## Research Article

# Flexible Diodes/Transistors Based on Tunable p-n-Type Semiconductivity in Graphene/Mn-Co-Ni-O Nanocomposites

Lihong Su [,](https://orcid.org/0000-0001-8889-4624) **1,2** Zhou Yang,**<sup>3</sup>** Xitong Wang,**1,2** Ziao Zou,**<sup>1</sup>** Bo Wang,**1,4** Gary Hodes,**<sup>5</sup>** Ninghui Chang,**<sup>1</sup>** Yongyong Suo,**1,4** Zhibo Ma,**1,6** Haoxu Wang,**3,7** Yucheng Liu,**<sup>3</sup>** Junping Zhang,**<sup>1</sup>** Shuanhu Wang,**1,8** Yuefei Li,**1,2** Fengxia Yang,**1,2** Jixin Zhu,**<sup>9</sup>** Fei Gao,**<sup>3</sup>** Wei Huang<sup>(D[,](https://orcid.org/0000-0001-7004-6408) I, 9</sup> and Shengzhong Liu<sup>n)</sup><sup>[3](https://orcid.org/0000-0002-6338-852X)</sup>

<sup>1</sup>School of Chemistry and Chemical-Engineering, Northwestern Polytechnical University, Xi'an, 710129 Shaanxi, China<br><sup>2</sup>Donaguan Sanhang Civil Military Integration Innovation Institute Donaguan, 53300 Cuangdong, China  $<sup>2</sup>$ Dongguan Sanhang Civil-Military Integration Innovation Institute, Dongguan, 52300 Guangdong, China</sup>

 ${}^{3}$ Laboratory of Applied Surface and Colloid Chemistry, Ministry of Education; Shaanxi Key Laboratory for Advanced Energy Devices; Shaanxi Engineering Lab for Advanced Energy Technology; Institute for Advanced Energy Materials; School of Materials Science and Engineering, Shaanxi Normal University, Xi'an 710119, China <sup>4</sup>

 ${}^{4}$ School of Aeronautics, Northwestern Polytechnical University, Xi'an, 710072 Shaanxi, China

 $5$ Department of Materials and Interfaces, Weizmann Institute of Science, Rehovot 76100, Israel

<sup>6</sup>Key Lab of Micro/Nano Systems for Aerospace, Ministry of Education, Northwestern Polytechnical University, Xi'an, 710129 Shaanxi, China

 $^7$ University of Queensland, Australian Institute for Bioengineering & Nanotechnology, Nanomaterials Centre, St. Lucia, Qld, Australia

<sup>8</sup>School of Physical Science and Technology, Northwestern Polytechnical University, Xi'an, 710129 Shaanxi, China<br><sup>9</sup>Institute of Flexible Electronics, Northwestern Polytechnical University, Xi'an, 710129 Shaanxi, China Institute of Flexible Electronics, Northwestern Polytechnical University, Xi'an, 710129 Shaanxi, China

Correspondence should be addressed to Lihong Su; hlshong@nwpu.edu.cn, Wei Huang; iamwhuang@nwpu.edu.cn, and Shengzhong Liu; liusz@snnu.edu.cn

Received 9 February 2021; Accepted 11 August 2021; Published 13 October 2021

Copyright © 2021 Lihong Su et al. Exclusive Licensee Science and Technology Review Publishing House. Distributed under a Creative Commons Attribution License (CC BY 4.0).

We report a novel Mn-Co-Ni-O (MCN) nanocomposite in which the p-type semiconductivity of Mn-Co-Ni-O can be manipulated by addition of graphene. With an increase of graphene content, the semiconductivity of the nanocomposite can be tuned from p-type through electrically neutral to n-type. The very low effective mass of electrons in graphene facilitates electron tunneling into the MCN, neutralizing holes in the MCN nanoparticles. XPS analysis shows that the multivalent manganese ions in the MCN nanoparticles are chemically reduced by the graphene electrons to lower-valent states. Unlike traditional semiconductor devices, electrons are excited from the filled graphite band into the empty band at the Dirac points from where they move freely in the graphene and tunnel into the MCN. The new composite film demonstrates inherent flexibility, high mobility, short carrier lifetime, and high carrier concentration. This work is useful not only in manufacturing flexible transistors, FETs, and thermosensitive and thermoelectric devices with unique properties but also in providing a new method for future development of 2D-based semiconductors.

