



Elasticity of individual protocadherin 15 molecules implicates tip links as the gating springs for hearing

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Hair cells, the sensory receptors of the inner ear, respond to mechanical forces originating from sounds and accelerations. An essential feature of each hair cell is an array of filamentous tip links, consisting of the proteins protocadherin 15 (PCDH15) and cadherin 23 (CDH23), whose tension is thought to directly gate the cell's transduction channels. These links are considered far too stiff to represent the gating springs that convert hair bundle displacement into forces capable of opening the channels, and no mechanism has been suggested through which tip-link stiffness could be varied to accommodate hair cells of distinct frequency sensitivity in different receptor organs and animals. Consequently, the gating spring's identity and mechanism of operation remain central questions in sensory neuroscience. Using a high-precision optical trap, we show that an individual monomer of PCDH15 acts as an entropic spring that is much softer than its enthalpic stiffness alone would suggest. This low stiffness implies that the protein is a significant part of the gating spring that controls a hair cell's transduction channels. The tip link's entropic nature then allows for stiffness control through modulation of its tension. We find that a PCDH15 molecule is unstable under tension and exhibits a rich variety of reversible unfolding events that are augmented when the Ca²⁺ concentration is reduced to physiological levels. Therefore, tip link tension and Ca²⁺ concentration are likely parameters through which nature tunes a gating spring's mechanical properties.

auditory system | entropic stiffness | hair cell | optical trap | vestibular system

Mechanically gated ion channels are ubiquitous. In addition to underlying our senses of hearing, balance, and touch, they are involved in the regulation of such processes as muscle extension, blood pressure, pulmonary inflation, and visceral distension. These channels are opened and closed through the action of gating springs, elastic elements that are tensioned by mechanical stimulation and in turn communicate stress to the molecular gates of the respective channels. Gating springs accordingly store mechanical energy and use it to regulate channels' open probabilities. For bacterial mechanoreceptors, which respond to osmotic stress, the cellular membrane itself serves as a gating spring (1). The ubiquitous Piezo channels of vertebrates extend three membrane-embedded arms that likely act as gating springs by flexing in response to membrane stretching (2, 3). Other mechanosensitive channels, such as NOMPC (TRPN1) in *Drosophila*, appear to be gated by the tension in elastic ankyrin domains (4).

Gating springs were first posited for hair cells of the vertebrate inner ear, the sensors of the auditory and vestibular systems (5–7). Each hair cell is surmounted by a hair bundle—a cluster of erect, actin-filled processes termed stereocilia—that is deflected by mechanical stimulation. However, the identity of the gating springs in these cells has remained controversial. A plausible candidate discovered soon after the gating-spring hypothesis was advanced is the tip link. Extending approximately 150 nm between the tip of each stereocilium and the side of its longest neighbor, the tip link is positioned to sense the shear between

stereocilia when a hair bundle is deflected (Fig. 1 *A* and *B*) (8–10). The tip link is a dimer of parallel dimers, comprising two PCDH15 molecules joined at their amino termini to a pair of CDH23 molecules through a “handshake,” whose stability depends on the presence of bound Ca²⁺ ions (Fig. 1*C*) (11).

The mechanical properties of hair bundles imply a gating-spring stiffness (12, 13) on the order of 1 mN·m⁻¹ (1 pN/nm). However, electron microscopy images suggest that the tip link is relatively rigid (14), and crystallographic studies and molecular dynamics simulations of the relevant cadherins support a stiffness 50-fold greater than that measured (15). Therefore, it has been posited that most of a gating spring's elasticity resides at a tip link's two attachments, rather than within the link itself. However, it is possible that the tip link elongates in response to force by the straightening of its lateral thermal undulations, the working principle of an entropic spring (16). To clarify the identity of the hair cell's gating spring, we examined the elastic properties of a tip link protein.

Results

Mechanical Characteristics of PCDH15. The mechanical properties of a protein can be tested by tethering it between two surfaces, applying a force that pulls the surfaces apart, and measuring the protein tether's elongation (17). To directly explore the stiffness of a tip link, we used a high-precision optical trap to determine

Significance

Our hearing depends on mechanosensitive channels in hair cells of the inner ear. Experiments suggest that each channel is opened by a “gating spring,” an elastic element that conveys displacement of a hair bundle to the channel. Appropriate stiffness of the gating spring permits the discrimination of different sound amplitudes; if the spring is too stiff, then a faint sound will elicit the same response as a loud sound, opening all of a cell's channels. Although the tip link—a fine molecular filament—might be the gating spring, its properties have remained controversial. Using high-precision optical tweezers, we demonstrate that the mechanical properties of a tip link protein correlate with those of a gating spring in vivo.

