# **Divalent Ion Trapping Inside Potassium Channels of Human T Lymphocytes**

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ABSTRACT Using the patch-clamp whole-cell recording technique, we investigated the influence of external  $Ca^{2+}$ ,  $Ba^{2+}$ ,  $K^+$ ,  $Rb^+$ , and internal  $Ca^{2+}$  on the rate of  $K^+$  channel inactivation in the human T lymphocyte-derived cell line, Jurkat E6-1. Raising external Ca<sup>2+</sup> or Ba<sup>2+</sup>, or reducing external K<sup>+</sup>, accelerated the rate of the K<sup>+</sup> current decay during a depolarizing voltage pulse. External  $Ba^{2+}$  also produced a use-dependent block of the  $K<sup>+</sup>$  channels by entering the open channel and becoming trapped inside. Raising internal  $Ca^{2+}$  accelerated inactivation at lower concentrations than external  $Ca^{2+}$ , but increasing the  $Ca^{2+}$  buffering with BAPTA did not affect inactivation. Raising  $[K^+]_o$  or adding Rb<sup>+</sup> slowed inactivation by competing with divalent ions. External  $Rb<sup>+</sup>$  also produced a use-dependent removal of block of  $K^+$  channels loaded with  $Ba^{2+}$  or  $Ca^{2+}$ . From the removal of this block we found that under normal conditions  $\sim$ 25% of the channels were loaded with Ca<sup>2+</sup>, whereas under conditions with 10  $\mu$ M internal Ca<sup>2+</sup> the proportion of channels loaded with  $Ca^{2+}$  increased to ~50%. Removing all the divalent cations from the external and internal solution resulted in the induction of a nonselective, voltage-independent conductance. We conclude that  $Ca<sup>2+</sup>$  ions from the outside or the inside can bind to a site at the  $K^+$  channel and thereby block the channel or accelerate inactivation.

# INTRODUCTION

Voltage-dependent  $K^+$  channels in human T lymphocytes inactivate almost completely over a period of hundreds of milliseconds during depolarizing voltage steps (DeCoursey et al., 1984; Matteson and Deutsch, 1984). The inactivation rate is sensitive to the presence of divalent ions in the bathing solution; elevated external  $Ca<sup>2+</sup>$ or added lanthanum increase, whereas manganese and cobalt decrease, the rate of inactivation (DeCoursey et al., 1985). One possible explanation for these results is that  $Ca<sup>2+</sup>$  or lanthanum can gain access to a site that inactivates the channel, whereas manganese and cobalt may prevent  $Ca<sup>2+</sup>$  from gaining access to this site. Increasing  $[Ca^{2+}]$  also reduces the  $K^+$  conductance of T cells (Bregestovski et al., 1986) and accelerates inactivation of  $K^+$  channels of B cells (Choquet et al., 1987). Thus, both external and internal  $Ca^{2+}$  regulate the effective number of  $K^+$  channels.

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The purpose of the experiments reported here is to define mechanisms of ion interaction with  $K<sup>+</sup>$  channels. The similarity of effects when raising external or internal  $[Ca<sup>2+</sup>]$  suggests a common site of action. In some nerve preparations, inactivation of voltage-gated  $Ca^{2+}$  channels depends on  $Ca^{2+}$  entry through open  $Ca^{2+}$ channels and intracellular accumulation, rather than on membrane potential alone, a phenomenon known as  $Ca^{2+}$ -dependent inactivation of the  $Ca^{2+}$  channel (for review see Eckert and Chad, 1984). An analogous mechanism in which sparingly permeant  $Ca^{2+}$  ions pass through the K<sup>+</sup> channel and lead to inactivation might  $\alpha$  account for the effects of external di- and polyvalent cations on the rate of  $K^+$  channel inactivation described above. We therefore investigated whether the model of  $Ca^{2+}$ -dependent inactivation of  $Ca^{2+}$  channels could be also applied to K<sup>+</sup> channel inactivation of human T lymphocytes. Our experiments suggest that rather than modifying a site accessible to intracellular  $Ca^{2+}$  buffers, divalent ions can become trapped inside K<sup>+</sup> channels. Since  $[Ca^{2+}]$ <sub>i</sub> is known to rise shortly after mitogen addition, the interaction of ions within the  $K<sup>+</sup>$  channel may be relevant to signal transduction mechanisms. Some of the results have been reported in preliminary communications (Grissmer and Cahalan, 1987, 1988).

### METHODS

The experiments were carried out on single cells of a human T lymphocyte-derived cell line, Jurkat E6-1, using the whole-cell recording mode of the patch<lamp technique (Hamill et al., 1981). All experiments were done at room temperature (22-26°C).

#### *Solutions*

The cells under investigation were bathed in normal Ringer solution. External solutions containing varying concentrations of  $Ca^{2+}$ ,  $Ba^{2+}$ ,  $K^+$ , and  $Rb^+$  are listed in Table I. The chamber volume contained  $\sim$ 250  $\mu$ l of solution which could be totally exchanged during the recordings by bath perfnsion within 15-20 s. The patch-pipette usually contained 160 mM K-aspartate with  $[Ca<sup>2+</sup>]$  buffered to  $10<sup>-8</sup>$  M (see Table I). In some experiments the pipette solution contained higher  $[Ca<sup>2+</sup>]$ .

#### *Data Acquisition*

The holding potential was adjusted in all experiments to  $E = -80$  mV. The patch-clamp amplifier (either List L/M-EPC 7, Adams and List Associates, Ltd., Great Neck, NY; or Axon Instruments, Inc., Axopatch, Burlingame, CA) was used in the voltage-clamp mode without series resistance compensation. Electrodes were pulled from Accufill 90 Micropets (Becton, Dickinson & Co., Parsippany, NJ) in three stages, coated with Sylgard (Dow Coming Corp., Midland, MI) and fire polished to resistances, measured in the bath, of  $2-7$  M $\Omega$ . K<sup>+</sup> channel inactivation was not different if hard glass was used as pipette glass (compare Cota and Armstrong, 1988; Furman and Tanaka, 1988).

In all experiments, the command input of the patch-clamp amplifier was controlled by a computer (PDP 11/73) via a digital-analog converter, and membrane currents were recorded at a bandwidth of 2 kHz. Correction for linear leakage and capacitative currents was achieved by analog subtraction and by digital subtraction of an appropriately scaled mean current associated with eight pulses delivered from a hyperpolarized potential. All potentials were corrected for the liquid junction potential that develops at the tip of the pipette if the pipette solution is different from that of the bath. The liquid junction potential between the normal internal (pipette) and external (bath) solution was  $-13$  mV.



External $Cl^-$ anion (pH 7.4)	Na <sup>+</sup>	$\mathbf{K}^+$	$Ca2+$	$Mg^{2+}$	<b>HEPES</b>	<b>EGTA</b>	$Ba2+$	$Rb$ <sup>+</sup>
Normal ringer $Ca2+$	160	4.5	$\overline{2}$	l	5			
"0"	160	4.5		$\boldsymbol{3}$	5	1		
0.5	160	4.5	0.5	2.5	5			
4.6	156	4.5	4.6	$\mathbf{I}$	5			
12	144	4.5	12.4	$\mathbf{I}$	5			
54	80	4.5	53.8	$\mathbf{I}$	5			
106		4.5	105.7	$\mathbf{1}$	5			
$\mathbf{K}^+$								
$\pmb{0}$	164.5		$\overline{2}$	1	5			
2.2	162.2	2.2	$\overline{2}$	ı	5			
10	154.5	10	$\overline{\mathbf{2}}$	ı	5			
20	144.5	20	$\overline{\mathbf{2}}$	ı	5			
40	124.5	40	$\overline{\mathbf{2}}$	$\mathbf{1}$	5			
80	84.5	80	$\mathbf 2$	1	5			
133	31.5	133	$\overline{\mathbf{2}}$	$\mathbf{I}$	5			
140	24.5	140	$\overline{\mathbf{2}}$	$\mathbf{1}$	5			
160	4.5	160	$\overline{\mathbf{2}}$	1	5			
Rb <sup>+</sup> -Ringer $Ba2+$	4.5		$\overline{2}$	$\mathbf{1}$	5		¢	160
0.3	160	4.5		2.7	5		0.3	
$\overline{2}$	160	4.5		$\mathbf{1}$	5		2.0	
10	144	4.5	$\boldsymbol{2}$	$\mathbf{I}$	5		10.4	
54	80	4.5	$\boldsymbol{2}$	$\mathbf{1}$	5		53.8	
106		4.5	$\overline{\mathbf{2}}$	$\mathbf{1}$	5		105.7	
Internal								
(pH 7.2)	$\mathbf{K}^+$	$Ca2+$		$Mg^{2+}$	<b>HEPES</b>	<b>EGTA</b>	$Cl^-$	$Asp^-$
$[Ca^{2+}]_i = 10^{-8} M$	160	0.1		$\overline{2}$	5	1.1	4.2	160
$[Ca2+]$ <sub>i</sub> - 10 <sup>-5</sup> M	160	1.1		$\bf 2$	5	1.1	4.2	160
$[Ca^{2+}]_i = 2 \times 10^{-5}$	160	2.0		I	5		166	

