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Chemically reactive MHD micropolar nanofluid flow with velocity slips and variable heat source/sink

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The two-dimensional electrically conducting magnetohydrodynamic flow of micropolar nanofluid over an extending surface with chemical reaction and secondary slips conditions is deliberated in this article. The flow of nanofluid is treated with heat source/sink and nonlinear thermal radiation impacts. The system of equations is solved analytically and numerically. Both analytical and numerical approaches are compared with the help of figures and tables. In order to improve the validity of the solutions and the method convergence, a descriptive demonstration of residual errors for various factors is presented. Also the convergence of an analytical approach is shown. The impacts of relevance parameters on velocity, micro-rotation, thermal, and concentration fields for first- and second-order velocity slips are accessible through figures. The velocity field heightens with the rise in micropolar, micro-rotation, and primary order velocity parameters, while other parameters have reducing impact on the velocity field. The micro-rotation field reduces with micro-rotation, secondary order velocity slip, and micropolar parameters but escalates with the primary order velocity slip parameter. The thermal field heightens with escalating non-uniform heat sink/source, Biot number, temperature ratio factor, and thermal radiation factor. The concentration field escalates with the increasing Biot number, while reduces with heightening chemical reaction and Schmidt number. The assessment of skin factor, thermal transfer, and mass transfer are calculated through tables.

List of symbols

c, d	Constants
κ	Vertex viscosity
B_0	Magnetic field strength
μ	Dynamic viscosity
ρ	Density
σ	Electrical conductivity
ρc_p	Heat capacitance
k	Thermal conductivity

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Mr	Micro-rotation parameter
$j = \nu/c$	Micro-inertia density
$\alpha = \kappa/\mu$	Material parameter
h_f	Heat transfer coefficient
h_s	Mass transfer coefficient
K_n	Knudsen number
λ	Molecular free path
M	Magnetic factor
M_r	Microrotation parameter
C_r	Chemical reaction parameter
γ, δ	Velocity slip factors
C_f	Skin friction
Re_x	Reynolds number
$\tilde{h}_f, \tilde{h}_g, \tilde{h}_\theta, \tilde{h}_\phi$	Auxiliary factors
Sh	Sherwood number
x, y	Coordinates
u, v	Velocity components
D	Diffusion coefficient
K_1	Reaction rate
N	Micro-rotation velocity
T	Temperature
T_s	Temperature at the surface
T_∞	Ambient temperature
C	Concentration
C_s	Surface concentration
C_∞	Ambient concentration
A^*, B^*	Non-uniform heat source and sink parameters
P, Q	Constants
a	Momentum coefficient
α	Micropolar factor
Pr	Prandtl number
Rd	Non-linear thermal energy factor,
Bi_1, Bi_2	Biot numbers
λ	Stretching factor
Nu	Nusselt number
$a_i (i = 1 - 9)$	Constants in general solution
Res	Residual error
C_s	Couple stress

The uses of nanotechnology in the field of engineering and technologies are electric circuits, solar cells, food processing, chemical sensors, batteries, fuels, ethanol, etc. Such applications and advantages have endorsed the researchers' interest in the field of nanotechnology. Nowadays, the investigators are analyzing a new subclass of nanotechnology called nanofluids. Initially, the nanofluid was determined by Choi¹. Rashidi et al.² explored the heat transfer analysis of nanofluid. Bahiraei and Hangi³ investigated the transfer of heat by nanofluids with a magnetic influence. Ghasemian et al.⁴ tested the nanofluid heat transfer with alternating and constant magnetic field. Ellahi et al.⁵ intentioned the nanofluid stream of the boundary layer. Xuan et al.⁶ assessed heat propagation in the nanofluid flow. Sheikholeslami et al.⁷ probed the Coulomb force influence on the heat transfer of a nanofluid. Alsabery et al.⁸ presented the nanofluid flow in an inclined enclosure with a porous medium. Sheikholeslami⁹ probed the flow of nanofluid in porous media. Sheikholeslami¹⁰ analyzed the nanofluid flow with a magnetic field effect. Hassan et al.¹¹ observed the flow of nanofluid in a porous medium with a convective heat exchange. Dawar et al.¹² studied the flow of nanofluid with thermal energy and energy source. Khan et al.¹³ tested the convective boundary layer nanofluid flow through mass and heat distribution. Sheikholeslami et al.¹⁴ monitored the thermal conduction to a magnetic nanofluid. Shah et al.¹⁵ tested the nanofluid thin film flow with a nonlinear thermal radiation. Shah et al.¹⁶ looked at the coupled stress nanofluid flow with the Cattaneo heat design. Sohail et al.¹⁷ presented the ferrofluid videography valuation in drug targeting. Dawar et al.¹⁸ inspected the nanofluid flow with thermophoretic and Brownian motion influences. Majeed et al.¹⁹ probed the nanofluid heat transmission with heat source and thermal energy effects. Dawar et al.²⁰ investigated MHD nanofluid with dissipation impact. Sheikholeslami et al.²¹ numerically checked the flow of ferrofluid using porous media. Sajjad et al.²² offered the flow of nanofluid with thermophoretic and Brownian motion influences using Darcy-Forchheimer relation. Dawar et al.²³ probed the thin film nanofluid flow with thermal energy using the Darcy-Forchheimer relation. Alamri et al.²⁴ monitored the flow of nanofluid in a porous medium with slip conditions. Ahmad et al.²⁵ offered the couple stress nanofluid flow with the Cattaneo heat model using the Darcy-Forchheimer relation.

