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Accounting for Cooperativity in the Thermotropic Volume Phase Transition of Smart Microgels

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Abstract: A full quantitative description of the swelling of smart microgels is still problematic in many cases. The original approach of Flory and Huggins for the monomer–solvent interaction parameter χ cannot be applied to some microgels. The reason for this obviously is that the cross-linking enhances the cooperativity of the volume phase transitions, since all meshes of the network are mechanically coupled. This was ignored in previous approaches, arguing with distinct transition temperatures for different meshes to describe the continuous character of the transition of microgels. Here, we adjust the swelling curves of a series of smart microgels using the Flory–Rehner description, where the polymer–solvent interaction parameter χ is modeled by a Hill-like equation for a cooperative thermotropic transition. This leads to a very good description of all measured microgel swelling curves and yields the physically meaningful Hill parameter ν . A linear decrease of ν is found with increasing concentration of the cross-linker N,N'-methylenebisacrylamide in the microgel particles p(NIPAM), p(NNPAM), and p(NIPMAM). The linearity suggests that the Hill parameter ν corresponds to the number of water molecules per network chain that cooperatively leave the chain at the volume phase transition. Driven by entropy, ν water molecules of the solvate become cooperatively "free" and leave the polymer network.

Keywords: Flory–Rehner theory; microgel; Hill cooperativity; Flory–Huggins parameter; swelling behavior; NNPAM; NIPAM; NIPMAM

1. Introduction

In recent years, thermoresponsive microgels have received considerable interest in soft matter science due to their enormous application potential [1-9]. They can be used, for example, as carriers for enzymes [10,11], as drug delivery systems [12,13], as responsive surface coatings [14], or in catalysis [15–17]. The ability to take up and release chemicals at different temperatures results from their temperature-dependent swelling behavior [18–23]. The swelling behavior depends on the chosen monomer, and the related lower critical solution temperature (LCST) of the obtained polymer varies for different acrylamides [6,24]. The best-studied systems are certainly *N*-isopropylacrylamide- [25] and vinylcaprolactam-based [26] microgels. These microgels deswell at approximately 33 °C. Alternative monomers are N-n-propylacrylamide (LCST \approx 22 °C) [27] and Nisopropylmethacrylamide (LCST \approx 45 °C) [28]. Microgels based on these monomers show a volume phase transition (VPT) upon heating with a volume phase transition temperature (VPTT) close to the LCST of the respective linear polymer. Above this temperature, the microgels are collapsed. This swelling behavior is described by the Flory-Rehner theory [29–35]. In some cases, this description is quantitative; in other cases, it becomes quantitative assuming a Matrioshka-like cross-linker gradient [36]. However, recent super-resolution fluorescence microscopy results indicate that the microgels have a rather extended highly cross-linked core which is less heterogeneous than expected, as well as a rather fuzzy shell [37,38]. This is in line with some slightly older studies with



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Copyright: © 2021 by the authors. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (https:// creativecommons.org/licenses/by/ 4.0/). neutrons [39,40] which also show a nearly constant network density in about 90% of the particle volume. Moreover, often the behavior can only be predicted qualitatively [19,41,42]. Hence, there is still a need to achieve a better description of the swelling of microgels. The Flory–Huggins theory uses a lattice model for the calculation of the system partition function [29]. It is assumed that every lattice site has the same volume, so that the solvent molecules and the monomer units occupy an equal volume. The interaction between solvent and polymer network is described by the Flory–Huggins interaction parameter χ . For the calculation of χ , the interactions only with the nearest neighbor are considered. Since in some cases the experimental data cannot be rationalized with this approach, several attempts have been made to modify the Flory–Huggins theory [18]. Erman and Flory suspected that higher-order interactions must also be considered, because in the collapsed state of the gels the packing density is very high [43]. Therefore, they proposed to model the interaction parameter by a series expansion with respect to the swelling ratio. However, in this approach in most cases the virial coefficients do not have a physically meaningful interpretation. Such series expansions were successfully used to fit the swelling curves of different microgels (e.g., [19,44–47]).

