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OPEN One-step Synthesis of Few-layer **WS2 by Pulsed Laser Deposition**

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Atomically thin tungsten disulfide (WS₂) has attracted much attention in recent years due its indirect**to-direct band gap transition, band gap tunability, and giant spin splitting. However, the fabrication of atomically thin WS2 remains largely underdeveloped in comparison to its structural analogue MoS2. Here we report the direct fabrication of highly crystalline few-layer WS2 on silver substrates by pulse laser deposition at the relatively low temperature of 450°C. The growth takes places by conventional epitaxy, through the** *in-situ* **formation of nearly lattice-matching Ag2S on the silver surface. Intriguingly,** it was observed that the resulting film was composed of not only the usual semiconducting 2H-WS₂ structure but also the less common metallic 1T-WS₂. Modifications of the synthesis parameters allow **for control over the crystalline quality, film thickness and crystal phase composition of the resulting WS2 film.**

Among the many materials that have the capacity to form 2-dimensional (2D) layers, much attention has been devoted to the lamellar transition metal dichalcogenides (TMDC)¹⁻⁴. This class of materials, in which a plane of metal atoms are sandwiched between two planes of chalcogen atoms, exhibits interesting and diverse properties, ranging from insulators to semiconductors to metals^{[2](#page-7-1)}. As a typical example of a TMDC, tungsten disulphide (WS2) demonstrates a unique combination of structural, electronic, optical, mechanical, chemical, and thermal properties that have been the focus of various studies. Its electronic band gap for example, undergoes an indirect (1.4 eV) to direct (2 eV) (2 eV) (2 eV) transition when its size is reduced from bulk to a single layer². The material has been used in fluorescent emitters⁵ and field effect transistors^{[6](#page-7-3)}, for photovoltaics^{[7](#page-7-4)} and photocatalysis⁸, and is regarded as highly valuable for high power applications such as solid state batteries^{[9](#page-7-6)} and supercapacitors¹⁰. Monolayer WS_2 also possesses strong spin-orbit induced electronic band splitting and spin-valley coupling^{[11](#page-7-8),[12](#page-7-9)}, in addition to a band structure that can be tuned by variation in strai[n13.](#page-7-10) Until recently, most of the attention has been focused on the semiconducting phase of WS_2 and MoS_2 , both of which possess a prismatic coordination for the metal atom (2H). This 2H-structure is generally produced by physical exfoliation or vapour growth methods. In contrast, the metallic 1T-phase, with an octahedral coordination for the metal atom, has been synthesised primarily through lithium intercalation. Both 1T-WS₂ and 1T-MoS₂ layers have recently been shown to be very efficient hydrogen evolution electrocatalysts¹⁴⁻¹⁶.

Unlike related systems such as $MoS₂$, techniques for effective preparation of 2D WS₂ remains largely underdeveloped. There are currently two main methods of preparing ultrathin WS_2 : (1) top-down exfoliation and (2) bottom-up substrate growth. Exfoliation can be mechanical^{17–18}, chemically assisted¹⁹ (e.g. sonication in a good solvent) or purely chemical^{[14](#page-7-11),[20](#page-7-14)} (intercalation with e.g. lithium), but the end results remain similar: high quality flakes can be prepared but there is very little control over their size, shape or nature of their edges. On the other hand, the substrate growth technique, commonly some form chemical vapour deposition²¹⁻²³ (CVD), is capable of producing large area monolayers with high crystallinity and good control over flake shape. However, many of these substrate growth techniques utilize different solid precursors heated to high temperatures in the range of 750–1000 °C and require long growth times. Although pulsed laser deposition (PLD) is also a type of bottom-up substrate growth technique, it is unique in that it is a purely physical method. Among its advantages is the ability to grow high quality films, to ablate any material and to obtain a stoichiometric transfer of target material onto the substrate, which is especially useful in the case of composite materials such as WS₂. While PLD has found success in the fabrication of 2D materials such as few-layer graphene²⁴ and MoS_2^{25} , it remains underdeveloped compared to exfoliation and CVD methods.

