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A visible-light phototransistor based on the heterostructure of ZnO and TiO₂ with trap-assisted photocurrent generation†

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Visible-light phototransistors have been fabricated based on the heterojunction of zinc oxide (ZnO) and titanium oxide (TiO₂). A thin layer of TiO₂ was deposited onto the spin-coated ZnO film *via* atomic layer deposition (ALD). The electrical characteristics of the TiO₂ layer were optimized by controlling the purge time of titanium isopropoxide (TTIP). The optimized TiO₂ layer could absorb the visible-light from the sub-gap states near the conduction band of TiO₂, which was confirmed *via* photoelectron spectroscopy measurements. Therefore, the heterostructure of TiO₂/ZnO can absorb and generate photocurrent under visible light illumination. The oxygen-related-states were investigated *via* X-ray photoelectron spectroscopy (XPS), and the interfacial band structure between TiO₂ and ZnO was evaluated *via* ultraviolet photoelectron spectroscopy (UPS). Oxygen-related states and subgap-states were observed, which could be used to generate photocurrent by absorbing visible light, even with TiO₂ and ZnO having a wide bandgap. The optimized TiO₂/ZnO visible-light phototransistor showed a photoresponsivity of 99.3 A W⁻¹ and photosensitivity of 1.5 × 10⁵ under the illumination of 520 nm wavelength light. This study provides a useful way to fabricate a visible-light phototransistor based on the heterostructure of wide bandgap oxide semiconductors.

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Introduction

Thin-film transistors (TFTs) based on oxide semiconductors have been commercialized in numerous electronics industries due to excellent electrical characteristics, such as high field effect mobility, low off-current, and high on/off ratio. Oxide semiconductors can be used as a large-scale, flexible, and highly transparent electronics.^{1–5} Furthermore, oxide semiconductors have been a potential candidate for an active channel material of photosensors.^{6–8} For these reasons, research on UV phototransistors based on oxide semiconductors has been conducted.^{9–11} The device with a wide bandgap oxide semiconductor can generate a photocurrent under the illumination of high photon energy, such as UV light. However, it is hard to absorb the low

photon energy, such as visible light, and generate a photocurrent due to the wide bandgap (>3.0 eV).

To generate photocurrent under visible light illumination using an oxide semiconductor, it is essential to introduce additional absorbing layers on the device, such as two-dimensional materials, metal nanoparticles, organic materials, and quantum dots.^{12–18} However, heterojunction structure phototransistors have drawbacks such as poor uniformity, interface traps, and defect sites between other absorbing layers and oxide semiconductors.¹⁹ In particular, traps at the interface deteriorate electrical characteristics with high-off current, low on/off ratio, and roll-off effects.²⁰ To resolve these problems, heterostructures based on oxide semiconductors have been suggested instead of organic and other nanomaterials on the wide bandgap oxide semiconductor.^{21,22}

In this study, we introduced a heterojunction structure based on the oxide semiconductors of titanium oxide (TiO₂) and zinc oxide (ZnO) to fabricate a visible-light phototransistor. A TiO₂ layer was deposited onto a spin-coated ZnO *via* atomic layer deposition (ALD) under various purge time periods of titanium isopropoxide (TTIP) to find an optimal absorbing TiO₂ layer. The optimized TiO₂ layer on ZnO could absorb the visible light and generate photo-excited electrons from oxygen-related-states *via* trap-assisted photoelectron generation. These photo-excited electrons could be transferred from TiO₂ into the

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active channel of ZnO. To investigate the origin of the trap-assisted photoelectron generation, X-ray photoelectron spectroscopy (XPS) and ultraviolet photoelectron spectroscopy (UPS) measurements were conducted. Therefore, we suggest an efficient heterostructure composed of TiO₂/ZnO for the visible-light phototransistor based on oxide semiconductors.

Result and discussion

Fig. 1a shows a schematic of the visible-light phototransistor with the heterostructure composed of TiO₂/ZnO on the SiO₂/Si substrate and the ALD process of TiO₂. To optimize the process-condition of the TiO₂ layer as a stable photo-active channel, process controls for the optimal sub-gap states of TiO₂ were conducted according to different purge time periods (0.5, 10, and 20 s) of titanium isopropoxide (TTIP). Table S1† summarizes the saturation field-effect mobilities and electron concentrations. The field effect mobility of TiO₂/ZnO TFTs was higher than that of the ZnO TFT due to the bilayer of TiO₂/ZnO.²³ Also, electron concentration shows a similar tendency with the field effect mobility. Transfer curves of TiO₂/ZnO TFTs under various purge time periods of TTIP were measured to confirm the electrical characteristics, as shown in Fig. 1b. TiO₂ (purge_{TTIP} 0.5 s)/ZnO TFT showed poor transfer curve characteristics since the purge process time of TTIP was short and organic residues composed of the carbon groups (C) remained

