advanced microfluidics applications for various purposes, such as sorting, mixing, focusing, positioning, and aggregating biological cells or microparticles.^{15,21–24} Microstreaming trapping of very large particles has also demonstrated the advantage of hydrodynamic manipulation over competing micro trapping techniques such as optical or dielectrophoretic trapping.^{16,25,26}

In a recent study, the size-dependent migration of microparticles in three-dimensional microbubble flows was experimentally demonstrated. It was shown that microparticles of finite size migrate with high reproducibility in a size-dependent manner and are trapped at specific spatial positions. Three regimes with qualitatively different migration behavior were observed for different particle sizes, promising applications for particle separation, trapping, or sorting using microbubble streaming.¹⁶

In this study, we present a label-free, sheath-free, and inexpensive automated method for size-dependent particle removal *via* ultrasound-driven microbubble streaming. This method is designed to selectively capture and remove microparticles from the main fluid with minimal loss of the main fluid in microchannels. By utilizing a combination of continuous Poiseuille flow and microbubble streaming, microparticles and biological cells of various sizes can be automatically separated in a two-step process. Initially, larger particles are captured in proximity to the microbubble and subsequently ejected at a pre-determined time, or gathered in a collection chamber for subsequent analysis. Our design maximizes the capacity of microbubble flow by orienting the cavity (*i.e.*, the microbubble) perpendicular to the microchannel wall. This configuration enables the flow to achieve the largest transverse distance from the bubble, thereby drawing larger particles into the core of the counter-rotating vortices. As a result, our technique does not require flow preconditioning such as flow focusing, sheath flow, or the addition of buffer flows. Furthermore, this method can be performed at higher flow rates. In contrast to many other techniques that rely on a continuous underflow of either the main fluid or buffer to remove target particles from the system, resulting in wastage of the main fluid (such as blood), our approach holds the target particles as the flow passes through

the stream and ejects them in an instant, resulting in significantly reduced waste of the main fluid. Our proposed approach uses a controllable microstreaming of microbubbles, enabling the separation of particles of different sizes within a single channel, unlike other techniques that require different channel sizes for different pairs of particles.⁹ Since there are no moving parts and the system relies solely on the microstreaming principle, the microchip system does not require any special cleaning or maintenance procedures, which increases its longevity and robustness for extended periods of operation.

Toward the study's conclusion, the study involved quantifying the statistical parameters concerning the maximum particle retention within microbubble streaming during successive activation cycles of a piezoelectric device. The purpose of these statistical analyses is to optimize the operational timeframe in a periodic mode under various conditions, encompassing distinct particle sizes, flow rates, and seeding densities. By determining the maximum number of particles of varying sizes that the microbubbles can trap, the system can be fine-tuned to accommodate specific operational scenarios effectively. This approach ensures enhanced efficiency and adaptability for diverse particle characteristics and hydrodynamic conditions.

2 Experimental setup

In this section, the procedure for preparing the solution and the experimental setup will be explained. The experimental setup comprises three essential components, namely a microfluidic system, an optical setup, and a control system. To minimize the likelihood of experimental errors, these components are mounted on a vibration-damped table.

2.1 Sample preparation

In this research, different sizes of polystyrene microspheres (2, 5, 10, and 15 μm) were utilized. These microspheres were obtained from Microparticles GmbH (Germany) and were suspended in an aqueous solution with 23.8 w-w% glycerol. The purpose was to achieve neutral buoyancy for the particles by adjusting their density to match that of the surrounding medium. To ensure uniform distribution, the microspheres were thoroughly mixed. This step was crucial to ensure random and consistent entry of the particles into the channel.

2.2 Microfluidic system

Fig. 1 depicts a schematic illustration of the experimental arrangements comprising microchannels and flow control systems. The experimental investigations were conducted within a polydimethylsiloxane (PDMS) microchannel possessing dimensions of height (H) equal to 100 μm , width (W) of 500 μm , and length (L) extending to 20 mm. This microchannel is additionally characterized by adjacent side pits having dimensions of width (w) measuring 80 μm and length (h) measuring 200 μm , situated at a distance of 300 μm from the discharge channel. The discharge channel itself possesses a width of 60 μm and is controlled with a microvalve. According to previous

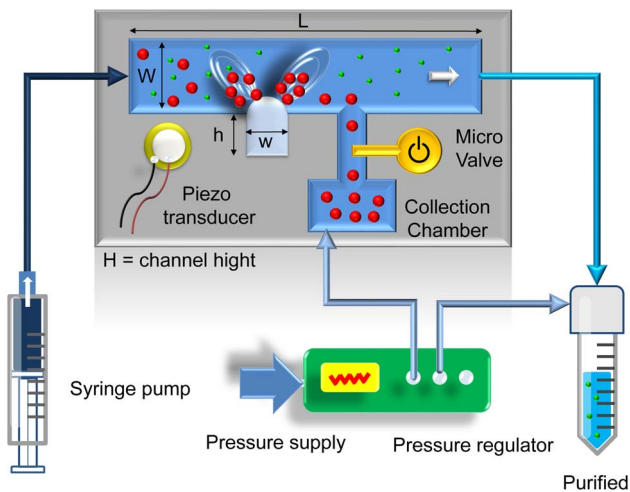


Fig. 1 Illustration of a microfluidic setup with precise control over flow rate and pressure. This system ensures stable conditions by carefully regulating flow parameters through a syringe pump and a pressure controller. It facilitates continuous and automated size-selective particle removal, allowing captured particles to be directed to a waste outlet or a designated collection chamber.

studies,^{20,22,27–29} the width of the microbubble cavity was chosen to be not too wide to avoid bubble detachment from the cavity and not too narrow to form a very small bubble, resulting in relatively weak streaming with minimal impact on the flow field.