In this work, we describe the synthesis of highly crystalline few-layer WS₂ on Ag substrates using PLD at a relatively low temperature of 450 °C. Ag metal was selected as substrate for the ease of the ease of producing

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Figure 1. (a) Raman spectra of as-grown samples fabricated on Ag and sapphire. The WS₂ film on sapphire was fabricated at 200 mJ and 10 s. The left and right dashed lines indicate the positions of the 2LA (M) and A_{1g} phonon modes in bulk WS₂ respectively. (**b**) Multi-peak Lorentzian fitting of Raman bands in the Ag sample fabricated at 100 mJ and 10 s, with bulk WS₂ included for comparison. (c) Raman spectrum of as-deposited WS₂ on quartz with varying Ag buffer thickness. All samples were synthesized using the parameters of 100 mJ and 10 s. (**d**) Optical transmittance of a quartz substrate with Ag buffer layer of 8 nm, before and after WS₂ film growth. Inset: Photograph of the as-grown WS_2 sample and bare quartz.

Laser Energy (mJ)	Deposition Time	$A_{1g}(\Gamma)$ (cm ⁻¹)	$E_{2g}^{!}(\Gamma)(cm^{-1})$	$2LA (M) (cm-1)$	$\frac{I_{2LA}}{I_{A_{1g}}}$	Estimated Number of Layers
50	10	418	351	355	0.94	\sim 2
100	10	418.7	350.6	355	0.75	\sim 3
200	10	419.1	350.1	354.9	0.58	$4 - 5$
.50	30	419.5	350.2	355	0.51	>5
Bulk		420	350.4	355.2	0.46	Bulk

Table 1. Peak position for the Raman modes A1g (Γ), 2LA(M) and E! 2g (Γ), as well as the intensity ratio of A1g (Γ) to 2LA(M) for the as-deposited samples on Ag.

metal-semiconductor contacts without film transfer. Additionally, the introduction of a metal support can substantially alter the H binding energy of 2D TMDCs²⁶, enabled by charge transfer from the substrate to the overlying film and through strong interactions at the interface. Metal supported WS_2 could thus potentially be used as novel catalysts for hydrogen production. Surprisingly, using PLD to grow WS₂ does not produce the expected semiconducting prismatic 2H structure. Instead, it creates metallic WS₂ films with the distorted octahedral 1T-WS₂ structure. The structure and properties of these films were explored through a combination of Raman spectroscopy, photoluminescence (PL) measurement, x-ray photoelectron spectroscopy (XPS), high-resolution transmission electron microscopy (HRTEM), and powder x-ray diffraction (PXRD).

Results and Discussion

[Figure 1a](#page-1-0) shows the Raman spectra of the as-grown WS_2 film on Ag excited under ambient conditions, with the spectrum of bulk WS₂ and a WS₂ sample grown on insulating sapphire substrates included for comparison. The spectrum of the sapphire sample was normalized to remove the peaks from the substrate. At this excitation wavelength of 514.5nm, the spectrum reveals many second-order Raman peaks, in addition to the first-order phonon modes. The strongest peak at ~352 cm⁻¹ can be resolved by multi-peak Lorentzian fitting into three individual contributions at 343, 351, and 355 $\rm cm^{-1}$ as shown in [Fig. 1b](#page-1-0) for the sample fabricated at 100 mJ and 10 seconds. These modes are assigned to the in-plane vibrational $E'_{2g}(M)$ mode, the second-order mode of longitudinal acoustic phonon 2LA (M), and the in-plane vibrational $E'_{2g}(\Gamma)$ mode respectively¹². The other peaks at 418 and 343 cm⁻¹ are attributed to the out-of-plane A^1_{g} mode and the in-plane vibrational E^1_{2g} (M) mode. Studies have previously shown that Raman characterization can provide unambiguous and nondestructive identification of the thickness of WS₂; the A_{1g} (Γ) mode softens while both the 2LA (M) and E¹_{2g} (Γ) modes present a subtle red-shift with a decreasing number of layers^{5,27}. In particular, at 514.5 nm laser excitation, the WS₂ spectrum reveals a striking increase in the intensity ratio of the 2LA (M) to $A_{1g}(\Gamma)$ phonon modes due to a double resonance process²⁷. [Table 1](#page-1-1) summarizes the frequency for the three main Raman modes $A_{1g}(\Gamma)$, 2LA(M) and $E'_{2g}(\Gamma)$, as well as the intensity ratio for the two most intense peaks in our pulsed laser fabricated samples on Ag. Based on these values, it can be concluded that atomically thin WS₂ films (\leq 5 layers) can be formed at any laser energy within the range of 50–200 mJ as long as the deposition time was capped at a maximum of 10 seconds. Compared to chemically derived $\text{WS}_2^{21,23},$ $\text{WS}_2^{21,23},$ $\text{WS}_2^{21,23},$ our samples exhibited Raman peaks with broader full width at half maximums (FWHM), such that it becomes difficult to distinguish the phonon modes in the 260-330 cm⁻¹ range without multi-peak Lorentzian fitting. This

indicates the presence of defects arising from disorder in the atomic arrangement of the WS₂ film. In addition, the samples fabricated at a laser energy of 50mJ demonstrate much lower signal intensity, possibly indicating only partial crystallization or a large concentration of defects. As such, even though the thinnest films are formed at the lowest energy, for an optimal balance between film thickness and crystalline quality, a laser energy of 100mJ would be most ideal to grow highly crystalline few-layer $WS₂$.