on the TiO₂ film during the ALD-process.²³ These organic materials cause a current path with an increase in the off current at a negative voltage range and $\sim 10^3$ on/off ratio.^{21,24,25} In addition, to determine the origin of the unstable transfer curve characteristics of TiO₂/ZnO TFTs, the chemical states of TiO₂ films were analyzed with various purge time periods of ALD-TiO₂ (0.5, 10, and 20 s). Fourier transform infrared (FT-IR) spectroscopy was performed to investigate the chemical states related to the organic residues of TiO₂ on ZnO films. As shown in Fig. S1,† peaks were observed at 1375 and 1019 cm⁻¹, which were attributed to the chemical bonds of TiO₂ (-CH₃ and C-O, respectively).²⁶ With the decrease in the purge time of ALD-TiO₂, the intensity of the peaks increased due to the extra carbon groups of TiO₂ during the ALD-process. In particular, an increased intensity of the TiO₂ (purge_{TTIP} 0.5 s) film was observed since excess carbon groups were left on the TiO₂. However, TiO₂ (purge_{TTIP} 10 s and 20 s)/ZnO TFTs showed relatively stable transfer characteristics with low leakage current ($\sim 10^{-11}$ A) at the negative voltages and 10^6 on/off ratio compared to the TiO₂ (purge_{TTIP} 0.5 s)/ZnO TFT. These were the reasons why the ligands composed of carbon groups easily broke loose from TTIP, and the ratio of the bonds between Ti and O₃ increased in 10 and 20 s of the TTIP_{purge} time. Therefore, TiO₂ (>purge_{TTIP} 10 s) films are required to decrease the organic residues related to the carbon groups and to enhance the transfer characteristics of the TiO₂/ZnO phototransistors.

Fig. 2 shows the photoresponse characteristics of ZnO and TiO₂/ZnO TFTs with each TTIP purge time at various wavelengths ($\lambda = 635, 520, 450,$ and 405 nm, $P = 4.5$ mW cm⁻²). The transfer curve of ZnO TFT showed photoresponse under the illumination of various wavelengths, as shown in Fig. 2a. It enabled the detection of UV light with 405 nm due to the wide bandgap of ZnO and the migration of photo-electrons from the valence band to the conduction band of ZnO.²⁷ However, it is difficult to absorb visible light, such as 520 and 635 nm, owing to the low photon energy.²⁸ Moreover, an increased photocurrent and negative shift in the threshold voltage (V_{th}) were observed at the wavelength of 450 nm despite insufficient photon energy. Solution-processed ZnO had numerous defects related to the oxygen vacancies (V_O). Ionized V_O^+ and V_O^{2+} act as shallow donor states, leading to the generation of unexpected photoexcited charge carriers in the bandgap of ZnO by the photon energy.^{29,30} Fig. 2b shows the transfer curve of TiO₂ (purge_{TTIP} 0.5 s)/ZnO TFT under numerous wavelengths. An inferior photoresponse was observed at the visible light since the device showed a high leakage current path with unstable transfer characteristics. TiO₂ (purge_{TTIP} 10 s)/ZnO TFT showed a remarkable change in the photoresponse at the 520 nm wavelength light with an increased I_{photo}/I_{dark} of $\sim 10^5$ and -9.8 V of negative shift V_{th} , as shown in Fig. 2c. Photo-excited electrons were generated from the oxygen-related-states of TiO₂ and then transferred into the conduction band of ZnO.³¹ Also, much increased photocurrent was observed at the wavelengths of 450 and 405 nm, more than that of 520 nm due to the additionally generated photoexcited electrons from ZnO and TiO₂. In Fig. 2d, a small photocurrent and negative shift V_{th} were observed at the wavelength of 520 nm since the TTIP purge

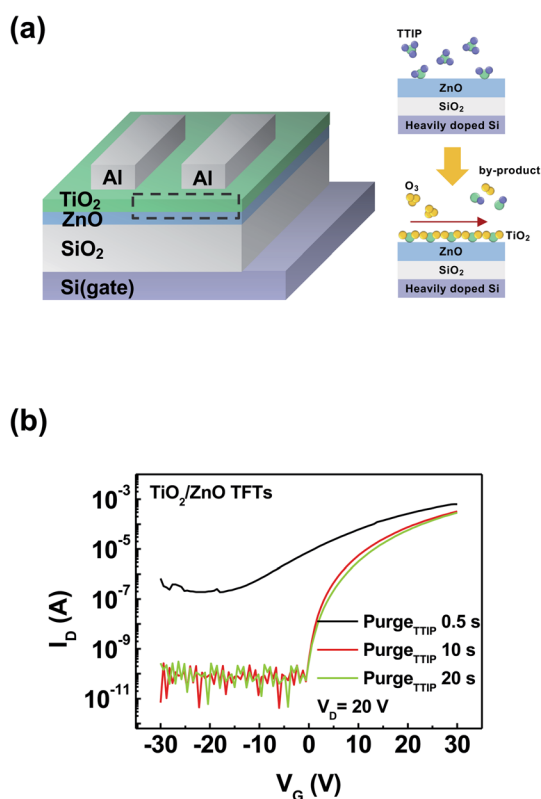


Fig. 1 (a) Schematic of the phototransistor with a heterostructure of TiO₂ and ZnO with the ALD process of TiO₂. (b) Transfer curves of the TiO₂/ZnO TFTs according to the change in the TTIP purge time (0.5, 10, and 20 s).

