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# **UV-Induced Charge-Transfer States in Short Guanosine-Containing DNA Oligonucleotides**

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Charge transfer has proven to be an important mechanism in DNA photochemistry. In particular, guanine (dG) plays a major role as an electron donor, but the photophysical dynamics of dG-containing charge-transfer states have not been extensively investigated so far. Here, we use UV pump (266 nm) and picosecond IR probe (~5-7 μm) spectroscopy to study ultrafast dynamics in dG-containing short oligonucleotides as a function of sequence and length. For the pure purine oligomers, we observed lifetimes for the charge-transfer states of the order of several hundreds of picoseconds, regardless of the oligonucleotide length. In contrast, pyrimidine-containing dinucleotides d (GT) and d(GC) show much faster relaxation dynamics in the 10 to 30 ps range. In all studied nucleotides, the charge-transfer states are formed with an efficiency of the order of ~50%. These photophysical characteristics will lead to an improved understanding of DNA damage and repair processes.

irradiation. This repair is explained by a charge transfer from the

of dG-containing short oligonucleotides. We applied transient

UV-pump/mid-IR probe techniques to characterize short-lived

intermediate states. We identified transient charge-transfer

states through their vibrational modes and determined lifetimes

and population as a function of sequence and oligonucleotide

length. The experimental results are discussed in the context of

In this study, we investigated the photophysical properties

transiently formed charge-transfer state d(G°+A°-).[5a]

#### Introduction

UV irradiation induces a selection pressure to which DNA molecules have been exposed for several billions of years on the surface of the earth.<sup>[1]</sup> It is assumed that the selection of the four canonical DNA mononucleotides is related to photostability provided by the ultrashort lifetimes of the excited states. [2] As information carrier, DNA contains the individual canonical nucleotides in strands of defined sequences. This geometric arrangement drastically alters the photophysics. In particular, the interaction between adjacent nucleotides opens alternative reaction channels such as the formation of dinucleotide lesions or of UV-induced charge-transfer states.[3]

Charge-transfer states are highly reactive and can, on the one hand, cause DNA damage.[4] On the other hand, chargetransfer states have been shown to promote the repair of DNA lesions.<sup>[5]</sup> The repair activity strongly depends on the sequence where in particular dG-containing sections are a key factor to self-repair, as dG acts as a strong electron donor. [6] Recently it has been shown that the cyclobutane thymine-dimer lesion (T=T) in the sequence d(GAT=T) is strongly repaired under UV

**Results and Discussion** 

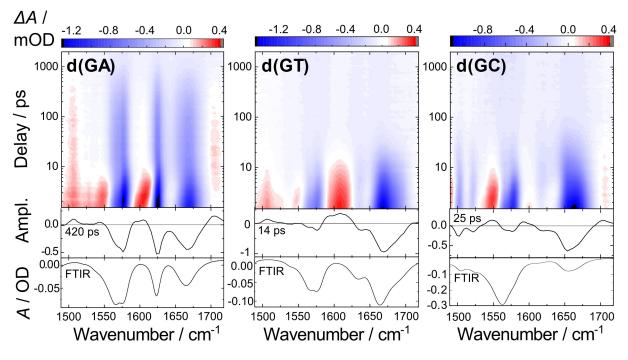
DNA-photoactivity.

In a first set of experiments, we investigated the dinucleotides d (GA), d(GT) and d(GC) by using transient UV-pump and mid-IR probe spectroscopy (Figure 1). The upper part shows the transient absorption difference spectra  $\Delta A$  (color coded) plotted on a logarithmic timescale. Negative spectral bands (blue) are due to depopulation of the initial ground state and occur at wavenumbers corresponding to the bands found in the FTIR absorption spectra (Figure 1, bottom). Positive bands (red) are due to induced absorption, characteristic for excited states and photoproducts and are used, to identify the nature of the intermediates. At early times, within the first 10 ps, the absorption dynamics are dominated by excited state relaxation and vibrational cooling. Subsequent absorption dynamics can be assigned to charge-transfer states and are discussed here in detail. Typically, the data can be modeled by a global fit with two time constants  $\tau_1$  (in the 5 ps range, relaxation of excitation and vibrational cooling) and  $\tau_{\rm 2}$  (related to the slower charge transfer dynamics). The time constants for all investigated samples are given in Table 1. The fit amplitudes, that is, the decay-associated difference spectra (DADS) related to  $\tau_2$  allow the spectroscopic characterization of the corresponding intermediate state (Figure 1).