The classical Flory–Rehner theory describes thermoresponsive (or thermotropic) volume phase transitions in gels without accounting for the potentially inhomogeneous distribution of the average degree of polymerization N_{Gel} in microgel particles. In thermoresponsive microgels, continuous transitions are observed nearly exclusively, which is attributed to this inhomogeneous distribution of the average degree of polymerization N_{Gel} . This observation was until now explained by different transition temperatures for individual meshes, leading to a continuous character of the VPT [20]. However, what has not yet been considered is that the cross-linking leads to a cooperativity of the VPT, since all meshes of the network are mechanically coupled.

Leite et al. [48] proposed a Hill-like model for χ that yields physically comprehensible parameters and takes into account the cooperativity of the VPT. Moreover, this approach does not rely on the degree of heterogeneity of the microgels.

In our present study, the interaction parameter χ of the Flory–Rehner theory is modeled with this Hill-like model, and is used to describe the swelling behavior of p(NNPAM), p(NIPAM), and p(NIPMAM) homopolymer microgels with different cross-linker concentrations.

2. Theory

For isotropic swelling, the polymer volume fraction ϕ is related to the hydrodynamic radius $R_{\rm H}$ of the polymer particle as:

$$\frac{\phi}{\phi_0} = \left(\frac{R_{\rm H,0}}{R_{\rm H}}\right)^3 \tag{1}$$

where $R_{H,0}$ is the hydrodynamic radius of the particle and ϕ_0 is the polymer volume fraction in the collapsed state, respectively. The hydrodynamic radius of the particle R_H is given by the Flory–Rehner equation:

$$\ln\left[1-\phi_0\left(\frac{R_{\rm H,0}}{R_{\rm H}}\right)^3\right] + \phi_0\left(\frac{R_{\rm H,0}}{R_{\rm H}}\right)^3 + \chi\phi_0^2\left(\frac{R_{\rm H,0}}{R_{\rm H}}\right)^6 + \frac{\phi_0}{N_{Gel}}\left[\frac{R_{\rm H,0}}{R_{\rm H}} - \frac{1}{2}\left(\frac{R_{\rm H,0}}{R_{\rm H}}\right)^3\right] = 0$$
(2)

where N_{Gel} is the average degree of polymerization of a polymer chain, or the number of segments between two cross-linking points. The hydrodynamic radius of the particle $R_{\rm H}$ is a free variable, which can be found by solving Equation (2). In 1996, Hino and Prausnitz suggested a modified version of Equation (2) which is more suitable for heterogeneous gels [49]. However, due to the considerations mentioned in the Introduction, we will not use their approach in the present work.

Note that the N_{Gel} and the so-called number of segments per chain N_{Seg} are described by the same relationship $N_{Gel} = N_{Seg} = V_0 \phi_0 N_A / (v_s N_C)$, where N_A is the Avo-

gadro's constant, v_s the molar volume of solvent, for water $v_s = 18 \text{ cm}^3/\text{mol}$, $V_0 = (4/3)\pi R_{\text{H},0}^3$ and N_{C} the number of chains in the polymer network [18,30]. Here, we use N,N'-methylenebisacrylamide (BIS) as cross-linker. There are two reactive vinyl units for each BIS molecule. Hence, each cross-linker molecule connects pairs of chains [30]. Therefore the number of chains N_C in the polymer network is given by $N_C = 2N_{BIS}$, where N_{BIS} is the number of BIS molecules in the network. The total number of monomers N_M (sites of lattice) occupied by the polymer network is given by [18]:

$$N_M = \frac{\phi_0 V_0 N_A}{v_s}.$$
(3)

If the nominal amount [BIS] has been incorporated, N_{BIS} is given by:

Ν

$$J_{BIS} = \frac{N_M[\text{BIS}]}{100 \text{ mol}\%} \tag{4}$$

where [BIS] is the nominal amount of the cross-linker in mol%. With these notations we can apply $N_{Gel} = N_M / N_C$ and:

$$N_{Gel} = \frac{50 \text{ mol}\%}{[\text{BIS}]}.$$
(5)

 N_{Gel} vs. [BIS] for the different microgel particles is presented in Figure 1. In all cases, N_{Gel} decreases almost hyperbolically with the increasing concentration of BIS.