## 1. Introduction

Modern transistor manufacturing is based on p-n interface phenomena developed in semiconductor science and technology. The migration and separation of electrons/holes in electronic devices are dependent on carrier lifetimes in the devices. Quantum tunneling at the contact interface only involves <2 nm in interface thickness, but the leakage current caused by this tunneling can result in unreliability of the circuit in the manufacture of transistors below 10 nm. Over the past 70 years, metal/semiconductor contacts were divided into two types: Schottky contact and ohmic contact. The formation of a Schottky barrier is based on the mechanism of carrier electron diffusion and thermionic emission, and its properties are

<span id="page-1-0"></span>

Figure 1: The electrical and chemical properties of complex films with different graphene to MCN ratio. (a) The resistance of 10 *<sup>μ</sup>*m thick single-layer films at 25°C as a function of different graphene content. (b) Deconvolution of Mn XPS peaks into Mn<sup>3+</sup> and Mn<sup>4+</sup> showing how the  $Mn^{3+}$ :  $Mn^{4+}$  ratio increases with an increase in the graphene content.

determined ideally by the work functions of metal and semiconductor and, in practice, often also by interface defects or reaction [[1](#page-6-0), [2\]](#page-6-0). Materials with fixed chemical composition can be p-type, n-type, or intrinsic (or compensated) semiconductors [\[3, 4\]](#page-6-0). We found that the work function of negative temperature coefficient (NTC) p-type Mn-Co-Ni-O (MCN) semiconductor nanoparticles is about 0.27 eV higher than that of graphene, which allows electron transfer from graphene to MCN at the contact interface. With the increase of graphene content, the contact between them changes the conductivity of the composite from p-type through electrically neutral to n-type. The contact barrier is dominated by the near-zero electron mass and the volume effect of the nanometer scale MCN particles (Figure [S1\)](#page-6-0). There is no leakage of tunneling current base on electronic quantum-tunneling-effect in nanoscale components because of the electron mobile near Dirac point. The development of these new semiconductor composite materials, with inherent flexibility, has broad prospects for the design and manufacture of flexible nanoelectronic components, including design and manufacture of chips, transistors, thermoelectric and Hall components, and sensors [\[5](#page-6-0)–[20](#page-6-0)].

## 2. Results

Due to graphene's extremely high electron mobility velocity and zero bandgap characteristics, it is difficult to realize it in semiconducting form without doping or changing the graphene structure. In the past 70 years, contacts between conductors and semiconductors have been divided into ohmic and Schottky contacts. Generally, ohmic contacts have no barrier or only a small one ( $\langle a \rangle$  few hundred mV) and are high recombination contacts. The conductivity of MCN at room temperature

is about  $10^{-6}$  S/m and that of graphene is about  $10^{6}$  S/m. We measured the MCN work function (WF) and that of graphene by two methods—Kelvin probe (KP) and UPS—and take the average between the two measurements [\(SI, F](#page-6-0)). The average WF of graphene is about 4.43 eV while that of the MCN as a ceramic film is about 4.7 eV, (each ±0.05 eV). Therefore, the MCN WF is only ~0.27 eV higher than that of graphene which should result in a close-to-ohmic junction.

We carried out both resistivity and Hall effect measurements (the latter for determining conductivity type and carrier concentration). The contact geometry is shown in Figure [S1](#page-6-0) of the [SI](#page-6-0) and explained in the corresponding text in the [SI](#page-6-0); measurements are made both across the films and through it. This directionality may have a strong effect on measured conductivity type and carrier concentrations. Figure 1(a) shows the resistance values of  $10 \mu m$  thick single-layer films as a function of graphene content (similar graphs for different film thicknesses are shown in Figure [S4\)](#page-6-0). At low graphene concentrations, holes dominate the conductivity. The surface resistance of the composite film decreases nonlinearly (on a semilog plot) when the graphene content increases from zero to about 12% and does not change much when the graphene content increases beyond 12%, when the graphene phase becomes continuous and the film properties are like a conductor. In the intermediate region, the film changes from weak p-type through intrinsic (when a balance is reached between electrons and holes in the film) to increasing n-type as the graphene concentration is increased and finally a conductor dominated by the graphene.

Hall effect measurements show that the MCN-graphene composite films (10 *μ*m thick) change from p-type through



FIGURE 2: Schematic composition diagram of (a) graphene and MCN nanoparticle composite. (b) Holes in MCN are filled by electrons from graphene. Reducing  $Mn^{4+}$  to  $Mn^{3+}$  and  $Mn^{2+}$  and forming holes in the graphene. (c) With high enough graphene concentration, the graphene's electron dominates the n-type conductivity.

intrinsic type to n-type with an increase in the graphene content from 1% to about 13%. The trend of change of resistance and conductivity type of films with different thicknesses is the same, but the critical content of graphene needed to give different conductivity types varies according to thickness. Although the graphene content is less than MCN, the most conductive component in the composite materials is graphene because of its extremely large surface area as well as its much higher conductivity.