The data for d(GA) are presented in the left part of Figure 1. The spectral transients can be fitted by exponential functions

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**Figure 1.** Transient IR absorption difference spectra of three dG-containing dinucleotides after excitation at 266 nm. Top: Absorption changes plotted in contour representation (red: positive, blue: negative absorption changes). Bottom: Fitted decay-associated difference spectra (DADS) corresponding to intermediate states with the indicated time constants. The inverted ground state spectra (FTIR) are shown for comparison.

**Table 1.** Excited-state lifetimes and quantum yields of the charge separated states. All data are fitted with a multi-exponential decay model, including 3 time constants for GC ( $\tau_0$ =1.1±0.4 ps) and 2 time constants for all other samples, following a global fitting routine. <sup>[14]</sup> The quantum yields are estimated from the amplitudes of the absorption bleach in the DADS spectra (for details see the Supporting Information).

Sample	$\tau_1$ /[ps]	$ au_2$ /[ps]	фст/[%]
d(GA)	4±1	420 ± 120	42±20
d(GT)	$5\pm3$	$14\pm5$	$66 \pm 35*$
d(GC)	$4\pm1$	$25\pm 9$	$54 \pm 30*$
d(AG)	$4\pm1$	$280\pm160$	$32\pm15$
d(GAAG)	$5\pm 2$	$500\pm180$	$56\pm30$
d(GAGA)	$5\pm 2$	$630\pm200$	$55\pm30$
d(GA) <sub>10</sub>	$11\pm4$	$500\pm140$	$63\pm30$

\* For GT and GC the short decay time  $\tau_2$  imposes additional difficulties for a quantitative determination of the quantum yields  $\phi_{CT}$ . Here the yields should be taken only as a rough estimate.

with time constants of  $\tau_1 = 4 \pm 1$  ps and  $\tau_2 = 420 \pm 120$  ps. The three negative bleaching bands at 1576, 1627 and 1673 cm<sup>-1</sup> are striking. The wave numbers correspond to the positions in the stationary FTIR absorption spectrum of d(GA) (Figure 1 bottom left). The two bands at 1576 and 1673 cm<sup>-1</sup> are due to the bleaching of guanine, and the band at 1627 cm<sup>-1</sup> is due to the bleaching of adenine absorption after excitation. The features found at early delay times (< ca. 5 ps)

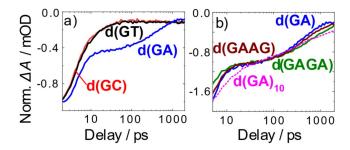
correspond to absorption changes expected from the cooling of a vibrationally hot ground-state. The final recovery of the ground state absorption occurs on the 400 ps timescale. The DADS of the long-lived  $\tau_2$ =420 ps component is plotted in Figure 1 left, center. Here the positive peaks at 1608 cm<sup>-1</sup> and

1704 cm $^{-1}$  are due to the absorption of the guanine radical cation. $^{[3c,9]}$  The weak bands around 1510 and 1543 cm $^{-1}$  correspond to an adenine radical anion band (calculated spectra in Figure S2). In agreement with previous studies we can assign the long-lived component to the charge-transfer state  $d(G^{\bullet+}A^{\bullet-}).^{[3c]}$ 

Results from the dinucleotide d(AG) are shown in Figure S1. They agree well with those obtained for d(GA). The marker bands in the DADS are the same. The observed time constants are  $\tau_1 = 4 \pm 1$  ps and  $\tau_2 = 280 \pm 160$  ps. The somewhat shorter time constant  $\tau_2$  is in agreement with previous studies in the UV and visible range. [2c,3b,10]

The d(GG) dinucleotide is not treated in details since at a concentration of 5 mM, required for transient IR measurements, the UV/Vis spectrum deviates from the spectrum at low concentrations (Figure S4).<sup>[11]</sup>