Figure 1. The average degree of polymerization between cross-links N_{Gel} vs. concentration of N_rN' -methylenebisacrylamide (BIS) in p(NNPAM) (blue), p(NIPAM) (green), and p(NIPMAM) (red) particles. The parameter N_{Gel} was calculated using the Flory–Rehner Equation (2) with the Hill-like Equation (6) for the interaction parameter χ . The theoretical N_{Gel} (black) was calculated with Equation (5). To better see the hyperbolic dependence between N_{Gel} and [BIS], the theoretical N_{Gel} is shown in the inset. More details will be discussed later.

However, since the BIS cross-linkers in the microgel are not homogeneously distributed [20], Equation (5) describes the average degree of polymerization N_{Gel} only qualitatively, as discussed in the Results and Discussion.

Hill-Like Equation

It is obvious that the VPT of strongly cross-linked microgels must be cooperative, since the meshes of the network are mechanically coupled. One can imagine that in a network with an inhomogeneous distribution of N_{Gel} , a polymer chain first collapses locally at volume phase transition and water molecules leave the chain. Since the col-

lapsed chain is mechanically coupled to the neighboring chains, the VPT is induced in the neighboring chains.

To describe this cooperativity, we use the Hill model for a cooperative aggregation to model the Flory–Huggins interaction parameter χ [50]. The dependence of χ on the temperature *t* is modeled by the Hill-like function for a cooperative thermotropic structure transition with linear baseline:

$$\chi(t) = \chi_0 + a(t - t_a) + b \frac{t_{rel}^{\nu}}{t_{rol}^{\nu} + K}$$
(6)

where χ_0 is the value of the χ parameter at $t = t_a$, t_a is the first temperature data point, and t_e is the last (end) temperature point of the data set, a is the slope of the baseline, b is the dimensionless amplitude parameter of the Hill transition, K is the half-saturation constant, ν is the Hill parameter, and $t_{rel}(t)$ is the relative temperature given by:

$$t_{rel}(t) = \frac{t - t_a}{t_e - t_a}.$$
(7)

Note that the relative temperature $t_{rel}(t)$ is used as an analog of polymer concentration. Therefore, the concentration dependence of the original Flory–Huggins parameter is preserved. The Hill parameter ν is the stoichiometric coefficient of the reaction:

1

$$PS_{\nu} \rightleftharpoons P + \nu S. \tag{8}$$

The symbols *P* and νS denote the states of the polymer and solvent after the VPT, respectively, and the symbol PS_{ν} denotes the aggregate state at the onset of the VPT. In contrast to a lyotropic transition, for the thermotropic transition the concentration of species is replaced by the relative temperature t_{rel} changing in the region $0 \le t_{rel}(t) \le 1$. If $\chi > 0.5$, the polymer–solvent binding is no longer energetically favorable. The gel particles collapse at VPT (thermo-shrinking gels). If $\chi < 0.5$, the free energy of the binding decreases and the microgels are swelling [18].

