Monte Carlo-based beam quality and phantom scatter corrections for solid-state detectors in ⁶⁰Co and ¹⁹²Ir brachytherapy dosimetry

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Beam quality correction, $k_{QQ_0}(r)$, for solid-state detectors diamond, LiF, Li₂B₄O₇, Al₂O₃, and plastic scintillator are calculated as a function of distance, r, along the transverse axis of the 60Co and 192Ir brachytherapy sources using the Monte Carlobased EGSnrc code system. This study also includes calculation of detector-specific phantom scatter correction, $k_{phan}(r)$, for solid phantoms such as PMMA, polystyrene, solid water, virtual water, plastic water, RW1, RW3, A150, and WE210. For ⁶⁰Co source, $k_{OO}(r)$ is about unity and distance-independent for diamond, plastic scintillator, $\text{Li}_2 \tilde{B}_4^0 O_7$ and LiF detectors. For this source, $k_{QQ_0}(r)$ decreases gradually with r for Al₂O₃ detector (about 6% smaller than unity at 15 cm). For ¹⁹²Ir source, $k_{OO_2}(r)$ is about unity and distance-independent for Li₂B₄O₇ detector (overall vāriation is about 1% in the distance range of 1–15 cm). For this source, $k_{OO}(r)$ increases with r for diamond and plastic scintillator (about 6% and 8% larger than unity at 15 cm, respectively). Whereas $k_{QQ_0}(r)$ decreases with r gradually for LiF (about 4% smaller than unity at 15 cm) and steeply for Al₂O₃ (about 25% smaller than unity at 15 cm). For ⁶⁰Co source, solid water, virtual water, RW1, RW3, and WE210 phantoms are water-equivalent for all the investigated solid-state detectors. Whereas polystyrene and plastic water phantoms are water-equivalent for diamond, plastic scintillator, Li₂B₄O₇ and LiF detectors, but show distance-dependent $k_{phan}(r)$ values for Al₂O₃ detector. PMMA phantom is water-equivalent at all distances for Al_2O_3 detector, but shows distance-dependent $k_{phan}(r)$ values for remaining detectors. A150 phantom shows distance-dependent $k_{phan}(r)$ values for all the investigated detector materials. For ¹⁹²Ir source, solid water, virtual water, RW3, and WE210 phantoms are water-equivalent for diamond, plastic scintillator, Li₂B₄O₇ and LiF detectors, but show distance-dependent $k_{phan}(r)$ values for Al₂O₃ detector. All other phantoms show distance-dependent $k_{nhan}(r)$ values for all the detector materials.

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I. INTRODUCTION

¹⁹²Ir and ⁶⁰Co sources are used in high-dose-rate (HDR) brachytherapy.⁽¹⁻⁴⁾ Dosimetry of a brachytherapy source is generally carried out using various solid-state detectors. The response of the detector is required to be corrected for absorbed dose energy dependence, when it is not water-equivalent. Although water is recommended as the reference medium for dosimetry of

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brachytherapy sources,^(5,6) different solid phantoms are also used to overcome practical problems such as waterproofing and precise positioning of detectors. In a previously published article,⁽⁷⁾ relative absorbed-dose energy response corrections, R, were reported for different solid-state detectors for ¹⁶⁹Yb and ¹²⁵I brachytherapy sources. The study also included the influence of solid phantom materials such as PMMA (polymethylmethacrylate) and polystyrene on R. In another study,⁽⁸⁾ the values of R were reported for different radiochromic films for high-energy brachytherapy sources such as ⁶⁰Co, ¹³⁷Cs, ¹⁹²Ir, and ¹⁶⁹Yb in liquid water, PMMA, and polystyrene phantom materials. In a recently published article by Selvam et al.,⁽⁹⁾ beam quality corrections (which is the inverse of R) were reported for different solid-state detectors at ¹³⁷Cs energy. In addition, detector-specific phantom scatter corrections for different solid phantoms were also reported in their study. The purpose of this study is to calculate the beam quality corrections for solid-state detectors for ⁶⁰Co and ¹⁹²Ir brachytherapy sources. This study also includes detector-specific phantom scatter corrections for different solid phantoms for these sources. The investigation of phantom scatter also includes water as a detector material. The EGSnrc-based⁽¹⁰⁾ user-codes DOSRZnrc and FLURZnrc ⁽¹¹⁾ are used in the study.