The measured osmolarity of the solutions was 290-320 mosmol. Asp<sup>-</sup>, aspartate-; HEPES, N-2-hydroxyethylpiperazine-N'-2-ethanesulfonic acid; EGTA, ethylene glycolbis( $\beta$ -aminoethylether)N,N,N',N'-tetraacetic acid. [Ca<sup>2+</sup>]<sub>i</sub> was calculated assuming a dissociation constant for EGTA and  $Ca<sup>2+</sup>$  at pH 7.2 of  $10<sup>-7</sup>$  M according to Portzehl et al., 1964; in some experiments the intracellular  $Ca^{2+}$  buffering was increased by including 55 mM 1,2-bis(2-aminophenoxy)-ethane-N,N,N',N'-tetraacetic acid (BAPTA) and 5 mM CaCl<sub>2</sub> (free  $[Ca^{2+}]_t = 10^{-8}$  M) in the pipette solution. All concentrations are in millimolar.

# *Analysis*

Exponentials were fitted to the decay of the  $K<sup>+</sup>$  current as described in more detail by Cahalan et al. (1985) by inputting values through potentiometers. Attempts to fit the decay of the  $K^+$  current in high  $[K^+]$ <sub>o</sub> with a single exponential failed to describe the time course of the  $K<sup>+</sup>$  channel inactivation sufficiently. The fit, however, was improved by including a certain amount of baseline offset. This is consistent with the idea that there are  $K<sup>+</sup>$  channels that do not inactivate during a depolarizing pulse, or that there is always a certain number of open  $K<sup>+</sup>$  channels during depolarizations due to an equilibrium between the rate of activation and inactivation. Even with this fitting procedure, including an offset, the decay of the  $K^+$  current could not be adequately characterized, especially during the first 300-500 ms after the onset of the depolarizing pulse. Therefore, the decay of the  $K^+$  current in high  $[K^+]_o$  was described with a double exponential function.

# RESULTS

# *External Ca*<sup> $2+$ </sup> and Ba<sup> $2+$ </sup> Enhance Inactivation

Raising  $[Ca^{2+}]_o$  reversibly increases the rate of  $K^+$  channel inactivation in Jurkat cells, as illustrated by  $K^+$  currents in Fig. 1. Raising  $\left[Ca^{2+}\right]_0$  from 2 to 12 mM  $Ca^{2+}$ consistently reduced the peak current amplitude  $(I_{Kpeak}$ , from 1,314 to 1,200 pA in the cell of Fig. 1), as well as speeding the  $K<sup>+</sup>$  current decay. The values for the inactivation time constants  $\tau_h$ , obtained by single exponential fitting (see Methods), in normal Ringer solution before and after the  $Ca^{2+}$  treatment were 230 and 226



ms, respectively;  $\tau_h$  in Ringer solution with 12 mM Ca<sup>2+</sup> was 129 ms. While  $I_{Kpeak}$  in 12 mM Ca<sup>2+</sup> Ringer is reduced to ~90% of the control value,  $\tau_h$  is almost 1.8 times faster than in normal Ringer solution. Adding external  $Ba<sup>2+</sup>$  results in similar effects on the time course of inactivation, but also induces a long-lasting "use-dependent" block of  $K^+$  currents, as described below.

Raising  $[Ca<sup>2+</sup>]_{o}$  may change the negative surface potential of the cell by neutralizing surface charges, thereby shifting the voltage dependence of  $K^+$  channel gating to more depolarized potentials. To assess the surface potential effect,  $I_K(E)$  curves were measured and the corresponding  $g_k(E)$  curves were calculated. These two curves are shown in Fig. 2, A and B. The  $g<sub>K</sub>(E)$  curve in 12 mM Ca<sup>2+</sup> is shifted by 12 mV to more depolarized potentials, and  $g_{Kmax}$  is reduced to 93% of the control values before and after the  $Ca<sup>2+</sup>$  treatment.

To quantitate the influence of raising  $[Ca^{2+}]_0$  on the inactivation of the K<sup>+</sup> channel, we measured  $K^+$  currents associated with pulses of different amplitudes. The inactivation time constants,  $\tau_h$ , shown in Fig. 3 were slightly voltage dependent in normal Ringer, confirming earlier measurements on human T lymphocytes (see Fig. 6 of Cahalan et al., 1985). Raising Ca<sup>2+</sup> shifts the voltage dependence of  $\tau<sub>h</sub>$  by an GRISSMER AND CAHALAN Lymphocyte *K<sup>+</sup> Channel Inactivation* 



FIGURE 2. Shift in  $K^+$  channel activation in 12  $Ca^{2+}$  Ringer. (A) Current-voltage relations before ( $\Box$ ), during ( $\triangle$ ), and after (X) application of 12  $Ca<sup>2+</sup> Ringer. [Ca<sup>2+</sup>] = 10<sup>-8</sup> M.$ Peak  $K^+$  currents obtained with 2-s depolarizing pulses are plotted against the absolute membrane potential, E, in mV. The lines through the points were drawn by eye for clarity.  $(B)$  Peak  $K^+$  conductance-voltage relations before ( $\square$ ), during ( $\triangle$ ) and after (X) application of  $12 \text{ Ca}^{2+}$  Ringer.

The lines through the points were fitted with the Boltzmann equation:

 $g_K(E) = g_{K \text{ max}} / \{1 + \exp[(E_n - E) / k]\}^4$ 

with the following parameters:



Same cell as in Fig. 1.



FIGURE 3. Inactivation time constants  $\tau_h$  of K<sup>+</sup> currents before ( $\Box$ ), during  $(\triangle)$ , and after  $(X)$  application of 12 Ca<sup>2+</sup> Ringer.  $[Ca^{2+}]_i = 10^{-8}$  M. Pulses were given every minute to allow complete recovery from inactivation. The time constants were obtained by fitting a single exponential to the decay of the  $K<sup>+</sup>$  currents. The  $\bullet$  were obtained by correcting the  $\tau_h$  values measured in 12 Ca<sup>2+</sup> for a shift along the voltage axis determined from the shift of the conductance-voltage curve in  $12 \text{ Ca}^{2+}$  compared with normal Ringer solution. The lines through the points were drawn by eye for clarity. Same cell as in Figs. 1 and 2.

amount comparable to the shift of activation shown in Fig. 2. In addition to this shift,  $\tau_h$  is reduced in 12 mM Ca<sup>2+</sup>-Ringer. The combination of both effects causes crossover of the control curve and the curve obtained in  $12 \text{ mM } Ca^{2+}$ . To correct for surface potential effects, the  $\tau_h$  values obtained in 12 mM Ca<sup>2+</sup> (labeled "12 Ca uncorrected") were shifted to more hyperpolarized potentials by the same amount that the  $g_k(E)$  curve was shifted to more depolarized potentials (12 mV) when  $[Ca^{2+}]$ <sub>o</sub> was raised from 2 to 12 mM ("12 Ca corrected"). The reduction of  $\tau_h$  was then determined by calculating the ratios of the control  $\tau_h$  and the corrected  $\tau_h$  in 12 mM  $Ca<sup>2+</sup>$ . At  $E = 10$  mV the reduction was to 52% of the control value, whereas at  $E = 50$  mV the reduction was to 57% of the control value. This means that, aside from a shift, the  $Ca<sup>2+</sup>$  effect on the inactivation time constant shows very little, if any, voltage dependence.