3. Results and Discussion

The hydrodynamic radius $R_{\rm H}$ as a function of temperature t calculated using the Flory–Rehner Equation (2) with the Hill-like Equation (6) for χ was fitted to the swelling curves at different concentrations of BIS in the range of $2.5 \leq [BIS]/mol\% \leq 15.0$. Here we studied microgel particles of p(NNPAM) (Figures S1 and S2 in the Supplementary Materials), p(NIPAM) (Figures S3 and S4 in the Supplementary Materials), and p(NIPMAM) (Figures S5 and S6 in the Supplementary Materials). The parameters ϕ_0 , N_{Gel} , K, and ν were obtained as fitting parameters; see Tables S1-S3 (also given in the Supplementary Materials). The polymer volume fraction ϕ_0 in the collapsed state was practically independent of the BIS concentration for all microgel types. The average value of ϕ_0 for p(NNPAM) was 0.75 ± 0.02 , for p(NIPAM) it was 0.72 ± 0.02 , and for p(NIPMAM) it was 0.72 ± 0.03 . These values are still controversial since several small-angle neutron scattering (SANS) studies gave values between 0.4 and 0.6 at the particle center [39,51] for ϕ_0 . However, for a direct comparison, the p(r) density functions obtained from SANS (or SAXS) need to be integrated over the microgel radius [39]. Moreover, the difference might be due to different weighting of the contribution of the outer fuzzy regions in SANS compared to photon correlation spectroscopy (PCS). In PCS-based studies, a value of $\phi_0 = 0.8$ has been reported [19,35,42,46,52–54]. In Equation (6), the parameters a, K, b, and χ_0 were also independent of BIS concentration. Parameters *a*, *K*, *b*, and χ_0 were fitted for one swelling curve within a series and the values obtained from this first fit and were subsequently used as input parameters and held constant for all other fits within the series. Fitting of the data was performed with the software Mathcad Prime 6.0. As an example, Figure 2 shows a selection of the swelling curves fitted with Equation (2), and Table 1 lists the



respective results of the fits for p(NNPAM), p(NIPAM), and p(NIPMAM) microgels with the respective cross-linker concentrations of 5 mol%, 10 mol%, and 15 mol%.

Figure 2. Hydrodynamic radius $R_{\rm H}$ vs. temperature *t* at different concentrations of the crosslinker BIS in p(NNPAM) (**a**–**c**), p(NIPAM) (**d**–**f**), and p(NIPMAM) (**g**–**i**) particles. The fit of the hydrodynamic radii $R_{\rm H}$ (solid curve) was calculated using the Flory–Rehner Equation (2) with the Hill-like Equation (6) for the interaction parameter χ . A nearly perfect fit to the experimental data was achieved for all measured BIS concentrations.

Table 1. Parameters resulting from the fit of swelling curves with the Hill-like Equation (6) and the volume phase transition temperature (VPTT) determined from the inflection point of the $R_{\rm H}(t)$ -swelling curve.

	[BIS]/mol%	<i>K</i> /10 ⁻⁴	ν	$N_{\rm Gel}$	$t_{0.5}/^{\circ} C$	VPTT/°C
p(NNPAM)	5.00	0.057	14.0	153.5	23.2	23.2
	10.0	20.00	8.16	43.28	24.1	23.9
	15.0	260.0	4.77	19.43	24.4	24.0
p(NIPAM)	5.00	0.307	22.7	213.8	35.4	35.7
	10.0	20.00	15.6	57.40	36.8	37.1
	15.0	230.0	10.5	27.50	37.8	38.6
p(NIPMAM)	5.00	8.000	32.1	209.1	48.3	48.0
-	10.0	80.00	20.2	60.67	47.6	47.8
	15.0	220.0	13.8	22.40	46.1	46.4

The statistical significances of the fit parameters ν , K, and N_{Gel} were estimated using the chi-square functions $(chi)^2$, Equation (S.1), and the normalized partial derivatives of the $(chi)^2$ (see Figure S7 and explanations in the Supplementary Materials). The three parameters ν , K, and N_{Gel} have similarly high sensitivities to small deviations around the minimum. Such sensitivities allow a reliable determination of the fit parameters. In almost all fits, very small values for $(chi)^2$ were obtained. For instance, for p(NNPAM) the values of $(chi)^2$ were on average 6.0 with 61 data points of the swelling curve, for p(NIPAM) $(chi)^2 = 1.0$ with 41 data points, and for p(NIPMAM) $(chi)^2 = 7.4$ with 32 data points.

Only the $R_{\rm H}(t)$ -swelling curves of p(NNPAM) at [BIS]= 2.5 mol% and 5 mol% show a partly discontinuous VPT. These curves could not be satisfactorily described by the Flory-Rehner theory with the Hill-like model for the interaction parameter χ (Figure 2 and S1).