The microchannels were made through the utilization of the conventional soft lithography technique, following a methodology similar to that detailed by Wang *et al.*³⁰ for efficient fabrication. However, in this particular investigation, we implemented the projection photolithography method suggested by Ostmann *et al.*³¹ The inclusion of this innovative photolithography technique within the soft lithography process has not only optimized fabrication but has also led to cost savings. Specifically, we accomplished this by creating a larger-scaled photomask (*e.g.*, 10 times larger) using a standard printer on transparent paper. Subsequently, the pattern on the mask is reduced to the actual size through the optical path, eliminating the requirement for a costly quartz or film photomask. Finally, the microchannel was bonded to a glass slide previously covered with a 1 mm-thick PDMS film by functionalizing the surfaces with a corona plasma treater (Elveflow, France) so that all channel walls were composed of PDMS. In order to uphold experimental reproducibility, the consistent size of microbubbles is maintained by controlling the pressure differential between the interior of the channel and the surrounding ambient pressure, as suggested by Volk *et al.*³² This control is accomplished by managing the flow of the aqueous sample solution into the channel using a syringe pump and a pressure regulator, as depicted in Fig. 1 (for further details, refer to Bakhtiari *et al.*^{15,33}).

2.3 Optical setup

The microfluidic chip was placed on a motorized stage with three-dimensional movement capabilities, which is integrated

into an upright Zeiss AxioImager.Z22 microscope (Germany). The microscope is equipped with a dichroic filter and a 10× objective lens (EC Plan Neouar 10×/0.3 M27, Zeiss, Germany) that is connected to an sCMOS camera (pco.edge 5.5, USA). Fluorescent particles from Microparticles GmbH, Germany, specifically polystyrene particles labeled with red-fluorescent dye (PS-FluoRed), were utilized in this study. The excitation and emission wavelengths for these particles were 530 nm and 607 nm, respectively. Epifluorescence microscopy was chosen over bright-field microscopy due to its superior signal-to-noise ratio (SNR). Epifluorescence microscopy involves illuminating fluorescent particles using a continuous-wave laser or high-power LED along the optical path, while the camera captures the emitted fluorescent light from the particles. Subsequently, particle detection is achieved through straightforward filtering and thresholding techniques.

2.4 Control system setup

Lab-VIEW (National Instrument, USA) is used to simultaneously control the entire process of image acquisition, particle detection, particle counting, and switching the output channel to discharge the desired particles in live mode by a specially developed routine. The control system consists of a function generator (GW INSTRUK AFG-2225, Good Will Instrument Co., Ltd, Taiwan), an amplifier (Krohn Hite 7500, USA), and an oscilloscope (Teledyne LeCroy HDO6104, Germany) to send the predefined electrical signal to the piezoelectric transducer. A National Instruments data acquisition device (DAQmx USB-6002, USA) controls the output channel by switching an electrical microvalve (see Fig. 2).

2.4.1 Method. It was shown that the position of particles in a steady microbubble streaming strongly depends on the

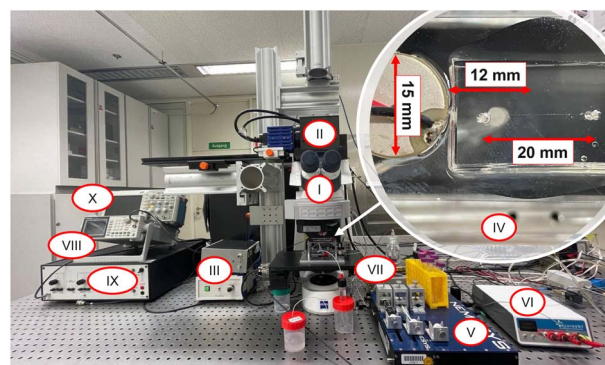


Fig. 2 The experimental configuration comprises three main components: a microfluidic system, an optical arrangement, and a control system. The optical setup includes an upright microscope equipped with a 10× objective lens (I), a camera (II), and light sources (III). Within the microfluidic setup, there is a microchip housing a piezoelectric transducer, situated 12 mm away from a microbubble cavity within a 20 mm long microchannel (IV). The flow in the channel is controlled by a syringe pump (V), and a pressure controller (VI) is employed to maintain bubble size equilibrium. Flow to the exit or collecting chamber is regulated by an electric microvalve (VII). The piezoelectric element is activated through a function generator (VIII), an amplifier (IX), and an oscilloscope (X) under the control of LabVIEW.

