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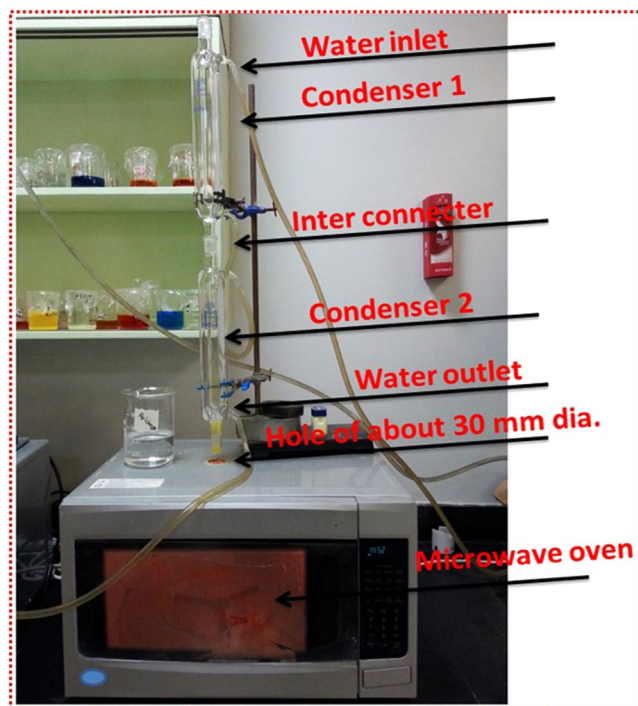
## Tailoring the structural, morphological, optical and dielectric properties of lead iodide through $\text{Nd}^{3+}$ doping

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Hexagonal single crystal nanosheets of  $\text{Nd}^{3+}$  doped  $\text{PbI}_2$  were effortlessly synthesized via microwave-assisted technique under a power of 700W and in a duration of 15 minutes with a homogeneous morphology. X-ray diffraction, energy dispersive X-ray spectroscopy, scanning electron microscope, FT-Raman, UV-Visible, photoluminescence and dielectric measurement were employed to study the product. High purity, single phase and presence of  $\text{Nd}^{3+}$  doping was confirmed. SEM study confirms the formation of nanorods and single crystal nanosheets of very few nanometers in size. Robust vibrational analysis has been carried out and the observed bands are assigned to the vibration modes of  $E_2^1$ ,  $A_1^1$ ,  $A_1^2$ ,  $2E_2^1$  and  $2E_1^1$ , respectively. These bands are red-shifted when compared to the corresponding bulk values which indicate relaxed nanostructure formation and occurrence of confinement effect. The thickness of the synthesized single crystal nanosheets are found to be in the range of ~20 to 30 nm. The energy band gap was calculated and found to be 3.35, 3.34, 3.42 and 3.39 eV for pure, 1, 3 and 5%  $\text{Nd}^{3+}$  doped lead iodide, respectively. The clear blue luminescence has been observed at 440 nm and 466 nm when excited at 250 nm and 280 nm respectively. Dielectric and ac electrical conductivity was also measured and discussed.

In recent past, nanostructured semiconductor materials receive the colossal attention from scientists and researchers around the world due to their wide range of applications in the field of environmental and energy applications, gas-sensing, field-emission, radiation detection, solar cell, and optoelectronic devices<sup>1–5</sup>. In particular, Lead iodide ( $\text{PbI}_2$ ) is a large band gap (2.3 eV) p-type semiconductor material and possess wide range of applications such as in active matrix flat panel imagers, room temperature radiation detection, especially for the low energy X-ray spectrometer, mammography energy range detection and nuclear particle detectors, photoconductors, photo-detectors, photovoltaic, co-precipitation sensors, biological labeling and diagnostic, light emitting diodes etc.<sup>6,7</sup>. In past decade, various research and development (R&D) activities were carried out on one-dimensional nanostructured materials based on varying forms including nanocrystals, nanorods, nanobelts, nanoribbons, nanowires, nanotubes etc. due to their novel properties and impending applications in nano devices<sup>8,9</sup>. However, only few reports are available on the synthesis of  $\text{PbI}_2$  nanostructures including  $\text{PbI}_2$  nanorods or nanorods like particles by inverse micro-emulsion, microwave, hydrothermal and ultrasonic methods<sup>3,5,10,11</sup>,  $\text{PbI}_2$  nanocrystals by colloidal, hydrothermal and sol-gel route<sup>12,13</sup>,  $\text{PbI}_2$  nanoparticles by inverse micelles<sup>14</sup>,  $\text{PbI}_2$  single crystalline nanobelts by hydrothermal route<sup>15</sup>, and  $\text{PbI}_2$  quantum dots and clusters by chemical route<sup>16–18</sup>. Lately, the fabrications of nanosheets, in general, are commanding much attention from varying bodies of scientific communities and institutions because of their promising potentials in many advanced and futuristic technologies<sup>19–21</sup>. Hence, and by the same token, our research group has recently synthesized the  $\text{Gd}^{3+}$  and  $\text{Cs}$  doped  $\text{PbI}_2$  nanosheets by hydrothermal and microwave routes<sup>5,22,23</sup>. Embarking on the same research direction, the microwave-assisted rapid synthesis of Neodymium ( $\text{Nd}^{3+}$ ) doped  $\text{PbI}_2$  single crystal nanosheets with uniform morphology is going to be addressed and discussed in this article. In fact, synthesizing nano/thin films of key semiconductor materials like  $\text{ZnO}$ ,  $\text{TiO}_2$ ,  $\text{SnO}_2$  with  $\text{Nd}^{3+}$  doping has resulted in modified and interesting