The situation of stretching performs a significance role in boundary layer flow examination due to its remarkable results in the area of engineering and industries for instance, paper production, polymer engineering, metallic beds cooling, wire drawing, hot rolling, plastic sheets extraction, glass formation, etc. Preceding the stretching phenomenon, Crane²⁶ was the discoverer who introduced the fluid flow over the extending plate. Hayat et al.²⁷ examined the micropolar fluid flow. Najib et al.²⁸ investigated the stagnation point flow with chemical reaction. Babu et al.²⁹ scrutinized the micropolar fluid flow's stagnation point with suction impact. Soid et al.³⁰ presented

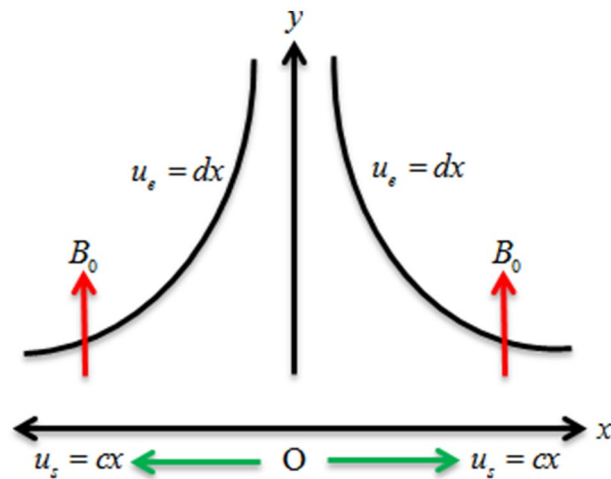


Figure 1. Geometry of the problem.

the heat transmission of a fluid flow. Makinde et al.³¹ probed the boundary layer nanofluid flow with a magnetic field. Mabood et al.³² scrutinized the nanofluid flow with melting absorption/generation impacts.

Rarefaction influences should always be assumed in order to accurately design a microsystem. Rarefaction impacts become huge as the distinctive length of a system moves toward the molecular mean free pathway of the liquid inside the framework. The Knudsen number, K_n , characterized as the proportion of the molecular mean free way to the trademark length of the framework, is regularly used to exactly measure the impacts of rarefaction³³. For continuum flow, $K_n \leq 0.01$ rarefaction impacts are immaterial and traditional preservation conditions are utilized. $0.1 \leq K_n \leq 10.0$, extremely infrequent flows are part of the transitional system. $K_n \geq 10.0$, the open molecular system. The presumption that because a liquid rigidly adhere to a solid boundary and this called no-slip boundary condition was shown to be ineffective in a variety of cases, like: the high molecular liquid flows, micro-channel flows or dynamics of thin fluids. Slip conditions are very important for the nanofluids studies. Eggs yolk, oil, and liquid combination, grease, polymer solutions, etc. are the application containing slip conditions influences. A substantial majority of models were discussed to characterize the slip that actually happens at a solid boundary. In 1827, Navier³⁴ presented the fluid motion through a sheet with a slip condition. Fang et al.³⁵ investigated the flow of viscous fluid in a second-order slip state. Beg et al.³⁶ tested the magnetic flow in slippery conditions. Martin and Boyd³⁷ tested the convection in a boundary layer flow with slip condition. Ibrahim et al.³⁸ proposed a thermal transfer of a magnetohydrodynamic micropolar fluid with a second-order slip state. Mabood et al.³⁹ explored the heat transmission of stagnation point flow with second-order slip conditions. Other relevant analyses are mentioned in^{40–43}. Lund et al.⁴⁴ presented the MHD micropolar fluid flow with joule heating and viscous dissipation influences over exponentially shrinking sheet. Yasmin et al.⁴⁵ investigated the thermal and mass transmission in MHD micropolar fluid flow over a stretching surface. Kumar et al.⁴⁶ investigated the MHD micropolar fluid flow over a stretching sheet with heat flux model. The MHD stagnation point flow of micropolar fluid over a convective surface with nonlinear radiation influence was analyzed by Kumar et al.⁴⁷. Kumar et al.⁴³ analyzed the MHD first and second orders slips flow of micropolar fluid over a convective surface. Kumar et al.^{48,49} investigated the micropolar fluid flow with heat flux model under the influence of magnetic field, thermophoresis and Brownian motion using coagulated and stretching surfaces. Further studies of Kumar et al. can be seen in^{50–53}.