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at high loading rates (27). Even in the absence of an acoustic stimulus, tip links experience a constant resting tension (28) that varies with the frequency sensitivity of a hair cell, from 5 pN at 1 kHz to 34 pN at 4 kHz, with possibly even greater tensions at higher frequencies (13). In each tip link, two copies of PCDH15 act in parallel, and each copy assumes one-half of the tension of the entire link, approximately 3 pN up to at least 17 pN. During stimulation from a moderately loud sound, with a 60-dB sound pressure level, a hair bundle is displaced (29) by <10 nm. Such a stimulus stretches the gating springs (12) by up to 1.4 nm, superimposing an oscillation of only a few piconewtons onto the resting tension. Therefore, in our experiments we had to apply slowly changing forces to explore the influence of resting tension on the mechanical properties of PCDH15.

Unless noted otherwise, we applied force ramps to the single molecule tethers at a loading rate of $130 \text{ pN}\cdot\text{s}^{-1}$ by linearly increasing the spring constant of the stimulus trap (*SI Appendix, Fig. S1*). The dynamic forces during normal hearing at a frequency of 1 kHz vary much faster, at roughly $5,000 \text{ pN}\cdot\text{s}^{-1}$. Our loading rate represented a compromise between slow force application and our desire to collect a statistically relevant number of extension-relaxation cycles for each molecule in a reasonable amount of time. For each cycle, we ramped the force up to 60 pN to cover the entire range of physiologically relevant tensions, then returned it at the same rate to a holding level of 2–4 pN. Depending on the Ca^{2+} concentration, we adjusted the holding level and duration to allow the protein to refold domains after many, but not all, cycles. The chosen loading rate likewise led to unfolding events in only a subset of cycles. With these parameter choices, we were able to trap the protein in a given conformational state for several extension-relaxation cycles, allowing us to precisely characterize the mechanics of each state.

Representative examples of individual force-extension relationships for Ca^{2+} concentrations of 3 mM, 20 μM , and 0 M feature hockey stick-like shapes, as expected for the extension of a biopolymer in a heat bath (Fig. 2 *A–C*) (30). As we quantify below, this functional shape indicates that entropic effects dominate PCDH15's elastic response. Abrupt, stepwise extensions, or “rips,” in the force-extension relationship correspond to structural changes of the protein under force. In contrast to typical single-molecule experiments, under our loading conditions, PCDH15 never fully unfolded during the extension phase of the stimulus; therefore, we frequently observed extensional structural changes even during the relaxation phase of the stimulus (B_U ; Fig. 2*B*). For each Ca^{2+} concentration, a set of conformational changes led to a modulated occupation of the force-extension state space, which we visualize by overlapping hundreds of extension-relaxation cycles for one representative molecule apiece (Fig. 2 *D–F*).

Conformational Changes. Rips in the force-extension relations at physiological forces suggest that PCDH15 exists in different structural states during hearing. At a saturating Ca^{2+} concentration, the conformational states accessible to the protein are limited; the state-space heatmap reveals only two major configurations (Fig. 2 *A* and *D*). The second of these configurations is further divided into two substates separated by a difference of only a few nanometers in contour length. State 1 reflects the extensibility of fully folded PCDH15 (*SI Appendix, Note S1*). By fitting a polymer model to the force-extension relations for five molecules, we find that state 2 arises from a combination of two classes of conformational changes, leading to mean \pm SEM elongations of $4.0 \pm 0.2 \text{ nm}$ and $15.8 \pm 0.7 \text{ nm}$ (respectively, distributions of A_U and B_U in Fig. 2*G* and Table 1). The structural origin of these conformational changes is unknown and difficult to determine owing to the large size of the protein; however, we can rule out the unfolding of entire cadherin domains as the origin of the unfolding events. The length of the

folded peptide in each of the 11 cadherin domains ranges from 94 to 123 aa, with a mean of 104 residues. At a contour length of 0.39 nm per residue (31), the unfolding of each cadherin domain is expected to augment the contour length by approximately 36 nm, an elongation of 40.6 nm minus 4.5 nm to account for loss of the folded cadherin domain. The observed contour length changes of the elongations A_U and B_U therefore represent protein rearrangements less extreme than the complete unfolding of a cadherin domain. At a saturating Ca^{2+} concentration and physiological forces, we never observed length changes in the wild-type protein compatible with full unfolding.

We next asked whether there are unique structural features in PCDH15 that give rise to the elongations A_U and B_U , or whether several different conformational changes, each with a similar contour length change, underlay the observed distributions. In most of the extension-relaxation cycles, we did not observe more than one of either class of events (*SI Appendix, Fig. S2*). The rare occasions in which several events A_U or B_U were detected in a single trace were not reproducible across trials or protein molecules and thus could be explained as the rupture of nonspecific interactions between the protein and either of the confining surfaces. We conclude that a single, unique structural alteration of PCDH15 is responsible for event A_U , whereas a distinct structural change results in event B_U , precluding the occurrence of several events of either type within the same cycle. Interestingly, the force sensitivity of events A_U is much narrower than that of events B_U (Fig. 2*J*). From these distributions, we determined the statistical dependence of both classes. An event A_U generally follows an event B_U ($P < 0.05$ and $P < 0.1$ for three and two of the five tested molecules, respectively). It is plausible that both structural changes resulted from the same cadherin domain, with elongation B_U leading to a destabilization that facilitated elongation A_U .