The swelling curves in all microgel types show that the VPT became increasingly smeared as the cross-linker concentration increased. This behavior is also evident in the curve of the interaction parameter as a function of temperature $\chi(t)$ (Figure 3). Wu et al. suggested that this phenomenon can be explained by the inhomogeneous radial distribution of the degree of polymerization N_{Gel} within the gel particles [20]. They assumed that each polymer chain between two cross-linkers has its own volume phase transition temperature. Accordingly, if the distribution of N_{Gel} is very narrow, the VPT is discontinuous. This is the situation for some macroscopic gels [20,55–57]. According to Wu et al., larger N_{Gel} have a lower VPTT compared to the VPTT at smaller N_{Gel} [20]. Additionally, the microgels showed a sharper VPT at a low cross-linker concentration, suggesting that the distribution of N_{gel} was narrow. With an increase of the cross-linker concentration, N_{Gel} decreased, so, in line with Wu et al., the VPTT should increase [20]. However, the increase of the VPTT was only observed for p(NNPAM) and p(NIPAM) (Table 1). Contrary to p(NNPAM) and p(NIPAM), the VPTT of p(NIPMAM) decreased with the increase of the cross-linker concentration (Table 1).



Figure 3. Interaction parameter χ vs. temperature *t* at different concentrations of BIS in p(NNPAM) (**a**–**c**), p(NIPAM) (**d**–**f**), and p(NIPMAM) (**g**–**i**) particles. The interaction parameter χ was calculated with the Hill-like Equation (6).

3.1. Average Degree of Polymerization N_{Gel}

The values of N_{Gel} obtained from the analysis of swelling curves are systematically larger than given by Equation (5) (Figure 1). If N_{Gel} was radius-independent in the microgel particles, the experimentally obtained values of N_{Gel} could be correctly described by Equation (5). The radial dependence of N_{Gel} in p(NNPAM), p(NIPAM), and p(NIPMAM) particles was also observed by Bergmann et al. and Wrede et al. using super-resolution fluorescence microscopy [37,38].

3.2. Half-Saturation Constant K

The half-saturation constant is given by:

$$K = \left(\frac{t_{0.5} - t_a}{t_e - t_a}\right)^{\nu}.$$
(9)

Assuming that half of the water molecules have left the gel at the volume phase transition, *K* is then the concentration of PS_{ν} at the VPTT. Hence, the half-temperature $t_{0.5}$ corresponds to the VPTT. Tables S1–S3 clearly show that the $t_{0.5}$ and VPTT were almost identical in all cases studied here.

3.3. Hill Parameter v

For all microgels studied in the present work, ν linearly decreased with the increase of [BIS] (Figure 4). The results suggest that ν represents the number of water molecules per segment with a certain length cooperatively leaving the segment at the VPT (Figure 5).



Figure 4. The Hill parameter ν vs. concentration of BIS in p(NNPAM) (blue), p(NIPAM) (green), and p(NIPMAM) (red) particles. The Hill parameter ν has been calculated with the Hill-like Equation (6) for the interaction parameter χ . The dashed lines are guides to the eye.



Figure 5. Schematic representation of the water molecules attached to an arbitrary polymer segment. It is assumed that the water molecules bound to the monomer units E leave the segment at the VPT while the water molecules bound to the cross-linker D remain at their position because the cross-linker has no LCST. Segment A contains only one cross-linker. In segment B the concentration of the cross-linker is two times higher and in segment C it is four times higher than in segment A. The number of water molecules that leave the microgel per segment at the VPT should decrease linearly with the increase of the cross-linker concentration.

Since the cross-linker molecules have no LCST, we assume that the water molecules bound to the cross-linker do not leave the segment at the VPT. Indeed, Figure 4 suggests that the number of water molecules cooperatively leaving the segment decreased linearly with the increasing cross-linker concentration.