Although insulating substrates such as sapphire are preferred substrates for growing 2D materials for characterization of electronic properties, it was found that such substrates do not produce good quality WS₂ films in pulsed laser synthesis. [Figure 1a](#page-1-0) shows that despite the use of higher laser energies of 200mJ to fabricate the sapphire sample, its Raman spectrum resembles the partially crystalline and highly defective spectrum of the Ag sample fabricated at 50 mJ and 10 s. At lower energies of 100 and 50 mJ, the characteristic peaks of 2H-WS, are completely absent. This obstacle can be however, overcome by sputtering a very thin layer of the Ag buffer on the substrate. Our investigations reveal that the minimum thickness that the Ag layer can have before it begins to significantly affect the crystalline quality of the overlying WS_2 film is ~8 nm ([Fig. 1c\)](#page-1-0). These WS_2 films were grown on quartz, which not only permits the direct synthesis of crystalline few-layer WS₂ on insulating substrates without the need for transfer techniques, but also allows the fabrication of samples that are capable of transmitting visible light. [Figure 1d](#page-1-0) shows that the triple-layered WS₂ film on quartz has an average optical transmittance of 71% in the wavelength range between 400 and 800nm. Reducing the buffer layer to such minimal thickness also results in *in-situ* consumption of most, if not all, of the pure Ag metal to form Ag₂S (a phase which was detected in the TEM, XPS and XRD analysis). This was deduced from the increase in optical transmission in the sample after deposition of WS₂, also shown in [Fig. 1d.](#page-1-0) The high optical transmittance of our WS₂ films on quartz means that they are suitable for solar energy applications. The presence of the Ag₂S phase is also beneficial as the material appears to be a promising solar absorbing material with its narrow band gap of $\sim 0.9 \text{ eV}^{28}$ and its unique combination of properties such as light absorbance in the near-infrared spectral regions²⁹.

High resolution transmission electron microscope (TEM) images of our pulsed laser fabricated samples are shown in [Fig. 2](#page-2-0), and reveal the stacking of WS₂ (002) layers with an interplanar spacing of 0.62 nm on top of silver. [Figure 3a](#page-3-0) provides direct evidence of the successful formation of double-layered WS₂ in the sample fabricated at 50 mJ and 10 seconds, with the first WS₂ layer forming covalent bonds to the previously mentioned Ag₂S phase that develops *in-situ* on the surface of the Ag buffer. This bonding between the two layers manifests as an indistinct boundary at the interface contrary to the sharp interfaces of 2D layered materials grown by van der Waals epitaxy. The formation of the Ag_2S phase is supported by the presence of lattice fringes with interplanar spacings of 0.27 nm located underneath the WS₂ film that can be ascribed to the (120) plane of monoclinic acanthite Ag₂S [JCPDS #14-0072]. This Ag₂S phase was also previously found to similarly promote the growth of ultrathin MoS₂ by PLD through lattice matching and conventional epitaxy^{[25](#page-7-17)}. [Figure 3b](#page-3-0) shows the change in the WS₂ film when the deposition time is increased to 30 seconds. The number of layers have increased to >5, verifying the results of the Raman measurements.

To characterize the chemical nature and bonding state of WS_2 on Ag metal, X-ray photoelectron spectroscopy (XPS) was employed. [Figure 3a](#page-3-0) depicts the W 4f and S 2p core level XPS scans for the WS₂ film on silver. Two doublets are present in the S 2p spectra, one occurring at 162.3 and 163.5 eV, consistent with the S^{2−} species of WS₂, while the second pair located at 161.1 and 162.4 eV can be assigned to the S $2p_{3/2}$ and S $2p_{1/2}$ peaks of Ag₂S. The presence of both doublets confirms the successful growth of WS₂ as well as the *in-situ* formation of Ag₂S in the samples. For the W 4f spectra, there is a small shoulder at 35.6 eV corresponding to the W^{6+} state that shows the formation of tungsten oxide in the as-deposited films. This oxide can arise in part from surface oxidation of the WS_2 target, and in part from unreacted W atoms remaining on the substrate surface after completion of the WS_2 film. These unreacted W atoms are always produced because some S atoms are initially consumed in the making