We found the protein in state 1 at the beginning of many extension-relaxation cycles and concluded that there is a high probability of refolding of both event A_U and event B_U between cycles. Indeed, we routinely detected refolding events A_F during the relaxation phase of our protocol (Fig. 2 *A* and *G*), but rarely observed refolding events B_F (*SI Appendix, Fig. S3*). The latter events likely occurred only at very low forces, for which the slight shortening was lost in Brownian noise (*SI Appendix, Fig. S4*).

We next reduced the Ca^{2+} concentration to the physiological value of 20 μM and exposed a tethered PCDH15 protein to the same force protocol. The extension-relaxation cycles showed mean conformational changes identical to those of the previously observed classes (A_U : $3.9 \pm 0.4 \text{ nm}$; B_U : $15.0 \pm 0.6 \text{ nm}$; $n = 8$) (Fig. 2*H*). At this Ca^{2+} concentration, however, an additional class of unfolding events emerged with a mean contour length change of $35 \pm 1 \text{ nm}$ (C_U in Fig. 2 *B* and *H*), in excellent agreement with the elongation expected for unfolding of an entire cadherin domain (32). At holding forces of 2–4 pN, the refolding of cadherin domains was a slow process and occurred on a time scale of seconds, in line with other proteins that feature Ig-like motifs (33). In some extension-relaxation cycles, we observed the successive unfolding of several cadherin domains (*SI Appendix, Fig. S5*). Because unfolding of any of the 11 extracellular cadherin domains should increase the contour length by a similar amount, we could neither assign unfolding events to particular domains nor elucidate the sequence in which the domains unraveled.

The heatmap of all force-extension relationships originated from a mixture of the unfolding events A_U , B_U , and C_U and their respective refolding events. The contour lengths that gave rise to the annotated states 1–8 (Fig. 2*E*) were a consequence of the unfolding of up to three cadherin domains in series with up to one unfolding event of type B_U . Structural changes of types A_U and A_F , which are clearly visible in the heatmap for a Ca^{2+} concentration of 3 mM as a subdivision of state 2 (Fig. 2*D*), are