The change in  $Mn^{3+}$ :  $Mn^{4+}$  ratio on addition of graphene to MCN can be seen clearly from XPS measurements (Figure [1](#page-1-0)(b) and Table [S1\)](#page-6-0) where the ratio increases from 2.75 (pure MCN, 0% graphene) to 3.92 (5% graphene) and 5.23 (10% graphene). At ~15% graphene concentration, the Mn signal is no longer clearly visible (Figure [S3b](#page-6-0)), meaning the MCN is essentially completely covered with a multilayer of graphene at that graphene concentration. The data show that the MCN nanoparticle adsorption on the graphene includes both physical and chemical effect.

The (Hall) mobility of the graphene/MCN was measured and found to be  $91.4 \text{ cm}^2/\text{V} \cdot \text{s}$  for 3.5% graphene and 55.0 cm<sup>2</sup>/V·s for 6.5% graphene content (see Table [S1,](#page-6-0) average values). At low graphene concentrations (graphene below 7.5%), the composite shows p-type conductivity while between 7.5% and 12%, it exhibits n-type conductivity (for 10% graphene, the mobility is  $0.9 \text{ cm}^2/\text{V}$  with n-type conductivity). When the graphene content is increased beyond 12.5%, the graphene phase becomes continuous and the properties are like a conductor. The carrier type and concentration of the composite films can be modulated by the graphene MCN ratio and can vary over a wide range. Note that all tests are completed by pressing composite contact particles into thin films at 5-30MPa with different thicknesses. In addition to the graphene/MCN ratio, the specific conductivity data is also related to the formation pressure and thickness of the thin films. Although there will be fluctuation errors in the test data, the change trend of each sample is the same.

## 3. Discussion

A contact potential difference will occur between any two solid interfaces with different WFs and explains diode rectification. The width of the depletion region determines the thickness of the contact interface. Electron tunneling associated with the contact interface between conductor and semiconductor in general microelectronic devices only involves a distance < 2 nm. The carrier electrons in crystals will be scattered  $10^{12}$ - $10^{13}$  times per second, but because of the high electron mobility in graphene, electron scattering is strongly reduced. However, the electron space density of graphene is so large that diffusion is not the main effect [[18, 20](#page-6-0)–[24\]](#page-6-0) and the electrons behave like a two-dimensional electronic fluid [[25](#page-6-0)].

The Schottky contact barrier model is widely used to explain the interface carrier transport mechanism, which is generally based on thermionic emission. Although electron diffusion also occurs, the main contact mechanism between graphene and MCN is dominated by quantum tunneling because of the zero mass electrons of graphene (Figure 2) [[22](#page-6-0)].

The formula of quantum tunneling transmittance is

$$
T_{\text{RECTe}} \approx \frac{16E(V_0 - E)}{V_0^2} \exp\left(-\frac{2a}{\hbar} \sqrt{2m_e(V_0 - E)}\right), \quad (1)
$$

where the electronic mass =  $9.10956 \times 10^{-31}$  kg and the reduced Planck constant  $\hbar = 1.05457266 \times 10^{-34}$  j.s. From the work function difference between MCN and graphene, the potential well difference can be assumed to be  $V_0 - E \approx$ 0.3 eV approximately. Assuming a potential well  $a = 1$  nm, for general electronic calculation, we get

$$
T_{\text{RECTe}} \approx 4.8 \times 10^{-19} \times \exp\left(-4.45\right) \frac{E}{V_0^2}.
$$
 (2)

The electron velocity is 1/300 of the speed of light, so the

<span id="page-3-0"></span>

FIGURE 3: SEM images of cross-sections of films after different multilayer gradient stacking. (a) From left to right, the content of MCN nanometer powder decreases gradually, while the conductivity increases. The darker part of the film has a higher content of graphene and the light parts are the plastic films. (b) From left to right, the content of MCN powder increases gradually; then, on passing the broken green center line, it decreases gradually. (c) From left to right, the content of MCN nanometer powder decreases gradually; then, on passing the broken green line, it increases gradually.

effective dynamic mass,  $m = m_0 / \sqrt{1 - v^2/c^2}$ ,  $m_{\text{Ge}} = m - m_0$ ,  $m_{\text{Ge}}$ is approximately  $5.5 \times 10^{-6}$  *m*<sub>0</sub>.