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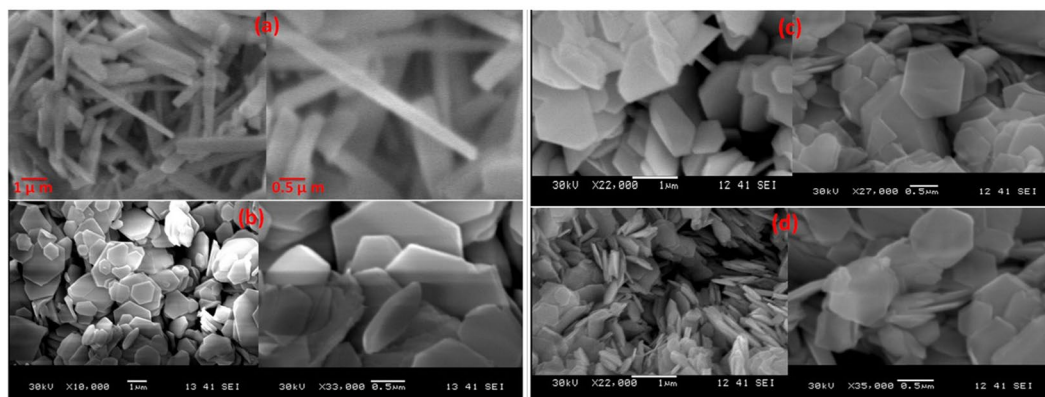
**Figure 1.** Used modified microwave system for nanosynthesis.

properties as reported elsewhere<sup>24–27</sup>. Therefore, it would be of interest to scientists and engineers alike to investigate the effect of Neodymium ( $\text{Nd}^{3+}$ ) doping on the important  $\text{PbI}_2$  semiconductor materials. Moreover, as per the authors' knowledge on the available literature concerning  $\text{PbI}_2$ , this would be the first report on the microwave-assisted rapid synthesis of Neodymium ( $\text{Nd}^{3+}$ ) doped  $\text{PbI}_2$  single crystal nanosheets with uniform morphology. It is well known that the microwave-assisted synthesis offers a method to resolve the problem in which microwave provides a rapid and volumetric heating at higher reaction rate and better selectivity and very short reaction time in comparison of conventional heating<sup>20,28,29</sup>. The synthesized single crystal nanosheets of  $\text{Nd}^{3+}$  doped  $\text{PbI}_2$  were characterized by X-ray diffraction (XRD), scanning electron microscope (SEM)/energy dispersive X-ray spectroscopy (EDXS), FT-Raman and UV-Visible spectroscopies, Photoluminescence (PL) and impedance spectroscopy measurements. The obtained results are discussed, and the conclusion was drawn as well.

## Experimental details

**Chemicals.** The analytical grade materials of Lead acetate ( $\text{Pb}(\text{C}_2\text{H}_3\text{O}_2)_2$ ), sodium iodide (NaI), Sodium dodecyl sulfate (SDS), Poly(vinyl alcohol) (PVA), Neodymium (III) nitrate were purchased from CDH, Alfa Aesar, and Sigma Aldrich and used for the titled materials synthesis without further purification.

**Microwave-assisted synthesis.** In a typical synthesis of pure and  $\text{Nd}^{3+}$  doped  $\text{PbI}_2$  nanorods and nanosheets the following steps were taken: (I) 0.5 M lead acetate (18.967 g) was dissolved in 50 ml double distilled water in a highly cleaned cylinder, once it was dissolved and clear solution appears, 50 ml SDS (10 g/ml) and 50 ml of PVA (45 g/liter) solutions were added one by one as surfactant and continuous stirring was performed at a stable magnetic stirrer fixed at 1000 rpm to get homogeneous solution at room temperature. The use of the surfactant is justified as it plays a key role in the synthesis of well-organized nanostructures and various kinds of surfactant have been used extensively in the past for the same purpose by many researchers and scientists<sup>5,30,31</sup>. (II) 1 M Sodium Iodide (14.988 g) was taken in another beaker and 50 ml distilled water used to dissolve it through stirring. Within 30 minutes, the homogeneous solutions are achieved, and finally, the prepared solutions were mixed together and stirred continuously. The temperature was set to  $60 \pm 2^\circ\text{C}$  during the whole synthesis process and within few seconds the solution becomes yellow which indicates the formation of  $\text{PbI}_2$ . Preparing the other three solutions for doping purpose, the same latter procedure was used. (III) During the synthesis of  $\text{PbI}_2$ , the different concentrations of Neodymium (III) nitrate [i.e. 1, 3 and 5 wt% of  $\text{Nd}^{3+}$  doping] were dissolved in 25 ml double distilled water in three different beakers and added to the  $\text{PbI}_2$  solutions at the same temperature. Finally, all the prepared solutions were reassign into cylindrical glass vessels for microwave-assisted synthesis. For microwave-assisted synthesis, a domestic microwave system of frequency 2.54 GHz bought from LG (Model No. MS5246VR/00) has been used, and the programming was done for 15 minutes at the fixed power of 700 W and afterward the microwave was automatically set off for natural cooling. The domestic microwave system was modified by making a hole at the center of the top of the furnace, and two condensers were used on the furnace and connect with the solution contained cylinder as shown in Fig. 1. The washing of the finally prepared yellow



**Figure 2.** SEM micrographs of (a) pure and (b) 1 wt%, (c) 3 wt% and (d) 5 wt% Nd doped  $\text{PbI}_2$  nanorods and single crystals nanosheets.