particle size. While sufficiently small particles behave as passive tracers and follow the streamlines, larger particles become entrapped in more rapid and smaller loops located in close proximity to the microbubble's surface.¹⁶ Here, we extract the larger particles from a two-phase flow with particles of different sizes by a superposition of Poiseuille flow and microbubble flow, which is generated automatically or periodically by the activation and deactivation of piezoelectric elements. This approach employs microstreaming to induce the entrainment and confinement of larger particles in the proximity of the microbubble surface across the entire width of the microchannel while allowing smaller particles to escape through the counter-rotating vortices and remain in the Poiseuille flow towards the outlet. Upon attaining a desired concentration of larger particles or when space limitations arise, the piezoelectric is deactivated, propelling the cloud of larger particles downstream by the main flow near the wall, where a side channel is automatically opened at a distance of 300 μm and promptly closed again to discharge the larger particles into a waste or collection chamber for further investigation. The operational duration, microstreaming activation, and the time between opening and closing of the exit channel depend on the main flow rate, the concentration of larger particles, and channel geometry, thus allowing adaptation of the system as periodic, manual, or automatic, based on the conditions. This study evaluates critical parameters, such as the maximum number of particles trapped per run, and the effects of particle size and mainstream velocity on this limit.

2.4.2 Advanced control system. We have formulated algorithms using LabVIEW that possess diverse functionalities to monitor and control the system while it operates continuously in live mode. Maintaining constant microstreaming is crucial for achieving optimum efficacy and guaranteeing the replicability of outcomes. This is accomplished by controlling the microbubble penetration depth and preserving its stability over prolonged durations through the use of a control system. This system simultaneously regulates the pressure differential

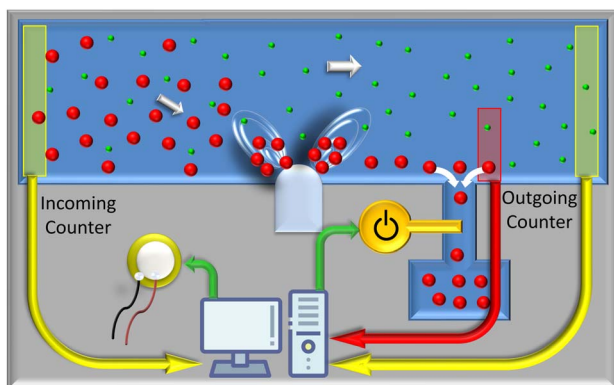


Fig. 3 Automated particle collection principle: our custom LabVIEW control system manages various functions, including particle counting, piezo transducer and microvalve control, pressure regulation, and flow management. The yellow regions of interest (ROIs) are responsible for particle counting, while the red ROI serves as a sensor to detect particle leaks and trigger valve activation.

between the interior of the microchannel and its surrounding environment, which is approximately 6–8 mbar in this particular investigation (see in ref. 32). To ensure the efficient removal of larger particles with minimal fluid waste, a virtual detector (measuring 20×200 micrometers) is placed downstream of the outlet (Fig. 3). The detector is set to detect the presence of larger particles in a predefined region of interest (ROI) as soon as they pass through. Once a larger particle is detected, the outlet is opened for a short period of time (1 second), allowing for rapid removal of the particle to the waste or collector chamber. The quantification of confined particles is achieved through live mode by employing analogous regions of interest (ROIs) to automatically detect and tally larger particles entering and exiting the upstream and downstream sections of the microbubble. The ROIs are narrow in dimension (20 μm) and extend across the entire width of the microchannel, with the ability to exclusively identify and quantify particles that exceed a specific size threshold (10 μm). In Fig. 4, a comprehensive illustration of the size-selective particle removal process in live mode is presented. During this process, a blend of microparticles is injected into the microchannel. This mixture comprises a significantly higher quantity of 2 micrometer particles

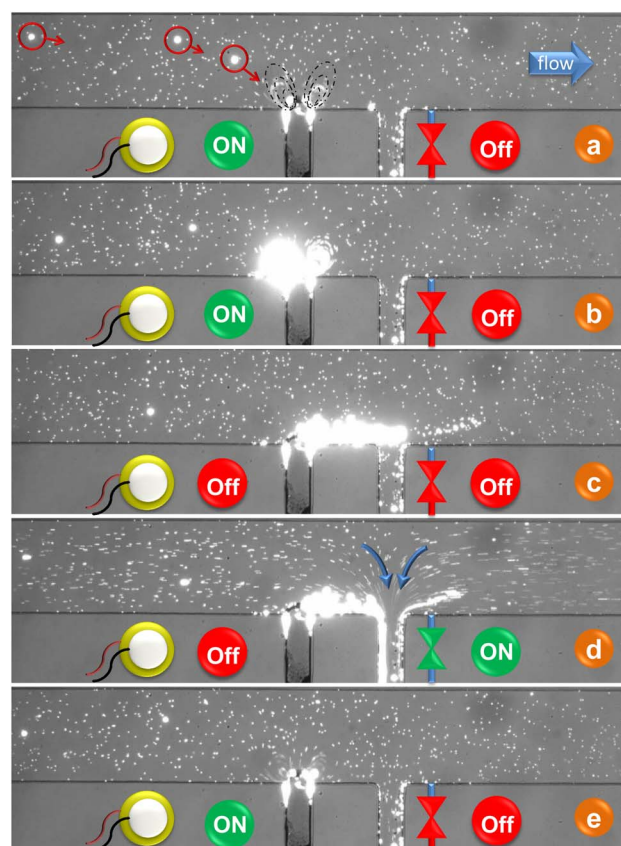


Fig. 4 Sequential visualization of size-selective particle removal in the microchannel: (a) entrapment of 10 micrometer particles by counter-rotating vortices, (b) augmentation of larger particle concentration, (c) deactivation of the piezo transducer positions larger particles near the exit channel, (d) controlled microvalve opening initiates swift particle suction, (e) repetition of these steps for continuous operation.