When the potential well width  $a = 1$  nm,

$$
T_{\text{RECTe}} \approx 4.8 \times 10^{-19} \times \exp(-0.0104) \frac{E}{V_0^2}.
$$
 (3)

(3)/(2) is 85.6 times. If  $a = 3$  nm, 5 nm, and 10 nm, the velocity of graphene electrons entering MCN will drop very rapidly and the effective mass of electrons will increase, but  $T_{RECTe}$  increases by 310, 334, and 1.0 times at the different values of *a*. Although these assumptions include errors, it can be found that the tunneling transmittance at the contact interface between graphene electrons and MCN can be  $\sim 10^2$ times bigger than that of electrons from most other materials (SI tunneling effect calculation part).

However, the real size of MCN nanoparticles is about 10 nm, the actual action width scale of graphene electron tunneling action is above the MCN 10 nm diameter, and the electrons migrate into the whole volume structure of the MCN nanoparticles. For the various Mn valences (4+, 3+, and 2+) in the MCN structure, electrons fill holes resulting in an overall reduction in Mn oxidation state. As discussed above, our XPS data show that the  $Mn^{3+}/Mn^{4+}$  ratio increases from 2.8 to 5.2 or more as the graphene concentration is increased.

The MCN by itself is p-type as we have already noted. Upon addition of a small amount of graphene ("small" meaning that there is much less than complete coverage of MCN by graphene), the graphene will donate electrons to the high WF MCN (the high electron mobility of monolayer graphene is expected to favor electron transfer to the MCN, particularly where tunneling occurs). This results in two main effects: that of electrons donated to the MCN and holes

remaining in the graphene. The effect of electrons on the MCN is to reduce  $Mn^{4+}$  to  $Mn^{3+}$  (as shown in Figure [1\(](#page-1-0)b)/Table [S2\)](#page-6-0). This would be expected to make the MCN even more p-type. However, there are various Ni, Co, and O species that are also present, so it is difficult to predict what happens overall to the MCN, although increasing p-type would be expected in most cases. As the graphene coverage increases and a percolating threshold through graphene forms, conductivity should increase. The conductivity type of the composite in this situation can still be p-type due to the holes in the graphene. As the graphene concentration continues to increase, graphene layers not in direct contact with MCN will conduct by electrons and the conductivity changes to n-type. This is of course very simplified since graphene not in contact with MCN can donate electrons to underlying graphene layers that contained holes; however, the overall picture remains correct. Finally, when the graphene concentration is high enough, the composite will conduct essentially as if it were pure multilayer graphene.

The MCN/graphene composite has many unique characteristics [[3, 26\]](#page-6-0). By pressing composite powders of two different MCN: graphene ratios to form films and stacking the composite films, a flexible p-n junction could be formed due to their different majority carrier types. This was extended to p-n-p or n-p-n transistors by combining three films. Unlike traditional semiconductor devices, electrons are excited from the filled band into the empty band at the graphene Dirac points. These semiconductor devices are based on the two-dimensional electronic fluid characteristics of graphene, where the electrons are very mobile on the twodimensional surface of the graphene. Potential devices made of this composite thin film include diodes, triodes, and ultraviolet sensors.

<span id="page-4-0"></span>As an example, Figure [3\(](#page-3-0)a) shows the structure and Figure 4 shows an I-V plot of a p-n junction. If two of these p-n junction films are superimposed back-to-back, an n-p-n transistor configuration can be formed (Figure [3\(](#page-3-0)b)) and analogously for a p-n-p configuration (Figure  $3(c)$ ) ([SI,](#page-6-0) Figure [S5\)](#page-6-0). As a flexible material, the actual gradient films can be assembled into more complex flexible structures and devices of various shapes. This points to a new potential flexible substitute for silicon integrated circuit components.

Because the electron movement in the composite involves tunneling and the graphene's electron transition process is extremely rapid [[25](#page-6-0), [27](#page-6-0)–[32\]](#page-7-0), transistors manufactured by the composite material can be expected to operate at high frequencies. The carrier mobility of the p-type composite increases with temperature in the range of -50-150° C. This is determined by the properties of MCN itself, and the negative temperature coefficient is proportional to the MCN content. The electron mobility is basically temperature independent in this range because most carriers move on the surface of graphene and the mobility of graphene itself is little affected by temperature. It is found that due to the ultrahigh mobility characteristics of graphene, compared with pure MCN, the composite carrier mobility can be modulated and increased by  $10^2$ -10<sup>3</sup> times, and carrier concentration can be adjusted between  $10^{10}$  and  $10^{19}$ /cm<sup>3</sup>.