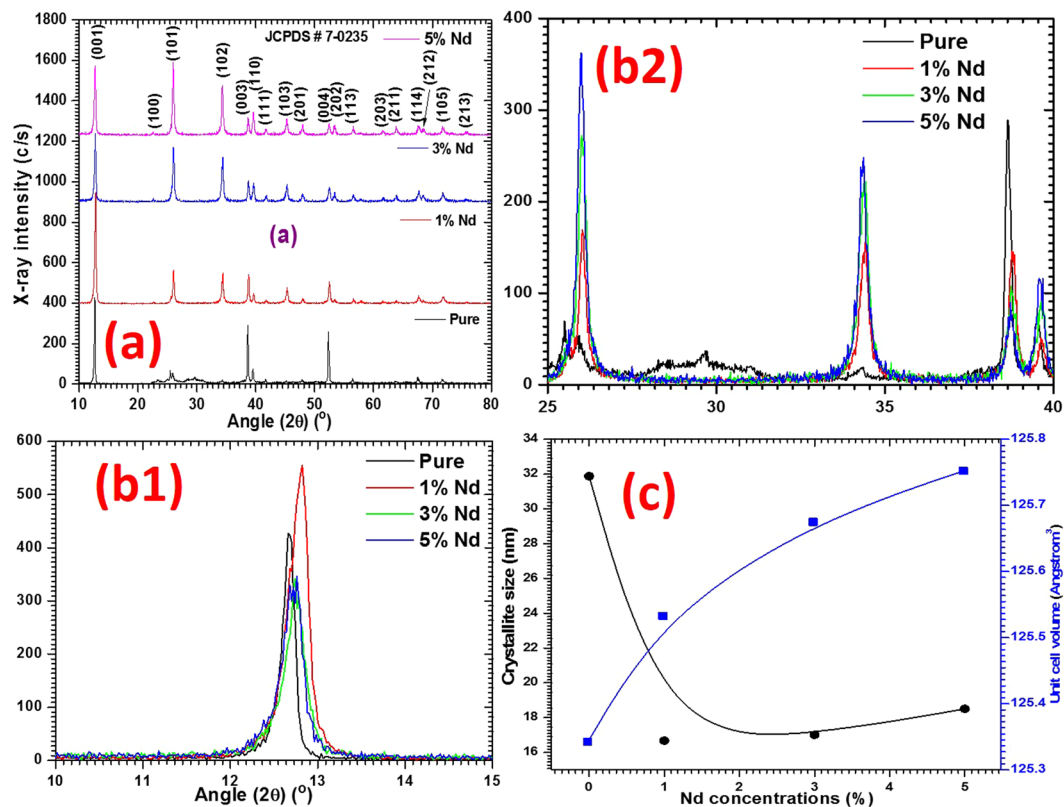
color salts with double distilled water was done many times during the filtration processes, and the obtained powder was dried in an oven at  $100^\circ\text{C}$  for 24 h.

## Results and Discussion

**Morphological analysis.** Figure 2(a–d) shows the collected SEM micrographs of the pure and 1, 3, and 5 wt.%  $\text{Nd}^{3+}$  doped  $\text{PbI}_2$  nanorods and nanosheets, respectively, captured at different resolutions/scales. As it is clearly visible from the Fig. 2(a), the morphology of the pure  $\text{PbI}_2$  is nanorods of diameters in the range of  $\sim 70$  to  $100$  nm and length of few microns. However, the morphology of  $\text{PbI}_2$  has been remarkably changed from nanorods (as the case of the pure  $\text{PbI}_2$ ) to single crystal nanosheets of hexagonal shape (as it is the case of the  $\text{Nd}^{3+}$  doped  $\text{PbI}_2$ ) of thickness in the range of  $\sim 20$  to  $30$  nm and width of about few nano- as well as micro-meters shown in Fig. 2(b,c and d). Apparently, Nd ions are playing the role of the heterogeneous seeds for the growing of the nanosheets. From SEM micrographs it is revealed that the synthesized nanostructures are of well-defined nanorods for pure  $\text{PbI}_2$  and single crystals nanosheets for  $\text{Nd}^{3+}$  doped  $\text{PbI}_2$  of homogeneous morphology. The scanning electron microscopy results confirms the strong effect of  $\text{Nd}^{3+}$  doping on the morphology of  $\text{PbI}_2$  nanostructures. However no major change in morphology or shape i.e. hexagonal structure, at higher concentration of  $\text{Nd}^{3+}$  doping has been observed, which means that by varying  $\text{Nd}^{3+}$  doping content within  $\text{PbI}_2$  the single crystal nanosheets with the same hexagonal morphology the  $\text{PbI}_2$  material can still be achieved. The synthesized  $\text{Nd}^{3+}$  doped  $\text{PbI}_2$  nanosheets may provide the exceptional optoelectronic properties needed in many modern devices and applications<sup>21</sup>.

**X-ray diffraction analysis.** The phase structure of the synthesized product (i.e. pure and 1, 3, 5 wt.%  $\text{Nd}^{3+}$  doped) was examined by powder X-ray diffraction study. The recorded X-ray patterns are depicted in Fig. 3. The sharpness of the diffraction peaks proved good crystallinity within the fabricated nanostructures. The close inspection of X-ray diffraction pattern of the synthesized nanostructures confirms the high purity of the product as there is no peak due to other impurities like:  $\text{PbO}$ ,  $\text{PbOH}$  and  $\text{Pb}(\text{OH})_2$  was observed. This was further confirmed by EDX/SEM mapping as well as FT-Raman studies presented in the forthcoming sections.