compared to 10 micrometer particles, with a concentration ratio of 200 to 1, respectively. The flow direction is from the left to the right. In the initial step, as shown in Fig. 4a, the piezo transducer is activated while the discharge valve remains closed. Notably, this action generates counter-rotating vortices that draw the particles in close proximity to the microbubble. The vortices effectively confine the 10 micrometer particles, allowing them to remain trapped, while the 2 micrometer particles can freely pass through the channel without any obstruction. Subsequently, in Fig. 4b, the system is activated to enhance the concentration of the larger particles. This step is heavily reliant on the main flow velocity and the seeding ratio between the larger and smaller particles. Moving on to Fig. 4c, the piezo transducer is deactivated, allowing approximately four seconds for larger particles to reach the correct position adjacent to the exit channel. At the precise moment when the larger particles are positioned near the disposal channel, the microvalve is opened briefly, typically for a duration of 1 second. This controlled opening creates a pressure differential that induces a rapid suction of the particles into the exit channel. The level of this pressure difference can be finely adjusted using a pressure controller, granting precise control over the suction flow to the exit channel. Finally, this entire cycle repeats continuously, as depicted in Fig. 4e. Operational programming provides users with a wide range of choices to customize the system to suit various conditions. The system can function without the need for a visual interface, relying solely on predefined and adjustable time intervals. These intervals are determined based on factors such as flow rate, particle ratio, or the desired level of filtering or separation precision.

3 Results and discussion

3.1 Flow characterization

Characterizing fluid flow under various conditions is imperative for precise system programming, ensuring optimal operational efficiency. In the initial phase, we conducted a General Defocusing Particle Tracking (GDPT)³⁴ analysis to characterize the combination of microbubble streaming and Poiseuille flow. To this end, we tracked 15 micrometer particles within a combined flow involving Poiseuille flow (left to right) and microbubble streaming generated by the continuous oscillation of a piezoelectric transducer, with varied voltage amplitudes of 70, 35, and 20 V_{pp} at a frequency of 33 kHz. The selection of 15 micrometer particles was deliberate, aligning with their critical size and higher Stokes number compared to their 2, 5, and 10 micrometer counterparts. This choice ensures that successful manipulation of the larger particles towards the microbubble logically extends to the smaller particles (2, 5, and 10 micrometers) with lower Stokes numbers, implying similar or faster dragging toward the bubble. As depicted in Fig. 5, the flow field is partitioned into distinct flow topologies by the combination of microbubble streaming and Poiseuille flow, resulting in two primary regions separated by a dividing line or separatrix. The streamlines and particles located above this separatrix line bypass the counter-rotating vortices and exit downstream of the bubble at nearly the same level as upstream of the bubble. On

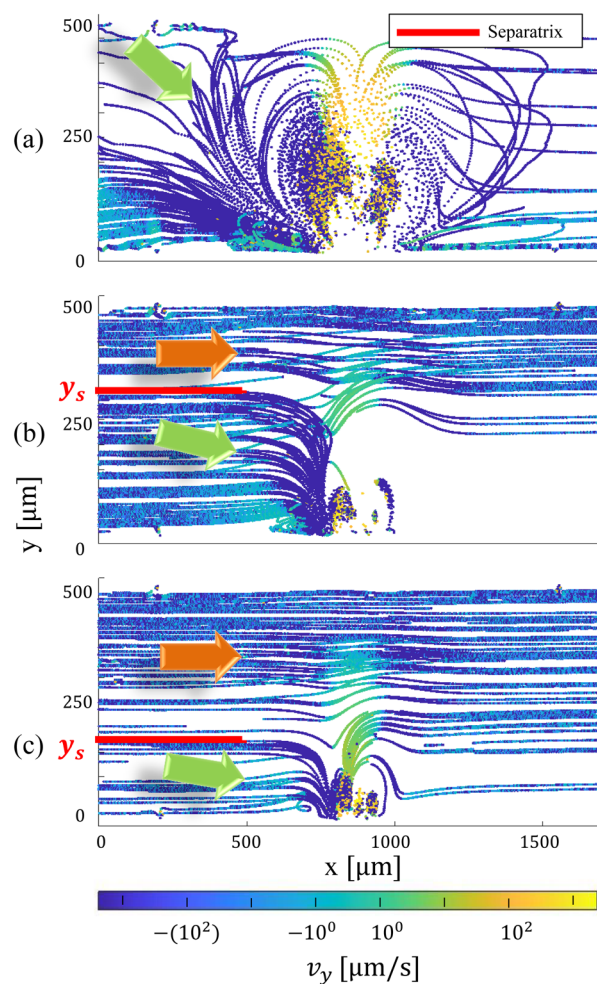


Fig. 5 Tracking 15 micrometer particles within a Poiseuille flow (left to right) accompanied by microbubble streaming induced through continuous excitation of the piezoelectric transducer at frequencies of 33 kHz. Cases (a)–(c) correspond to excitation amplitudes of 70, 35, and 20 V_{pp} respectively. In this analysis, it is evident that case (a) has the capability to attract and concentrate incoming particles near the microbubble, whereas cases (b) and (c) exhibit a decrease in amplitude, creating a bypass for particles situated above the red separatrix line where $y_{sc} < y_{sb}$ (y_s represents the y -value of the separatrix line).