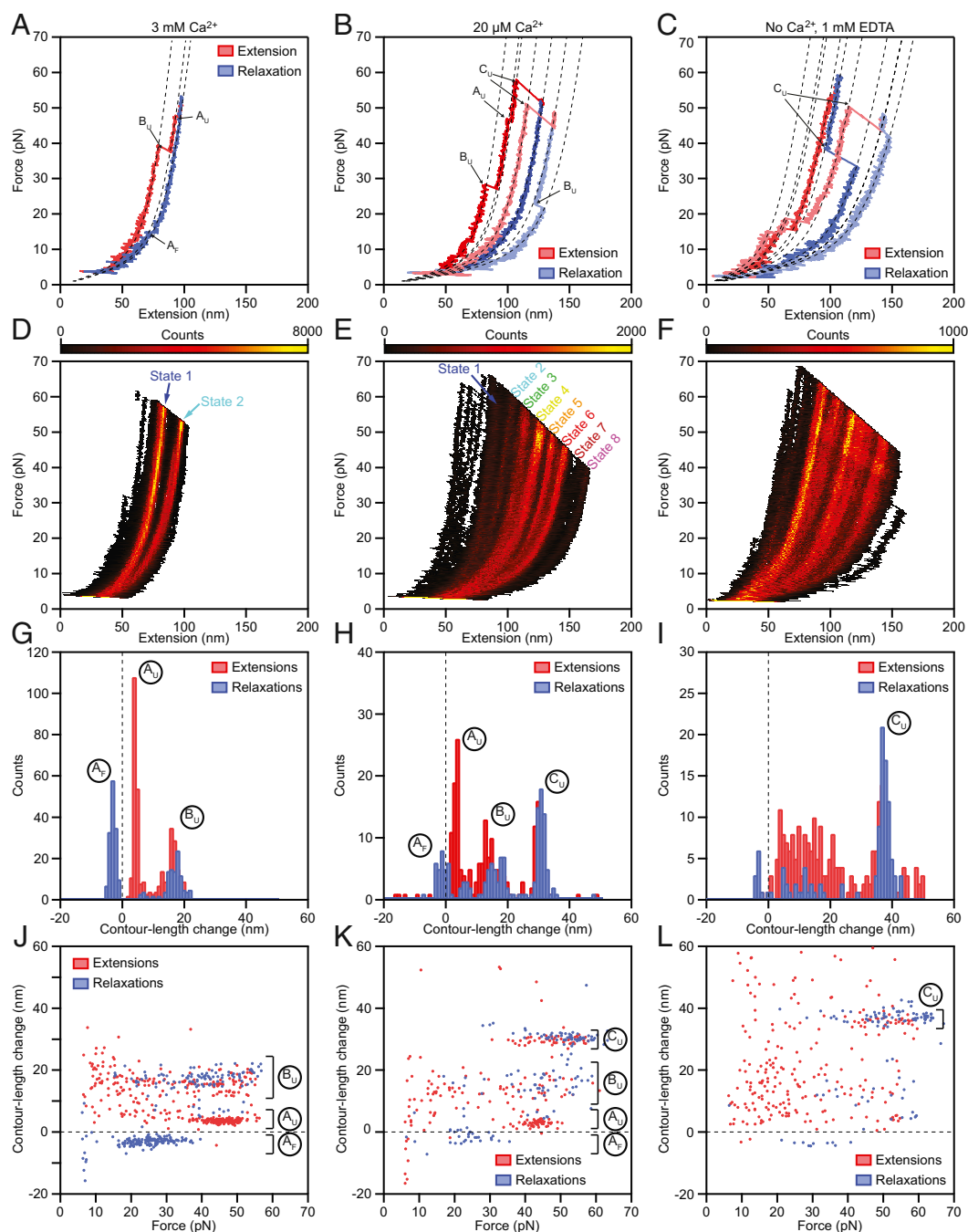


Fig. 2. Force-extension measurements of PCDH15 monomers. (A) At a Ca²⁺ concentration of 3 mM, individual force-extension cycles show two distinct classes of abrupt elongations, the unfolding events A_U and B_U, as well as refolding events of class A_F. The dashed lines represent fits to the trajectories by a protein model. (B) Reducing the Ca²⁺ concentration to 20 μM elicits an additional class of unfolding events, C_U, corresponding to the unfolding of entire cadherin domains. (C) In the absence of Ca²⁺, unclassifiable structural changes occur in conjunction with the well-defined events C_U. (D–F) Heatmaps displaying all the force-extension cycles for a single representative molecule at each Ca²⁺ concentration. The data were binned into pixels of 1 nm × 0.1 pN. A much smaller portion of the state space is accessible at a Ca²⁺ concentration of 3 mM than at a Ca²⁺ concentration of 20 μM or in the absence of Ca²⁺. The heatmaps illustrate which contour lengths of the tethered molecules occurred with increased likelihood during our force-loading protocol as an average over both extensions and relaxations. Prominent regions of elevated occupancy are labeled states 1–8, in which state 1 corresponds to the fully folded protein and state 8 results from the unfolding of three cadherin domains in series with one event of type B_U. (G–I) Histograms of the contour length changes of all abrupt elongations verify that these rips can be grouped into classes A_F, A_U, B_U, and C_U at Ca²⁺ concentrations of 3 mM and 20 μM. In the absence of Ca²⁺, most of the contour length changes are more broadly distributed. (J–L) Plots of the contour length change of every rip against the force at which that event occurred revealing the force distributions of each class of structural change. Note that the extensions never completely unfolded a PCDH15 molecule, so elongations could occur even during the relaxation phases. Because the contour lengths observed for the folded protein correspond to the length of monomeric PCDH15 in series with its molecular anchors, the measured extensions exceed those expected for the protein alone. Our analysis corrects for this influence, resulting in values in excellent agreement with the known structure of the protein (Table 1). All force-extension cycles were sampled at intervals of 10 μs and smoothed to a temporal resolution of 1 ms. The waiting times between cycles were 0.2 s for a Ca²⁺ concentration of 3 mM, 2 s for a Ca²⁺ concentration of 20 μM, and 4 s in the absence of Ca²⁺. The number of cycles recorded was 500 for a Ca²⁺ concentration of 3 mM and 200 for a Ca²⁺ concentration of 20 μM and in the absence of Ca²⁺.

Table 1. Material property values of PCDH15

[Ca ²⁺]	<i>b</i> , nm	<i>l</i> _{linker} , nm	<i>k</i> _{folded} , mN/m	Δ <i>l</i> _A , nm	Δ <i>l</i> _B , nm	Δ <i>l</i> _C , nm
3 mM	2.9 ± 0.5	1.4 ± 0.5	9 ± 4	4.0 ± 0.2	15.8 ± 0.7	—
20 μM	3.1 ± 0.4	1.3 ± 0.5	11 ± 5	3.9 ± 0.4	15.0 ± 0.6	35 ± 1