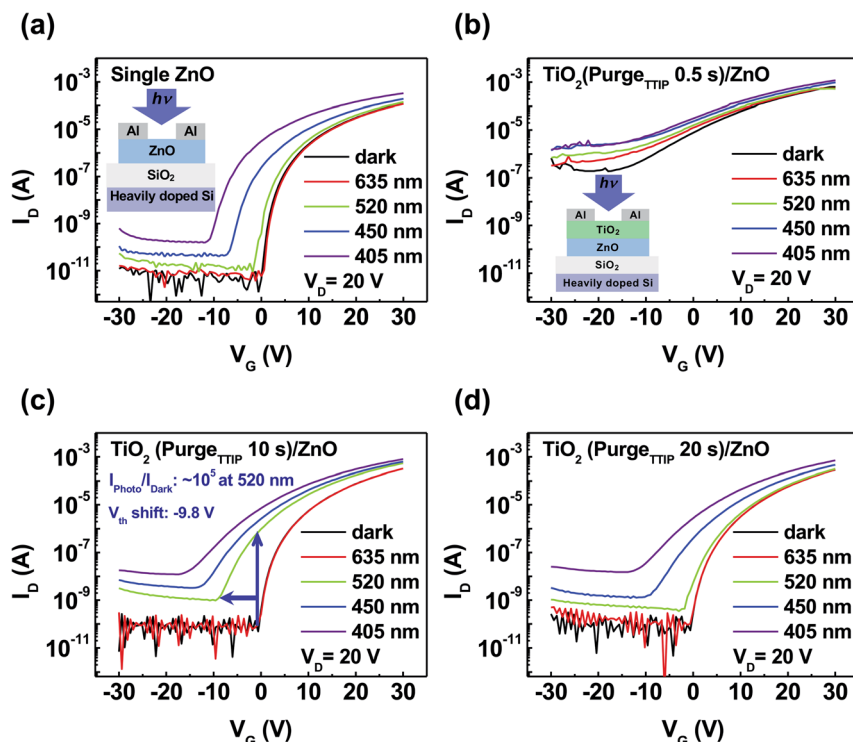


Fig. 2 Transfer curve characteristics with the illumination of various wavelengths of light for (a) ZnO phototransistor and TiO_2/ZnO phototransistors according to the different TTIP purge time periods of (b) 0.5, (c) 10, and (d) 20 s at $V_D = 20$ V.

time of 20 s enabled the separation of the ligands and carbon residues from TTIP and enhanced the bonds between Ti atoms and O atoms to form a stable TiO_2 film. These results indicate that the TiO_2/ZnO visible light phototransistor was optimized at the TTIP purge time of 10 s.

To investigate an interfacial electronic band structure at the optimized TiO_2 (purge_{TTIP} 10 s)/ZnO interface, UPS measurements were conducted, as shown in Fig. 3a. UPS spectra of the secondary electron cutoff (SEC) region and valence region were obtained. The incident energy of the UV source (He I) was 21.2 eV. E_C is the conduction band maximum energy level; E_{vac} is the vacuum level energy. The work function was obtained from the fitting of the UPS spectra. The work functions of ZnO and TiO_2 (purge_{TTIP} 10 s) were 4.07 eV and 4.41 eV, respectively. The energy level between Fermi energy (E_F) and valence band maximum (VBM) of ZnO and TiO_2 (purge_{TTIP} 10 s) were 3.02 and 3.06 eV, respectively. The band gaps ($E_g = E_C - E_V$) of ZnO and TiO_2 (purge_{TTIP} 10 s) were calculated from the Tauc's plot based on the UV-Vis spectroscopy measurement, as shown in Fig. S2.† The band gaps of ZnO and TiO_2 (purge_{TTIP} 10 s) were 3.24 and 3.30 eV, respectively. Fig. 3b shows the band alignment between ZnO and TiO_2 (purge_{TTIP} 10 s). The small band offset (0.02 eV) between the conduction band of ZnO and TiO_2 (purge_{TTIP} 10 s) was measured. A favorable interfacial band alignment between ZnO and TiO_2 (purge_{TTIP} 10 s) was achieved, and efficient photo-excited electron transport from TiO_2 (purge_{TTIP} 10 s) into the ZnO layer could be expected.