the contrary, particles located beneath the separatrix ingress into the perigee, a confined region characterized by a minimal separation between the bubble's exterior surface and the separatrix. For a given geometry, protrusion depth of the bubble, and resonant frequency, an increase in mean Poiseuille velocity causes the separatrix to move downward towards the bubble, while an increase in the amplitude of the bubble excitation enlarges the size of the vortices, which in turn levels the separatrix towards the wall opposite the bubble. On the other hand, when the relative strengths of Poiseuille flow and microbubble streaming are increased, it leads to an augmentation in the perigee, which is the closest point to the bubble interface.²¹ This perigee value plays a decisive role in the migration of particles toward the inner torii of the counter-rotating vortices. In essence, particles that exceed the perigee size are subsequently guided and confined within the inner torus.¹⁶

In order to enhance particle entrapment efficacy, it is imperative to achieve an optimal proportion between the Poiseuille flow rate and the magnitude of microstreaming. This equilibrium ensures that the resultant vortices adequately cover the complete width of the microchannel and makes the perigee sufficiently minimized to facilitate the effective capture of focused particles. Since particles larger than 2 μm encounter greater inertial forces, the research endeavors were replicated for particle dimensions measuring 5, 10, and 15 μm . This was undertaken to validate the notion that particles situated near the far wall of the microbubble could still be effectively drawn towards the perigee. The results of these experimental iterations vividly illustrate that, with an average Poiseuille flow rate of 200 $\mu\text{m s}^{-1}$, the microstreaming induced by the excited piezo transducer—operating at an amplitude of 66 V_{pp} and a frequency of 33 kHz within the microchannel—exhibits the ability to aggregate particles of all sizes towards the perigee. Notably, it succeeds in ensnaring particles measuring 5, 10, and 15 μm within the inner torii.

In cases where it is challenging to enhance amplification due to bubble detachment or when the influence is insufficient to overcome the dominance of Poiseuille flow, alternative strategies may be explored. These alternatives include employing deeper microchannels at constant flow velocities through the use of higher flow rates or introducing multiple microbubbles on opposing channel walls. These approaches enable the system to effectively function in wider microchannels, offering potential solutions to overcome operational limitations.

3.2 High-frequency μPTV

Under specific conditions, especially when target particles are rare, continuous system activation may not be necessary. Instead, the system can be designed to initiate its operation automatically upon detecting a single particle within the first Region of Interest (ROI). To ensure the effective detection of particles and minimize the risk of missing them, precise placement of the ROI within the camera's field of view, along with a thorough understanding of the response time between activating the piezoelectric element and redirecting the particle, becomes crucial. In this investigation, we employ high-frequency microparticle tracking velocimetry (μPTV) to scrutinize the dynamic shift in the flow field, transitioning from pure Poiseuille flow to the combination of Poiseuille flow (left to right) along with the introduction of microbubble streaming upon activation and deactivation of a piezoelectric element. High-Frequency Micro-Particle Tracking Velocimetry (μPTV) at 24 000 frames per second was conducted to track 2 micrometer particles in this combined flow field. The experimental setup, as described in 2.3, was utilized for high-frequency (μPTV), with the exception of employing a Phantom v2640 ONYX high-speed camera operated by Davis 10 (LaVision, Germany). Subsequently, the GDPT method was applied for post-processing the images. Fig. 6 offers a visual representation of our high-frequency μPTV analysis, spanning 3000 frames (lasting 125 ms, depicted in black) before and 10 000 frames (416 ms, illustrated in color) after initiating bubble excitation *via* the piezoelectric element.

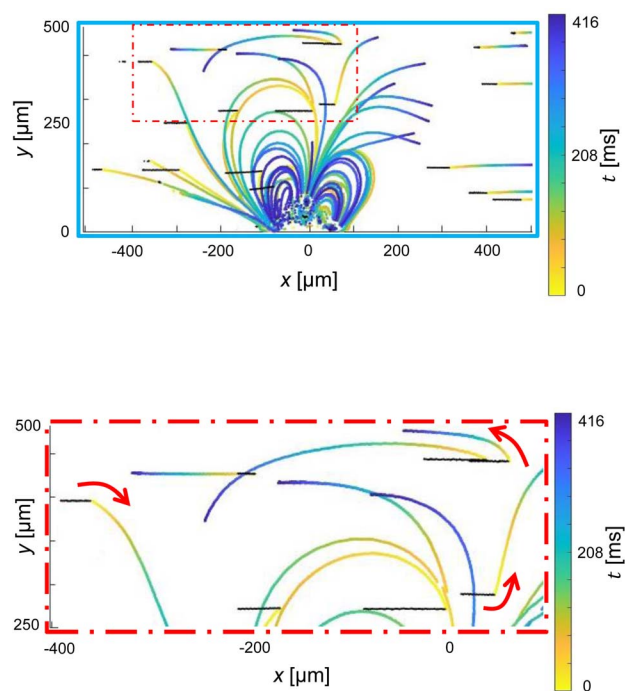


Fig. 6 High-Frequency Micro-Particle Tracking Velocimetry (μPTV) at 24 000 frames per second: tracking 2-micrometer particles in the presence of Poiseuille flow (left to right) and microbubble streaming. (a) Presents the entire experimental field of view, while (b) provides an enlarged view (indicated by the red dashed line rectangle) to enhance the visibility of finer details.