of the Ag₂S phase. The oxide is believed to form sort of capping layer over the WS₂ film when it is first exposed to air, and does not negatively affect the quality the film itself as the calculated stoichiometric ratio of S : W is very close to the ideal value of 2. It was also observed that the main tungsten doublet peak can be deconvoluted into two separate pairs. The second doublet (red curve) occurs at binding energies of 32.7 and 34.8 eV, which corresponds well with the W⁴⁺ species of highly crystalline $2H\text{-}WS_2^{30}$. The doublet at 32.1 and 34.2 eV (blue curve) on the other hand lies between the binding energies of metallic tungsten (W^{0+}) and W^{4+} , and appears to be due to a partially sulfided, intermediate W^{x+} state. However, calculation of the stoichiometric ratio of S: W atoms reveal that both W tungsten species (red and blue curves) have a S: W ratio of \sim 2, indicating that the W 4f doublet located at 32.1 and 34.2 eV is also due to a W⁴⁺ state. The negative shift of peak binding energies by 0.6 eV is thus believed to be attributed to the formation of the 1T-phase of WS₂, consistent with its metallic nature, and is comparable to with previous studies on $1T-MX_2$ materials^{[14](#page-7-11),[31](#page-7-24)}.

The formation of the 1T-phase can also be established by closer inspection of the Raman spectra, as well as by powder X-ray diffraction (XRD). As noted by Voiry *et al*¹⁴, the Raman spectra of 1T-WS₂ display the additional modes of J_1 , J_2 and J_3 that are attributed to the superlattice structure of the distorted 1T-phase. However, these Raman modes are only readily observable when the concentration of 1T-WS₂ dominates over the 2H-phase. [Figure 3b](#page-3-0) shows the Raman results from two of our as-deposited samples with widely differing concentrations of 1T-WS₂. It is apparent that the J₁, J₂ and J₃ modes are noticeable only in the sample fabricated at 50 mJ and 10 s due to the much higher 1T-to-2H ratio of 1.67. In contrast, the sample fabricated at 100 mJ and 10 s only has a 1T-to-2H ratio of 0.73, and correspondingly the typical Raman peaks of 1T-WS₂ have almost completely vanished. As-deposited samples were also annealed at varying temperatures to test the stability of the 1T-phase. Included for comparison in [Fig. 3b](#page-3-0) is a sample annealed at 300°C for 30minutes, wherein a complete phase transformation from 1T to 2H-WS₂ was achieved as revealed by XPS [\(Fig. 3c](#page-3-0)). The Raman bands of 2H-WS₂ in this sample appear more

Figure 4. Powder X-ray diffractograms of WS₂ nanosheets (a) annealed at 300 °C and (b) as-deposited at 50 mJ and 10 s. (c) Bulk diffraction peaks of 2H-WS₂.

distinct; in particular the combination mode of 2LA- E^2_{2g} at ~320 cm⁻¹ becomes visibly noticeable. Nonetheless, the peaks still lack the sharpness of CVD grown WS₂ sheets, and there is only a slight increase in the peak intensities while the FWHMs remained largely unchanged. This indicates that there is some degree of atomic disorder present even in the pure 2H-phase, and the crystalline quality of PLD grown WS₂ on Ag is not quite comparable to CVD grown WS₂ layers. The fact that the 1T to 2H phase transition can be induced with thermal annealing indicates that the 1T-WS₂ produced in this work is metastable. Indeed, it was observed that partial 1T to 2H transition begins to occur even at temperatures as low as 100 °C, which would be detrimental if the 1T-phase was intended for use at high temperatures. The ability to prepare metastable 1T-WS₂ layers at temperatures higher than the 1T to 2H transition is rather surprising. However, this is likely to be simply due to the short deposition times and fast cooling rate, such that the 1T-structure does not have sufficient time to relax to the more stable 2H-phase.