To understand the effect of  $\text{Nd}^{3+}$  doping on the crystalline structure of the  $\text{PbI}_2$ , we have compared the X-ray diffraction patterns of pure and doped nanostructures as shown in Fig. 3(b1 & b2). The figure clearly indicates that the increase in Nd/Pb atomic ratio due to doping is having major influence on the preferential growth directions, specially along (101), of the nanostructures. Moreover, the relative peaks positions are slightly shifted in the case of doped  $\text{PbI}_2$  compared to the pure material. The intensity of peaks (101), (102), (103) and many others is found to be noticeably enhanced with doping. However, it can be noticed that there is a noticeable reduction in the peaks intensities, specially, for the (003) and (004) planes, and no additional peak due to impurity was observed, which confirms the absence of any free Neodymium nitrate in the prepared nanostructures. Furthermore, comparing the XRD patterns within Fig. 3(b1 & b2), a slight shift in the peak positions is observed in the doped nanostructures. This shift in the peaks position provides a clear indication of the change taking place within the lattice parameters of the grown nanostructure and lattice size due to Nd doping effect. The variation in lattice size of  $\text{PbI}_2$  suggests that  $\text{Nd}^{3+}$  ions can be incorporated into the crystalline matrix of the  $\text{PbI}_2$  structure either in substitutional or interstitial positions. To find out the nature of  $\text{Nd}^{3+}$  doping in the  $\text{PbI}_2$  lattice, we have calculated the lattice parameters (a, c) and grain size as a function of  $\text{Nd}^{3+}$  doping concentrations in  $\text{PbI}_2$  are presented in Table 1; the result is in close agreement with the standard literature values on  $\text{PbI}_2$  [JCPDS #7-0235]. The possible error in the refined parameters is totally based on the experimental procedure and environmental conditions and it may be around  $\pm 0.001 \text{ \AA}$ . However, in the current work we have taken care about the condition to perform the X-ray measurements and also the refinement has been carried out many times. The nanostructures prepared by microwave-assisted route possess hexagonal phase of 2H-polytype of  $\text{PbI}_2$  with space group  $p-3m1$  (164). The crystallite/grain size of the fabricated nanorods and nanosheets was calculated using Scherer's formula:  $D = \frac{k\lambda}{\beta \cos\theta}$ , where,  $k = 0.9$  and  $D$  is an average crystallite size,  $\lambda$  is X-ray wavelength ( $0.15406 \text{ nm}$ ), and  $\beta$  is full width at half maximum in radian. Furthermore, the values of lattice strain ( $\epsilon$ ) and dislocation density ( $\delta$ ) were



**Figure 3.** (a) X-ray diffractions patterns, (b1 & b2) presentation of shifting in diffraction peaks and (c) Dependence of crystallite size of and unit cell volume as a function of  $\text{Nd}^{3+}$  doping.

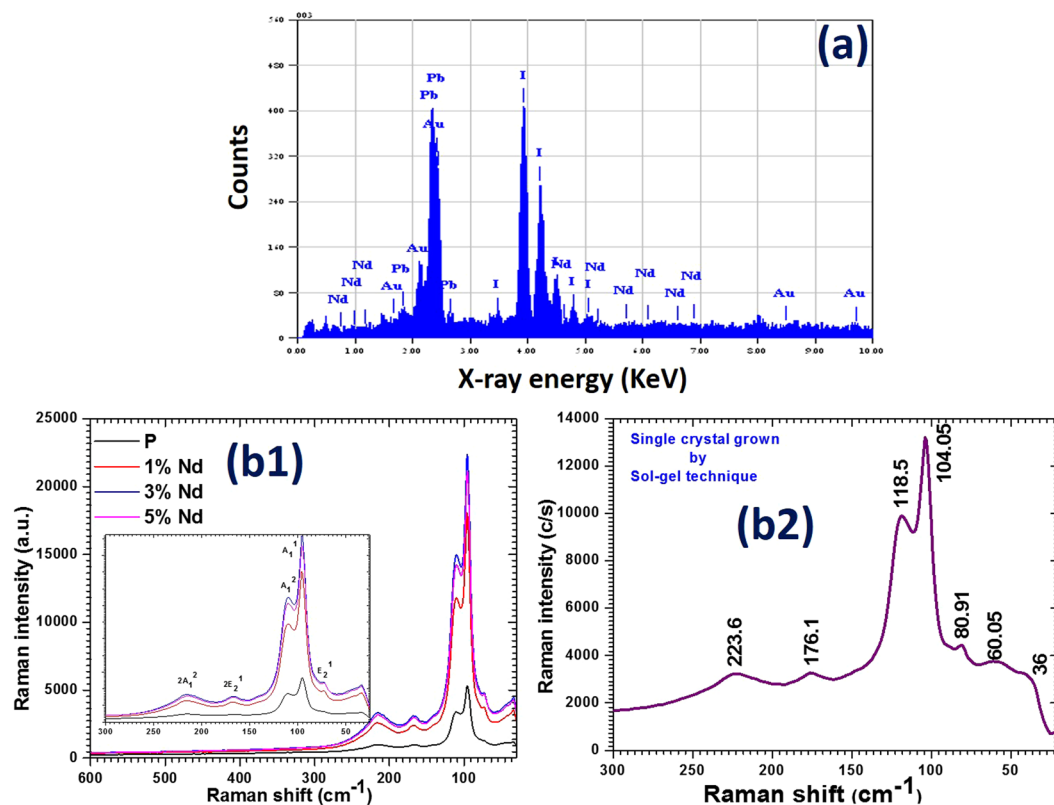
Samples	a = b (Å)	c (Å)	V(Å) <sup>3</sup>	D (nm)	$\epsilon \times 10^{-4} \text{ lin}^{-2} \text{ m}^{-4}$	$\delta \times 10^{15} \text{ lines m}^{-2}$
JCPDS- (07-0235)	4.557	6.979	125.511	—	—	—
pure_PbI <sub>2</sub>	4.554	6.978	125.342	31.871	0.110	0.0114
1 wt% Nd_PbI <sub>2</sub>	4.557	6.981	125.533	16.679	0.208	0.0410
3 wt% Nd_PbI <sub>2</sub>	4.559	6.983	125.675	17.008	0.204	0.0394
5 wt% Nd_PbI <sub>2</sub>	4.561	6.983	125.752	18.509	0.187	0.0333