The parameters of a polymer model (*SI Appendix*) were as follows: *b*, length of each stiff segment of the folded protein; *l*_{linker}, contour length of each flexible linker between the stiff segments; *k*_{folded}, enthalpic stiffness of the folded protein; Δ*l*_A, contour length increase, type A; Δ*l*_B, contour length increase, type B; Δ*l*_C, contour length increase, type C. The parameter values were determined from fits of the model to the force-extension curves in Fig. 2. The values for the folded protein are largely independent of the Ca²⁺ concentration. Δ*l*_C is in good agreement with the expected contour-length change per unfolded cadherin domain. Fitted parameter values are given as mean ± SEM for five experiments at a Ca²⁺ concentration of 3 mM and for eight experiments at a Ca²⁺ concentration of 20 μM.

not apparent in the state space for a Ca²⁺ concentration of 20 μM. It is possible that there were Ca²⁺-dependent sub-nanometer changes in the contour length that averaged out small effects of events A_U and A_F in the heatmap. We again investigated a potential sequential dependence of unfolding classes by their force distributions (Fig. 2*K*). Even though events of type A_U followed events of type B_U at a Ca²⁺ concentration of 3 mM, at a physiological Ca²⁺ concentration these structural changes in eight molecules were independent of one another (*P* > 0.1 for each molecule). Moreover, events of types A_U and B_U were independent of events of type C_U, the unfolding of entire cadherin domains.

We next investigated the mechanics of PCDH15 in the absence of Ca²⁺. Representative force-extension relationships featured a plethora of conformational changes (Fig. 2*C* and *I*), many of which could no longer be clearly grouped into any of the classes A_{U,F} and B_U. Events with a mean of 37 ± 2 nm (for five molecules) continued to characterize a well-defined class C_U. The heatmap of all extension-relaxation cycles had a structure reminiscent of that at a Ca²⁺ concentration of 20 μM (Fig. 2*F*). The structure in the absence of Ca²⁺ arose from the unfolding of a discrete number of cadherin domains in series with the unclassifiable shorter structural changes that likely represented the partial unfolding of one or more domains. This lack of well-defined short structural changes was also evident from the force distribution of the observed rips (Fig. 2*L*).

Stiffness of PCDH15. It is unknown whether tip link cadherins are completely or only partially folded during normal hearing. Therefore, we investigated the stiffness not only of folded PCDH15, but also that of conformational states with a progressively greater number of unfolded domains. The total stiffness of PCDH15 comprises both enthalpic and entropic components, whose contributions we quantified by fitting the force-extension relationships with a model of the protein. A freely jointed chain (34) formed by the 11 folded cadherin domains, in series with a worm-like chain (35) representing the 10 unstructured linker regions, captured the protein's disorder under tension and reflected its entropic extensibility. The protein's enthalpic extensibility, which likely originates from deformations of its secondary and tertiary structures rather than from unfolding, was modeled through a Hookean spring constant. We included an additional worm-like chain to model any unfolded portions of the protein (*SI Appendix*, Fig. S6).

Because the unfolded polypeptide chains and the linker regions are structurally similar, we modeled both with the same persistence length. Fits to the data for 13 molecules yielded a mean *l*_{peptide} value of 0.49 ± 0.04 nm. For folded PCDH15 at a Ca²⁺ concentration of 3 mM, we found a mean length of 2.9 ±

0.5 nm for each of the 11 solid segments of the chain, a mean length of 1.4 ± 0.5 nm for each of the 10 flexible linkers between the solid segments, and a mean enthalpic spring constant of 9 ± 4 mN·m⁻¹ for the Hookean stiffness of the protein (*n* = 5 molecules) (Table 1). The full length of a solid segment combined with its associated linker region was 4.3 ± 0.7 nm, in excellent agreement with the value of 4.5 nm per cadherin repeat from crystal structures of cadherin domains (22). Much to our surprise, these values did not change when the Ca²⁺ concentration was decreased to 20 μM, the physiological level in the cochlea. An elevated Ca²⁺ concentration stabilizes cadherin domains against unfolding but does not augment the stiffness of the folded protein. The stiffness predicted by our model is in good agreement with the slopes of the different states in the state-space heatmap (Fig. 3).

Across all states and Ca²⁺ concentrations, the measured and predicted stiffness of the protein is much smaller than its enthalpic stiffness of ~10 mN·m⁻¹. The additional compliance is entropic, arising from the thermal motion of the individual cadherin domains and from thermal undulations in the interdomain linker regions and unfolded polypeptide chains. When PCDH15 is tensed, this thermal kinking is smoothed out, and the protein elongates. The progressive unfolding of domains further softens the protein by introducing additional disordered polypeptide chains (Fig. 3*B*). At high forces, most thermal bends have been straightened, and the enthalpic elasticity begins to dominate the protein's response. Importantly, we find that for physiological tensions, the protein's response to force is determined by entropic elasticity. The protein's stiffness approaches its enthalpic value only for unphysiologically high tensions (Fig. 3 and *SI Appendix*, Fig. S7).

Unfolding of Cadherin Domains Under Forces Relevant for Hearing.