The key explanation for the considerable attention paid to the analysis of micropolar fluid flows is the uses and applications in manufacturing processes, including: animal blood, liquid crystal solidification, bath metal plate cooling, suspension and colloidal solutions, polymer fluid extrusion, and exotic lubricants. In order to present the current work in the field of micropolar fluids, we present the flow of micropolar nanofluid over an extending sheet in the presence of first and second orders velocity slip conditions with chemical reaction. The system of equations is solved analytically and numerically. The effects of developed factors on the nanofluid flow are presented through graphs and deliberated their features.

Problem formulation

We considered the electrically accompanying magnetohydrodynamic flow of micropolar fluid over an extending sheet with chemical reaction and secondary slips conditions. The nanofluid flow is treated with heat source/sink and nonlinear thermal radiation. The nanofluid flow is considered in 2D coordinates system. The x -axis is considered along the nanofluid flow and y -axis is considered vertical to the nanofluid flow. The strength of the magnetic field is taken vertically to the nanofluid flow. The velocities are $u_s = cx$ and $u_e = dx$ where $c > 0$ and $d > 0$ are constants as expressed in Fig. 1.

The proposed model leads to the following equations^{54–56}:

$$\frac{\partial v}{\partial y} + \frac{\partial u}{\partial x} = 0, \tag{1}$$

$$\rho \left(v \frac{\partial u}{\partial y} + u \frac{\partial u}{\partial x} \right) = \kappa \left(\frac{\partial^2 u}{\partial y^2} + \frac{\partial N}{\partial y} \right) + \mu \frac{\partial^2 u}{\partial y^2} - \sigma B_0^2 u, \tag{2}$$

$$\rho j \left(v \frac{\partial N}{\partial y} + u \frac{\partial N}{\partial x} \right) = -\kappa \left(2N + \frac{\partial u}{\partial y} \right) + \Gamma \frac{\partial^2 N}{\partial y^2}, \tag{3}$$

$$\rho c_p \left(v \frac{\partial T}{\partial y} + u \frac{\partial T}{\partial x} \right) = k \frac{\partial^2 T}{\partial y^2} - \frac{\partial q_f}{\partial y} + q'''', \tag{4}$$

$$u \frac{\partial C}{\partial x} + v \frac{\partial C}{\partial y} = D \frac{\partial^2 C}{\partial y^2} - K_1 (C - C_\infty), \tag{5}$$

where Γ , q_f and q'''' are defined as

$$\begin{aligned} \Gamma &= \left(\mu + \frac{\kappa}{2} \right) j = \left(1 + \frac{\alpha}{2} \right) \mu j, \quad q_f = -\frac{4\sigma^*}{3k^*} \frac{\partial T^4}{\partial y} = -\frac{16\sigma^*}{3k^*} T^3 \frac{\partial T}{\partial y}, \\ q'''' &= \frac{k(T_s - T_\infty)u_s}{xv} \left(A^* f' + B^* \frac{(T - T_\infty)}{(T_s - T_\infty)} \right), \end{aligned} \tag{6}$$

Equation (4) is reduced as:

$$\rho c_p \left(v \frac{\partial T}{\partial y} + u \frac{\partial T}{\partial x} \right) = \frac{16\sigma^*}{3k^*} \frac{\partial}{\partial y} \left(T^3 \frac{\partial T}{\partial y} \right) + k \frac{\partial^2 T}{\partial y^2} + \frac{k(T_s - T_\infty)u_s}{xv} \left(A^* f' + B^* \frac{(T - T_\infty)}{(T_s - T_\infty)} \right), \tag{7}$$

with boundary conditions⁵⁶:

$$\begin{aligned} u = u_s + u_{\text{slip}}, \quad v = 0, \quad N = -Mr \frac{\partial u}{\partial y}, \quad -\frac{\partial T}{\partial y} = \frac{h_f}{k} (T_s - T), \quad -\frac{\partial C}{\partial y} = \frac{h_s}{D} (C_s - C) \quad \text{at } y = 0, \\ u \rightarrow u_e, \quad N \rightarrow 0, \quad T \rightarrow T_\infty, \quad C \rightarrow C_\infty \quad \text{as } y \rightarrow \infty. \end{aligned} \tag{8}$$