## 4. Conclusions

In conclusion, we have developed a method to effectively regulate p- and n-type semiconductivity of graphene-MCN nanocomposite, making it possible to easily fabricate graphene-based functional semiconductor devices. Unlike traditional semiconductor devices, transistor charge transport is dominated by highly mobile electron movement at the Dirac points on the 2-dimensional surface of graphene. It is anticipated that availability of both p- and n-type graphene-based materials will spur development to fabricate a new generation of semiconductor devices.

## 5. Materials and Methods

MCN nanometer powders (Figures [5\(](#page-5-0)a) and [5](#page-5-0)(d)) were synthesized by mixing  $Mn^{2+}$ ,  $Co^{2+}$ , and  $Ni^{2+}$  salts with aqueous ammonia and the precipitate fired at 900–1000 K in air for several hours. MCN inorganic semiconductor devices can work stably for a long time with high precision, sensitivity and reliability at high temperatures (up to 780° C), high voltage, and mechanical stress. The transition metal oxide  $Mn<sub>3</sub>$  $_{x}Co_{2-y}NiO_{8-x-y}$  (MCN,  $0 < x + y < 0.35$ ) has a spinel structure. The radii of  $Mn^{4+}$ ,  $Mn^{3+}$ , and  $Mn^{2+}$  ions are close to those of  $Co^{3+}$ ,  $Co^{2+}$ , and  $Ni^{3+}$  ions, respectively, which conforms to the Hume-Rothery rule. The MCN nanocrystal structure is compact and stable. MCN is usually a p-type semiconductor with negative temperature coefficient of resistivity in the range of -2%/K to -6%/K. At room temperature, the carrier concentration of these oxides is about 10<sup>9</sup>/cm<sup>3</sup>. It has very low carrier mobility, about 10<sup>-5</sup> cm<sup>2</sup>/V·s; its intrinsic conductivity is close to that of insulators [\[18](#page-6-0)–



Figure 4: I/V curve of a gradient multilayer thin film p-n junction.

[20\]](#page-6-0), and the negative temperature coefficient depends on the chemical composition of MCN (SI). Co and Ni are used as accepting or doping, and high-temperature sintering forms MCN solid solution with nonstoichiometric defects, which increases the oxygen hole density in the material. Due to Jahn-Teller distortion, some  $Mn^{3+}$  ions become  $Mn^{4+}$ . In the normal temperature equilibrium state, MCN nanocrystals actually have fixed ionic chemical compositions of  $Mn^{4+}/Mn^{3+}/Mn^{2+}$  at room temperature (25°C):  $Mn^{4+}$  and  $Mn^{3+}$  ions account for more than 80% and are the main components. The size of the individual MCN nanoparticles is about 10 nm (the maximum size of agglomerated particles does not exceed 60 nm, Figure [5](#page-5-0)(d)), and the specific surface area of the MCN nanoparticles is about  $10 \text{ m}^2/\text{g}$ . The MCN crystal structure takes  $Mn_4Mn_8O_{16}$  as the basic unit crystal cell [[33](#page-7-0), [34\]](#page-7-0). It can be calculated that particles with the 10 nm MCN account for more than 50% of the interfacial unit crystal cells.

Single-layer graphene (Figure [5](#page-5-0)(e)) was produced by a micromechanical stripping method (>99% is single-layer graphene, the apparent bulk density is below <0.001  $g/cm^3$ ) [\[21\]](#page-6-0). The specific surface area of mechanically exfoliated graphene is about 2630 m<sup>2</sup>/g. Single-layer graphene has an electron energy band structure containing Dirac cones with a zero bandgap. This band structure leads to an effective mass of the two-dimensional surface electrons in graphene that tends to zero, resulting in their very fast velocity of nearly 1/300 of the speed of light. One of its striking properties is the very high mobility ~ up to 150,000 cm<sup>2</sup>/V·s (for single-layer graphene at room temperature) and good electrical conductivity. Current research on semiconducting graphene focuses on heterojunction formation and elemental doping. The direct realization of functional devices by utilizing the Dirac cone characteristics is a topic which is of interest to many researchers [\[6](#page-6-0), [19, 27](#page-6-0)–[32](#page-7-0), [35](#page-7-0)–[42\]](#page-7-0).