XPS was performed to investigate the oxygen-related-states of TiO_2 according to numerous TTIP purge time periods (0.5, 10, and 20 s). Fig. S3a† shows the XPS survey spectra of ZnO, TiO_2

layers with variation in the purge time periods of TTIP. In the XPS spectra, the O 1s, Zn 2p, and Ti 2p peaks were intensively observed. However, Zn 2p was negligibly observed in the TiO_2/ZnO samples since the TiO_2 film was deposited on the ZnO film. The O 1s peak was analyzed to confirm the oxygen-related-states of the TiO_2 layers. As shown in Fig. 4, the O 1s peaks are

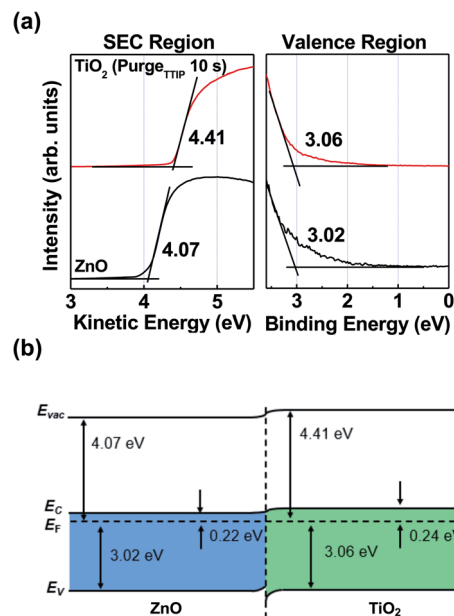


Fig. 3 (a) UPS spectra, and (b) schematic of the band alignment between the ZnO and TiO_2 (TTIP 10 s) films.

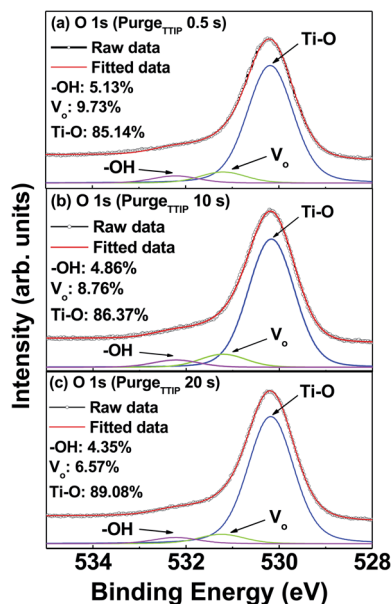


Fig. 4 O 1s spectra of TiO₂ on the ZnO film with different TTIP purge time periods of (a) 0.5, (b) 10, and (c) 20 s.

deconvoluted into three peaks at 530.1, 531.2, and 532.2 eV, which are attributed to the Ti–O bond of the lattice oxygen, V_O in the TiO₂ layer, and hydroxides species (–OH) on the film surface, respectively.³² With the increase in the purge time of TTIP, the Ti–O bonds at the lattice oxygen increased, and the relative area ratios were 85.14, 86.37, and 89.08% at the TTIP purge time periods of 0.5, 10, and 20 s, respectively.³³ With the increase in the TTIP purge time, the relative area ratio of V_O decreased (9.73, 8.76, and 6.57% at TTIP purge time periods of 0.5, 10, and 20 s, respectively), indicating that subgap-states related to V_O exist at TiO₂.³⁴ As shown in Fig. S3b–d,† the Ti 2p peaks are deconvoluted. The peaks at 463.4 and 457.8 eV are assigned to Ti³⁺ 2p_{1/2} and Ti³⁺ 2p_{3/2}, respectively. Also, the peaks at 464.9 and 458.8 eV are assigned to Ti⁴⁺ 2p_{1/2} and Ti⁴⁺ 2p_{3/2}, respectively.³⁵ Table 1 summarizes the binding energy and area ratio of Ti 2p_{1/2} and Ti 2p_{3/2}. According to the TTIP purge time periods (0.5, 10, and 20 s) of TiO₂, the area ratios of Ti³⁺ at Ti 2p_{1/2} are 48.65, 47.82, and 46.50%, and the area ratios of Ti³⁺ at Ti 2p_{3/2} are 3.29, 2.47, and 0.91%. As the purge time of TTIP increased, the area ratio of Ti³⁺ decreased, which corresponded to oxygen vacancies with Ti³⁺ from deficient bonding between Ti and O. However, the area ratio of Ti⁴⁺ gradually increased.³⁶ Therefore, it indicates that the TiO₂ with the TTIP purge time of 10 s was a more useful film than that of the TTIP purge time of 0.5 s for the visible-light phototransistor, as

shown in Fig. 2b. The optimized oxygen-related defects by V_O in the bandgap of TiO₂ could excite photo-generated electrons in TiO₂ by the illumination of visible light.^{37–39} Therefore, it was found that the process of 10 s of TTIP purge time during TiO₂ deposition was appropriate for a visible-light phototransistor.