FIGURE 4. The change  $p$  of the inactivation time constant  $\tau_h$  in varying concentrations of  $Ca<sup>2+</sup>$  ( $\square$ ) or  $Ba<sup>2+</sup>$  ( $\triangle$ ). All external solutions contain 4.5 mM  $[K^+]_{\alpha}$ .  $[Ca^{2+}]_{\alpha} = 10^{-8}$  M. p is the normalized change of  $\tau_{\rm h}$ , defined as

$$
p = (\tau_{\rm h} - \tau_{\rm hmin})/(\tau_{\rm hmax} - \tau_{\rm hmin})
$$

with  $\tau_{\text{hmax}} = 430$  ms and  $\tau_{\text{hmin}} = 65$ ms. Divalent ion concentration  $[X^{2+}]_o$  is plotted on a logarithmic scale. The bars on the points reflect the standard deviation, each point reflects the mean of at least three measurements. The curve through the points was calculated assuming that  $p$  is proportional to  $K_d$  $([X^{2+}]_{o} + K_{d})$  with  $K_{d} = 3.5$  mM for  $Ca^{2+}$ , and  $K_d = 1.5$  mM for  $Ba^{2+}$ .

Fig. 4 shows the change in  $\tau_h$  (p, as defined in the legend) in Ringer solutions with different  $[Ca^{2+}]_{\text{o}}$  or  $[Ba^{\overline{2}+}]_{\text{o}}$ . The results indicate that raising  $[Ca^{2+}]_{\text{o}}$  from  $10^{-9}$  M to 100 mM accelerates the rate of inactivation more than fivefold. The change is normalized such that  $p = 1$  corresponds to the maximum inactivation time constant,  $\tau_{\text{hmax}}$ , of 430 ms;  $p = 0$  corresponds to the minimum inactivation time constant,  $\tau_{\text{hmin}}$ , of 65 ms. The estimate of  $\tau_{\text{hmax}}$  was obtained by measuring  $\tau_{\text{h}}$  in Ringer solutions with  $\left[\text{Ca}^{2+}\right]_{o} = 10^{-9}$  M and it was found to be 422  $\pm$  17 ms (n = 3), whereas  $\tau_{\text{hmin}}$  was determined by fitting the calculated curve to the measured data points.  $\tau_{\rm hmin}$  was the same for either Ca<sup>2+</sup> or Ba<sup>2+</sup>. It is concluded that Ca<sup>2+</sup> or Ba<sup>2+</sup> can bind with apparent  $K_d$ 's of 3.5 and 1.5 mM to a site with one-to-one stoichiometry (Hill plots of these curves yield a slope of 1.03 and 1.01, respectively), and thus influence  $K<sup>+</sup>$  channel inactivation.

In addition to accelerating the rate of inactivation,  $Ba<sup>2+</sup>$  ions result in progressive "use-dependent" block of  $K^+$  current. Fig. 5 A shows a  $K^+$  current in normal Ringer associated with a 1-s test pulse to 40 mV from a holding potential of  $-100$  mV. Immediately after recording this trace, the bath solution was changed to 10 mM  $Ba<sup>2+</sup>$  and  $K<sup>+</sup>$  currents elicited by identical test pulses were recorded. The first trace in 10 mM  $Ba<sup>2+</sup>$  shows an acceleration in the decay of the K<sup>+</sup> current, whereas the peak current amplitude is hardly affected, which is similar to the effect of raising external  $Ca<sup>2+</sup>$  shown in Fig. 1. Subsequent test pulses result in a progressive reduction of the peak current amplitude. This type of "use-dependent block" is similar to the effects of external  $Ba^{2+}$  on the delayed outward  $K^+$  currents in muscle fibers (Werman and Grundfest, 1961; Sperelakis et al., 1967), squid axons (Armstrong and Taylor, 1980), and myelinated nerve fibers (Arhem, 1980). The peak  $K^+$  current



FIGURE 5. Block of K<sup>+</sup> currents by 10 mM Ba<sup>2+</sup>. (A) K<sup>+</sup> currents associated with 1-s depolarizing test pulses to 40 mV from a holding potential of  $-100$  mV in normal Ringer solution (control) and immediately after application of a Ringer solution containing 10 mM  $Ba^{2+}$ .  $[Ca<sup>2+</sup>]_{<sub>1</sub>} = 10<sup>-8</sup>$  M. The interval between test pulses was 53 s, long enough to allow complete recovery from inactivation in normal Ringer solution. (B) Peak current amplitudes of the  $K^+$ currents  $(I<sub>Kpeak</sub>)$ , shown in A, at different times after the solution change from normal Ringer to 10 mM Ba<sup>2+</sup>. Cell 603.

amplitudes are plotted in Fig.  $5 B$ . A single exponential fit to these points yields an apparent time course for blocking the  $K^+$  channels by  $Ba^{2+}$ . This apparent time course strongly depends on the interval between test pulses (data not shown), suggesting that  $Ba^{2+}$  can block only if the channel is open.

Fig. 6 illustrates that removal of the  $Ba^{2+}$  block also depends on channel opening. After changing the bath solution from 10 mM  $Ba<sup>2+</sup>$  to normal Ringer solution (shown at the arrow), the  $Ba^{2+}$  block of  $K^+$  current is slowly removed as the depolarizing pulses are continued; peak  $K^+$  current recovers almost to the original level,  $-1,400$  pA. The time constant for washing out the Ba<sup>2+</sup> using this particular protocol was ~80 s. If, on the other hand, pulses are not delivered after washing out the  $Ba<sup>2+</sup>$ , the channels remain blocked, as illustrated in the right-hand side of the figure. This time, after establishing the  $Ba^{2+}$  block in 10 mM  $Ba^{2+}$  and then restoring the bath solution to normal Ringer solution, the membrane potential was held constant at  $-80$  mV for 10 min without pulses. When the same sequence of depolarizing pulses was then delivered, the time course for washing out the  $Ba^{2+}$  was the same as the one on the left side, except for the 10-min delay in which channels remained closed.  $Ba<sup>2+</sup>$  can therefore be trapped within the closed channel and cannot dissociate from the blocking site until after the channel opens. The time course for washing out the  $Ba<sup>2+</sup>$  also depends very strongly on the interval between pulses (data not shown), suggesting that the longer the channels are held open, the faster the washout. This mechanism is similar to Ba<sup>2+</sup> block of delayed outward  $K^+$  current in squid axons and myelinated nerve fibers (Armstrong and Taylor, 1980; Eaton and Brodwick, 1980; Armstrong et al., 1982; Woll, 1982) and of voltage-dependent  $Ca^{2+}$ -



FIGURE 6. Removal of the  $Ba^{2+}$  block. K<sup>+</sup> currents were elicited with the same pulse protocol as described (legend of Fig. 5), and peak  $K^+$  currents ( $I_{K_{peak}}$ ) were plotted at different times after the solution change from normal Ringer to 10 mM Ba<sup>2+</sup>. [Ca<sup>2+</sup>]<sub>i</sub> = 10<sup>-8</sup> M. The bath solutions were changed from 10 mM  $Ba<sup>2+</sup>$  to normal Ringer solution at the arrows. (A) The block of the  $K^+$  current is removed without delay, and the peak  $K^+$  current recovers almost to the original level before the  $Ba^{2+}$  treatment if depolarizing pulses are given every 53 s. Cell 604. (B) No recovery of  $I_{\text{Kpeak}}$  occurred if the membrane was not depolarized. After 10 min in normal Ringer solution the membrane was depolarized with the same pulse protocol used for monitoring the  $Ba^{2+}$  washout in A. Cell 603.

activated  $K^+$  channels from skeletal muscle (Vergara and Latorre, 1983; Miller, 1987; Miller et al., 1987).