II. MATERIALS AND METHODS

A. Radioactive sources

The brachytherapy sources investigated in this study are BEBIG HDR ⁶⁰Co (model Co0.A86; Eckert & Zielger BEBIG BmbH, Berlin, Germany)⁽³⁾ and HDR ¹⁹²Ir (model MicroSelectron; Elekta, Stockholm, Sweden).⁽²⁾ In the Monte Carlo calculations, two gamma lines 1.17 MeV and 1.33 MeV are considered for the ⁶⁰Co source. For ¹⁹²Ir source, the spectrum is taken from the literature.⁽¹²⁾

B. Detector and phantom materials

The detector materials investigated in this study are diamond, LiF, $\text{Li}_2\text{B}_4\text{O}_7$, plastic scintillator, and Al_2O_3 . The solid phantom materials investigated are PMMA, polystyrene, solid water, virtual water, plastic water, RW1, RW3, A150, and WE210. The atomic composition and density details of RW1 and virtual water phantoms are taken from the published studies.^(13,14) Remaining phantom data are taken from the study by Seco and Evans⁽¹⁵⁾

C. Beam quality and phantom scatter corrections

As described in the published study by Selvam et al.,⁽⁹⁾ beam quality correction, $k_{QQ_0}(r)$, and phantom scatter correction, $k_{phan}(r)$, can be calculated at a brachytherapy beam quality Q for solid-state detectors by using the following relations:

$$k_{\underline{QQ}_0}(r) = \frac{\left[D_{w,\underline{Q}}(r)/D_{\det,\underline{Q}}(r)\right]}{\left[D_{w,\underline{Q}_0}/D_{\det,\underline{Q}_0}\right]}$$
(1)

$$k_{phan}(r) = \left[D_{\det,Q}(r) / D_{\det,phan,Q}(r) \right]$$
⁽²⁾

Here, $D_{w,Q}(r)$ and $D_{det,Q}(r)$ are the absorbed dose to water and absorbed dose to detector in liquid water at a distance, r, along the transverse axis of the photon emitting brachytherapy source of beam quality Q (in the present study, it is ¹⁹²Ir or ⁶⁰Co), respectively. $D_{w,Q}(r)$ and $D_{det,Q}(r)$ are the absorbed dose to water and absorbed dose to detector in water at the reference beam quality Q₀ (realistic ⁶⁰Co teletherapy beam), respectively. $D_{det,Phan,Q}(r)$ is the absorbed dose to detector in the solid phantom at r from the brachytherapy source of beam quality Q.

For the calculation of $k_{QQ_0}(r)$, the values of water-to-detector dose ratio at Q_0 (denominator of Eq. (1)) are taken from the published article.⁽⁷⁾ Note that $k_{phan}(r)$ converts absorbed dose to detector at r for the brachytherapy source (of beam quality Q) in a solid phantom to absorbed dose to detector in liquid water phantom at the same r. Numerator of $k_{QQ_0}(r)$ corrects for the difference in the energy absorption properties of water and detector at brachytherapy beam quality Q at r, and the denominator of $k_{QQ_0}(r)$ corrects for the same, but at reference beam quality Q₀.

D. Monte Carlo calculations

Dose ratios of water-to-detector at beam quality Q (numerator of Eq. (1)) are based on the FLURZnrc user-code⁽¹¹⁾ as described in the published study.⁽⁷⁻⁹⁾ In the Monte Carlo calculations, the source is positioned at the centre of a 40 cm diameter × 40 cm height cylindrical phantoms (liquid water and solid phantoms). The photon fluence spectrum is scored along the transverse axis of the source (r = 1-15 cm) in 0.5 mm high and 0.5 mm thick cylindrical shells. The fluence spectrum is converted to collision kerma-to-water and collision kerma-to-detector materials by using the mass-energy absorption coefficients of water and detector materials, respectively.⁽¹⁶⁾ Note that the FLURZnrc⁽¹¹⁾ simulations also provide fluence-weighted mean energy of photons, E_{fl}. Up to 10⁹ photon histories are simulated. The 1 σ statistical uncertainty on the calculated absorbed dose and collision kerma values are about 0.2%. The statistical uncertainties on the calculated values of $k_{QQ_0}(r)$ and $k_{phan}(r)$ are less than 0.5%. The Monte Carlo parameters used in the calculations are same as that used in the earlier work.^(8,9)