UV-visible spectroscopy measurements were conducted to investigate the photoresponse by subgap-states under visible light, as shown in Fig. 5a. ZnO and TiO₂ (purge_{TTIP} 10 s)/ZnO films were measured at the wavelength ranging from 350 to 800 nm. Compared to the ZnO film, the intensity of absorbance remarkably increased at the TiO₂ (purge_{TTIP} 10 s)/ZnO film under visible light (λ = 400 to 700 nm) by the subgap-states of the TiO₂.⁴⁰ Moreover, Fig. 5b shows the UPS spectra of ZnO and TiO₂ (purge_{TTIP} 10 s)/ZnO films. The difference in the intensity indicates that the subgap-states near the valence band edge in the TiO₂ (purge_{TTIP} 10 s)/ZnO film exist more than those of the ZnO film in the range from 0 to 3.5 eV.^{41,42} These results suggest that subgap-states in the TiO₂ (purge_{TTIP} 10 s)/ZnO film can induce visible light absorption. Therefore, it indicates that electrons at the subgap-states and oxygen-related states can be trapped near the CBM and VBM levels of the TiO₂ film. As shown in Fig. 5c, these trapped-electrons can generate photoexcited electrons by trap-assisted generation under low photon energy, such as the visible light. Then, photoexcited electrons near the CBM of TiO₂ can be smoothly transferred into ZnO with a favorable band alignment between ZnO and TiO₂ due to the minimal difference in the CBM energy level (0.08 eV). As a result, the increased number of the photoexcited electrons by trapped electrons at the subgap-states of the TiO₂ causes the photocurrent under the visible light.

Fig. 6a shows the photoresponsivity of ZnO and TiO₂ (purge_{TTIP} 10 s)/ZnO phototransistors under the exposure at the wavelength of 520 nm as a function of V_G and V_D = 20 V, which was plotted using the equation:

$$\text{Photoresponsivity} = \frac{(I_{\text{light}} - I_{\text{dark}})/A_{\text{pt}}}{P/A_{\text{pd}}} = \frac{J_{\text{ph}}}{P}$$

where I_{light} is the current in the phototransistor under the exposure of visible light, I_{dark} is the dark current, P is the power of the incident light, A_{pt} is the product of the channel width and thickness, A_{pd} is the spot size of the laser source, J_{ph} is the photocurrent density, and P is the incident laser power density.¹⁵ The maximum photoresponsivity was calculated as 99.3 A W⁻¹ at the TiO₂ (purge_{TTIP} 10 s)/ZnO phototransistor, and the laser power was kept under 4.5 mW cm⁻². Also, the photoresponsivity at the TiO₂ (purge_{TTIP} 10 s)/ZnO phototransistor showed superior photoresponse towards the ZnO

Table 1 XPS analysis results for the Ti 2p core level of the TiO₂/ZnO film

| | Ti 2p _{1/2} [eV (%)] | | Ti 2p _{3/2} [eV (%)] | |
|--|--------------------------------|------------------|--------------------------------|------------------|
| | Ti ³⁺ | Ti ⁴⁺ | Ti ³⁺ | Ti ⁴⁺ |
| TiO ₂ (purge _{TTIP} 0.5 s)/ZnO | 463.38 (48.65) | 464.92 (51.35) | 457.83 (3.29) | 458.83 (96.71) |
| TiO ₂ (purge _{TTIP} 10 s)/ZnO | 463.38 (47.82) | 464.90 (52.18) | 457.83 (2.47) | 458.80 (97.53) |
| TiO ₂ (purge _{TTIP} 20 s)/ZnO | 463.37 (46.50) | 464.89 (53.50) | 457.75 (0.91) | 458.77 (99.09) |

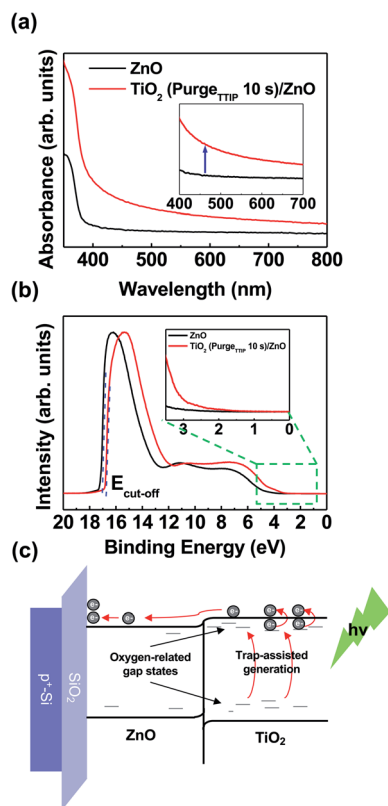


Fig. 5 (a) Absorption spectra of the ZnO and TiO₂ (TTIP 10 s)/ZnO films. The inset shows the absorption spectra from 400 to 700 nm. (b) UPS spectra of the ZnO and TiO₂ (TTIP 10 s)/ZnO films. (c) Schematic of the photoexcited charge transport mechanism at the interface between ZnO and TiO₂ (TTIP 10 s) with the illumination of light.