# *Internal Ca*<sup>2+</sup> Mimics External Ca<sup>2+</sup>

Bregestovski et al. (1986), as well as Choquet et al. (1987), suggested that internal  $Ca<sup>2+</sup>$  may play a role in regulating the number of  $K<sup>+</sup>$  channels capable of being activated. In further experiments, we increased the intracellular  $Ca<sup>2+</sup>$  concentration and monitored the inactivation time constant,  $\tau_h$ , at +40 mV after breaking into the cells with pipettes filled with either  $[Ca^{2+}]_i = 10^{-8}$  M,  $10^{-5}$  M, or 2 mM, as illustrated in Fig. 7. With pipettes containing  $10^{-8}$  M Ca<sup>2+</sup> (squares) there is a small reduction of  $\tau_h$  10 min after breaking into the cell. This is probably due to the gradual disappearance of a junction potential between the pipette solution and the cytoplasm (Marty and Neher, 1983). When the patch pipettes were filled with  $\left[Ca^{2+}\right]_i =$  $10^{-5}$  M there was a much more pronounced effect of  $\tau_h$  after breaking into the cell (triangles). After 10 min, the  $\tau_h$  values with  $[Ca^{2+}]_i = 10^{-5}$  M correspond to a 1.7fold increase in the rate of inactivation, which was somewhat less than the threefold decrease in half-decay time reported by Choquet et al. (1987), and possibly due to the different test pulse voltages used. If the patch pipettes contained 2 mM  $[Ca^{2+}]_{i}$ , inactivation accelerated even further after breaking into the cell (circles). After 5 min, the  $\tau_h$  values with  $[Ca^{2+}]_i = 2$  mM corresponded to a fivefold increase in the



FIGURE 7.  $K^+$  currents measured in normal Ringer solution at different times  $t$  after breaking into the cell with varying  $[Ca<sup>2+</sup>]$  in the pipette solution. (A)  $K<sup>+</sup>$  currents were associated with depolarizing pulses to  $+40$  mV from a holding potential of  $-80$  mV at time 0, 53 s, and 106 s (time 0 means ~20 s after reaching the whole-cell configuration). The pipette solution contained 2 x 10<sup>-3</sup> M Ca<sup>2+</sup>. (B)  $\tau<sub>h</sub>$  values were obtained by fitting single exponentials to the decay of the K<sup>+</sup> currents obtained from current records similar to the one shown in A.  $\tau_h$ values were normalized to the  $\tau<sub>h</sub>$  values measured at time 0, meaning ~15 to 45 s after reaching the whole-cell configuration. The squares represent ten experiments with  $\left[\text{Ca}^{2+}\right]_{i} = 10^{-8}$ M, the triangles represent nine experiments with  $[Ca^{2+}]_i = 10^{-5}$  M, the circles represent four experiments with  $[Ca^{2+}]_i = 2$  mM. The lines through the points were drawn by eye for clarity.

rate of inactivation. With such a high internal  $Ca<sup>2+</sup>$  concentration it was impossible to record longer than 6 min after breaking into the cell.

If the binding site for  $Ca^{2+}$  that influences the rate of  $K^+$  channel inactivation is located within the membrane field, and  $Ca<sup>2+</sup>$  is able to reach this site from either side of the membrane, raising  $[Ca^{2+}]$  may alter the voltage dependence of the effect of  $[Ca^{2+}]_o$  on  $\tau_h$ . To assess this we did experiments of the type shown in Figs. 1-3, but with  $[Ca^{2+}]_i = 10^{-5}$  M. The  $\tau_h$  values obtained in normal Ringer with  $[Ca^{2+}]_i =$  $10^{-5}$  M were plotted against the membrane potential, E, as shown in Fig. 8 (open squares). Raising  $[Ca^{2+}]_o$  from 2 to 12 mM with  $[Ca^{2+}]_i = 10^{-5}$  M shifted the  $I_K(E)$ and  $g_K(E)$  curves of the K<sup>+</sup> current by 10 mV to more depolarized potentials (not shown). Therefore, it was necessary to correct the  $\tau<sub>h</sub>$  values measured with  $\left[Ca^{2+}\right]_0 =$ 12 mM and  $\left[\text{Ca}^{2+}\right]_{i} = 10^{-5}$  M (filled triangles), in a manner similar to that shown in Fig. 3. The reduction of  $\tau_h$  M Ca<sup>2+</sup> was then determined by calculating the ratios of  $\tau_h$  with normal Ringer outside and the corrected  $\tau_h$  in Ringer with 12 mM Ca<sup>2+</sup> outside. At  $E = 10$  mV, the reduction was to 63% of the control value, whereas at  $E = 50$  mV the reduction was to ~72% of the control value. This means that the effect of raising  $[\text{Ca}^{2+}]_o$  on the inactivation time constant shows little voltage dependence. Comparing the effect of increased  $\left[\text{Ca}^{2+}\right]_0$  on  $\tau_h$  with  $\left[\text{Ca}^{2+}\right]_i = 10^{-8}$  or  $10^{-5}$ M yields two results. First, the voltage dependence of increased  $[Ca<sup>2+</sup>]_{o}$  on the inactivation time constant is steeper with  $[Ca^{2+}]_i = 10^{-5}$  M compared with  $10^{-8}$  M. Second, the change in  $\tau_h$  produced by raising  $[Ca^{2+}]_o$  is smaller when  $[Ca^{2+}]_i = 10^{-5}$  M compared with 10<sup>-8</sup> M, which suggests competition between  $[Ca^{2+}]_i$  and  $[Ca^{2+}]_o$  for the same site. In addition, this series of experiments provides further evidence that the lymphocyte K<sup>+</sup> channel is not activated by intracellular  $Ca^{2+}$ , as we did not



FIGURE 8. Inactivation time constants  $\tau_h$  of K<sup>+</sup> currents in a Jurkat cell in normal Ringer  $( \Box )$  and in 12  $Ca^{2+}$  Ringer (v) with  $[Ca^{2+}]_i = 10^{-5}$ M in both cases. The  $\blacktriangledown$  values were obtained by correcting the  $r<sub>h</sub>$  values measured in 12 Ca<sup>2+</sup> ([Ca<sup>2+</sup>]<sub>i</sub> = 10<sup>-5</sup> M) for a shift along the voltage axis due to a change in surface potential, a correction similar to the one in Fig. 3. The lines through the points were drawn by eye for clarity. Cell 133.

observe any significant shift of the activation parameters to more hyperpolarized potentials with  $[Ca^{2+}]_i = 10^{-5}$  M, compared with  $[Ca^{2+}]_i = 10^{-8}$  M.

If there is a binding site located on the cytoplasmic side of the membrane, which can be accessed by  $Ca^{2+}$  ions from the outside moving through the channel, then increasing the intracellular  $Ca^{2+}$  buffering may reduce the effect of extracellular  $Ca<sup>2+</sup>$ , by analogy with  $Ca<sup>2+</sup>$ -dependent inactivation of  $Ca<sup>2+</sup>$  channels (Eckert and Chad, 1984). However, increasing cytoplasmic Ca<sup>2+</sup> buffering by including 55 mM BAPTA and 5 mM Ca<sup>2+</sup> in the pipette (free  $[Ca^{2+}]_i = 10^{-8}$  M) did not change the effect of  $Ca^{2+}$  on  $\tau_{hf}$  (data not shown). These results suggest either that the site at which  ${[Ca<sup>2+</sup>}$  binds is inside the K<sup>+</sup> channel and not accessible to cytoplasmic  $Ca<sup>2+</sup>$ buffers, or that there are two sites, one external and one internal.