In contrast to the pure purine dinucleotides d(GA) and d (AG) the pyrimidine containing dinucleotides d(GT) and d(GC) display a much faster recovery of the ground state (Figure 1). The data can be fitted again with a sum of exponentials d(GT):  $\tau_1 = 5 \pm 3$ ,  $\tau_2 = 14 \pm 5$  ps, d(GC):  $\tau_0 = 1.1 \pm 0.4$  ps,  $\tau_1 = 4 \pm 1$  ps,  $\tau_2 = 25 \pm 9$  ps. The transient spectral changes on early times can again be assigned to vibrational cooling. However, the time constant  $\tau_2$  of the slower process is found to be much shorter for the pyrimidine containing dinucleotides. This is evident when the temporal evolution of the ground state bleach at the guanine band (at 1576 cm<sup>-1</sup>) is plotted in Figure 2a for d(GT), d (GC) and d(GA). In the pyrimidine containing dimers the decay  $\tau_2$  is accelerated by a factor of about 10. The DADS for d(GT) and d(GC) (central part of Figure 1) clearly show that the



**Figure 2.** Time dependence of the IR absorption changes at the guanine ground-state band around 1576 cm $^{-1}$  after excitation at 266 nm. a) dG-containing dinucleotides (normalized at the local minimum around 1.5 ps). b) dG- and dA-containing oligonucleotides (normalized at a delay time t=51 ps).

process  $\tau_2$  not only represents the recovery of the ground state bands but also exhibits the characteristic positive marker bands for charge-transfer states. For d(GC) positive amplitudes are found at 1704, 1592 (dG<sup>•+</sup>) and 1549 cm<sup>-1</sup> (dC<sup>•-</sup>).<sup>[3c]</sup> The lifetime of the charge-transfer state of d(GC) also agrees with d(CG) reported in a study by Takaya and coworkers.<sup>[3b]</sup> For d(GT) the guanine cation absorption at 1608 cm<sup>-1</sup> is clearly seen while a shoulder appears at > 1704 cm<sup>-1</sup>.

The broad positive band in the DADS of d(GT) around 1608 cm<sup>-1</sup> may be caused by the overlap of the bands of the dT<sup>•-</sup> anion and the dG<sup>•+</sup> cation.<sup>[3d]</sup> The corresponding calculated spectra are shown in Figure S2. The comparison of measured absorption changes (DADS) with the calculated band positions (see the Supporting Information) clearly shows, that the longer-lived components in all studied dinucleotides are

radical pair states where the electron is transferred from the electron donor dG to the acceptor dA, dT or dC. The radical pair state has a lifetime in the picosecond range (15 to 600 ps). The efficiency of this charge transfer reaction can be estimated from the amplitudes of the absorption bleach in the DADS spectra (for details see the Supporting Information). The efficiency of charge transfer ( $\phi_{CT}$ ) amounts to around 50%. Details are given in Table 1 for the different oligonucleotides. Although chargetransfer states are formed in all dinucleotides with similar amplitudes, there is a strong acceleration of the charge recombination in the pyrimidine containing dinucleotides. This effect can be described by Marcus theory. [3b,12] The charge recombination typically takes place in the Marcus inverted regime, where a lower thermodynamic driving force is related with a faster charge recombination. Since the driving force, related to the redox potential difference is lower for the purinepyrimidine than for the purine-purine recombination, [6a,12] the lifetimes  $\tau_2$  should be shorter for d(GC), d(GT), which is observed experimentally (Table 1).

In a second set of experiments, we address the influence of oligonucleotide length on the formation and lifetime of charge-separated states. For this purpose, we choose the two tetranucleotides d(GAGA), d(GAAG) and the alternating 20-mer sequence d(GA)<sub>10</sub> as samples (Figure 3). The over-all spectral and temporal features resemble those of the d(GA) dimer: The band positions are very similar. The absorption transients at later times also involve biexponential kinetics with the slow components in the 500 ps range. The corresponding absorption features are again indicative for the presence of radical pair states. Figure 2b shows the normalized absorption transients at the ground state absorption band of dG (ca. 1576 cm<sup>-1</sup>) for the

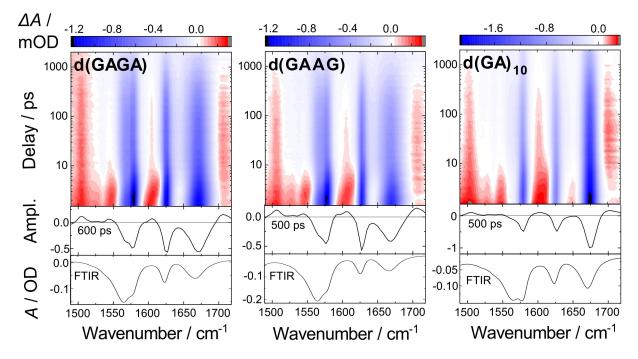


Figure 3. Transient IR absorption difference spectra of three dG- and dA-containing oligonucleotides after excitation at 266 nm. Top: Experimental data in contour representation (red: positive, blue: negative absorption changes). Bottom: Fitted decay associated difference spectra (DADS) corresponding to intermediate states with the indicated time constants. The inverted ground state spectrum (FTIR) is shown for comparison.