III. RESULTS & DISCUSSION

A. Fluence-weighted mean energy, E_{fl}

Tables 1 and 2 present the values of $E_{\rm fl}$ as a function of r for ⁶⁰Co and ¹⁹²Ir sources in various phantoms, respectively. $E_{\rm fl}$ decreases with distance due to degradation in the photon energy after scattering. The degree of decrease depends on the type of the phantom as well as the type of source. For the ⁶⁰Co source, the decrease in $E_{\rm fl}$ is higher in PMMA, polystyrene, and A150 phantom as compared to other phantoms. For example, $E_{\rm fl}$ decreases from 1.134 MeV to 455 keV in PMMA, from 1.146 MeV to 486 keV in polystyrene, and from 1.140 MeV to 481 keV in A150 phantom when the distance is increased from 1 cm to 15 cm. For phantoms such as water, RW1, RW3, and solid water, $E_{\rm fl}$ decreases from about 1.15 MeV to 520 keV in the above distance range. For the virtual water and WE210 phantoms, $E_{\rm fl}$ decreases from

TABLE 1. Monte Carlo-calculated values of fluence-weighted mean energy for different phantoms presented as a function of distance r along the transverse axis of the BEBIG 60 Co source.

Distance, r (cm)	Water	PMMA	Polystyrene	Plastic water	RW1	RW3	Virtual water	Solid water	A150	WE210
1	1.149	1.134	1.146	1.152	1.151	1.149	1.150	1.149	1.140	1.152
2	1.057	1.026	1.049	1.063	1.061	1.056	1.058	1.058	1.036	1.061
3	0.972	0.927	0.958	0.982	0.974	0.969	0.973	0.972	0.943	0.980
4	0.896	0.842	0.877	0.913	0.900	0.891	0.899	0.897	0.861	0.905
5	0.832	0.770	0.807	0.851	0.831	0.824	0.835	0.830	0.792	0.841
6	0.774	0.709	0.746	0.796	0.774	0.766	0.778	0.776	0.731	0.785
7	0.723	0.656	0.693	0.751	0.726	0.715	0.728	0.726	0.681	0.735
8	0.682	0.612	0.649	0.712	0.683	0.672	0.687	0.684	0.637	0.694
9	0.644	0.575	0.608	0.678	0.646	0.636	0.650	0.648	0.602	0.659
10	0.612	0.546	0.576	0.648	0.614	0.603	0.620	0.615	0.569	0.626
11	0.585	0.517	0.550	0.623	0.587	0.577	0.594	0.589	0.543	0.599
12	0.563	0.495	0.528	0.603	0.564	0.556	0.571	0.567	0.522	0.578
13	0.547	0.477	0.511	0.585	0.548	0.539	0.555	0.549	0.503	0.559
14	0.532	0.462	0.498	0.573	0.533	0.523	0.542	0.534	0.490	0.547
15	0.520	0.455	0.486	0.562	0.523	0.512	0.530	0.524	0.481	0.536

Distance, r (cm)	Water	PMMA	Polystyrene	Plastic water	RW1	RW3	Virtual water	Solid water	A150	WE210
1	0.325	0.320	0 324	0.327	0.325	0 324	0.325	0.325	0.321	0.326
2	0.295	0.285	0.292	0.299	0.295	0.294	0.296	0.294	0.288	0.296
3	0.270	0.257	0.265	0.276	0.270	0.269	0.271	0.270	0.262	0.272
4	0.249	0.234	0.242	0.258	0.249	0.247	0.250	0.250	0.240	0.252
5	0.233	0.216	0.223	0.242	0.232	0.229	0.233	0.233	0.223	0.235
6	0.218	0.200	0.208	0.230	0.217	0.215	0.220	0.219	0.208	0.221
7	0.206	0.188	0.195	0.220	0.205	0.203	0.208	0.207	0.196	0.210
8	0.197	0.178	0.184	0.210	0.195	0.193	0.198	0.197	0.186	0.200
9	0.189	0.169	0.175	0.203	0.187	0.185	0.190	0.189	0.178	0.192
10	0.182	0.162	0.168	0.196	0.179	0.177	0.184	0.183	0.171	0.186
11	0.175	0.156	0.161	0.192	0.173	0.171	0.178	0.176	0.165	0.179
12	0.170	0.152	0.156	0.186	0.168	0.166	0.172	0.172	0.160	0.174
13	0.166	0.147	0.152	0.183	0.164	0.162	0.169	0.168	0.156	0.171
14	0.163	0.145	0.148	0.180	0.161	0.159	0.165	0.165	0.153	0.167
15	0.161	0.142	0.147	0.178	0.159	0.156	0.163	0.163	0.151	0.165

TABLE 2. Monte Carlo-calculated values of fluence-weighted mean energy for different phantoms presented as a function of distance r along the transverse axis of the ¹⁹²Ir source.

about 1.150 MeV to 530 keV and, in the case of plastic water, from 1.152 MeV to 562 keV in the above distance range.