Graphene and MCN semiconductor powders were weighed according to different mass ratios, (G: MCN between 1:5 and 1:200 mass ratios). The packing density of graphene,  $\sim 0.001$  g/ml, is about 1,000 times that of MCN

<span id="page-5-0"></span>

FIGURE 5: The Raman data, scanning electron microscopy (SEM), and transmission electron microscope images (JSM-7600F instrument and JEM-F200, JEOL Beijing Co., Ltd., Japan). (a) Raman spectrum of pure MCN. (b) Raman spectrum of 6.5% graphene+MCN composite film. (c) Raman spectrum of 10% graphene+MCN composite film. (d) SEM image of MCN nanometer powder. (e) TEM image of graphene. (f) TEM image of MCN nanoparticles coated with graphene (G : MCN mass ratio 1 : 20~25).

nanopowder (about 1.1 g/ml). MCN powder was added to graphene slowly, and the mixture was ground in a planetary mill for over 36 hours. The average thickness of graphene covering the MCN particles varied from 0.35 to 60 nm depending on the graphene concentration (Figure  $4(f)$ ). The composite powder was pressed onto polytetrafluoroethylene (Teflon) film to form a 3 mm square thin film. Raman spectroscopy of these films showed that when the graphene content was less than 8%, the main spectral characteristics of MCN were observed, but the intensity decreased significantly compared with pure MCN (Figures 5(a) and 5(b)). When the graphene content reached 10%, the spectral characteristics of multilayer graphene between MCN nanoparticles were displayed (Figure  $5(c)$ ). The composite film is protected by two layers of Teflon membrane, showing excellent bending flexibility (Figure [S6](#page-6-0)). Repeated bending is for more than 500 times, and the fluctuation of performance test data is within 3%.

## Data Availability

All data is available in the main text or the supplementary materials. All materials used in the analysis must be available in some form to any researcher for purposes of reproducing or extending the analysis. Please contact the corresponding author [hlshong@nwpu.edu.cn](mailto:hlshong@nwpu.edu.cn).

#### Conflicts of Interest

The authors declare that there is no conflict of interest regarding the publication of this article.

## Authors' Contributions

Profs. L. Su and Frank Liu conceived the idea and designed the experiments. Z. Zou, XT. Wang, YC. Liu, HX. Wang, Z. Ma, YF. Li, FX. Yang, YY. Suo, NH. Chang, and JX. Zhu conducted the experiments. Dr. L. Su, Z. Yang, SH. Wang, B. Wang, JP. Zhang, F. Gao, W. Huang, G. Hodes, and SZ. Liu contributed equally to the writing of the manuscript.

#### Acknowledgments

Participation and assistance in this research work also include the following: Mr. Junyou Li, Ake Zuoli, Mr. Li An, and Ms. Fangyuan Yan (School of Chemistry and Chemical Engineering, Northwestern Polytechnical University); Ms. Lan Zheng, Min Chen, Yanling Ai, Yanmin Zhang, Yanling Zhu, and Xuefang Shi (Material Analysis Center, Northwestern Polytechnical University); Prof. Mingzhen Wang (Dept. of Material and Chemical Engineering, Shaanxi Normal University); and Prof. Youlong Xu (Xi'an Jiaotong University). This work was supported by the

<span id="page-6-0"></span>Shaanxi Fundamental Education Research Funds for the Universities (2017e065201102) and the Special Fund for Support by 2019 Guangdong Special Funds (2019B090904007).

## Supplementary Materials

Figure S1: XRD pattern of MCN nanometer powder. Figure S2: schematic of the contacts for resistance and Hall measurements. Table S1: the electrical properties of composite films. Figure S3: XPS wide range spectra of (a) pure MCN film. (b) 15% graphene/MCN composite film. Little Mn, Co, or Ni is seen in this spectrum showing high coverage by graphene. Table S2:  $Mn^{3+}/Mn^{4+}$  contents and ratios for different films. Figure S4: the test curve of MCN negative temperature coefficient. Figure S5: (a, b) typical input and output characteristics of gradient film triode. (c, d) The output characteristic curve of an n-p-n gradient film composite triode (common emitter circuit). IC is the collector current, VOUT is the output voltage, and Ib is the base current. Figure S6: photographs of multilayer gradient film devices under different degrees of bending. [\(Supplementary](https://downloads.spj.sciencemag.org/research/2021/9802795.f1.zip) [Materials\)](https://downloads.spj.sciencemag.org/research/2021/9802795.f1.zip)