different d(GA)-containing oligonucleotides. In this plot, the qualitative similarity of the transients becomes apparent. In details the curves show some weak differences, pointing to different decay times (results given in Table 1). For d(GA)<sub>10</sub> the observed differences can be assigned to structural inhomogeneities. The data on the oligonucleotides d(GAGA), d(GAAG), and d(GA)<sub>10</sub> show a weak trend towards longer lifetimes  $\tau_2$  of the radical pair states as compared to those of the dinucleotides d(GA) and d(AG). For d(GAGA) we find 630  $\pm$  200 ps, for d(GAAG) 500  $\pm$  180 ps and for d(GA)<sub>10</sub> 500  $\pm$  140 ps.

## **Conclusions**

Time-resolved IR spectroscopy was used to study photophysical dynamics of dG-containing di- and oligonucleotides. After excitation in the UV at 266 nm we found transient absorption changes pointing to excitation decay and vibrational cooling on the 5 ps timescale. Additional slower decaying intermediates were observed in all samples and assigned to charge-transfer states by comparison with the literature and by density functional calculations. While the lifetimes of the charge-transfer states of d(GT) and d(GC) were around 20 ps, the pure purine containing oligonucleotides showed lifetimes in the 300 to 600 ps range. We found yields of the charge-transfer states in the range of 50%. The observations clearly show that transient charge-transfer states are efficiently formed in dG-containing oligomers. The known reactivity and the high formation yield make charge-transfer states important promoters for secondary reactions in DNA or in its direct surroundings.

### **Experimental Section**

The oligonucleotides were purchased from biomers.net (UIm, Germany). The lyophilized and HPLC purified samples were dissolved in 50 mM phosphate buffer (KH $_2$ PO $_4$ , Na $_2$ HPO $_4$ ) in D $_2$ O (Sigma Aldrich). The oligo concentrations were kept at ~10 mM per base. The pump probe experiments were performed in a flow cell with BaF $_2$  windows (sample thickness 100  $\mu$ m) under ambient oxygen conditions at a temperature of about 23 °C.

The basic principles of the pump probe setup for time-resolved experiments has been described previously. [15] The light source in the pump and probe setup was a Ti:Sa based laser amplifier system (Libra, Coherent) with an output pulse duration of ~113 fs at 800 nm and a repetition rate of 1 kHz. The third harmonic (266 nm) of the fundamental was used for the excitation (pump).<sup>[15c]</sup> The excitation energy at the sample position was 2.2 µJ with a spot diameter of 150 µm FWHM. The pump pulse was transmitted through a fused silica rod of 25 cm length to increase its pulse duration to about 1 ps. The broadband mid-IR light between 5-7 μm (probe) was generated by a cascade of two optical parametric amplifiers (OPA) producing signal and idler pulses at 1400 and 1800 nm and difference frequency mixing in a silver thiogallate crystal (AgGaS<sub>2</sub>). Pump and probe pulses were focused and spatially overlapped in the sample. After passing the sample the probe light was spectrally dispersed (Chromex 250IS, Bruker) and detected on a 64-channel MCT array (IR-0144, Infrared Systems Development). All experiments were performed under magic angle conditions. The analysis of the transient data uses global fitting where the uncertainty of the time-constants was determined by an exhaustive search procedure. The transient and stationary IR spectroscopy and the density functional theory calculations follow the procedures described by Bucher et al. [3c,16]

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#### Conflict of Interest

The authors declare no conflict of interest.