For ¹⁹²Ir source, decrease in E_{ff} is higher for PMMA, A150, and polystyrene phantoms as compared to other phantoms. E_{ff} decreases from about 320 keV to 140 keV when the distance is increased from 1 cm to 15 cm. For phantoms such as water, WE210, virtual water, and solid water, E_{ff} decreases from about 325 keV to 160 keV in the above distance range. For RW1 and RW3 phantoms, $E_{\rm fl}$ decreases from about 325 keV to 156 keV in the above distance range. In the case of plastic water phantom, E_{ff} decreases from 327 keV to 178 keV when the distance is increased from 1 cm to 15 cm.

B. Beam quality correction, $k_{QQ_0}(r)$ Table 3 presents the values of $k_{QQ_0}(r)$ for the ⁶⁰Co and ¹⁹²Ir sources, respectively. For Li₂B₄O₇ detector, $k_{QQ_0}(r)$ is about unity, and is independent of r for both the sources. For the ⁶⁰Co source, $k_{QQ_0}(r)$ is about unity and distance independent for diamond, plastic scintillator, and LiF detectors. Whereas for the ¹⁹²Ir source, $k_{QQ_0}(r)$ increases gradually about 6% and 8%

TABLE 3. Beam quality correction, $k_{QQ_0}(r)$, presented for diamond, Al_2O_3 , $Li_2B_4O_7$, LiF, and plastic scintillator detectors as a function of distance r along the transverse axis of ⁶⁰Co and ¹⁹²Ir sources.

Distance, r	Diamond		Al ₂ O ₃		$Li_2B_4O_7$		LiF		Plastic Scintillator	
(cm)	⁶⁰ Co	¹⁹² Ir	⁶⁰ Co	¹⁹² Ir	⁶⁰ Co	^{192}Ir	⁶⁰ Co	¹⁹² Ir	⁶⁰ Co	¹⁹² Ir
1	1.000	1.004	0.998	0.973	1.000	1.000	1.000	0.996	1.001	1.017
2	1.001	1.008	0.996	0.955	1.000	1.001	0.999	0.994	1.001	1.022
3	1.001	1.012	0.992	0.935	1.000	1.001	0.999	0.991	1.002	1.027
4	1.002	1.016	0.989	0.913	1.000	1.002	0.998	0.987	1.003	1.031
5	1.003	1.021	0.984	0.892	1.000	1.003	0.998	0.984	1.003	1.037
6	1.003	1.026	0.980	0.870	1.000	1.003	0.997	0.980	1.004	1.043
7	1.004	1.031	0.975	0.849	1.000	1.004	0.996	0.977	1.005	1.048
8	1.005	1.036	0.970	0.830	1.000	1.005	0.996	0.973	1.006	1.055
9	1.006	1.041	0.965	0.813	1.001	1.006	0.995	0.970	1.007	1.061
10	1.007	1.045	0.960	0.797	1.001	1.006	0.994	0.967	1.008	1.067
11	1.008	1.050	0.956	0.783	1.001	1.007	0.994	0.964	1.009	1.071
12	1.009	1.054	0.952	0.770	1.001	1.008	0.993	0.961	1.010	1.075
13	1.009	1.057	0.949	0.760	1.001	1.008	0.993	0.959	1.011	1.078
14	1.010	1.060	0.946	0.752	1.001	1.009	0.992	0.957	1.012	1.083
15	1.010	1.062	0.944	0.746	1.001	1.009	0.992	0.956	1.012	1.084

larger than unity for diamond and plastic scintillator, but decreases about 4% smaller than unity for LiF detector with r over the distance range of 1–15 cm. For Al₂O₃ detector, $k_{QQ_0}(r)$ decreases with r gradually about 6% and steeply about 25% smaller than unity for ⁶⁰Co and ¹⁹²Ir sources respectively, in the above distance range.