phototransistor, which could not detect the visible light due to the wide bandgap of ZnO. Fig. 6b shows the photosensitivity of ZnO and TiO₂ (purge_{TTIP} 10 s)/ZnO phototransistors under the exposure at the wavelength of 520 nm at a function of V_G and $V_D = 20$ V, which was plotted using the equation:

$$\text{Photosensitivity} = \frac{(I_{\text{light}} - I_{\text{dark}})}{I_{\text{dark}}}$$

where I_{light} is the current in the phototransistor under the exposure of visible light, and I_{dark} is the dark current.¹⁹ The photosensitivity of the TiO₂ (purge_{TTIP} 10 s)/ZnO phototransistor was superior to that of the ZnO phototransistor under visible light illumination. The maximum photosensitivity of the TiO₂ (purge_{TTIP} 10 s)/ZnO phototransistor was 1.5×10^5 . To confirm the photomodulation characteristics of the TiO₂ (purge_{TTIP} 10 s)/ZnO phototransistor, the on/off photoresponse characteristic was measured at an illumination of a periodic light signal of 520 nm with 0.5 Hz, $V_G = -15$ V, and $V_D = 20$ V. As shown in Fig. 6c, the device periodically responded to the periodic light illumination. The result indicates that the visible light can excite trapped-electrons from the subgap-states and oxygen-related-states to the CBM of TiO₂.²¹ The result from the experiment indicates a feasible way to fabricate a visible-light-sensitive phototransistor based on the heterostructure composed of a wide bandgap oxide semiconductor with oxygen-related-states.

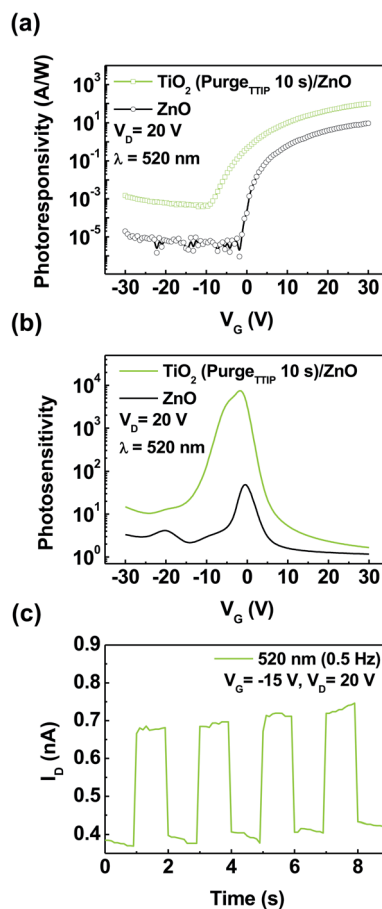


Fig. 6 (a) Photoresponsivity and (b) photosensitivity of the ZnO and TiO₂ (TTIP 10 s)/ZnO phototransistors at $V_D = 20$ V under the visible light illumination ($\lambda = 520$ nm, $P = 4.5$ mW cm⁻²). (c) Photoresponse characteristics of the TiO₂ (TTIP 10 s)/ZnO phototransistors under 0.5 Hz periodic illumination of 520 nm wavelength light.

Conclusion

A visible light phototransistor based on the heterostructure of ALD-TiO₂ and a spin-coated ZnO has been fabricated. The device generated a photocurrent under the exposure visible light due to the oxygen-related-states and subgap-states of TiO₂. To investigate the origin of photocurrent generation, UPS, XPS, and UV-visible spectroscopy measurements were performed. The oxygen-related-states enabled the generation of a photocurrent under visible light illumination on the device even when the active channel materials have a wide bandgap. The favorable band alignment between ZnO and TiO₂ (purge_{TTIP} 10 s) was confirmed by the measurement of UPS for the efficient charge transfer between TiO₂ and ZnO. Therefore, the TiO₂ (purge_{TTIP} 10 s)/ZnO phototransistor exhibited a photoresponsivity of 99.3 A W⁻¹, photosensitivity of 1.5×10^5 , and $I_{\text{photo}}/I_{\text{dark}}$ ratio of $\sim 10^5$. Moreover, it was hard to observe a persistent photocurrent effect during the modulation of the device with periodic visible light illumination. The result provides a useful way to fabricate a visible light phototransistor based on the heterostructure of wide bandgap oxide semiconductors.

Experimental

ZnO solution synthesis

A ZnO solution was synthesized using 0.0813 g of ZnO powder (Sigma Aldrich, 99.999%) and 12 mL ammonium hydroxide solvent (Alfa Aesar, 99.9%). Then, the ZnO solution was stirred for 30 min in the ambient atmosphere. The ZnO solution was refrigerated for 5 h to increase the solubility of the ZnO powder in the ammonium hydroxide solvent.⁴³

Fabrication of the device

To fabricate the ZnO phototransistor with an ALD-TiO₂, a ZnO film was spin-coated onto a SiO₂ (100 nm)/Si substrate at 3000 rpm for 30 s. Then, the ZnO film was annealed at 300 °C for 1 h under ambient atmosphere to remove the ammonium hydroxide solvent. Subsequently, a TiO₂ film was deposited onto the ZnO film *via* the ALD process with titanium isopropoxide (TTIP) and ozone. The ALD system was operated at an Ar flow rate of 150 sccm and 300 cyc as well as process variations to purge time of TTIP. The process was maintained at 250 °C of the substrate temperature. Finally, 100 nm-thick aluminum electrodes for source/drain were deposited onto the substrate using a thermal evaporator with a shadow mask. The shadow mask was a channel with a length (*L*) of 100 and width (*W*) of 1000 μm.