**Table 1.** Lattice parameters of pure and  $\text{Nd}^{3+}$  doped  $\text{PbI}_2$  nanorods and nanosheets.

also calculated from the following relations:  $\epsilon = \frac{\beta \times \cos\theta}{4}$ , and  $\delta = \frac{15\epsilon}{a \times D}$ , respectively. The calculated values of D,  $\epsilon$  and  $\delta$  are given in Table 1. The variation of grain size and unit cell volume is shown in Fig. 3(c), from which it can be concluded that the unit cell volume is found to be increasing with increasing  $\text{Nd}^{3+}$  concentrations and hence lattice parameters a and c. Such variation can be explained with help of Vegard's law, which suggests that when the incorporation process takes place substitutionally or interstitially within a crystal its lattice parameters are going to be shrinking or expanding, respectively<sup>32</sup>. In the present case, knowing the ionic radii of  $\text{Nd}^{3+}$  is 0.11 nm and  $\text{Pb}^{2+}$  is 0.133 nm, the lattice parameters of  $\text{PbI}_2$  are found to be expanding due to doping which suggests the incorporation of  $\text{Nd}^{3+}$  ions within the interstitial sites of the crystalline matrix of the  $\text{PbI}_2$ . We have not observed any amorphization even at higher concentration doping of  $\text{Nd}^{3+}$ , which means that no disorientation was occurred in the prepared  $\text{PbI}_2$  nanostructures due to doping.

**EDX and vibrational analyses.** Figure 4(a) shows the recorded energy dispersive X-ray spectroscopy (EDXS) spectra for  $\text{Nd}^{3+}$  doped  $\text{PbI}_2$  nanosheets. As it is clear from figure that the peaks of  $\text{Nd}^{3+}$  are present in all the doped products. There is an extra energy peak due to Au in the pattern which was sputter on the samples for this measurement. Moreover, the SEM mapping was also carried out (see Fig. 1S supplementary data) to learn about the distribution of dopant inside the parent matrix and shows that the Nd doping has been taken places homogeneously in  $\text{PbI}_2$ . Figure 4(b1) shows the recorded FT-Raman spectrum of the microwave-assisted synthesized pure (nanorods) and  $\text{Nd}^{3+}$  doped  $\text{PbI}_2$  nanosheets. From the figure, it is clear that no vibrational frequency is present in the Raman spectrum due to any impurity, which again proves that the grown material is high-quality nanostructure. The Raman peaks intensity is found to be enhanced in the case of the nanosheets structures at 1 wt% and 3 wt%  $\text{Nd}^{3+}$  doping, however at 5 wt%  $\text{Nd}^{3+}$  doping it was found to be reduced, but it is still higher than that of the pure  $\text{PbI}_2$ .

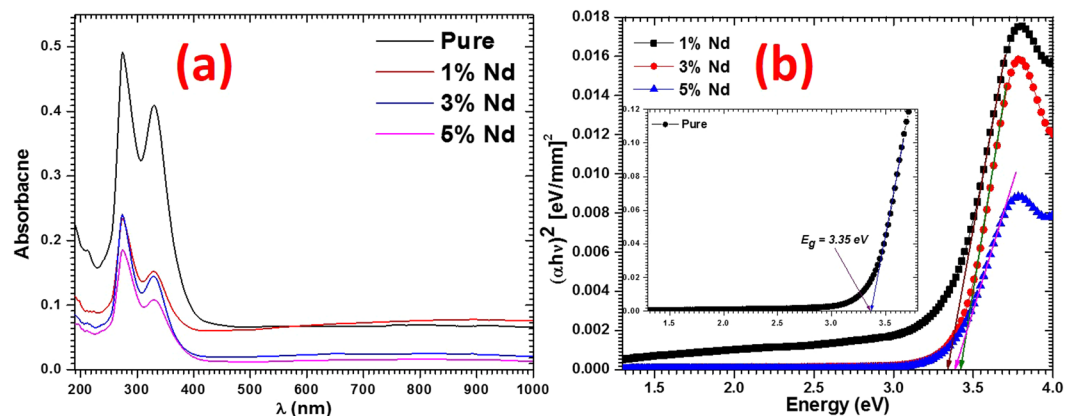




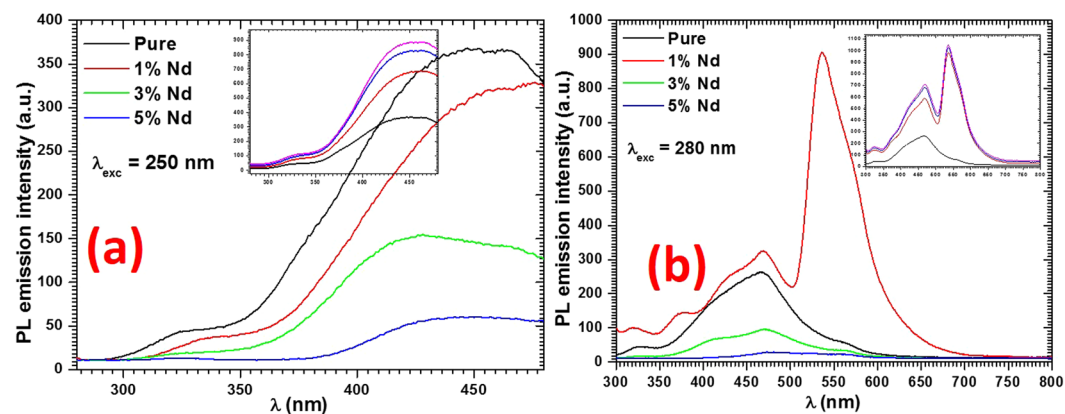
**Figure 4.** EDXS spectra for doped (a) and Raman spectra for pure and doped nanostructures (b1) and for single crystal (b2).