C. Phantom scatter correction, $k_{phan}(r)$

Table 4 presents the summary of $k_{phan}(r)$ results for diamond, Al₂O₃, Li₂B₄O₇, LiF, and plastic scintillator detectors in the investigated phantom materials for the ⁶⁰Co and ¹⁹²Ir sources, respectively. In this table, phantoms which are water-equivalent (i.e., $k_{phan}(r)$ is unity) at all distances (1–15 cm) are designated as "Yes". "No" implies that the phantoms show distance-dependent $k_{phan}(r)$ values. For such phantoms results are discussed below.

TABLE 4. Summary of $k_{phan}(r)$ results presented for diamond, Al₂O₃, Li₂B₄O₇, LiF, and plastic scintillator detectors for the ⁶⁰Co and ¹⁹²Ir sources, respectively. In this table, "Yes" implies the phantom is water-equivalent (i.e., $k_{phan}(r)$ is unity) at all distances (1–15 cm) along the transverse axis of the sources. "No" implies that the phantoms show distance-dependent $k_{nhan}(r)$ values (figure number is shown in parenthesis).

Phantom	Diamond / Plastic Scintillator		Al ₂ O ₃		Li ₂ E	$B_4 O_7$	LiF	
Materials	⁶⁰ Co	¹⁹² Ir	⁶⁰ Co	¹⁹² Ir	⁶⁰ Co	¹⁹² Ir	⁶⁰ Co	¹⁹² Ir
PMMA	No (Fig. 2)	No (Fig.5)	Yes	No (Fig.5)	No (Fig.2)	No (Fig.5)	No (Fig.2)	Yes
Polystyrene	Yes	No (Fig.6)	No (Fig.1)	No (Fig.6)	Yes	No (Fig.6)	Yes	No (Fig.6)
Plastic water	Yes	No (Fig.8)	No (Fig.1)	No (Fig.8)	Yes	No (Fig.8)	Yes	No (Fig.8)
RW1	Yes	No (Fig.7)	Yes	No (Fig.7)	Yes	No (Fig.7)	Yes	No (Fig.7)
RW3	Yes	Yes	Yes	No (Fig.4)	Yes	Yes	Yes	Yes
Virtual water	Yes	Yes	Yes	No (Fig.4)	Yes	Yes	Yes	Yes
Solid water	Yes	Yes	Yes	No (Fig.4)	Yes	Yes	Yes	Yes
A150	No (Fig. 3)	No (Fig.9)	No (Fig.3)	No (Fig.9)	No (Fig.3)	No (Fig.9)	No (Fig.3)	No (Fig.9)
WE210	Yes	Yes	Yes	No (Fig.4)	Yes	Yes	Yes	Yes

C.1 ⁶⁰Co source

Phantoms such as solid water, virtual water, RW1, RW3, and WE210 are water-equivalent (i.e., $k_{phan}(r)$ is unity) at all distances (1–15 cm) for all the solid-state detectors (maximum deviation from unity is about 1% at 15 cm for Al₂O₃ detector in solid water, RW1, and RW3). Polystyrene and plastic water phantoms are water-equivalent at all distances for all the detectors (with a maximum deviation of about 1% from unity for LiF), other than Al₂O₃. Figure 1 presents the distance-dependent $k_{phan}(r)$ values for the Al₂O₃ detector in plastic water and polystyrene phantoms. PMMA is water-equivalent at all distances for Al₂O₃ detector (larger than unity by about 1% at 15 cm), whereas $k_{phan}(r)$ increases with r for remaining detector materials, including water (see Fig. 2). In this phantom, $k_{phan}(r)$ values are comparable for diamond, plastic scintillator, Li₂B₄O₇, LiF, and water detectors at all distances. For Al50 phantom, $k_{phan}(r)$ values are comparable for the detectors, including water (see Fig. 3). For this phantom, $k_{phan}(r)$ values are comparable for the detectors diamond, plastic scintillator, Li₂B₄O₇, LiF, and water detectors diamond, plastic scintillator, Li₂B₄O₇, LiF. and water at all distances for Simple for the detectors diamond, plastic scintillator, Li₂B₄O₇, LiF. and water at the detectors diamond, plastic scintillator, Li₂B₄O₇, LiF. and water at the detectors diamond, plastic scintillator, Li₂B₄O₇, LiF. and water at the detectors diamond, plastic scintillator, Li₂B₄O₇, LiF. and water at the detectors diamond, plastic scintillator, Li₂B₄O₇, LiF.

all distances, with a maximum value of about 1.045 at 15 cm. For Al_2O_3 , the maximum value of $k_{phan}(r)$ is 1.027 at 15 cm.