Elevated tension not only increases the stiffness of PCDH15, but also increases the likelihood that entire cadherin domains unfold. Do cadherin domains unfold during normal hearing? If so, do they refold under physiological conditions, or could tip links with persistently unstructured regions exist *in vivo*?

To determine the unfolding rate of cadherin domains under physiological tensions, we transformed the force distributions of type C_U unfolding events into unfolding rates as a function of constant force (Fig. 4) (27). For a given unfolding event, we could not determine which of the 11 cadherin domains had unfolded, so our result is an average over several or all of the domains. The transformation also assumed no cooperativity between the unfolding of individual domains. If the unfolding of one domain were to increase the probability that an adjacent domain would unravel, for example, then our approach would have systematically overestimated the rate at which fully folded tip links unfold.

We computed unfolding rates for Ca²⁺ concentrations of both 20 μM and 0 M (Fig. 4). As expected, the presence of Ca²⁺ stabilized cadherin against force-induced unfolding; the domains unfolded many times faster in the ion's absence than in its presence. With a physiological Ca²⁺ concentration and at 20 pN of tension, the upper range of physiological resting values, a single cadherin domain unfolded at a rate of approximately once every 100 s. A fully folded PCDH15 molecule—consisting of 11 domains—then unfolded a domain every 10 s. At even higher tensions, this rate rapidly increased to roughly 0.4 s⁻¹ for the unfolding of a single cadherin domain at a tension of 60 pN. Because unfolding events at forces of 10–20 pN were extremely rare and might have corresponded to transitions from molten-globule states (*SI Appendix*, Note S2), we were unable to reliably calculate force-dependent unfolding rates for this force range. Nevertheless, extrapolation of the available data suggests a rate of ~0.003 s⁻¹ at 10 pN of tension.

These results indicate that unfolding of cadherin domains would not occur within individual cycles of an auditory stimulus. For physiological forces, the unfolding rate is much too low to