As illustrated in Fig. 6a, the immediate alteration in particle displacements becomes evident after the initial cycle of bubble activation. Although the magnitude of these changes is smaller for particles located farther from the bubble's center (e.g., particles with $x_p < -400$ and $x_p > +400$ μm), they all initiate promptly, irrespective of their proximity to the vortex core. A closer look in Fig. 6b reveals that the newly adopted trajectories, categorized by their initial positions, exhibit perfect synchronization, devoid of any temporal delay. This tracking behavior remains consistent over an extended period, confirming the absence of any transition delay and the instantaneous establishment of the superposition pattern. Similarly, when transitioning from the superposition flow to pure Poiseuille flow, we observe an immediate cessation of lateral particle movement without delay. The particles maintain their respective lateral positions until they reach the end of the microchannel.

3.3 Particle trapping capacity

Fig. 7 illustrates the outcomes of particle trapping efficiency across varying particle sizes (5, 10, and 15 μm), using a consistent microbubble excitation frequency (33 kHz, 66 V_{pp}) and main flow velocity (200 $\mu\text{m s}^{-1}$). The quantification of entrapped particles was accomplished by tracking inflow and outflow. This involved counting particles entering upstream of the microchannel and subtracting those exiting downstream of the bubble. The best trapping performance was observed for 10

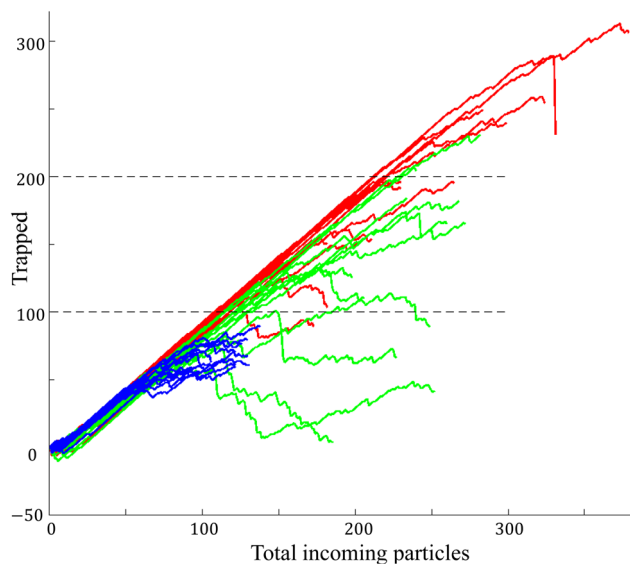


Fig. 7 The visual representation illustrates the efficiency of particle trapping as it relates to particle size (5, 10, and 15 μm) when subjected to uniform microbubble excitation at a frequency of 33 kHz with an amplitude of 66 V_{pp} . The experiment maintains a constant main flow velocity of 200 $\mu\text{m s}^{-1}$. Each line in the graph corresponds to a single experimental run. Notably, the trapping performance is significantly superior for 10 micrometer particles (depicted in red) when compared to both 5 micrometer particles (green) and 15 micrometer particles (blue).

micrometer particles with a success rate of 100 particles and negligible leakage. However, as the number of particles entering increased, the leakage rate gradually rose. This continued until an equilibrium was reached between entering and exiting particles. In some instances, higher particle concentrations led to notable particle–particle interactions within the trapping vortex, causing sudden leakage. Fig. 7 presents the highest capacity achievable for clean trapping, devoid of particle leakage. This capacity decreased to 70 particles for 5 micrometer particles and 60 particles for 15 micrometer particles. We identified two primary factors contributing to this reduction in efficiency for particles larger or smaller than 10 micrometers. Firstly, larger particles occupied space within the torus (vortex) at a faster pace, promoting quicker particle–particle interactions. Secondly, the proximity of 5 micrometer particles to the critical perigee size resulted in an increased likelihood of leakage as trapped particle concentration intensified within the vortices.