The values of ν in Figure 4 show that p(NIPMAM) bound about 40% more water compared to p(NIPAM) microgel particles. P(NIPAM) bound about 70% more water in relation to the p(NNPAM) segments. Hence, the calculated ν parameter nicely follows the differences in hydrophilicity of the three polymers. This observation is in line with the fact that the p(NNPAM) microgels were more hydrophobic than p(NIPAM) microgels, and p(NIPAM) microgels were more hydrophobic than p(NIPAM) microgel particles. The hydrophobicity of the gels can be seen from the VPTT. The lower the VPTT, the larger the hydrophobicity of the gels. The relatively small amount of water molecules leaving the microgel per segment at the VPT is at least qualitatively consistent with the results of recent molecular dynamics (MD) simulations [58–60].

4. Conclusions

We have shown that using the Flory-Rehner theory and the Hill-like model for the interaction parameter χ , the swelling behavior of the p(NNPAM), p(NIPAM), and p(NIPMAM) microgels with different BIS concentrations can be quantitatively described. The Hill-like model for χ provides a deeper insight into the volume phase transition of acrylamide microgels than the use of series expansions for the interaction parameter χ [1]. Interestingly, the original approach of Flory and Huggins for the monomer-solvent interaction does not apply successfully to some microgels [32]. The reason for this is obviously that the cross-linking enhances the cooperativity of the volume phase transitions because all meshes of the network are mechanically coupled. This was ignored in previous descriptions arguing with distinct transition temperatures for individual meshes to describe the broad continuous character of the transition of microgels. Moreover, we observed a linear relationship between the Hill parameter ν and the BIS concentration. This linearity suggests that the Hill parameter ν corresponds to the number of water molecules per network chain that cooperatively leave the chain at the volume phase transition. Additionally, we found that the Hill parameter increased with increasing microgel hydrophilicity. In the future, this approach will be applied to other acrylamide-based microgels, and to copolymer particles as well. Additional improvements might be possible by accounting for the surface charges stemming from the initiator.

5. Materials and Methods

5.1. Chemicals

N-n-propylacrylamide (NNPAM) was synthesized via a Schotten–Baumann reaction published by Hirano et al. [61]. For this reaction, acryloylchloride (Sigma-Aldrich Chemie GmbH, Munich, Germany; purity 98%), n-propylamine (Fluka, Sigma-Aldrich Chemie GmbH, Munich, Germany; purity 99%), triethylamine (Grüssing GmbH Analytika, Filsum, Germany; purity 99%), and methylenechloride (p.a.) were used as received. The obtained monomer NNPAM was washed with NaHCO₃ (10 wt%) and dried over MgSO₄. After filtration, the solvent was evaporated and the product was distilled in vacuum (115 °C, 10 mbar). N-isopropylacrylamide (NIPAM; Sigma-Aldrich Chemie GmbH, Munich, Germany; purity 97%) and N-isopropylmethacrylamide (NIPMAM; Sigma-Aldrich Chemie GmbH, Munich, Germany; purity 97%) were purified by recrystallization from hexane. The cross-linker N,N'-methylenebisacrylamide (BIS; Sigma-Aldrich Chemie GmbH, Munich, Germany; purity 99%), the initiator ammonium persulfate (APS; Sigma-Aldrich Chemie GmbH, Munich, Germany; purity \geq 98%), and pyrene (Sigma-Aldrich Chemie GmbH, Munich, Germany; purity \geq 99%) were used without further purification. For all experiments, purified water from an Arium pro VF system (Sartorius AG, Göttingen, Germany) was used.