# $K^+$  and  $Ca^{2+}$  Have Opposing Effects on  $K^+$  Channel Inactivation

Reducing  $[K^+]_0$  from 4.5 to 0 mM has almost the same effect on the  $K^+$  current as does increasing  $[Ca^{2+}]_o$ , as shown in Fig. 9 A. The time course of the decay of the

K<sup>+</sup> current is faster in 0 mM K<sup>+</sup> compared with normal Ringer solution, with  $\tau_h$ values of 200 and 285 ms, respectively. Thus, removing  $K^+$  from the Ringer accelerated the rate of inactivation by a factor of 1.4. Fig. 9 B illustrates the effect of increasing  $[K^+]$ , from 4.5 to 160 mM. The reduction of the  $K^+$  current amplitude in high  $[K^+]_o$  is due to a change in reversal potential  $E_K$  for potassium ions; thus, the driving force of these ions at  $+40$  mV is reduced. In addition to slowing channel closing during tail currents (Cahalan et al., 1985), raising  $[K^+]_o$  slows down inactivation. The time course of the inactivation in  $160$  mM K<sup>+</sup> Ringer could be well described by a double exponential with fast and slow inactivation time constants,  $\tau_{\text{hf}}$ and  $\tau_{\rm hs}$ . The value for  $\tau_{\rm hf}$  was 320 ms and was 1.6 times larger than the  $\tau_{\rm h}$  value



FIGURE 9. Effects of varying  $[K^+]_0$ on K<sup>+</sup> channel inactivation.  $[Ca^{2+}]_i =$  $10^{-8}$  M. (A) K<sup>+</sup> currents associated with 5-s depolarizing pulses to  $+40$ mV from a holding potential of  $-80$ mV in normal Ringer solution (4.5  $mM$  K<sup>+</sup>) and after application of a Ringer solution containing  $0 \text{ mM } K^+$ . Cell 119.  $(B)$  K<sup>+</sup> currents in a Jurkat cell in normal Ringer solution (4.5  $mM$  K<sup>+</sup>) and after application of a  $K^+$ -Ringer solution (160 mM  $K^+$ ).  $K^+$ outward currents associated with 3.8-s pulses to  $+40$  mV. Cell 80.

found in normal Ringer solution (200 ms). The value for  $\tau_{\text{hs}}$  was  $\sim$  4 s, a slow inactivation not seen in normal Ringer.

A dose-response curve demonstrating the effect of  $[K^+]_0$  on the fast component,  $\tau_{\rm hf}$  of K<sup>+</sup> channel inactivation is shown in Fig. 10. The change in the fast inactivation time constant was normalized as was done previously with  $Ca<sup>2+</sup>$ . A value of 1 corresponds to the maximum inactivation time constant of 430 ms. This  $\tau_h$  value was assumed to be the same as in the experiments with  ${[Ca<sup>2+</sup>}$ <sub>0</sub> = 10<sup>-9</sup> M. A value of 0 corresponds to the inactivation time constant of 200 ms and was obtained in experiments with  $[K^+]_o = 0$  mM as shown in Fig. 9 A. Fits of these data to the Hill equation yielded a slope of 1.1, which is consistent with the idea that only a single potassium ion is required to slow  $\tau_{hf}$ . A calculation of an adsorption isotherm to the doseresponse curve (line in Fig. 10) is in agreement with the data with a value of 13 mM for the apparent dissociation constant of  $K^+$  at its site of action.

Three explanations for the opposing influences of  $Ca^{2+}$  and  $K^{+}$  on inactivation are possible. (a) Two different receptors at the  $K^+$  channel interact allosterically;  $Ca<sup>2+</sup>$  bound to its receptor results in faster inactivation and K<sup>+</sup> bound to its receptor results in a loss of affinity of the  $Ca^{2+}$  receptor for  $Ca^{2+}$ . (b)  $Ca^{2+}$  can bind to a site where is can speed inactivation, and  $K^+$  either competes with  $Ca^{2+}$  for this site or prevents  $Ca<sup>2+</sup>$  from reaching it. (c) There are two separate and independent binding sites for Ca<sup>2+</sup> and K<sup>+</sup>; Ca<sup>2+</sup> binding will result in a faster inactivation, whereas K<sup>+</sup> binding will result in slower inactivation. The agreement of  $\tau_{\text{hmax}}$  with  $\left[\text{Ca}^{2+}\right]_{\text{o}} =$  $10^{-9}$  M and of  $\tau_{\text{hfmax}}$  with  $[K^+]_0 = 160$  mM (compare Figs. 4 and 10) argues against this idea of independent sites. Table II shows calculations of apparent dissociation constants of the  $Ca^{2+}$  site reaction (a) for the case of an affinity change at the  $Ca^{2+}$ binding site and (b) for the case of competition. Values of  $K_d$ , reflecting the change in affinity of the site for  $Ca<sup>2+</sup>$  under (a), change almost two orders of magnitude (compare  $K_d = 1$  mM for  $[K^+]_0 = 0$  mM with  $K_d = 79$  mM for  $[K^+]_0 = 160$  mM).



FIGURE 10. Dose-response curve of  $[K^+]$ <sub>o</sub> on the change  $p$  of the fast inactivation time constant  $\tau_{\rm hf}$  in Ringer solutions with  $[Ca^{2+}]_{o} = 2$  mM.  $[Ca^{2+}]_i = 10^{-8}$  M. p is the normalized change of  $\tau_{\text{hf}}$ , with  $\tau_{\text{hfmin}} = 200$  ms and  $\tau_{\text{hfmax}} = 430 \text{ ms}$ , as discussed in the legend of Fig. 4. The curve through the points was calculated assuming that  $p$  is proportional to  $K_{d}/([K^{+}]_{o} + K_{d})$  with  $K_{d} = 13$  mM.

This change in  $K_d$  seems to be too large to be explained by a loss in affinity of the  $Ca<sup>2+</sup>$  binding site for  $Ca<sup>2+</sup>$ . For the case of competition (b), the calculated dissociation constant of  $Ca^{2+}$  and its receptor, in the presence of the competing  $K^{+}$ , should be independent of  $[K^+]_{\alpha}$ , as indicated in Table II assuming that the true dissociation constant for the reaction of K<sup>+</sup> with the receptor, i.e.,  $K_{K}$ , is ~2 mM. From these calculations we favor the idea that  $Ca^{2+}$  and  $K^{+}$  compete for the same binding site.

### *External Rb + and Divalent Trapping*

The effects of replacing NaC1 and KCI in normal Ringer with RbCI were similar to those shown in Fig. 9 B in which  $[K^+]_o$  is increased from 4.5 to 160 mM. Interestingly, however, the first trace, after the application of Rb<sup>+</sup>-Ringer, was consistently different in the time course of activation from the second and following traces in  $Rb^+$ -Ringer, as illustrated in Fig. 11 A. One can see that the  $K^+$  current associated with the first test pulse to  $+40$  mV in Rb<sup>+</sup>-Ringer activates very rapidly (within 5 ms) to 767 pA, then activates further very slowly to reach a peak after 100 ms. It then inactivates with a time course similar to the slow inactivation seen in 160 mM  $K^+$ -Ringer. The  $K^+$  currents associated with the second and following test pulses in  $Rb^+$ -Ringer do not show the slow activation; the peak  $K^+$  current is reached within  $10$  ms. The  $K^+$  current decays associated with the first, second, and subsequent test pulses in Rb<sup>+</sup>-Ringer are superimposable. In addition, tail currents in Rb<sup>+</sup>-Ringer are prolonged compared with normal, as observed previously (Cahalan et al,, 1985). We believe that the slow phase of activation represents the removal of divalent ions

			$K_d$ calculated for assumptions		
$[\text{Ca}^{2+}]_0$	$[K^+]$	n	$(a)^*$	$(b)$ <sup>1</sup>	
m <sub>M</sub>	mM		mM	mM	
2.0	0.0	5	1.0	1.0	
2.0	2.2	$\mathbf{2}$	1.7	0.8	
2.0	4.5	58	2.8	0.9	
2.0	10.0		3.5	0.6	
2.0	20.0	ł	7.4	0.7	
2.0	40.0	$\mathbf 2$	24.1	1.1	
2.0	80.0	$\boldsymbol{2}$	31.2	0.8	
2.0	144.4	6	79.1	1.1	
2.0	160	5	79.1	1.0	
0.5	4.5	2	3.7	1.1	
4.6	4.5	2	3.3	1.0	
12.4	4.5	4	3.7	1.1	
53.8	4.5	5	3.4	1.1	
105.7	4.5	3	3.3	1.0	

TABLE II *Calculated Dissociation Constant, K<sub>4</sub>, of Ca<sup>2+</sup> Binding Site* 