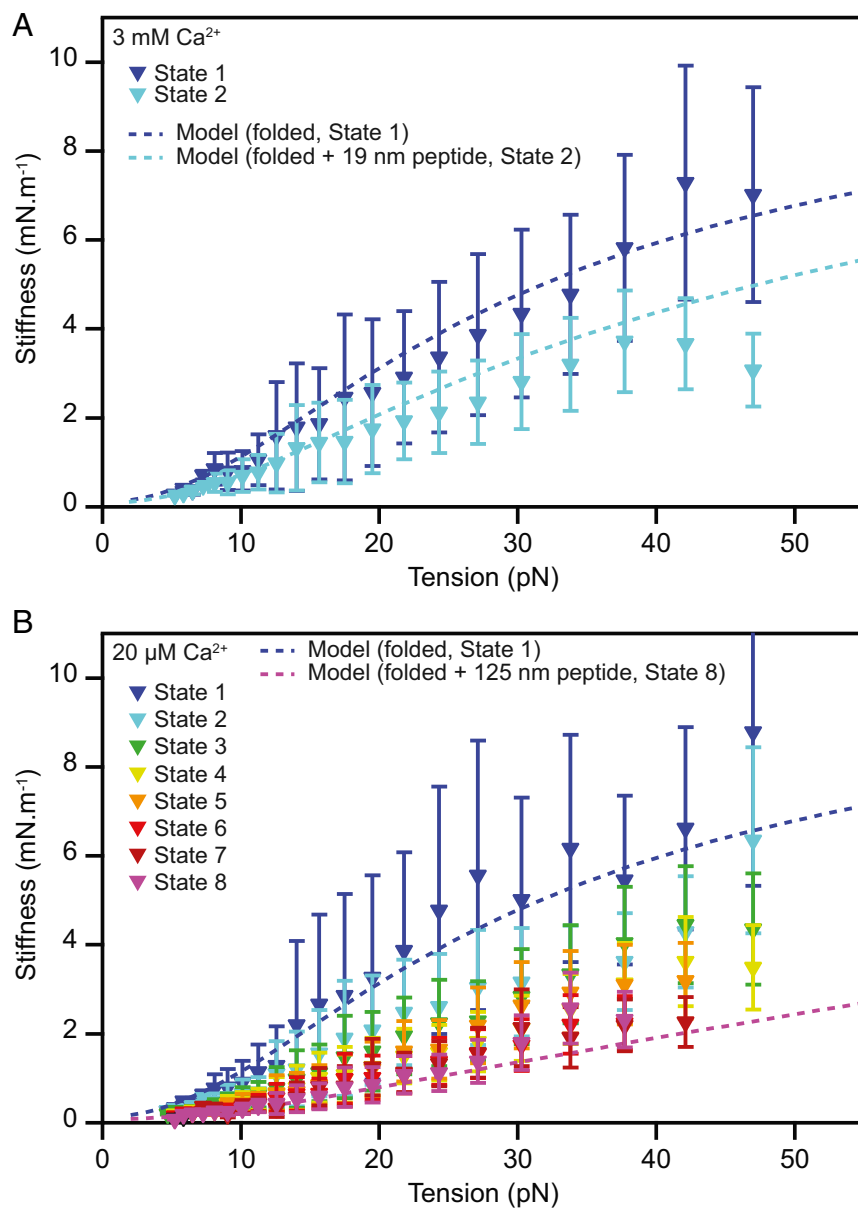


Fig. 3. Stiffness of monomeric PCDH15. (A) The stiffnesses of the different conformational states of PCDH15 at a Ca^{2+} concentration of 3 mM correspond to the slopes of the highly occupied regions of the state space in Fig. 2 D and E and are corrected for the stiffness of the molecular tags and anchors. The dark-blue dashed line represents the stiffness of our model of state 1, the fully folded protein, with the parameter values of Table 1 ($b = 3.0$ nm; $l_{\text{linker}} = 1.35$ nm; $k_{\text{folded}} = 10$ $\text{mN}\cdot\text{m}^{-1}$). Parameter values were averaged over both Ca^{2+} concentrations. The light-blue dashed line represents the model for state 2, with an additional 19-nm segment of unfolded protein with a persistence length of 0.49 nm representing the combined effect of the events A_U and B_U . (B) The corresponding data for a Ca^{2+} concentration of 20 μM capture a variety of unfolding events leading to states 2–8. The dark-blue dashed line represents a model of the fully folded protein (state 1); the pink dashed line depicts the modeled stiffness of the protein in state 8, with an unstructured 125-nm-long peptide to represent the unfolding of three cadherin domains in series with contour length changes of 15 nm and 4 nm. The experimental data are mean \pm SEM for five molecules at a Ca^{2+} concentration of 3 mM and for six molecules at a Ca^{2+} concentration of 20 μM .

follow stimuli with frequencies ranging from hundreds to thousands of hertz. However, provided that at the link's resting tension, a domain's unfolding rate exceeds its refolding rate, cadherin domains in a tensed tip link could exist in a permanently unfolded state. We never observed refolding of cadherin domains during any of the recorded extension-relaxation cycles, even though such events should be readily detectable at tensions exceeding 4 pN (*SI Appendix, Fig. S4*). Instead, refolding occurred only during the holding phase between successive cycles, provided that the holding force was <4 pN and the waiting time was on the order of several seconds.

We conclude that for tensions >20 pN per molecule, the unfolding rate—although very small—exceeds the refolding rate. Our data indicate that this is also the case for the force range of 4–20 pN, but owing to the possible influence of molten-globule states, we could not determine this with certainty. Future experiments with force jumps and constant forces might shed light on the existence and behavior of molten-globule states in this force regime. Our result suggests that some cadherin domains of tip links *in vivo* exist in a perpetually unfolded state. Such unfolded

states would decrease the protein's stiffness (Fig. 3B) and could be a mechanism by which a tip link softens even under high tension.

The critical force at which the unfolding and refolding rates of a cadherin domain are equal remains to be determined. However, the giant muscle protein titin, which has immunoglobulin folds similar to those of cadherin, exhibits a critical force of 5.4 pN (33). If cadherin domains feature a similar value, the resting tensions in low-frequency hair cells might be less than the critical force and bias the tip link's domains toward a fully folded state, whereas the tip link might remain in a partially unfolded state in hair cells sensitive to high frequencies.

Effect of a Mutation Associated with Hearing Loss. More than 100 mutations of the tip link cadherins cause hearing loss in humans (36). The deletion of residue V767 in EC7 (Fig. 5A) is of particular interest, for it leads to deafness—stemming from a deficit in the cochlea, with its low Ca^{2+} concentration—but not loss of function in the vestibular labyrinth, which has a higher Ca^{2+} concentration (37). This mutation evidently does not hinder

increasing its stiffness. Finally, the arrangement of a dimer of CDH23 in series with the PCDH15 dimer is expected to reduce the stiffness by approximately 70%. At very high tensions, when entropic effects are largely suppressed and enthalpy dominates, we estimate the stiffness of the full-length, dimeric tip link as $6 \text{ mN}\cdot\text{m}^{-1}$ (SI Appendix, Note S3), a value in good agreement with the stiffness of strongly tensioned gating springs in vivo (13). Provided that entropy further softens the tip link at lower tensions, our results suggest that the tip link cadherins are a major component of the gating spring for mechanotransduction in hair cells. The entropic contributions to the mechanics of dimeric PCDH15 and of the entire tip link remain to be determined.