Such high intensity in Raman peaks indicates that the synthesized single crystal nanosheets are more crystalline than that of the nanorods structure of  $\text{PbI}_2$ . The Raman bands for pure  $\text{PbI}_2$  (nanorods) are observed at  $\sim 73, 113, 96, 111, 164, 213 \text{ cm}^{-1}$  and for doped  $\text{PbI}_2$  (nanosheets) of 1 wt%  $\text{Nd}^{3+}$  at 73, 96, 110, 167, 216  $\text{cm}^{-1}$ , for 3 wt%  $\text{Nd}^{3+}$  at 73, 96, 110, 168, 215  $\text{cm}^{-1}$ , and for 5 wt%  $\text{Nd}^{3+}$  at 73, 96, 110, 168, 215  $\text{cm}^{-1}$ , and are assigned to vibration modes:  $E_2^1, A_1^1, A_1^2, 2E_2^1$  and  $2E_1^1$ , respectively. The observed bands confirm the formation of 2H- $\text{PbI}_2$  Polytypes nanorods as well as single crystal nanosheets. These bands are found to be red shifted towards the lower wavenumber as compared to their corresponding values of the bulk material; they are also comparable with other reports on  $\text{PbI}_2$  nanostructures<sup>33,34</sup>. Furthermore, we have also grown the single crystal of  $\text{PbI}_2$  by gel method and recorded the Raman spectrum as shown in Fig. 4(b2). It can be clearly seen that the vibration bands in the grown single crystal are observed at 36, 60.05, 80.91, 104.05, 118.5, 176.1 and 223.6  $\text{cm}^{-1}$ , which shows a clear shifting in the position of Raman bands compared to the fabricated nanostructures of  $\text{PbI}_2$ . The red shifting and broadening in the Raman bands of the synthesized both nanorods and single crystal nanosheets of  $\text{PbI}_2$  confirms the relaxed behavior of the grown nanostructure. Such red shifting and broadening can be explained in terms of the relaxation of the phonon wave vector  $q$ - vector, related to the lattice of the studied crystal within the framed theory of the Raman scattering in solid material. When the fundamental wave vector  $q \approx 0$ , Raman selection rule is relaxed for a predetermined size domain or nanostructure. Such relaxation initiates contribution of phonons away from the Brillouin zone center of the crystal, according to the Heisenberg uncertainty principle. The uncertainty in phonon wave vector goes around as  $\Delta q \approx 1/D$ , where  $D$  indicates the diameter or thickness of the synthesized nanostructures. This spatial confinement of optical phonon modes produces the red shifting and asymmetric broadening in active Raman peaks<sup>35,36</sup>.

**Optical studies.** *UV-Visible spectroscopic analysis.* The absorbances of pure and doped  $\text{PbI}_2$  nanostructures were recorded by preparing their colloidal solutions using methanol as solvent as shown in Fig. 5(a). The energy band gap was evaluated from the recorded absorption spectrum in a transmission mode using the Tauc's plot obtained from Tauc's relation as shown in Fig. 5(b). To find out band gap from this plot we have followed the standard procedure, in which the value of the direct band gap  $E_g$ , was determined to be correspond to the point of the intersection at  $(\alpha h\nu)^2 = 0$ , of the extrapolated straight portion of the curve  $(\alpha h\nu)^2$  vs.  $h\nu$ , with the abscissa axis, where  $h\nu$  and  $\alpha$  respectively measure the photon energy and absorption coefficient. For calculating  $\alpha$  we have used well-known Beer–Lambert relation,  $\alpha = 2.303A/d$ , where  $A$  is UV-visible absorbance and  $d$  is the path length of the quartz cuvette (10 mm). The optical band gap is found to be 3.35, 3.34, 3.42 and 3.39 eV for pure and for doped with 1 wt% Nd, 3 wt% Nd and 5 wt% Nd doped  $\text{PbI}_2$  nanostructures, respectively. The band gap value is not much effected by the 1 wt% Nd doping however, at 3 wt% and 5 wt% dopings it was found to be enhanced. It is clear that the values of band gaps of the  $\text{PbI}_2$  nanostructures are enhanced compared to that of bulk  $\text{PbI}_2$  value i.e. 2.3/2.27 eV<sup>37</sup>. The value of band gap for the fabricated nanostructures are enriched for about



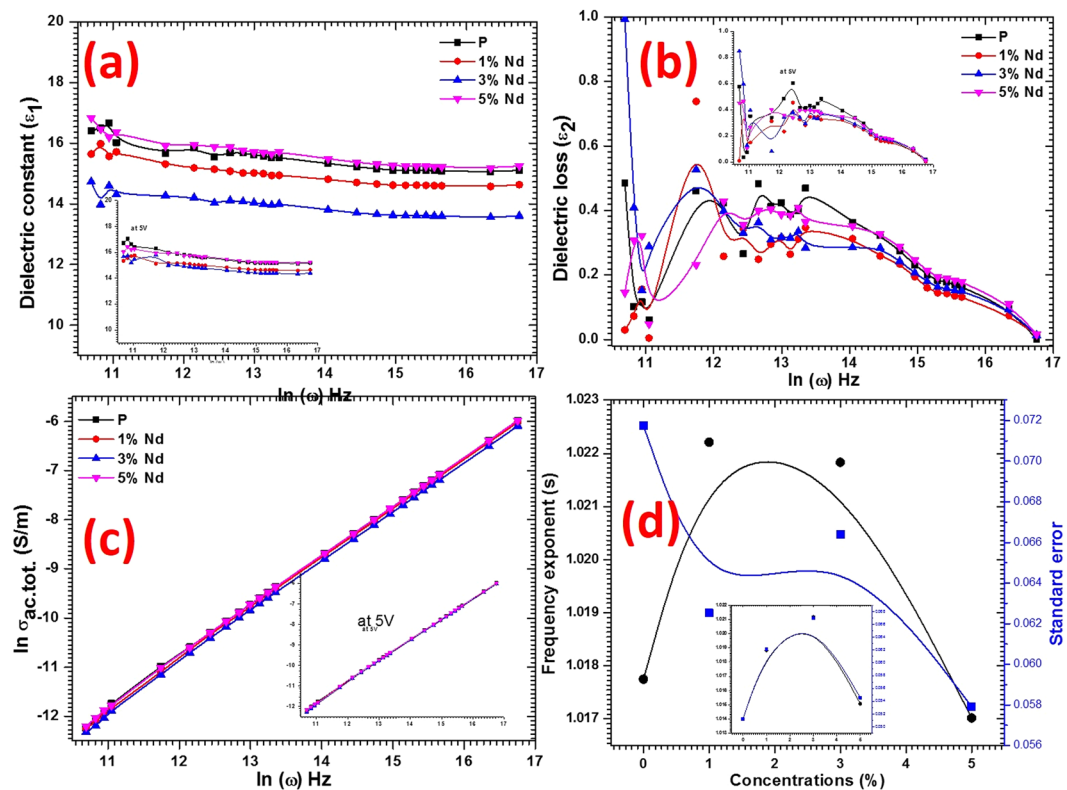
**Figure 5.** (a) Absorbance spectra and (b) energy band gap determination plots.