Apparent dissociation constant,  $K_d$ , of Ca<sup>2+</sup> site reaction calculated for the case of (a) exclusive change of affinity of the  $Ca<sup>2+</sup>$  binding site by K<sup>+</sup>, and for the case of (b) competition. The equations used are from Ulbricht and Wagner (1975) and Grissmer (1984) with little modification. The equations are:

$$
*K_{d} = [Ca^{2+}]_{0}/(1/p - 1)
$$

 $(p$  is defined in legend of Fig. 4) and

$$
{}^{t}K_{d} = [Ca^{2+}]_{0}/(1/p - 1)(1 + [K^{+}]_{0}/K_{K})
$$

 $K_{\rm K}$  is the dissociation constant for the reaction of  $K^+$  with the Ca<sup>2+</sup> binding site, assuming that there is no direct influence of  $K^+$  on the fast inactivation with  $K_K = 2$  mM.

from  $K<sup>+</sup>$  channels, for reasons described below. On this assumption, the difference between the first and second current traces after applying Rb<sup>+</sup>-Ringer represents the time course of unblocking. The peak amplitude of the difference current was 20% of the amplitude of the scaled second current trace, and the time course of the decay of this difference current was 72 ms. In a total of five experiments we obtained a relative peak amplitude of the difference current of  $24 \pm 3\%$ (mean  $\pm$  SD); the time constant of the decay of these difference currents was 75  $\pm$  6 ms.

Pretreatment with 10 mM BaCl<sub>2</sub> increased the amplitude of the difference cur-

rent between the first and second current traces in Rb+-Ringer to 76% of the peak, without changing the time constant of the decay of this difference current (70 ms), compared with that without  $Ba^{2+}$  pretreatment. This result, illustrated in Fig. 11 B, indicates that the fraction of channels containing trapped divalent ions is correlated with the amplitude of the difference current between the first and second current traces in Rb+-Ringer. It also suggests that the relative amplitude of the difference current under normal conditions (i.e., no  $Ba^{2+}$  pretreatment) reflects the percentage of channels that have  $Ca<sup>2+</sup>$  trapped inside.

To test whether the fraction of channels with trapped  $Ca<sup>2+</sup>$  can be modified, we varied external and internal  $[Ca^+]$ . Increasing  $[Ca^{2+}]_0$  to 106 mM  $Ca^{2+}$  accelerated inactivation, but changed neither the relative amplitude of the difference current



FIGURE 11.  $K^+$  currents associated with the first and second 2-s depolarizing test pulses to  $+40$  mV (holding potential  $-80$  mV) after changing to Rb<sup>+</sup>-Ringer. Pulse interval was 53 s. The bath solution was totally exchanged before delivering the pulses. *(Upper row)* Bath solution was changed from normal Ringer to Rb<sup>+</sup>-Ringer ( $(Ca^{2+})$  =  $10^{-8}$  M) in A, from a Ringer solution with 10 mM Ba<sup>2+</sup> to Rb<sup>+</sup>-Ringer ( $[Ca<sup>2+</sup>]_{i} = 10<sup>-8</sup>$  M) in B, and from normal Ringer to Rb<sup>+</sup>-Ringer ( $[Ca<sup>2+</sup>]$ <sub>i</sub> = 10<sup>-5</sup> M) in C; records were taken after the solution change. (Lower row) Difference current of the first and second current trace in Rb<sup>+</sup>-Ringer.

between the first and second current traces in Rb<sup>+</sup>-Ringer nor its decay. This suggests that  $Ca<sup>2+</sup>$  loaded from the outside during the pulse may have been released to the inside during the 1-min pulse interval between the last current trace in 106 mM  $Ca<sup>2+</sup>$  and the first trace in Rb<sup>+</sup>-Ringer. This led us to attempt to determine whether we could increase the number of channels that have  $Ca<sup>2+</sup>$  trapped inside by increasing  $[Ca^{2+}]_i$ . For this purpose we used patch-pipettes filled with  $[Ca^{2+}]_i = 10^{-5}$  M. After breaking into the cell, we waited at least 5 min to attain equilibrium between the pipette solution and the cell cytoplasm. The pipette solution was then changed to Rb+-Ringer. The difference between the first and second current trace obtained under these conditions is shown in Fig. 11  $C$ . The peak amplitude of the difference current was 57% of the second current trace; the decay of the difference current

was 71 ms. In a total of five experiments, we obtained  $52 \pm 8\%$  for the relative amplitude of the difference current; the time course for the decay of these difference currents was 76  $\pm$  4 ms. We conclude that it is possible to load the channels with  $Ca^{2+}$  from the inside, and that this  $Ca^{2+}$  can be released to the outside.

# *Effects of Removing All Divalent Cations*

To see whether the residual inactivation rate of  $K<sup>+</sup>$  channels (which corresponds to  $\tau_{\text{hmax}}$  in low external and internal  $[\text{Ca}^{2+}]$  is due to block by Mg<sup>2+</sup> ions, which would be analogous to the Mg<sup>2+</sup> block of inward rectifier  $K^+$  channels (Matsuda et al., 1987; Vandenberg, 1987), experiments were performed using internal solutions with no added  $Mg^{2+} ([Ca^{2+}]_i = 10^{-8} M)$ . After breaking into a cell we recorded normal inactivating  $K^+$  currents even after 15 min with normal Ringer outside. If we changed to a Ringer solution with no added Mg<sup>2+</sup> and  $\left[Ca^{2+}\right]_0 = 10^{-9}$  M, changes in the membrane current that were similar to those described by Armstrong and Lopez-Barneo (1987) occurred. We observed the following results. (a) The inward holding current increased. (b) When the voltage was stepped to  $E = +40$  mV from a holding potential of  $-80$  mV, there was an instantaneous jump in current from inward to outward. (c) The  $K^+$  current became smaller, with slower decay of the current (comparable with the experiments with  $\left[Ca^{2+}\right]_0 = 10^{-9}$  M), and disappeared at last. Extending these results, which are similar to those found previously on squid neurons (Armstrong and Lopez-Barneo, 1987), we found that: (d) the induction of the new "leakage" conductance could be prevented by 2 mM  $Mg^{2+}$  inside. (e) This conductance could be blocked by either 1 mM  $Ca<sup>2+</sup>$  or 1 mM Mg<sup>2+</sup> outside. (f) Ba<sup>2+</sup> could not only prevent the induction of this conductance but could also block it after its development with a  $K_d$  around 3 mM. (g) The induction of this conductance could not be prevented by totally blocking the  $K^+$  channels with 10 nM charybdotoxin, a potent blocker of the  $K^+$  channel (Sands et al., 1988), either before or during treatment with zero-divalent external and internal solutions. Although raising  $[Mg^{2+}]_o$  to 12 mM has little effect on the K<sup>+</sup> channel inactivation rate (data not shown), even 1 mM Mg<sup>2+</sup> can "protect" the K<sup>+</sup> channels while also preventing the increased leakage current seen in zero divalent. We are uncertain whether this conductance is related to the  $K^+$  channel protein, or, alternatively, represents a nonspecific conductance induced by zero divalent.

# DISCUSSION

It is well established that ions can alter the gating of voltage-dependent channels. One mechanism involves electrostatic attraction and binding of ions to negative surface charges at each side of the membrane (for review, see Hille, 1984), thereby influencing the voltage dependence of channel gating. Several observations, however, cannot be explained by surface potential theory alone. One example is that externally applied  $\text{Zn}^{2+}$  ions slow activation without altering the time course of deactivation in squid axons (Gilly and Armstrong, 1982). State-dependent binding mechanisms, in which the ion binds more strongly to a particular channel conformation, provide an explanation for effects of divalent ions on channel gating. A special case of state-dependent mechanisms involves ion binding to open channels. For example, occupancy of an open channel by either a permeant or blocking ion has been hypothesized to hinder channel closing for both acetylcholine receptors and  $K^+$ channels (Marchais and Marty, 1979; Swenson and Armstrong, 1981; Cahalan and Pappone, 1983; Matteson and Swenson, 1986; Armstrong and Matteson, 1986). The effects of permeant monovalent, divalent, and blocking ions on  $T$  cell  $K^+$  channels are further examples of state-dependent channel mechanisms. A novel aspect of the results reported here is that we have focused on the effects of ions on channel inactivation gating.