**Keywords:** charge transfer ⋅ DNA damage ⋅ guanine photophysics ⋅ picosecond IR spectroscopy

- a) S. Ranjan, R. Wordsworth, D. D. Sasselov, Astrophys. J. 2017, 843, 110;
   b) S. A. Benner, E. A. Bell, E. Biondi, R. Brasser, T. Carell, H. J. Kim, S. J. Mojzsis, A. Omran, M. A. Pasek, D. Trail, ChemSystemsChem 2019; c) S. J. Mojzsis, G. Arrhenius, K. D. McKeegan, T. Harrison, A. Nutman, C. Friend, Nature 1996, 384, 55–59.
- [2] a) B. Cohen, P. M. Hare, B. Kohler, J. Am. Chem. Soc. 2003, 125, 13594–13601; b) C. E. Crespo-Hernández, B. Cohen, P. M. Hare, B. Kohler, Chem. Rev. 2004, 104, 1977–2020; c) A. A. Beckstead, Y. Zhang, M. S. de Vries, B. Kohler, Phys. Chem. Chem. Phys. 2016, 18, 24228–24238.
- [3] a) G. W. Doorley, M. Wojdyla, G. W. Watson, M. Towrie, A. W. Parker, J. M. Kelly, S. J. Quinn, J. Phys. Chem. Lett. 2013, 4, 2739–2744; b) T. Takaya, C. Su, K. de La Harpe, C. E. Crespo-Hernández, B. Kohler, Proc. Natl. Acad. Sci. USA 2008, 105, 10285–10290; c) D. B. Bucher, B. M. Pilles, T. Carell, W. Zinth, Proc. Natl. Acad. Sci. USA. 2014, 111, 4369–4374; d) B. M. Pilles, D. B. Bucher, L. Liu, P. Gilch, W. Zinth, W. J. Schreier, Chem. Commun. 2014, 50, 15623–15626.
- [4] a) S. Marguet, D. Markovitsi, J. Am. Chem. Soc. 2005, 127, 5780-5781; b) M. Barbatti, A. C. Borin, S. Ullrich, Photoinduced Phenomena in Nucleic Acids II: DNA Fragments and Phenomenological Aspects, Vol. 356, Springer, 2014; c) G. Lin, L. Li, Angew. Chem. Int. Ed. 2010, 49, 9926–9929; d) D. Pörschke, Proc. Natl. Acad. Sci. USA 1973, 70, 2683-2686; e) S. Kumar, P. C. Joshi, N. D. Sharma, S. N. Bose, R. J. H. Davies, N. Takeda, J. A. McCloskey, Nucleic Acids Res. 1991, 19, 2841-2847; f) H. Görner, J. Photochem. Photobiol. B 1991, 10, 91-110; g) M. Gomez-Mendoza, A. Banyasz, T. Douki, D. Markovitsi, J.-L. Ravanat, J. Phys. Chem. Lett. 2016, 7, 3945-3948; h) G. S. Madugundu, J. R. Wagner, J. Cadet, K. Kropachev, B. H. Yun, N. E. Geacintov, V. Shafirovich, Chem. Res. Toxicol. 2013, 26, 1031-1033; i) F. Bergeron, F. Auvré, J. P. Radicella, J.-L. Ravanat, Proc. Natl. Acad. Sci. USA 2010, 107, 5528-5533; j) M. Münzel, C. Szeibert, A. F. Glas, D. Globisch, T. Carell, J. Am. Chem. Soc. 2011, 133, 5186-5189; k) T. Matsuura, Bioorganic Photochemistry: Photochemistry and the Nucleic Acids, Vol. 1, Editor: Harry Morrison, 1990, John Wiley & Sons Inc., New York
- [5] a) D. B. Bucher, C. L. Kufner, A. Schlueter, T. Carell, W. Zinth, J. Am. Chem. Soc. 2016, 138, 186–190; b) R. Szabla, H. Kruse, P. Stadlbauer, J. Šponer, A. L. Sobolewski, Chem. Sci. 2018, 9, 3131–3140; c) B. Giese, B. Carl, T. Carl, T. Carell, C. Behrens, U. Hennecke, O. Schiemann, E. Feresin, Angew. Chem. Int. Ed. 2004, 43, 1848–1851; d) S. Breeger, U. Hennecke, T. Carell, J. Am. Chem. Soc. 2004, 126, 1302–1303; e) C. Behrens, T. Carell, Chem. Commun. 2003, 1632–1633; f) A. Schwögler, L. T. Burgdorf, T. Carell, Angew. Chem. Int. Ed. 2000, 39, 3918–3920.
- [6] a) Y. Paukku, G. Hill, J. Phys. Chem. A 2011, 115, 4804–4810; b) Z. Pan, J. Chen, W. J. Schreier, B. Kohler, F. D. Lewis, J. Phys. Chem. B 2011, 116, 698–704; c) Z. Pan, M. Hariharan, J. D. Arkin, A. S. Jalilov, M. McCullagh,