#### References

- [1] W. Schottky, "Zur Halbleitertheorie der Sperrschicht- und Spitzengleichrichter," Zeitschrift für Physik, vol. 113, no. 5-6, pp. 367–414, 1939.
- [2] N. F. Mott, "Note on the contact between a metal and an insulator or semi-conductor," Mathematical Proceedings of the Cambridge Philosophical Society, vol. 34, no. 4, pp. 568–572, 1938.
- [3] S. M. Sze and K. K. Ng, Physics of Semiconductor Devices, John wiley & sons, 3rd edition, 2006.
- [4] J. Bardeen, "Surface states and rectification at a metal semiconductor contact," Physical Review, vol. 71, no. 10, pp. 717– 727, 1947.
- [5] Z. S. Wu, G. Zhou, L. C. Yin, W. Ren, F. Li, and H. M. Cheng, "Graphene/metal oxide composite electrode materials for energy storage," Nano Energy, vol. 1, no. 1, pp. 107–131, 2012.
- [6] F. Schwierz, "Graphene transistors," Nature Nanotechnology, vol. 5, no. 7, pp. 487–496, 2010.
- [7] R. T. Tung, "Recent advances in Schottky barrier concepts," Materials Science and Engineering: R: Reports, vol. 35, no. 1-3, pp. 1–138, 2001.
- [8] M. H. Guimaraes, H. Gao, Y. Han et al., "Atomically thin ohmic edge contacts between two-dimensional materials," ACS nano, vol. 10, no. 6, pp. 6392–6399, 2016.
- [9] L. Lin, L. Liao, J. Yin, H. Peng, and Z. Liu, "Building graphene p-n junctions for next-generation photodetection," Nano Today, vol. 10, no. 6, pp. 701–716, 2015.
- [10] J. Gu, I. She, and Y. Yue, "Micro/nanoscale thermal characterization based on spectroscopy techniques," ES Energy & Environment, vol. 9, no. 83, pp. 15–27, 2020.
- [11] G. Liu, G. Yao, J. Xu, and X. Yan, "Spatial decoupling of light absorption and scattering centers in plasmon-assisted bubble column evaporator for solar steam generation," ES Energy  $\phi$ Environment, vol. 9, no. 76, pp. 41–49, 2020.
- [12] G. Liu, Y. Chen, S. Gao, B. Zhang, R. W. Li, and X. Zhuang, "Recent advances in resistive switching materials and devices:

from memories to memristors," Engineered Science, vol. 4, no. 4694, pp. 4–43, 2018.