**Figure 6.** PL spectra for pure and doped PbI<sub>2</sub> nanostructures.

$1.07 \pm 0.08$  eV compared to bulk. This gives a strong indication about the confinement effect taking place within the prepared nanostructures<sup>38,39</sup>. From such observation, we can assume that the UV-Visible method for extracting the optical properties is a suitable and more reliable technique for studying semiconductor nanostructured materials, particularly PbI<sub>2</sub> since it is not changing color even when doped; PbI<sub>2</sub> has a deep yellow color. Hence we will not get proper results on nanostructured PbI<sub>2</sub> by the DR analysis, which provides the data from the materials surface only. In a previous study on doped PbI<sub>2</sub> in which we have reported band gap results using DR technique, one can see the discrepancies between DR values of the band gaps and the present results obtained by UV-Vis data<sup>5,22</sup>. The high value of the band gap of the nanostructured Nd-doped PbI<sub>2</sub> is a positive characteristic that will allow such material to be utilized in optoelectronic devices and it suggests that Nd-doped PbI<sub>2</sub> may be operated at higher voltages, frequency and temperature when compared to usual semiconductor materials. Due to such advantage of the fabricated nanorods and nanosheets of pure and Nd<sup>3+</sup>-doped PbI<sub>2</sub>, it may be used to fabricate a more powerful, cheaper and energy-efficient electrical system<sup>28,40–43</sup>.

**Photoluminescence (PL) analysis.** Figure 6(a) and (b) shows the measured PL emission spectra for the fabricated nanorods and single crystal nanosheets of PbI<sub>2</sub> at two excitation wavelengths, i.e., 250 and 280 nm. It may be mentioned here that the figures presented in the inset of figure (a) and (b) are for more clarity (cumulative) to get the right and more clear information. For PL measurement, the colloidal solutions of the fabricated nanostructures were prepared in methanol. It is evident from both figures that the emission spectra of pure and Nd-doped PbI<sub>2</sub> consist of two emission peaks. The emission spectra recorded at 250 nm excitation wavelength have two broad peaks, i.e., at  $\sim 325 \pm 2$  and  $440 \pm 2$  nm in pure and Nd<sup>3+</sup>-doped PbI<sub>2</sub>. However, when such materials are excited at 280 nm, there are three peaks which are at  $325 \pm 2$ ,  $466 \pm 2$  and  $536 \pm 2$  nm in pure and Nd<sup>3+</sup>-doped PbI<sub>2</sub>. In this PL spectra, some more bands are also observed as shoulder peaks at 409 nm and 563 nm in pure, while in Nd<sup>3+</sup>-doped nanostructures, some more bands are also found at  $373 \pm 1$  nm,  $409 \pm 2$  nm and  $577 \pm 2$  nm as shoulder peaks with high intensity compared to the pure one. Some of them confirm the presence of Nd<sup>3+</sup> doping in PbI<sub>2</sub>. However, more details regarding this dopant in the titled materials can be observed in the higher wavelength range as seen for the same doping in other materials<sup>27</sup>. In fact, more bands can be found for Nd<sup>3+</sup>, but that depends on the ability of detection of the instrument. The band observed at 325 nm is due to band edge emission and the band observed at 409 nm is corresponding to energy emitted due to free excitons,  $E_p$ , while the other is due to bound excitons,  $E_B$ <sup>38,44,45</sup>. It can be noticed that when the fabricated nanostructures are excited at 250 nm, a blue luminescence at  $\sim 440$  nm was observed. However, when excited at 280 nm



**Figure 7.** Plots for (a)  $\epsilon_1$ , (b)  $\epsilon_2$ , (c)  $\sigma_{ac.tot.}$  and (d)  $s$  values of pure and doped  $\text{PbI}_2$  nanostructures.

this band is observed at  $\sim 466$  nm in all of the prepared nanostructures. The PL intensity for 1%  $\text{Nd}^{3+}$  doped  $\text{PbI}_2$  is enriched remarkably, however for doped  $\text{PbI}_2$  at higher concentrations: the intensity of the PL spectra quenched down when excited at 280 nm. Such observations clearly indicates that there is a fundamental changes in the type and density of the defects within  $\text{PbI}_2$  due to doping, hence affecting the radiative and non-radiative recombination processes within the studied material due the increase in the increase of number of defects, dislocations, surface/interface and grain boundary states by the  $\text{Nd}^{3+}$  doping<sup>27</sup>.