3.4 Leaked particles: Z-directional distribution analysis

To gain deeper insight into the particle leakage phenomenon, we performed comprehensive volumetric microparticle tracking velocimetry (PTV) using defocus imaging.³⁵ This allowed us to accurately track the precise positions of particles in the z-axis as they exited the vortices and traversed the microchannel. In earlier stages of our research, both numerical simulations and experimental observations led to intriguing results: during the trapping phase, characterized by the absence of Poiseuille flow and particle leakage, larger particles consistently exhibited

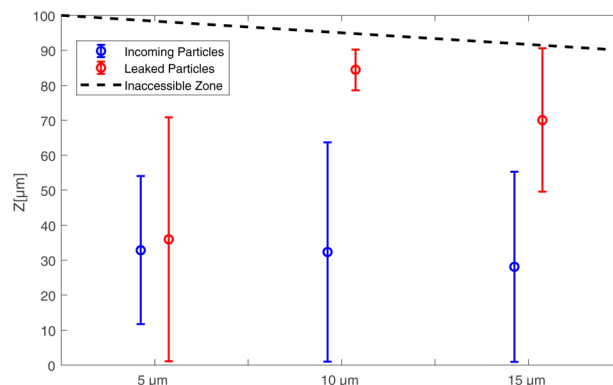


Fig. 8 This figure presents the mean and standard deviation of the z-axis distribution for incoming and leaked particles with diameters of 5, 10, and 15 micrometers. A dashed black line originating from the microchannel's bottom wall delineates the region inaccessible to particles due to their finite size. Incoming particles (depicted in blue) demonstrate a uniform z-direction distribution. Conversely, the 10 and 15 micrometer leaked particles (illustrated in red) exhibit a distinctive clustering pattern near the channel wall. In contrast, the 5 micrometer particles disperse evenly across the entire channel depth, without displaying the wall-focused behavior observed in larger particles.

a distinct migration pattern along the z-axis. This rapid migration occurred on the order of milliseconds and was strongly influenced by both particle size and flow velocity. The results of the migration of particles with diameters of 5, 10, and 15 μm in the microbubble flow show that 15 micrometer particles are focused closer to the wall, similar to 10 micrometer particles, and closer than 5 micrometer particles.¹⁶ Fig. 8 illustrates volumetric micro-Particle Tracking Velocimetry (micro-PTV) results pertaining to the z-axis spatial distribution of particles with diameters of 5, 10, and 15 micrometers. These measurements were obtained while subjecting the system to uniform microbubble excitation at a frequency of 33 kHz, with an amplitude of 66 V_{pp} , and maintaining a consistent main flow velocity of 200 $\mu\text{m s}^{-1}$. While incoming particles are initially uniformly distributed in both the y and z directions, we observe a distinct behavior. The 10 and 15 micrometer leaked particles exhibit a focused clustering near the channel wall, with a standard deviation of approximately 10–20 μm . In contrast, 5 micrometer particles do not display this wall-focused behavior; instead, they disperse throughout the channel depth. These observations suggest that during the particle trapping phase, the larger diameter particles, 15 and 10 μm , experience a more significant shear-induced lift force compared to wall-induced lift force. Consequently, they have a greater propensity to position themselves closer to the nearest channel wall, resulting in a common exit point along the z-axis. Furthermore, this exit process is faster for 15 micrometer particles than for 10 micrometer particles due to the quicker filling of the available space, which promotes faster particle–particle interactions.

4 Conclusions

In conclusion, this study has presented and validated an automated method for achieving Size-Selective Particle Depletion in

continuous microchannel flow. This method stands out for its label-free, sheath-free, and cost-effective nature. Particle separation and removal are accomplished using a PDMS microchannel and piezoelectric technology, offering a simple yet effective alternative to expensive devices like flow cytometers. The cost-effectiveness of microchip fabrication is further improved through the use of negative magnification during soft lithography. Moreover, the technique relies on actuated microbubbles without involving moving parts, resulting in a key advantage of minimal maintenance procedures. Notably, it effectively retains larger particles near microbubbles without necessitating the use of buffer flow, offering the option for subsequent removal or collection. Unlike conventional techniques, our approach keeps target particles within vortices near microbubbles while the primary fluid continues to flow through the microchannel. Consequently, particle ejection occurs rapidly, minimizing fluid wastage and preserving the primary fluid. Our analysis encompassed various critical aspects, including rigorous statistical examination, volumetric micro PTV flow characterization, high-frequency micro PTV for flow field observations, evaluation of particle trapping capabilities across different sizes with a proprietary algorithm, and exploration of z-axis particle distribution using volumetric micro PTV. The insights gained from this data have been instrumental in refining the system and optimizing its operational parameters to achieve peak efficiency under diverse conditions, accommodating varying particle sizes, flow rates, and seeding densities. Notably, our system's flexibility allows for customization to meet practical needs, accommodating different flow conditions, particle sizes, and particle distribution rates. This adaptability ensures its applicability in a wide range of scenarios, underscoring its potential for widespread utility and innovation in the field of microfluidics and particle manipulation.

Author contributions

Amirabas Bakhtiari: conceptualization, methodology, data acquisition and analysis, original draft writing. Christian J. Kähler: conceptualization, methodology, supervision, review and editing, funding acquisition.

Conflicts of interest

There are no conflicts to declare.

Acknowledgements

We acknowledge financial support by Universität der Bundeswehr München.