# *Divalent Effects on Inactivation: Block or Modulation*

We have shown that the lymphocyte  $K^+$  channel inactivation rate is strongly influenced by external  $Ca^{2+}$  (or  $Ba^{2+}$ ), as well as internal  $Ca^{2+}$ . There are two possible mechanisms that should be considered.  $Ca<sup>2+</sup>$  binding to a site in the channel may result in (or modulate the rate of) a conformational change that inactivates the channel. A second possibility is that block of the  $K^+$  channel by a  $Ca^{2+}$  ions is the inactivation mechanism. In this case, depolarization to open the channel would also enable entry and block by divalent ions; the slow recovery from inactivation would represent the unbinding of trapped  $Ca^{2+}$ . In either mechanism, if  $Ca^{2+}$  entry were decreased (a) by reducing  $[Ca^{2+}]_0$ , (b) by competition with an external permeant ion, or  $(c)$  by a  $K^+$  channel blocker, then inactivation should become slower. We now consider experimental evidence on each of these points. (a) As  $[Ca^{2+}]_o$  is reduced, the time constant of inactivation approaches a maximum value,  $t_{\text{hmax}}$ , of  $\sim$ 430 ms. This residual inactivation observed in essentially  $Ca^{2+}$ -free internal and external solutions may be due to an intrinsic inactivation mechanism, or due to block by another ion present, for example  $Mg^{2+}$ . The induction of "leak" conductance and disappearance of  $K^+$  conductance upon the removal of all divalents prevents a clear test of this hypothesis. (b) When  $[K^+]_o$  is elevated, inactivation becomes slower and incomplete, even during pulses lasting  $4 \text{ s}$  (Fig.  $9 \text{ B}$ ). Table II describes the interaction between  $[K^+]_o$  and  $[Ca^{2+}]_o$  on the fast component of inactivation; permeant monovalent ions may compete with  $Ca<sup>2+</sup>$  for entry into the channel. A slow component of inactivation with a time constant greater than  $t_{\text{hmax}}$  is revealed when  $[K^+]_o$  is raised. This may indicate that  $t_{\text{hmax}}$  is not due to an intrinsic inactivation mechanism. (c) External tetraethylammonium ( $[TEA^+]_0$ ) slows the time constant of inactivation as it reduces the current, resulting in crossover of the  $K^+$  currents in the presence and absence of  $TEA^+$  and a constant current integral during long depolarizing pulses as  $[TEA^+]_o$  is elevated (Grissmer and Cahalan, 1989). Similar crossover of K<sup>+</sup> current due to channel block, combined with slowing of inactivation, was also observed with external  $Mn^{2+}$  or  $Co^{2+}$  ions (DeCoursey et al., 1985). The kinetic effects of  $TEA<sup>+</sup>$  on inactivation were modeled by a scheme in which channels blocked by TEA<sup>+</sup> cannot inactivate (Grissmer and Cahalan, 1989). TEA<sup>+</sup> (or  $Mn^2$ <sup>+</sup> or  $Co^{2+}$ ) blocking the K<sup>+</sup> channel may prevent inactivation by preventing  $Ca^{2+}$ entry. If this is true, then the fact that TEA<sup>+</sup> results in  $t<sub>h</sub>$  values far greater than  $t<sub>h</sub>$ <sub>max</sub> might argue against an intrinsic  $t<sub>h</sub>$  of 430 ms, as discussed in (a) above, and in favor of divalent entry being a prerequisite for inactivation. Either divalent block or a conformational change induced by a divalent ion within the channel would be consistent with the evidence described above.

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The concentration dependence of the  $Ca^{2+}$  (or  $Ba^{2+}$ ) effect on inactivation is apparently inconsistent with a simple divalent-block hypothesis;  $t_{\text{hmin}}$  saturates at a minimum value of 65 ms as  $[Ca^{2+}]_0$  or  $[Ba^{2+}]_0$  is elevated. According to a simple divalent-block mechanism for inactivation, one would expect the time constant of inactivation to decrease monotonically as the concentration of divalent is raised, as is the case for internal  $Ba^{2+}$  block of squid axon  $K^+$  channels (Armstrong and Taylor, 1980; Eaton and Brodwick, 1980; Armstrong et al., 1982). The effects of both  $[Ca^{2+}]_o$  and  $[Ba^{2+}]_o$  on lymphocyte K<sup>+</sup> channel inactivation can be described by a simple 1:1 binding isotherm, with saturation at high  $[Ca<sup>2+</sup>]_{o}$  or  $[Ba<sup>2+</sup>]_{o}$  at the same value of  $t_{\text{hmin}}$  (Fig. 4). This result suggests that divalent binding to a site inside the channel induces a conformational change that inactivates the channel with a limiting rate corresponding to  $t_{\text{hmin}}$ . A more complicated blocking model involving extremely slow single filing of  $Ca^{2+}$  (or  $Ba^{2+}$ ) between an outermost saturable site and a blocking site within the channel might also account for  $t_{\text{hmin}}$ . Regardless of the exact mechanism, the acceleration of inactivation by divalent ions parallels the idea that channel closing (deactivation) is more rapid when channels contain  $Ca<sup>2+</sup>$  than when they contain monovalent ions, as proposed for squid axon  $K<sup>+</sup>$  channels (Armstrong and Matteson, 1986).

# *Blocking Model for*  $[Ca^{2+}]_o$  *and*  $[Ca^{2+}]_i$

Raising either internal or external  $[Ca<sup>2+</sup>]$  accelerates the rate of inactivation. Competition between internal and external  $Ca^{2+}$  ions for the same binding site within the channel may result in the diminished effect of  $[Ca^{2+}]_{\text{o}}$  on  $\tau_{\text{h}}$  when  $[Ca^{2+}]_{\text{i}}$  is raised (Figs. 3 and 8). We modeled the change in  $\tau_h$  by assuming that  $Ca^{2+}$  can block the open channel nearly irreversibly, and that  $Ca^{2+}$  may reach the blocking site from either side of the membrane. Under this assumption one can separate  $\tau_h$  due to an intrinsic inactivation process (measured in  $[Ca^{2+}]_o = 10^{-9}$  M and  $[Ca^{2+}]_i = 10^{-8}$  M) from  $\tau_{Ca}$  due to  $Ca^{2+}$  reaching the site. These  $\tau_{Ca}$  values measured in different  $[Ca^{2+}]_0$  and different  $[Ca^{2+}]_i$  were plotted against the absolute membrane potential E as shown in Fig. 12 (open squares). The voltage dependent of  $\tau_{Ca}$  changes from a small increase of  $\tau_{Ca}$  with more depolarization in  $[Ca^{2+}]_i = 10^{-8}$  M, to a marked reduction of  $\tau_{Ca}$  with depolarization in  $[Ca^{2+}]_i = 10^{-5}$  M. The lines through the points are least-squares fits to the data considering that the  $K<sup>+</sup>$  channel has one local free energy minimum (site) that can either be empty or occupied by  $Ca^{2+}$ .  $Ca^{2+}$ moves from either side of the membrane to the site over an intervening energy maximum (barrier). The data could be reasonably described by a free energy profile shown in Fig. 13, which illustrates a two-barrier one-site model. In this model the binding site for  $Ca^{2+}$  is almost at the extracellular side of the membrane, perhaps near the channel mouth. The reason for the very broad plateau of the inner barrier is the weak voltage dependence of  $\tau_{Ca}$  in all the solutions with different Ca<sup>2+</sup> concentrations. Alternatively, a second two-barrier one-site model, with the plateau at the outer barrier, is indistinguishable from the one shown in Fig. 13. The  $Ca<sup>2+</sup>$  binding site would then be about two thirds of the way across the electric field from the outside.