Unlike with 2H-WS₂, characterization of 1T-WS₂ by powder XRD is rather challenging because the 1T-phase does not exhibit a well-defined crystalline structure and well-resolved diffraction peaks³¹. This problem is exacerbated by the fact that our samples are all mixtures of 1T and 2H-WS₂. Nonetheless, we were still able to observe differences in the XRD diffractograms obtained before and after annealing. As shown in [Fig. 4,](#page-4-0) both the as-deposited and annealed samples have peaks corresponding to the known 2H-WS₂ pattern. In the as-deposited sample, the presence of a broad peak at 14.32° with the highest intensity reveals the preferential growth of WS₂ sheets along the (002) direction. In contrast, the XRD pattern of the annealed sample shows a more intense (100) peak, which suggests the growth of protrusion edges along the (100) direction. Compared with 2H-WS₂, the (002) peak of both the as-deposited and annealed sample is shifted to lower 2θ values, indicating a lattice expansion of (002) layers, i.e. 0.62nm compared to 0.616nm. The annealed sample also presents additional peaks corresponding to 2H-WS2, suggesting an improvement in crystallinity of the WS2 film that may be attributed to a 1T to 2H-phase transformation. There is furthermore a very minor peak around 19–20° that vanishes after thermal annealing. As this peak cannot be indexed to either 2H-WS₂, Ag, Ag₂S or even WO₃, and has previously been observed to be present in $1T\text{-}WS_2^{31}$, we believe that it can be attributed to the presence of the 1T-phase.

It was observed that the concentration of the 1T-WS₂ phase appears to decrease with longer deposition times and higher laser energies ([Fig. 5](#page-5-0)). In addition, the peak position of the W 4f doublet for 1T-WS₂ continues to shift to lower binding energies with increasing deposition times, an indication of increasingly metallic character. This phenomenon could possibly be due to bombardment of the growing film by energetic species in the pulsed laser deposition process, which may lead to the breaking of W-S bonds in 1T-WS₂. It is likely that with longer deposition times, the 1T-WS₂ phase would convert completely into metallic tungsten. Similarly, the formation of a metallic layer was previously observed during prolonged sputtering of WS_2^{-32} , another ion-assisted deposition technique. The breaking of the W-S bond appears to affect only 1T-WS₂ and not 2H-WS₂, as seen by the negligible negative shift in the W 4f peak position for the 2H polymorph. This is deduced to be a consequence of the higher

thermodynamic stability of 2H-WS₂, rendering it less likely to be affected by the bombardment of energetic species during film growth.

The presence of the 1T-phase also affects the photoluminescence (PL) measurements of our samples. Weak PL was observed in as-deposited WS₂ nanosheets, as expected from their partial metallic character. Enlarging the emission spectra allows us to see that the sample fabricated at 50mJ and 10 s presents a major peak at ~645nm and a shoulder peak at ~570nm ([Fig. 6\)](#page-6-0), corresponding to A and B excitonic transitions of the K point of the Brillouin zone. Curiously, the expected intensity and position dependence of the major PL peak with layer thickness can be observed even among the as-deposited samples despite the varying concentrations of 1T-WS₂. As shown in [Fig. 6,](#page-6-0) the major PL peak at ~640nm quickly diminishes in intensity and gradually red-shifts with increasing layer thickness, a trend that has been observed by other groups^{[12](#page-7-9),[22](#page-7-26),[33](#page-7-27)}. This is unusual as the thinnest sample also has the highest concentration of 1T-WS₂, which is expected to inhibit PL, and yet this sample gives the highest emission intensity. Because PL originates near the surface of a material, and is sensitive enough to be affected by surface

Figure 6. Photoluminescence spectra of as-grown samples fabricated on Ag substrates at $\lambda_{ex} = 488$ nm.

adsorbates, the fact that the PL spectra of as-deposited samples appears to be largely unaffected by the presence of 1T-WS₂ suggests that the 1T-phase is located away from the surface of the film. While the origin of 1T-WS₂ in our samples is as yet unclear, we believe that it can be attributed to the capacity of Ag atoms to act as an electron donor for W (owing to Ag having more valence electrons than W). The stabilization mechanism is therefore similar to chemically exfoliated $MOS₂$ and $WS₂$, wherein the 1T-phase is stabilized by substitutional doping of an electron donating atom. In such a case, the 1T-phase in our samples would be concentrated at the interface while 2H-WS₂ would dominate at the surface of the film, which could explain their unexpected PL results. Further experimental and theoretical studies would be required to verify these hypotheses.