Characterization and measurement of the film and device

A semiconductor parameter analyzer (HP 4145B) and probe station were used to investigate the transfer curves of the phototransistors. The optoelectrical characteristics were measured under dark conditions and exposure to the various illumination wavelengths (λ) of 405, 450, 520, 650, and 780 nm, with an illumination power of approximately 4.5 mW cm⁻². A UV-visible spectrometer (Jasco V-570) was used to record the absorbance spectra of the ZnO and TiO₂/ZnO film. XPS and UPS measurements were performed to investigate the interfacial electronic structure between ZnO and TiO₂ using a modified KRATOS AXIS NOVA system with an Al K α (1486.6 eV) source and a He-I line (21.22 eV) source at a base pressure of 5×10^{-9} torr.

Conflicts of interest

There are no conflicts to declare.

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Notes and references

- 1 K. Nomura, *Science*, 2003, **300**, 1269–1272.
- 2 H. Lee, K. S. Chang, Y. J. Tak, T. S. Jung, J. W. Park, W.-G. Kim, J. Chung, C. B. Jeong and H. J. Kim, *J. Inf. Disp.*, 2017, **18**, 131–135.
- 3 K. Nomura, H. Ohta, A. Takagi, T. Kamiya, M. Hirano and H. Hosono, *Nature*, 2004, **432**, 488–492.
- 4 S. Hong, J. W. Park, H. J. Kim, Y. Kim and H. J. Kim, *J. Inf. Disp.*, 2016, **17**, 93–101.
- 5 B. D. Ahn, H. S. Shin, G. H. Kim, J.-S. Park and H. J. Kim, *Jpn. J. Appl. Phys.*, 2009, **48**, 03B019.
- 6 A. Nathan, S. Lee, S. Jeon and J. Robertson, *J. Disp. Technol.*, 2014, **10**, 917–927.
- 7 W.-G. Kim, Y. J. Tak and H. J. Kim, *J. Inf. Disp.*, 2018, **19**, 39–43.
- 8 S.-E. Ahn, I. Song, S. Jeon, Y. W. Jeon, Y. Kim, C. Kim, B. Ryu, J.-H. Lee, A. Nathan, S. Lee, G. T. Kim and U.-I. Chung, *Adv. Mater.*, 2012, **24**, 2631–2636.
- 9 Z. Tao, X. Liu, W. Lei and J. Chen, *Mater. Lett.*, 2018, **228**, 451–455.
- 10 H. Choi, S. Seo, J.-H. Lee, S.-H. Hong, J. Song, S. Kim, S.-Y. Yim, K. Lee, S.-J. Park and S. Lee, *J. Mater. Chem. C*, 2018, **6**, 6014–6022.
- 11 D.-K. Kwon, Y. Porte, K. Y. Ko, H. Kim and J.-M. Myoung, *ACS Appl. Mater. Interfaces*, 2018, **10**, 31505–31514.
- 12 X. Xu, L. Yan, T. Zou, R. Qiu, C. Liu, Q. Dai, J. Chen, S. Zhang and H. Zhou, *ACS Appl. Mater. Interfaces*, 2018, **10**, 44144–44151.
- 13 H. Yoo, W.-G. Kim, B. H. Kang, H. T. Kim, J. W. Park, D. H. Choi, T. S. Kim, J. H. Lim and H. J. Kim, *ACS Appl. Mater. Interfaces*, 2020, **12**, 10673–10680.
- 14 J. E. Cho, J. Yu and S. J. Kang, *Curr. Appl. Phys.*, 2016, **16**, 1560–1563.
- 15 J. Yu, B. J. Kim, S. Park, I. K. Han and S. J. Kang, *J. Appl. Phys.*, 2018, **57**, 044001.
- 16 J. Yu, S. W. Shin, K.-H. Lee, J.-S. Park and S. J. Kang, *J. Vac. Sci. Technol., B: Nanotechnol. Microelectron.: Mater., Process., Meas., Phenom.*, 2015, **33**, 061211.
- 17 D. Kufer and G. Konstantatos, *ACS Photonics*, 2016, **3**, 2197–2210.
- 18 Y. S. Rim, K.-C. Ok, Y. M. Yang, H. Chen, S.-H. Bae, C. Wang, Y. Huang, J.-S. Park and Y. Yang, *ACS Appl. Mater. Interfaces*, 2016, **8**, 14665–14670.
- 19 B. J. Kim, S. Park, T. Y. Kim, E. Y. Jung, J.-A. Hong, B.-S. Kim, W. Jeon, Y. Park and S. J. Kang, *J. Mater. Chem. C*, 2020, **8**, 16384–16391.
- 20 Y. S. Rim, Y. M. Yang, S.-H. Bae, H. Chen, C. Li, M. S. Goorsky and Y. Yang, *Adv. Mater.*, 2015, **27**, 6885–6891.
- 21 J. Chung, Y. J. Tak, W.-G. Kim, B. H. Kang and H. J. Kim, *ACS Appl. Mater. Interfaces*, 2019, **11**, 38964–38972.
- 22 H. Yoo, W.-G. Kim, B. H. Kang, H. T. Kim, J. W. Park, D. H. Choi, T. S. Kim, J. H. Lim and H. J. Kim, *ACS Appl. Mater. Interfaces*, 2020, **12**, 10673–10680.
- 23 J. S. Kim, S. W. Cho, D. I. Kim, B. U. Hwang, Y. G. Seol, T. W. Kim and N.-E. Lee, *J. Nanosci. Nanotechnol.*, 2014, **14**, 8596–8601.
- 24 Y. S. Rim, W. H. Jeong, D. L. Kim, H. S. Lim, K. M. Kim and H. J. Kim, *J. Mater. Chem.*, 2012, **22**, 12491.
- 25 Y.-H. Kim, J.-S. Heo, T.-H. Kim, S. Park, M.-H. Yoon, J. Kim, M. S. Oh, G.-R. Yi, Y.-Y. Noh and S. K. Park, *Nature*, 2012, **489**, 128–132.