The similar effects raising internal and external  ${[Ca<sup>2+</sup>]}}$  on  $K<sup>+</sup>$  channel inactivation (Figs. 3 and 8), and the unblocking of divalent cations from  $K^+$  channels (Fig. 11),

are results consistent with  $Ca^{2+}$  ions being sparingly permeant through  $K^+$  channels, as suggested by the energy profile of Fig. 13. A calculation of the  $Ca^{2+}$  flux through an open  $K^+$  channel using the rate constants for leaving the  $Ca^{2+}$  binding site (obtained from the barrier heights and well depth given in legend of Fig. 13) shows that this Ca<sup>2+</sup> flux is ~0.6 ions/channel per s at 0 mV. Even with all  $K^+$  channels open, the maximum Ca<sup>2+</sup> flux in a human T cell with 400 channels would be  $\sim$ 250 ions/cell per s, at least 2 orders of magnitude smaller than estimates of  $Ca<sup>2+</sup>$  entry from tracer flux measurements which were reported by Metcalfe et al. (1980) and Hesketh et al. (1983) to be 24,000-130,000  $Ca<sup>2+</sup>$  ions/cell per s. The calculated flux



FIGURE 12. Time constants  $\tau_{Ca}$  for Ca<sup>2+</sup> reaching the site at the  $K^+$  channel at four different combinations of  $[Ca^{2+}]_o$ and  $[Ca^{2+}]_i$ . (A)  $[Ca^{2+}]_o = 2$ mM,  $[Ca^{2+}]_i = 10^{-8}$  M. (B)  $[Ca^{2+}]_o = 2$  mM,  $[Ca^{2+}]_i =$  $10^{-5}$  mM. (C)  $[Ca^{2+}]_{o} = 12$ mM,  $[Ca^{2+}]_i = 10^{-8}$  M. (D)  $[Ca^{2+}]_{o} = 12$  mM,  $[Ca^{2+}]_{i} =$  $10^{-5}$  M.  $\tau_{Ca}$  was obtained by dividing the time course of the decay of the  $K^+$  current with  $[Ca^{2+}]_o = 10^{-9}$  M by the time course of the decay of the  $K^+$ current in the different solutions. The newly created time course reflects the time course of  $Ca<sup>2+</sup>$  reaching its site at the channel. The lines through the points were least-squares fits to the data with values for the inner and outer barrier; the well and the electrical distances are given in legend of Fig. 13.

is based upon the blocking model with a static energy profile as shown in Fig. 13. As discussed above, the changes in inactivation may result from modulation of an intrinsic inactivation process as  $Ca^{2+}$  ions from either side of the membrane bind to a site within the membrane. We cannot estimate the flux in this case.

In single channel studies of the  $Ca^{2+}$ -activated K<sup>+</sup> channel, Neyton and Miller (1988b, c) described  $Ba^{2+}$  block from either the *cis* or *trans* side of the membrane in terms of a single site within the channel that can bind  $Ba^{2+}$ . In their model to account for interactions between  $Ba^{2+}$  and K<sup>+</sup>, three additional sites can bind K<sup>+</sup> ions (Neyton and Miller, 1988c), but when  $[K^+]_o$  is lowered a second Ba<sup>2+</sup> ion can also bind from the outside (C. Miller, personal communication). The asymmetric

concentration dependence and rates of  $Ba^{2+}$  block in the  $Ca^{2+}$ -activated K<sup>+</sup> channel are similar to the  $Ca^{2+}$  effect on lymphocyte  $K^+$  channel inactivation described here. Although a single-site model may account for the divalent effects on inactivation, we cannot rule out additional divalent binding sites.

# *Trapping of Divalent Ions*

The data presented in Figs. 5, 6, and 11 illustrate some of the dynamics of divalent trapping in the K<sup>+</sup> channel. Fig. 5 demonstrates that  $Ba^{2+}$  ions cannot enter the K<sup>+</sup> channel until it is opened. The channels must be opened in order for recovery from  $Ba<sup>2+</sup>$  block to occur (Fig. 6). The recovery rate strongly depends on the interval between test pulses, suggesting that if the channels could be held open longer, the recovery from  $Ba^{2+}$  block would be faster. However, the channels inactivate during depolarizing pulses, and so  $Ba^{2+}$  ions do not dissociate from all the channels during a single depolarizing pulse. If one could keep the channels open for a very long time



FIGURE 13. Energy profile for  $Ca^{2+}$ in the  $K^+$  channel. This energy profile was obtained by the least-squares fit shown in Fig. 12. The heights,  $B$ , and position of the barriers and the well are:  $B_i = 21 RT$ , well = -14.7  $RT$ ,  $B_0 = 15.7 RT$ , electrical distance of the well from the outside  $s_1 =$ 0.04, the length of the plateau of  $B_{\text{out}}$  $(p_{B2})$  is 0.52, and  $s_2/2$  is therefore 0.22.

by preventing inactivation one might see the real time course for  $Ba^{2+}$  leaving the channel with a single depolarizing test pulse. We succeeded in reducing the rate of inactivation by changing to high Rb<sup>+</sup>-Ringer, and we were able to measure the time course for  $Ba^{2+}$  and  $Ca^{2+}$  leaving the K<sup>+</sup> channel (see Fig. 11). The time course for leaving the  $K^+$  channel, reflected in the decay of the difference current shown in Fig. 11, is the same for either  $Ba^{2+}$  or  $Ca^{2+}$ . This similarity in the time course of the difference current may result from the dissociation of Rb ions being rate limiting; if  $Rb<sup>+</sup>$  binds to a site and can prevent the removal of a  $Ca<sup>2+</sup>$  or  $Ba<sup>2+</sup>$  ion, the time course of the difference current would be limited by the time course for Rb<sup>+</sup> leaving its site. This assumption could also explain the failure to observe a difference in the first and second current trace in  $160$  mM K<sup>+</sup> Ringer. Under normal conditions ~25% of all channels contain a  $Ca^{2+}$  ion. Surprisingly, raising  $[Ca^{2+}]_o$  did not increase this fraction, as indicated by the amplitude of the difference current between the first and second current trace in Rb+-Ringer. An explanation for this

behavior could be that  $Ca^{2+}$  is able to leave its binding site even if the channel is closed. The latter possibility is supported by the experiments with an increase in  $[Ca^{2+}]$ ; to  $10^{-5}$  M, which accelerates inactivation and increases the difference current to  $>50\%$ .

Several lines of evidence are presented in this paper which indicate that  $Ca<sup>2+</sup>$  and  $Ba^{2+}$  are able to enter open  $K^+$  channels from the outside. If the channel closes,  $Ca<sup>2+</sup>$  and  $Ba<sup>2+</sup>$  become trapped inside the channel.  $Ba<sup>2+</sup>$  remains trapped, whereas  $Ca<sup>2+</sup>$  is thought to cross the membrane to the inside during the interval between test pulses. Ba<sup>2+</sup> trapping has been described previously for the high conductance  $Ca^{2+}$ activated  $K^+$  channel from rabbit or rat muscle incorporated into planar lipid bilayers (Vergara and Latorre, 1983; Miller, 1987; Miller et al., 1987; Neyton and Miller, 1988a). The existence of a  $Ba^{2+}$ -trapping mechanism, as well as the previously described sensitivity to charybdotoxin of the  $Ca^{2+}$ -activated "maxi"-K<sup>+</sup> channel and the lymphocyte K<sup>+</sup> channel (Miller et al., 1985; Sands et al., 1988), implies structural similarities between these two channel types and raises the question of whether the  $K^+$  channel in T lymphocytes might be  $Ca^{2+}$  activated. Our experiments with high internal  $Ca^{2+}$  of up to 2 mM show no apparent activation of this channel; instead, raising  $[Ca^{2+}]$  speeds up inactivation. One conclusion from our experiments on  $Ba<sup>2+</sup>$  trapping is that divalent ions in the internal or external solution must wait for the channel to open before entering the pore. If the channel closes,  $Ba^{2+}$  is trapped inside the channel and remains trapped even in the absence of extracellular and intracellular  $Ba^{2+}$ . From the sidedness of the  $Ba^{2+}$  action, one has to conclude that both the Ca<sup>2+</sup>-activated K<sup>+</sup> channel and the voltage-gated K<sup>+</sup> channel in T lymphocytes have a gating mechanism at the external side that can shield the inside of the channel from externally applied divalent ions, as well as a gating mechanism at the cytoplasmic side of the membrane that shields the inside of the channel from internal blockers (cf. Armstrong, 1971).

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