FIG. 1. Phantom scatter correction, $k_{phan}(r)$, presented for Al₂O₃ detector in polystyrene and plastic water phantoms as a function of distance along the transverse axis of the BEBIG ⁶⁰Co brachytherapy source.



FIG. 2. Phantom scatter correction, $k_{phan}(r)$, presented for PMMA phantom as a function of distance along the transverse axis of the BEBIG ⁶⁰Co brachytherapy source. The values are presented for detector materials LiF, Li₂B₄O₇, diamond, plastic scintillator, and water.



FIG. 3. Phantom scatter correction, $k_{phan}(r)$, presented for A150 phantom as a function of distance along the transverse axis of the BEBIG ⁶⁰Co brachytherapy source. The values are presented for detector materials LiF, Li₂B₄O₇, diamond, plastic scintillator, Al₂O₃, and water.

C.2 ¹⁹²Ir source

Phantoms such as solid water, virtual water, RW3, and WE210 are water-equivalent in the distance range of 1–15 cm for all the detectors other than Al_2O_3 (with a maximum deviation of about 2% at 15 cm for solid water and RW3 phantoms). Figure 4 presents the distance-dependent $k_{phan}(r)$ values of Al_2O_3 detector for the above four phantom materials. For this detector, $k_{phan}(r)$ increases with r for solid water, virtual water, and WE210 phantoms and decreases with r for RW3 phantom. $k_{phan}(r)$ is comparable for solid water, virtual water, and WE210 phantoms.

PMMA is water-equivalent for LiF detector. Figure 5 presents $k_{phan}(r)$ values for all the detector materials other than LiF. For this phantom, $k_{phan}(r)$ decreases with r for Al₂O₃ detector (about 10% at 15 cm), whereas $k_{phan}(r)$ increases with r for all the other detectors. The degree of increase is higher for diamond detector and plastic scintillator (maximum deviation from unity at 15 cm is about 5% and 6%, respectively).

The phantoms polystyrene, RW1, plastic water, and A150 show distance-dependent $k_{phan}(r)$ values which are presented in Figs. 6 to 9. $k_{phan}(r)$ decreases with r for all the detector materials in polystyrene and RW1 phantoms (Figs. 6 and 7). However, the degree of decrease is higher for Al₂O₃ detector compared to all other detectors. For example, the value decreases to 0.821 and 0.960 at 15 cm for polystyrene and RW1 phantoms, respectively. For plastic water phantom, $k_{phan}(r)$ values increase with r for all the detector materials, including water (Fig. 8). The degree



FIG. 4. Phantom scatter correction, $k_{phan}(r)$, presented for Al₂O₃ detector in virtual water, solid water, RW3, and WE210 phantoms as a function of distance along the transverse axis of the ¹⁹²Ir brachytherapy source.



FIG. 5. Phantom scatter correction, $k_{phan}(r)$, presented for PMMA phantom as a function of distance along the transverse axis of the ¹⁹²Ir brachytherapy source. The values are presented for detector materials Li₂B₄O₇, diamond, plastic scintillator, Al₂O₃, and water.

of increase is higher for Al_2O_3 detector (about 20% larger than unity at 15 cm) compared to all other detectors (minimum deviation of about 5% from unity at 15 cm for diamond and plastic scintillator detector).



FIG. 6. Phantom scatter correction, $k_{phan}(r)$, presented for polystyrene phantom as a function of distance along the transverse axis of the ¹⁹²Ir brachytherapy source. The values are presented for detector materials Li₂B₄O₇, LiF, diamond, plastic scintillator, Al₂O₃, and water.



FIG. 7. Same as Fig. 6, but for RW1 phantom.



FIG. 8. Same as Fig. 6, but for plastic water phantom.