- 26 Z. Liu, Z. Jian, J. Fang, X. Xu, X. Zhu and S. Wu, *Int. J. Photoenergy*, 2012, **2012**, 8.
- 27 S. Park, B. J. Kim, S. J. Kang and N.-K. Cho, *J. Korean Phys. Soc.*, 2018, **73**, 1351–1355.
- 28 Y. S. Rim, B.-D. Ahn, J.-S. Park and H. J. Kim, *J. Phys. D: Appl. Phys.*, 2013, **47**, 045502.
- 29 Y. J. Tak, D. H. Yoon, S. Yoon, U. H. Choi, M. M. Sabri, B. D. Ahn and H. J. Kim, *ACS Appl. Mater. Interfaces*, 2014, **6**, 6399–6405.
- 30 H. S. Kang, J. S. Kang, J. W. Kim and S. Y. Lee, *J. Appl. Phys.*, 2004, **95**, 1246–1250.
- 31 Y.-C. Shen, C.-Y. Tung, C.-Y. Huang, Y.-C. Lin, Y.-G. Lin and R.-H. Horng, *ACS Appl. Electron. Mater.*, 2019, **1**, 783–788.
- 32 K. Shen, K. Wu and D. Wang, *Mater. Res. Bull.*, 2014, **51**, 141–144.
- 33 H.-H. Hsu, C.-H. Cheng, P. Chiou, Y.-C. Chiu, C.-Y. Chang and Z.-W. Zheng, *Solid-State Electron.*, 2014, **99**, 51–54.
- 34 B. H. Kang, W.-G. Kim, J. Chung, J. H. Lee and H. J. Kim, *ACS Appl. Mater. Interfaces*, 2018, **10**, 7223–7230.
- 35 J. Tian, Y. Leng, H. Cui and H. Liu, *J. Hazard. Mater.*, 2015, **299**, 165–173.
- 36 X. Zhang, H. Tian, X. Wang, G. Xue, Z. Tian, J. Zhang and Z. Zou, *Mater. Lett.*, 2013, **100**, 51–53.
- 37 N. Mintcheva, S. Yamaguchi and S. A. Kulinich, *Materials*, 2020, **13**, 719.
- 38 B. Bharti, S. Kumar, H.-N. Lee and R. Kumar, *Sci. Rep.*, 2016, **6**, 32355.
- 39 Y. Liu, J. Wang, P. Yang and K. Matras-Postolek, *RSC Adv.*, 2015, **5**, 61657–61663.
- 40 B. Ozden, C. Yang, F. Tong, M. P. Khanal, V. Mirkhani, M. H. Sk, A. C. Ahyi and M. Park, *Appl. Phys. Lett.*, 2014, **105**, 172105.
- 41 H. Tang, Y. Kishida, K. Ide, Y. Toda, H. Hiramatsu, S. Matsuishi, S. Ueda, N. Ohashi, H. Kumomi, H. Hosono and T. Kamiya, *ECS J. Solid State Sci. Technol.*, 2017, **6**, P365.
- 42 L. Gao, Y. Li, J. Ren, S. Wang, R. Wang, G. Fu and Y. Hu, *Appl. Catal., B*, 2017, **202**, 127–133.
- 43 J. J. Richardson and F. F. Lange, *Cryst. Growth Des.*, 2009, **9**, 2570–2575.