**Dielectric and ac electrical conductivity analyses.** For the dielectric functions and electrical conductivity studies, the capacitance (C), impedance (Z), and loss tangent ( $\tan\delta$ ) are measured in the high-frequency range from 7 kHz to 3 MHz at 0 and 5 V bias voltage. Using the obtained data the dielectric constant ( $\epsilon_1$ ) was calculated by the well-known formula:  $\epsilon_1 = \frac{C \times l}{\epsilon_0 \times A}$ , where  $l$  is the thickness of the sample,  $\epsilon_0$  is the vacuum space permittivity and  $A$  is the cross section area. Further, the dielectric loss ( $\epsilon_2$ ) was evaluated using the relation:  $\epsilon_2 = \tan \delta \times \epsilon_1$ . The variation of the calculated values of  $\epsilon_1$  and  $\epsilon_2$  as a function of frequency at the 0 bias voltage are shown in Fig. 7(a) and (b), respectively. The dielectric functions calculated at 5 V bias voltage are also shown in the inset of Fig. 7(a) and (b). It is evident from Fig. 7(a) that  $\epsilon_1$  values are almost constant in the entire testing frequency range for pure and  $\text{Nd}^{3+}$  doped  $\text{PbI}_2$  nanorods and single crystal nanosheets, which means that the electric field dipole does not follow the alternating field past a certain frequency for the studied material<sup>46–50</sup>. The high and stable value in the tested frequency range for the prepared nanostructures confirms their potential uses in optoelectronic devices applications. It can also be noticed from Fig. 7 that the values of  $\epsilon_1$  are found to be lower, when compared to the  $\epsilon_1$  values of the undoped  $\text{PbI}_2$ , due to  $\text{Nd}^{3+}$  doping up to 3 wt%, however at 5 wt%  $\text{Nd}^{3+}$  doping the dielectric constants tended to be higher than that of the pure  $\text{PbI}_2$ . The same type of trend was observed at the 5 V bias voltage. The average value of  $\epsilon_1$  at 0 bias voltage for pure, and 1%, 3% and 5 wt%  $\text{Nd}^{3+}$  doped  $\text{PbI}_2$  are found to be  $15.5 \pm 0.1$ ,  $14.5 \pm 0.1$ ,  $13.5 \pm 0.1$  and  $16.16 \pm 0.1$ , respectively, however these values are found to be enhanced and higher than  $14.5 \pm 0.1$  for all the grown nanostructures at 5 V bias voltage [see inset of Fig. 7(a)]. These values for the prepared nanostructures are found to be comparable as well as higher than the values reported for bulk  $\text{PbI}_2$ <sup>3,5,51</sup>. The reduction in the value of  $\epsilon_1$  at 1 and 3%  $\text{Nd}^{3+}$  doping is may be due to interfacial charge transfer. Such type of results were also reported previously for  $\text{Nd}^{3+}$  doped nanostructures<sup>52</sup>. Similar behavior was also observed for  $\epsilon_2$  and can be seen in Fig. 7(b). The low loss values confirmed that the prepared nanostructures contain lesser defects. The alternating current total conductivity ( $\sigma_{ac.tot.}$ ) was determined from the impedance measurement by the following relation:  $\sigma_{ac.tot.} = \frac{l}{Z \times A}$ . Figure 7 (c) shows the variation of calculated  $\sigma_{ac.tot.}$  with frequency. From figure it can be conclude that the value of  $\sigma_{ac.tot.}$  is increasing with increasing the frequency, and following the universal power law of frequency. It is also found to have similar trends as  $\epsilon_1$  and  $\epsilon_2$  functions. Furthermore, the well-know Jonscher law,  $\sigma_{ac.tot.} = \sigma_{dc} + B\omega^s$  (where  $\sigma_{dc}$  is the direct current conduc-

tivity,  $B$  is a constant and  $\omega$  is angular frequency and  $s$  is an exponent of frequency), was used to recognize the conduction behavior within the investigated material. The  $s$  values were noted for the grown nanostructures from the slope of linear part of the curve  $\ln \sigma_{ac}$  vs.  $\ln \omega$  [Fig. 7(c)] and presents in Fig. 7(d). The value of  $s$  is found to be increased first and then decreased with doping concentrations. Its average value is about 1, which confirms that the hopping mechanism of conduction in the studied material involves a translational motion with sudden carrier hopping within the prepared material nanostructures.

## Conclusions

Facile and rapid synthesis of well-defined and homogeneous morphology nanorods and single crystal nano-sheets of undoped and Nd-doped  $\text{PbI}_2$  have been achieved successfully using microwave-assisted method for the first time. The X-ray diffraction and FT-Raman analysis confirms the formation of hexagonal phase of  $2\text{H-PbI}_2$  polytypic. The presence of  $\text{Nd}^{3+}$  doping in the fabricated nanosheets of  $\text{PbI}_2$  was confirmed by EDX and EDX mapping analyses. SEM study confirms the formation of nanorods of diameter in the range of  $\sim 70$  to  $100$  nm for the pure and single crystal nanosheets with a thickness in the range of  $\sim 20$ – $30$  nm for  $\text{Nd}^{3+}$  doped  $\text{PbI}_2$  with homogeneous morphology. A red shift in the vibrational modes of the Raman spectra of the grown nanostructures of  $\text{PbI}_2$  was observed in comparison to the respective modes of the bulk material, which clearly indicates the formation of the more relaxed nanostructure within the grown lead iodide structures. The values of the band gap ( $E_g$ ) calculated from UV-Visible data are found to be 3.35, 3.34, 3.42 and 3.39 eV for the pure and for the 1%, 3%, 5% -  $\text{Nd}^{3+}$  doped  $\text{PbI}_2$  nanostructures, respectively. The PL spectra of the grown nanostructured  $\text{PbI}_2$  show a blue luminescence at  $\sim 440$  nm when excited at 250 nm and at  $\sim 466$  nm when excited at 280 nm in all fabricated samples. The average values of dielectric constant at 0 bias voltage are found to be  $\sim 15.5$ , 14.5, 13.5 and 16.16 for the pure, and for the 1%, 3% and 5 wt%  $\text{Nd}^{3+}$  doped  $\text{PbI}_2$  respectively. However, these values are found to be enhancing at 5 V bias voltage. The ac total electrical conductivity is also found to be enhancing with higher frequency. All obtained results suggest that the currently prepared single-crystal nanosheets of  $\text{PbI}_2$  are may be of great interest and potential for modern optoelectronic devices and applications.

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## Author Contributions

(Mohd. Shkir): Have synthesized the final product and written the manuscript. (S. AlFaify): Has contributed in performing the characterization of the final product and help in writing and modifying the language.

## Additional Information

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