From the above results, it appears that Ag is indeed capable of stabilizing the 1T-phase of WS₂, which is noteworthy as recent studies have shown that the 2H and 1T phase have matching lattices and coherent interfaces between domains of the two phases can be formed^{[34](#page-7-28)}. As 2H-WS₂ is a semiconductor and 1T-WS₂ a metal, localized phase stabilization of the 1T phase to form a 2H-1T hybrid structure presents a viable route to achieving unique electronic heterojunctions across a chemically homogeneous layer, which is advantageous for the fabrication of molecular electronic devices. Furthermore, studies have revealed that the 1T-WS₂ phase possesses higher reactivity and catalytic ability compared to the 2H-phase^{14,22}. With Ag also being catalytically active and considered a low-cost alternative to Pt³⁵, our findings open up the possibility of Ag stabilized 1T-WS₂ nanostructures as superior catalysts for the hydrogen evolution reaction. Even so, some challenges remain, not least of which is the impact of the Ag₂S phase on the intrinsic property of the WS₂ film. Due to the covalent bonds holding WS₂ to the buffer layer, it is unlikely that the Ag₂S phase can be removed without damaging the crystal structure of the WS_2 sheets. Thus any effort to incorporate pulsed laser deposited WS_2 into electronic devices must take into account the presence of this Ag₂S phase, and further fundamental studies on the exact electronic properties of the WS₂/ Ag2S material are necessary.

In summary, we have successfully fabricated crystalline few-layer WS₂ on Ag by pulsed laser deposition. The growth process involves the *in-situ* formation of a lattice matching Ag2S phase and that promotes the formation of crystalline WS₂ through conventional epitaxy rather than expected Van der Waals epitaxy. Samples fabricated at a laser energy of 100mJ and a deposition time of 10 seconds demonstrated the optimum balance between crystalline quality and film thickness. It was also found that crystalline WS₂ is able to form on Ag layers as thin as 8nm, producing samples with high optical transmittance in the visible range, which is advantageous for applications in optoelectronic devices and solar cells. In addition, the Ag substrate was observed to stabilize the 1T-polymorph of WS2, resulting in a hybrid 2H-1T structure. Further work would however be required to fully understand the factors behind the stabilization of the 1T phase on Ag before such controlled fabrication of such hybrid or even purely 1T-WS₂ structures can be achieved.

Methods

Synthesis. Ag films of approximately 500 nm thickness were sputtered onto normal doped $(n+)$ Si substrates and loaded in a KrF ($\lambda = 248$ nm) Lambda Physik excimer PLD system. During deposition, a WS₂ target, 99.9% (Inlab Supplies) was ablated at a laser frequency of 10Hz for a pulse duration of 25ns in a vacuum environment of 2×10^{-6} Torr. The target was rotated at a speed of 6 rpm, with the laser (spot size 1 mm³) ablating a circular outline of 2 cm in radius in order obtain a uniform film. The laser energies used for ablation was kept within the range of 50–200mJ, and the deposition time between 10–30 seconds. Substrate temperature was kept at 450 °C during deposition, and subsequently decreased at a controlled rate of 20 °C/min until the temperature reached 300 °C, thereupon natural cooling processes took over. The fabrication temperature of 450 °C was selected as it is the optimum for growth of ultrathin WS, films by PLD; any lower and the crystalline quality decreases because the atoms do not have sufficient thermal energy to rearrange themselves into a periodic configuration, higher temperatures only provide minimal increase in crystallinity, whereas temperatures >700°C promotes the reaction of tungsten with oxygen.

Characterization. WS₂ films were characterized by Raman spectroscopy performed in a Horiba MicroRaman HR Evolution System using an Argon laser beam with an excitation wavelength of 514.5nm.

The phonon mode from silicon at 520.6 cm⁻¹ was used for calibration. The Raman spectrum for bulk WS₂ was collected from the WS₂ target used. PL measurements were taken in a Perkin-Elmer fluorescence spectrometer LS 55 with excitation wavelength of 488nm. A JEOL JEM-2010F high resolution transmission electron microscopy (TEM) operated at 200kV was used to obtain bright field cross-section images of the samples. Surface composition was analyzed by x-ray photoelectron spectroscopy (XPS) using a Kratos Analytical Axis Ultra^{DLD} UHV spectrometer with a monochromatized Al Kα x-ray source (1486.6eV) scanning a spot size of 700μm by 300μm. Core-level XPS spectra were obtained by photoelectrons at a take-off angle of 90°, measured with respect to the sample surface at a vacuum of 5×10^{-9} Torr. The pass energy for the narrow scans were set at 10 eV.

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Author Contributions

T.A.J.L., D.H.C.C. and A.T.S.W. designed the experiments. The synthesis and characterization of samples was carried by T.A.J.L. All authors contributed to the manuscript preparation and discussion of results.

Additional Information

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