The stiffness of hair bundles has a complicated, nonmonotonic dependence on Ca^{2+} concentration (39, 40), an observation that may be interpreted as the interplay between strain-hardening of the tip links and the unfolding of cadherin domains. Low Ca^{2+} levels cause molecular motors to up-regulate tip link tension (41), which suppresses each link's thermal motion and increases its stiffness. Ca^{2+} also directly affects the rate at which cadherin domains unfold under force, however, resulting in states of PCDH15 with reduced stiffness (Fig. 3). At a tension of 20 pN and a Ca^{2+} concentration of 20 μM , individual cadherin domains unfold an order of magnitude more slowly than in the absence of Ca^{2+} . We never observed the unfolding of entire cadherin domains at a Ca^{2+} concentration of 3 mM. Thus, the Ca^{2+} concentration regulates both the tension on and the structure of PCDH15, the combined effect of which on the molecule's mechanics is hard to predict. Nevertheless, the sensitivity of PCDH15 to Ca^{2+} suggests that the variable concentrations of the ion in different receptor organs tunes the mechanical properties of tip links, and thus of hair bundles, to the organs' specific requirements. Even within a single organ, the guinea pig cochlea, the Ca^{2+} concentration increases fourfold along the tonotopic axis from base to apex (25). This gradient might adjust tip link stiffness to accord with the frequency response of the individual hair cells. Finally, hair cells can enhance the local Ca^{2+} concentration around their hair bundles through the activity of membrane Ca^{2+} pumps (42). This phenomenon raises the interesting possibility that the stiffness of tip links is modulated by the locally varying Ca^{2+} concentration in response to hair cell activity.

We found that the elastic properties of folded PCDH15 are surprisingly independent of the Ca^{2+} concentration. This result seems to contradict the impression conveyed by electron micrographs (43), in which cadherins transition from a disordered globular conformation to a rod-like chain of domains in the presence of progressively larger amounts of Ca^{2+} . Disordered states should make greater entropic contributions to PCDH15's elasticity than ordered states, a difference not apparent in our data. Note, however, that the divergence between electron-microscopic images and single-molecule data has also been observed for other proteins, such as titin (44). A possible explanation is that the configurations of proteins adsorbed to substrates are far from their equilibrium conformations, so that the variations in shape do not accurately capture the thermal motion in solution (45).

The displacement-transduction current relationships of hair bundles reveal that for unphysiologically large displacements in excess of 1 μm , gating springs can stretch (46) by at least 120 nm, a value thought to be incompatible with the extensibility of tip link cadherins. Consequently, it has been suggested that the gating spring's stiffness stems from the elasticity of the plasma membrane

or cytoskeleton into which the tip links insert (47). We have shown that invoking such sources of elasticity is unnecessary; cadherin domains in the tip link proteins can unfold under physiological stimuli, albeit at a low rate, and such unfolding events become very likely at high forces. Recall that physiological hair bundle displacements are on the order of 10 nm and augment the resting tension of each tip link by a few piconewtons. The stimulus in extreme deflections, which is 100-fold greater, is capable of causing the almost instantaneous unfolding of individual cadherin domains. Several unfolded cadherin domains could then account for the extreme tip link elongation of 120 nm.

In further support of domain unfolding, the length distribution of tip links in the bullfrog's hair bundles, as determined by electron-microscopic tomography, features two distinct classes (10) with means close to 110 nm and 170 nm. The unfolding of two cadherin domains per tip link monomer could account for this length difference. Such unfolding events could soften the gating spring at high resting tensions and protect both the tip link and the associated mechanotransduction machinery from damage during loud sounds.

In addition to the unfolding of entire cadherin domains, we also observed partial domain unfolding with contour length increases of 4 nm and 15 nm. Future single-molecule work is needed to elucidate the structural correlates of these conformational changes and to determine their role in hearing. Additional experiments are also needed to test the stiffness of PCDH15 dimers and of the full tip link and to confirm that domain unfolding occurs for those constructs and for tip links in vivo.

Methods

More detailed information is provided in SI Appendix, Materials and Methods.

Preparation of PCDH15 Monomers. Plasmids encoding murine PCDH15 were expressed in suspension-adapted HEK293 cells. To confine monomers in an optical trap, the amino terminus of the protein was labeled with an AviTag that permitted the enzymatic conjugation of biotin. The carboxyl terminus included a SpyTag that could form a covalent linkage with an immobilized SpyCatcher protein. The construct also encoded a signal peptide to promote secretion, short flexible linkers flanking the PCDH15 sequence, and a His tag for purification.

Optical Trapping and Measurements. SpyCatcher molecules were conjugated to 2- μm silicon dioxide pedestal beads that were covalently attached to a glass coverslip. PCDH15 monomers were washed into the sample chamber and allowed to attach covalently to the pedestal beads. To initiate an experiment, we optically trapped a streptavidin-coated 1- μm polystyrene bead and positioned it near a pedestal bead. Using a highly focused 852-nm laser beam, we applied stimulus forces to probe the molecule's mechanical properties. With a second, high-stability 1,064-nm laser, we monitored the position of the probe bead with a precision of ≤ 1 nm in each of three dimensions.

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