- G. C. Schatz, F. D. Lewis, *J. Am. Chem. Soc.* **2011**, *133*, 20793–20798; d) R. Chakraborty, D. Ghosh, *Phys. Chem. Chem. Phys.* **2016**, *18*, 6526–6533.
- [7] M. Banyay, M. Sarkar, A. Graslund, Biophys. Chem. 2003, 104, 477-488.
- [8] a) P. Hamm, S. Ohline, W. Zinth, J. Chem. Phys. 1997, 106, 519–529;
  b) W. J. Schreier, T. E. Schrader, F. O. Koller, P. Gilch, C. E. Crespo-Hernández, V. N. Swaminathan, T. Carell, W. Zinth, B. Kohler, Science 2007, 315, 625–629;
  c) W. J. Schreier, J. Kubon, N. Regner, K. Haiser, T. E. Schrader, W. Zinth, P. Clivio, P. Gilch, J. Am. Chem. Soc. 2009, 131, 5038–5039;
  d) W. J. Schreier, P. Gilch, W. Zinth, Annu. Rev. Phys. Chem. 2015, 66, 497–519;
  e) S. Laimgruber, W. J. Schreier, T. Schrader, F. Koller, W. Zinth, P. Gilch, Anaew. Chem. Int. Ed. 2005, 44, 7901–7904.
- [9] A. W. Parker, C. Y. Lin, M. W. George, M. Towrie, M. K. Kuimova, J. Phys. Chem. B 2010, 114, 3660–3667.
- [10] M. C. Stuhldreier, F. Temps, Faraday Discuss. 2013, 163, 173–188.
- [11] E. Balanikas, A. Banyasz, G. Baldacchino, D. Markovitsi, Molecules 2019, 24, 2347.
- [12] Y. Zhang, J. Dood, A. A. Beckstead, X.-B. Li, K. V. Nguyen, C. J. Burrows, R. Improta, B. Kohler, J. Phys. Chem. B 2015, 119, 7491–7502.
- [13] I. Kejnovská, J. Kypr, M. Vorlíčková, Chirality 2003, 15, 584-592.
- [14] H. Satzger, W. Zinth, Chem. Phys. 2003, 295, 287–295.
- [15] a) T. Schrader, A. Sieg, F. Koller, W. Schreier, Q. An, W. Zinth, P. Gilch, Chem. Phys. Lett. 2004, 392, 358–364; b) K. Haiser, B. P. Fingerhut, K. Heil, A. Glas, T. T. Herzog, B. M. Pilles, W. J. Schreier, W. Zinth, R.

- de Vivie-Riedle, T. Carell, *Angew. Chem. Int. Ed.* **2012**, *51*, 408–411; c) W. Schreier, **2008**.
- [16] M. J. Frisch, G. W. Trucks, H. B. Schlegel, G. E. Scuseria, R. M. A., C. J. R., J. Montgomery, J. A., T. Vreven, K. N. Kudin, J. C. Burant, J. M. Millam, S. S. Iyengar, J. Tomasi, V. Barone, B. Mennucci, M. Cossi, G. Scalmani, N. Rega, G. A. Petersson, H. Nakatsuji, M. Hada, M. Ehara, K. Toyota, R. Fukuda, J. Hasegawa, M. Ishida, T. Nakajima, Y. Honda, O. Kitao, H. Nakai, M. Klene, X. Li, J. E. Knox, H. P. Hratchian, J. B. Cross, V. Bakken, C. Adamo, J. Jaramillo, R. Gomperts, R. E. Stratmann, O. Yazyev, A. J. Austin, R. Cammi, C. Pomelli, J. W. Ochterski, P. Y. Ayala, K. Morokuma, G. A. Voth, P. Salvador, J. J. Dannenberg, V. G. Zakrzewski, S. Dapprich, A. D. Daniels, M. C. Strain, O. Farkas, D. K. Malick, A. D. Rabuck, K. Raghavachari, J. B. Foresman, J. V. Ortiz, Q. Cui, A. G. Baboul, S. Clifford, J. Cioslowski, B. B. Stefanov, G. Liu, A. Liashenko, P. Piskorz, I. Komaromi, R. L. Martin, D. J. Fox, T. Keith, M. A. Al-Laham, C. Y. Peng, A. Nanayakkara, M. Challacombe, P. M. W. Gill, B. Johnson, W. Chen, M. W. Wong, C. Gonzalez, J. A. Pople, Gaussian Inc., Wallingford, 2004.

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