Notes and references

- 1 S. Louhadj, N. Hammadi, S. Touhami, H. Louati, A. Hadjali, I.-E. Kimi and A. Tilmatine, *J. Electrostat.*, 2020, **105**, 103448.
- 2 F. Wang, Y. Zhao, T. Zhang, G. Zhang, X. Yang, Y. He, L. Wang and C. Duan, *J. Cleaner Prod.*, 2017, **165**, 452–457.
- 3 T. Spengler, M. Ploog and M. Schröter, *OR Spectrum*, 2003, **25**, 413–442.
- 4 L. Hansen, A. Wollmann, M. Weers, B. Benker and A. P. Weber, *Chem. Eng. Technol.*, 2020, **43**, 933–941.
- 5 E. N. Peleka, G. P. Gallios and K. A. Matis, *J. Chem. Technol. Biotechnol.*, 2018, **93**, 615–623.
- 6 R. T. Krivacic, A. Ladanyi, D. N. Curry, H. Hsieh, P. Kuhn, D. E. Bergsrud, J. F. Kepros, T. Barbera, M. Y. Ho, L. B. Chen, et al., *Proc. Natl. Acad. Sci. U. S. A.*, 2004, **101**, 10501–10504.
- 7 M. Toner and D. Irimia, *Annu. Rev. Biomed. Eng.*, 2005, **7**, 77–103.
- 8 M. E. Vincent, W. Liu, E. B. Haney and R. F. Ismagilov, *Chem. Soc. Rev.*, 2010, **39**, 974–984.
- 9 C. Liu, C. Xue, X. Chen, L. Shan, Y. Tian and G. Hu, *Anal. Chem.*, 2015, **87**, 6041–6048.
- 10 M. J. Kim, D. J. Lee, J. R. Youn and Y. S. Song, *RSC Adv.*, 2016, **6**, 32090–32097.
- 11 P. Deng, C.-J. Fu and Z. Wu, *RSC Adv.*, 2018, **8**, 35512–35520.
- 12 C. Devendran, N. R. Gunasekara, D. J. Collins and A. Neild, *RSC Adv.*, 2016, **6**, 5856–5864.
- 13 D. Yuan, S. H. Tan, Q. Zhao, S. Yan, R. Sluyter, N.-T. Nguyen, J. Zhang and W. Li, *RSC Adv.*, 2017, **7**, 3461–3469.
- 14 C. W. Shields IV, C. D. Reyes and G. P. López, *Lab Chip*, 2015, **15**, 1230–1249.
- 15 A. Bakhtiari and C. J. Kähler, *Microfluid. Nanofluid.*, 2022, **26**, 1–11.
- 16 A. Volk, M. Rossi, B. Rallabandi, C. J. Kähler, S. Hilgenfeldt and A. Marin, *Physical Review Fluids*, 2020, **5**, 114201.
- 17 N. Riley, *Annu. Rev. Fluid. Mech.*, 2001, **33**, 43–65.
- 18 P. Marmottant and S. Hilgenfeldt, *Nature*, 2003, **423**, 153–156.
- 19 M. Versluis, D. E. Goertz, P. Palanchon, I. L. Heitman, S. M. van der Meer, B. Dollet, N. de Jong and D. Lohse, *Phys. Rev. E: Stat., Nonlinear, Soft Matter Phys.*, 2010, **82**, 026321.
- 20 B. Rallabandi, C. Wang and S. Hilgenfeldt, *J. Fluid Mech.*, 2014, **739**, 57.
- 21 R. Thameem, B. Rallabandi and S. Hilgenfeldt, *Biomicrofluidics*, 2016, **10**, 014124.
- 22 C. Wang, B. Rallabandi and S. Hilgenfeldt, *Phys. Fluids*, 2013, **25**, 022002.
- 23 C. Wang, S. Jalikop and S. Hilgenfeldt, *arXiv*, 2010, preprint, arXiv:1010.3290, DOI: [10.48550/arXiv.1010.3290](https://doi.org/10.48550/arXiv.1010.3290).
- 24 Y. Gao, M. Wu, Q. Luan, I. Papautsky and J. Xu, *Lab Chip*, 2022, **22**, 805–813.
- 25 B. R. Lutz, J. Chen and D. T. Schwartz, *Anal. Chem.*, 2006, **78**, 5429–5435.
- 26 X.-B. Wang, Y. Huang, P. R. Gascoyne and F. F. Becker, *IEEE Trans. Ind. Appl.*, 1997, **33**, 660–669.
- 27 D. Ahmed, A. Ozcelik, N. Bojanala, N. Nama, A. Upadhyay, Y. Chen, W. Hanna-Rose and T. J. Huang, *Nat. Commun.*, 2016, **7**, 1–11.
- 28 A. R. Tovar and A. P. Lee, *Lab Chip*, 2009, **9**, 41–43.
- 29 D. Ahmed, X. Mao, B. K. Juluri and T. J. Huang, *Microfluid. Nanofluid.*, 2009, **7**, 727.
- 30 C. Wang, S. V. Jalikop and S. Hilgenfeldt, *Biomicrofluidics*, 2012, **6**, 012801.

- 31 S. Ostmann and C. J. Kähler, *Microfluid. Nanofluid.*, 2022, **26**, 24.
- 32 A. Volk, M. Rossi, C. J. Kähler, S. Hilgenfeldt and A. Marin, *Lab Chip*, 2015, **15**, 4607–4613.
- 33 A. Bakhtiari and C. J. Kähler, *Microfluid. Nanofluid.*, 2023, **27**, 1–9.
- 34 R. Barnkob, C. J. Kähler and M. Rossi, *Lab Chip*, 2015, **15**, 3556–3560.
- 35 M. Wu, J. W. Roberts and M. Buckley, *Exp. Fluids*, 2005, **38**, 461–465.