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In the case of A150 phantom, $k_{phan}(r)$ value increases with r for all the detector materials (maximum deviation of about 6% from unity at 15 cm for diamond detector) other than Al₂O₃ detector (Fig. 9). For Al₂O₃ detector, $k_{phan}(r)$ decreases from 0.997 (at 1 cm) to 0.978 (at 7 cm) and thereafter increases to unity at a distance of 15 cm. In order to verify this trend beyond 15 cm, auxiliary simulations are carried out using the FLURZnrc user-code⁽¹¹⁾ with larger dimensions (50 cm diameter × 50 cm height) of A150 and water phantoms, to calculate $k_{phan}(r)$ for r = 1–20 cm. Figure 10 compares $k_{phan}(r)$ values obtained in 50 cm diameter × 50 cm height and 40 cm diameter × 40 cm height phantoms for Al₂O₃ detector. Up to 15 cm $k_{phan}(r)$ values are comparable in both the phantom dimensions. For 50 cm diameter × 50 cm height phantom, $k_{phan}(r)$ reaches the value of 1.032 at r = 20 cm. To verify any possible influence of the detector is modeled as a 1 mm thick × 2 mm high cylinder. The values of $k_{phan}(r)$ are calculated along the transverse axis of the ¹⁹²Ir source for r = 1, 5, 10, 15, and 20 cm. The study shows that DOSRZnrc-based $k_{phan}(r)$ values are statistically identical to the corresponding FLURZnrc-based $k_{phan}(r)$ values.



FIG. 9. Same as Fig. 6, but for A150 phantom.



FIG. 10. Phantom scatter correction, $k_{phan}(r)$, presented for Al₂O₃ detector in 40 cm × 40 cm and 50 cm × 50 cm A150 phantoms. The calculations are based on the FLURZnrc user-code.

D. Influence of detector dimensions on detector response

The above-described FLURZnrc-based calculated values of $k_{phan}(r)$ and $k_{phan}(r)$ are based on the assumption that charged particle equilibrium exists and the presence of detector does not

affect the above corrections. In order to quantify the influence of detector thicknesses on the calculated response, auxiliary simulations are carried out in water phantom using the DOSRZnrc user-code.⁽¹⁰⁾ LiF, Li₂B₄O₇, plastic scintillator, and Al₂O₃ detectors are modeled as cylindrical shells of thicknesses 1 mm and height 2 mm, whereas diamond detector is modeled for two different thicknesses (0.2 mm and 0.4 mm) and height 2 mm. Absorbed dose and collision kerma to these detectors are calculated at r = 1 and 15 cm. For ¹⁹²Ir source, collision kerma and absorbed-dose values are statistically identical for all the detectors. For ⁶⁰Co source, collision kerma and absorbed-dose values are statistically identical for Al₂O₃, plastic scintillator, Li₂B₄O₇, and LiF detectors. In the case of diamond detector, the absorbed dose values are smaller by about 1% at 1 cm and about 1.5% at 15 cm, compared to the collision kerma values.

IV. CONCLUSIONS

Beam quality correction, $k_{QQ_0}(r)$, for solid-state detector materials such as diamond, plastic scintillator, LiF, Li₂B₄O₇, and Al₂O₃ are calculated as a function of distance along the transverse axis of the ⁶⁰Co and ¹⁹²Ir brachytherapy sources using the Monte Carlo-based EGSnrc code system. For ⁶⁰Co source, $k_{QQ_0}(r)$ is about unity and distance independent for diamond, plastic scintillator, Li₂B₄O₇, and LiF detector, and decreases gradually with r for Al₂O₃ (about 6% lesser than unity at 15 cm). For ¹⁹²Ir source, $k_{QQ_0}(r)$ is about unity and independent of distance for Li₂B₄O₇ detector. $k_{QQ_0}(r)$ increases with distance for diamond and plastic scintillator (about 6% and 8% larger than unity at 15 cm, respectively). $k_{QQ_0}(r)$ decreases gradually with r for LiF and steeply for Al₂O₃.

Phantom scatter correction, $k_{phan}(r)$, for various solid phantoms are calculated for the above detectors along the transverse axis of ⁶⁰Co and ¹⁹²Ir sources. For ⁶⁰Co source, phantoms such as solid water, virtual water, RW1, RW3, and WE210 are water-equivalent for all the investigated detectors. Polystyrene and plastic water phantoms are water-equivalent for diamond, plastic scintillator, Li₂B₄O₇, and LiF detectors, but shows distance-dependent $k_{phan}(r)$ values for Al₂O₃ detector. PMMA is water-equivalent at all distances for Al₂O₃, but shows distance-dependent $k_{phan}(r)$ values for remaining detectors. A150 phantom shows distance-dependent $k_{phan}(r)$ values for all the detector materials. For ¹⁹²Ir source, solid water, virtual water, RW3, and WE210 phantoms are water-equivalent for diamond, plastic scintillator, Li₂B₄O₇, and LiF detectors, whereas these phantoms show distance-dependent $k_{phan}(r)$ values for all the detector materials. For ¹⁹²Ir source, solid water, virtual water, RW3, and WE210 phantoms are water-equivalent for diamond, plastic scintillator, Li₂B₄O₇, and LiF detectors, whereas these phantoms show distance-dependent $k_{phan}(r)$ values, but the degree of dependence depends on the type of solid phantom and the detector. Li₂B₄O₇ detector shows $k_{phan}(r)$ values identical to that of water detector, and diamond detector shows $k_{phan}(r)$ values identical to that of plastic scintillator detector for all the investigated phantoms for ¹⁹²Ir and ⁶⁰Co sources.