5.2. Synthesis of the Homopolymer Microgels

The homopolymer microgels of NNPAM, NIPAM, and NIPMAM were synthesized via conventional precipitation polymerization without surfactant. All syntheses were performed in a 250 mL three-neck flask equipped with a reflux condenser, mechanical stirrer (210 rpm), and a nitrogen inlet. The monomer (11.05 mmol) and the cross-linker N,N'-methylenebisacrylamide (BIS) (2.5 mol%, 5.0 mol%, 6.75 mol%, 7.5 mol%, 8.75 mol%, 10.0 mol%, 11.25 mol%, 13.75 mol%, 15.0 mol% respective to the total monomer amount) were dissolved in 150 mL purified water and heated to 70 °C under continuous stirring and purged with nitrogen. After 1 h the polymerization was initiated by the addition of 2 mL of the 0.2 M solution of APS and left to proceed for 4 h at 70 °C. Subsequently, the solution was cooled to room temperature and stirred overnight. For purification, all samples were treated by four cycles of centrifugation, decantation, and redispersion in purified water using a JA-30.50 Ti Rotor in an Avanti J-30I centrifuge (Beckman Coulter, Brea, CA, USA) at 20,000 rpm and 25 °C.

5.3. Photon Correlation Spectroscopy (PCS)

The PCS measurements were performed on a custom-built fixed-angle setup (scattering angle θ : 60°) utilizing a He–Ne Laser (wavelength λ = 632.8 nm, 21 mW, Thorlabs, Newton, MA, USA) and two photomultipliers (ALV/SO-SIPD, ALV-GmbH, Langen, Germany) in a pseudo-cross-correlation configuration. The signal was correlated with an ALV-6010 multiple-tau correlator (ALV GmbH, Langen, Germany). Subsequently, the intensity–time correlation functions were converted to the field–time correlation function $g^1(t)$ and analyzed using the CONTIN software [62]. However, an analysis using a second-order cumulant function also leads to the same result within the exp. precision. The temperature was controlled via a thermostat (Phoenix II, Thermo Fisher Scientific, Waltham, MA, USA together with Haake C25P, Thermo Fisher Scientific, Waltham, MA, USA), and the sample was equilibrated for 25 min inside the decaline-filled refractive index matching bath. For each temperature, 5 consecutive measurements were performed. The obtained mean relaxation rates Γ of the $g^1(t)$ functions were converted to the hydrodynamic radius by

$$R_h = \frac{k_B T}{6\pi \eta \frac{\Gamma}{c^2}} \tag{10}$$

Here, k_B is the Boltzmann constant, η the solvent viscosity (water), T the temperature in Kelvin, and $q = \frac{4\pi n}{\lambda} \sin \frac{\theta}{2}$ the magnitude of the scattering vector. n is the refractive index of the solvent.

Supplementary Materials: The following are available online at https://www.mdpi.com/article/10 .3390/gels7020042/s1. Figure S1: Hydrodynamic radius $R_{\rm H}$ vs. temperature t, measured (points) and calculated (lines), at different concentrations of BIS in p(NNPAM) particles. Figure S2: Interaction parameter χ vs. temperature t at different concentrations of BIS in p(NNPAM) particles. Figure S3: Hydrodynamic radius $R_{\rm H}$ vs. temperature t, measured (points) and calculated (lines), at different concentrations of BIS in p(NIPAM) particles. Figure S4: Interaction parameter χ vs. temperature t at different concentrations of BIS in p(NIPAM) particles. Figure S5: Hydrodynamic radius $R_{\rm H}$ vs. temperature t, measured (points) and calculated (lines), at different concentrations of BIS in p(NIPAM) particles. Figure S5: Hydrodynamic radius $R_{\rm H}$ vs. temperature t, measured (points) and calculated (lines), at different concentrations of BIS in p(NIPMAM) particles. Figure S6: Interaction parameter χ vs. temperature t at different concentrations of BIS in p(NIPMAM) particles. Figure S7: Chi-Square $(chi)^2$ vs. the fit-parameters v, K and N_{Gel} of the $R_{\rm H}(t)$ -swelling curve from p(NIPAM) with a concentration of 2.5 mol% BIS. Table S1: Parameters resulting from the fit of the swelling curves of p(NIPAM) with the Hill-like Equation (6). Table S3: Parameters resulting from the fit of the swelling curves of p(NIPAM) with the Hill-like Equation (6).

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