- [13] S. D. Satpute, J. S. Jagtap, P. K. Bhujbal et al., "Mercurochrome sensitized  $ZnO/In_2O_3$  photoanode for dye-sensitized solar cell," ES Energy & Environment, vol. 9, no. 141, pp. 89–94, 2020.
- [14] M. A. Waghmare, N. I. Beedri, A. U. Ubale, and H. M. Pathan, "Fabrication and characterization of rose bengal sensitized binary  $TiO<sub>2</sub>-ZrO<sub>2</sub>$  oxides photoelectrode based dye-sensitized solar cell," Engineered Science, vol. 6, no. 1115, pp. 36–43, 2019.
- [15] C. Hou, W. Yang, X. Xie et al., "Agaric-like anodes of porous carbon decorated with  $MoO<sub>2</sub>$  nanoparticles for stable ultralong cycling lifespan and high-rate lithium/sodium storage," Journal of Colloid and Interface Science, vol. 596, pp. 396–407, 2021.
- [16] P. Xie, Y. Liu, M. Feng et al., "Hierarchically porous Co/C nanocomposites for ultralight high-performance microwave absorption," Advanced Composites and Hybrid Materials, vol. 4, no. 1, pp. 173–185, 2021.
- [17] S. Zhao and H. Wang, "An integrated H-type method to measure thermoelectric properties of two-dimensional materials," ES Energy & Environment, vol. 9, no. 36, pp. 59–66, 2020.
- [18] H. Y. Chiu, V. Perebeinos, Y. M. Lin, and P. Avouris,"Controllable p-n junction formation in monolayer graphene using electrostatic substrate engineering," Nano Letters, vol. 10, no. 11, pp. 4634–4639, 2010.
- [19] J. B. Kim, J. Li, Y. Choi, D. Whang, E. Hwang, and J. H. Cho, "Photosensitive graphene P-N junction transistors and ternary inverters," ACS applied materials & interfaces, vol. 10, no. 15, pp. 12897–12903, 2018.
- [20] H. J. Van Daal and A. J. Bosman, "Hall effect in CoO, NiO, and *α*-Fe<sub>2</sub>O<sub>3</sub>," *Physics Review*, vol. 158, no. 3, pp. 736–747, 1967.
- [21] K. S. Novoselov, Z. Jiang, Y. Zhang et al., "Room-temperature quantum Hall effect in graphene," Science, vol. 315, no. 5817, p. 1379, 2007.
- [22] L. Su, C. Wan, P. Yang et al., "Hybrid graphene oxide and NTC semiconductor material absorbs and transform light energy via a novel surface nanoscale plasmon mechanical," Plasmonics, vol. 11, no. 1, pp. 53–60, 2016.
- [23] K. S. Novoselov, A. K. Geim, S. V. Morozov et al., "Twodimensional gas of massless Dirac fermions in graphene," Nature, vol. 438, pp. 197–200, 2005.
- [24] L. A. Ponomarenko, R. V. Gorbachev, G. L. Yu et al., "Cloning of Dirac fermions in graphene superlattices," Nature, vol. 497, pp. 594–597, 2013.
- [25] H. Yang, J. Heo, S. Park et al., "Graphene barristor, a triode device with a gate-controlled Schottky barrier," Science, vol. 336, no. 6085, pp. 1140–1143, 2012.
- [26] A. I. Berdyugin, S. G. Xu, F. M. D. Pellegrino et al., "Measuring hall viscosity of graphene's electron fluid," Science, vol. 364, no. 6448, pp. 162–165, 2019.
- [27] L. H. Hess, M. V. Hauf, M. Seifert et al., "High-transconductance graphene solution-gated field effect transistors," Applied Physics Letters, vol. 99, no. 3, p. 033503, 2011.
- [28] G. Luongo, A. Grillo, F. Giubileo et al., "Graphene Schottky junction on pillar patterned silicon substrate," Nanomaterials, vol. 9, no. 5, p. 659, 2019.
- [29] W. Fu, L. Jiang, E. P. van Geest, L. M. Lima, and G. F. Schneider, "Sensing at the surface of graphene field-effect transistors," Advanced Materials, vol. 29, no. 6, p. 1603610, 2017.
- <span id="page-7-0"></span>[30] C. H. Lee, G. H. Lee, A. M. Van Der Zande et al., "Atomically thin P –N junctions with Van Der Waals heterointerfaces, " Nature Nanotechnology, vol. 9, no. 9, pp. 676 –681, 2014.
- [31] C. Li, M. T. Cole, W. Lei et al., "Highly electron transparent graphene for field emission triode gates," Advanced Functional Materials, vol. 24, no. 9, pp. 1218 –1227, 2013.
- [32] K. Xu, H. Lu, E. W. Kinder, A. Seabaugh, and S. K. Fullerton-Shirey, "Monolayer solid-state electrolyte for electric double layer gating of graphene field-effect transistors," ACS Nano, vol. 11, no. 6, pp. 5453 –5464, 2017.
- [33] L. Britnell, R. V. Gorbachev, R. Jalil et al., "Field-effect tunneling transistor based on vertical graphene heterostructures," Science, vol. 335, no. 6071, pp. 947 –950, 2012.
- [34] J. H. Chen, C. Jang, S. Xiao, M. Ishigami, and M. S. Fuhrer, "Intrinsic and extrinsic performance limits of graphene devices on SiO<sub>2</sub>," Nature Nanotechnology, vol. 3, no. 4, pp. 206 –209, 2008.
- [35] C. C. Chen, M. Aykol, C. C. Chang, A. F. Levi, and S. B. Cronin, "Graphene-silicon Schottky diodes," Nano Letters, vol. 11, no. 5, pp. 1863 –1867, 2011.
- [36] Y. Sun and J. A. Rogers, "Inorganic semiconductors for flexible electronics, " Advanced Materials, vol. 19, no. 15, pp. 1897 – 1916, 2007.
- [37] A. Allain, J. Kang, K. Banerjee, and A. Kis, "Electrical contacts to two-dimensional semiconductors," Nature materials, vol. 14, no. 12, pp. 1195 –1205, 2015.
- [38] P. Gallagher, C. S. Yang, T. Lyu et al., "Quantum-critical conductivity of the Dirac fluid in graphene," Science, vol. 364, no. 6436, pp. 158 –162, 2019.
- [39] A. Lucas, "An exotic quantum fluid in graphene," Science, vol. 364, no. 6436, p. 125, 2019.
- [40] R. Gy, G. D. C. Csete, A. N. Nolte, and I. M. Reaney, "Correlation between microstructure and conductance in NTC thermistors produced from oxide powders," Journal of the European Ceramic Society, vol. 19, no. 6, pp. 857 –860, 1999.
- [41] M. Vakiv, O. Shpotyuk, O. Mrooz, and I. Hadzaman, "Controlled thermistor effect in the system  $Cu_xNi_{1-x}$ ,  $Co_{2y}Mn_{2-y}O_4$ ," Journal of the European Ceramic Society, vol. 21, no. 10, pp. 1783 –1785, 2001.
- [42] S. K. Cheung and N. W. Cheung, "Extraction of Schottky diode parameters from forward current-voltage characteristics, " Applied Physics Letters, vol. 49, no. 2, pp. 85 –87, 1986.