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REFERENCES

- Karaiskos P, Angelopoulos A, Sakelliou L, et al. Monte Carlo and TLD dosimetry of an 192Ir high dose-rate brachytherapy source. Med Phys. 1998;25(10):1975–84.
- Daskalov GM, Loffler E, Williamson JF. Monte Carlo-aided dosimetry of a new high dose-rate brachytherapy source. Med Phys. 1998;25(11):2200–08.

- 3. Granero D, Perez-Calatayud J, Ballester F. Technical note: Dosimetric study of a new Co-60 source used in brachytherapy. Med Phys. 2007;34(9):3485–88.
- Ballester F, Granero D, Perez-Calatayud J, Casal E, Agramunt S, Cases R. Monte Carlo dosimetric study of the BEBIG Co-60 HDR source. Phys Med Biol. 2005;50(21):N309–N316.
- Rivard MJ, Coursey BM, DeWerd LA, et al. Update of AAPM Task Group No. 43 Report: A revised AAPM protocol for brachytherapy dose calculations. Med Phys. 2004;31(3):633–74.
- Nath R, Anderson LL, Luxton G, Weaver KA, Williamson JF, Meigooni AS. Dosimetry of interstitial brachytherapy sources: recommendations of the AAPM Radiation Therapy Committee Task Group No. 43. Med Phys. 1995;22(2):209–34.
- Selvam TP and Keshavkumar B. Monte Carlo investigation of energy response of various detector materials in 1251 and 169Yb brachytherapy dosimetry. J Appl Clin Med Phys. 2010;11(4):70–82.
 Subhalaxmi M and Selvam TP. Monte Carlo-based investigation of absorbed-dose energy dependence of radio-
- Subhalaxmi M and Selvam TP. Monte Carlo-based investigation of absorbed-dose energy dependence of radiochromic films in high energy brachytherapy dosimetry. J Appl Clin Med Phys. 2014;15(1):351–62.
- Selvam TP, Subhalaxmi M, Vishwakarma RS. Monte Carlo calculation of beam quality correction for solid-state detectors and phantom scatter correction at 137Cs energy. J Appl Clin Med Phys. 2014;15(1):339–50.
- Kawrakow I, Mainegra-Hing E, Rogers DW, Tessier F, Walters BRB. The EGSnrc Code System: Monte Carlo simulation of electron and photon transport. NRCC Report PIRS–701. Ottawa, ON: National Research Council of Canada; 2010.
- Rogers DW, Kawrakow I, Seuntjens JP, Walters BRB, Mainegra-Hing E. NRC user codes for EGSnrc. NRCC Report PIRS-702 (revB). Ottawa, ON: National Research Council of Canada; 2010.
- 12. Shirley VS. Nuclear data sheets for A = 192. Berkeley, CA: Lawrence Laboratory; 1991.
- Reniers B, Verhaegen F, Vynckier S. The radial dose function of low-energy brachytherapy seeds in different solid phantoms: comparison between calculations with the EGSnrc and MCNP4C Monte Carlo codes and measurements. Phys Med Biol. 2004;49(8):1569–82.
- Murphy MK, Piper RK, Greenwood LR, et al. Evaluation of the new cesium-131 seed for use in low-energy x-ray brachytherapy. Med Phys. 2004;31(6):1529–38.
- Seco J and Evans PM. Assessing the effect of electron density in photon dose calculations. Med Phys. 2006;33(2):540–52.
- Hubbell JH and Seltzer SM. Tables of x-ray mass attenuation coefficients and mass energy-absorption coefficients. Gaithersburg, MD: National Institute of Standards and Technology; 1995. Available from: http://www.nist.gov/ pml/data/xraycoef/