The secondary velocity slip model is considered as⁵⁶:

$$u_{\text{slip}} = \frac{2}{3} \lambda \left(\frac{3 - a l^2}{a} - \frac{3}{2} \left(\frac{1 - l^2}{K_n} \right) \right) \frac{\partial u}{\partial y} - \frac{1}{4} \lambda^2 \left(l^4 + \frac{2(1 - l^2)}{K_n^2} \right) \frac{\partial^2 u}{\partial y^2} = P \frac{\partial u}{\partial y} + Q \frac{\partial^2 u}{\partial y^2}, \tag{9}$$

Following the overhead equations, u , v are the velocity constituents along x - and y -coordinates, respectively, κ represents the vertex viscosity, B_0 is the magnetic field strength, μ indicates the dynamic viscosity, ρ designates the density, σ represents the electrical conductivity, ρc_p denotes the heat capacitance, k indicates the thermal conductivity, D specifies the diffusion coefficient, K_1 represents the reaction rate, N represents the micro-rotation velocity, Mr represents the micro-rotation parameter, $j = \nu/c$ represents the micro-inertia density, $\alpha = \kappa/\mu$ indicates the material parameter, T , T_s and T_∞ represent the temperature, temperature at the surface and temperature away from the surface, correspondingly, C , C_s and C_∞ represent the concentration, concentration at the surface and concentration away from the surface, correspondingly, h_f and h_s indicate the convective heat and mass transfer coefficients respectively, A^* and B^* are the non-uniform heat source and sink parameters respectively, P and Q are constants, K_n indicates the Knudsen number, $a(0 \leq a \leq 1)$ signifies the momentum coefficient, and λ indicates the molecular free path, $l = \min\left(\frac{1}{K_n}, 1\right)$ for all K_n .

The similarity transformations are defined as⁵⁴⁻⁵⁶:

$$\begin{aligned} \chi &= (cv)^{1/2} x f, \quad u = \frac{\partial \chi}{\partial y}, \quad v = -\frac{\partial \chi}{\partial x}, \quad N = cx \left(\frac{c}{v} \right)^{1/2} g, \\ T &= T_\infty (1 + (\theta_w - 1)\theta), \quad \theta_w = \frac{T_s}{T_\infty}, \quad \phi(\xi) = \frac{C - C_\infty}{C_s - C_\infty}, \quad \xi = \left(\frac{c}{v} \right)^{1/2} y, \end{aligned} \tag{10}$$

Equations (2), (3), (5) and (7) with boundary conditions (8) are reduced as:

$$(1 + \alpha) f'''' + f f'' - f'^2 + \alpha g' - M f' = 0, \tag{11}$$

$$\left(1 + \frac{\alpha}{2} \right) g'' + f g' - g f' - 2\alpha g - \alpha f'' = 0, \tag{12}$$

$$\theta'' + \text{Pr}f\theta' + A^*f' + B^*\theta + Rd \left\{ \begin{array}{l} \theta'' + \theta^3\theta''(\theta_w - 1)^3 + 3\theta^2\theta'^2(\theta_w - 1)^3 \\ + 3\theta^2\theta''(\theta_w - 1)^2 + 6\theta\theta'^2(\theta_w - 1)^2 \\ + 3\theta\theta''(\theta_w - 1) + 3\theta'^2(\theta_w - 1) \end{array} \right\} = 0, \tag{13}$$

$$\phi'' + \text{Sc}f\phi' - \text{Sc}C_r\phi = 0, \tag{14}$$

Subject to:

$$\begin{aligned} f = 0, f' = 1 + \delta f''' + \gamma f'', \quad g = -f'' M_r, \quad \theta' = -Bi_1(1 - \theta), \quad \phi' = -Bi_2(1 - \phi) \text{ at } \xi = 0, \\ f' \rightarrow \lambda, \quad g \rightarrow 0, \quad \theta \rightarrow 0, \quad \phi \rightarrow 0 \text{ as } \xi \rightarrow \infty \end{aligned} \tag{15}$$

Here M indicates the magnetic factor, α indicates the micropolar factor, Pr represents the Prandtl number, Rd signifies the non-linear thermal energy factor, M_r indicates the microrotation parameter, C_r is chemical reaction parameter, Bi_1 and Bi_2 represent the Biot numbers, γ and δ designate the velocity slip factors, and λ indicates the stretching factor which are defined as

$$\begin{aligned} M = \sigma B_0^2 / \rho c, \quad \text{Pr} = \frac{\mu c_p}{k}, \quad Rd = \frac{16\sigma^* T_\infty^3}{3kk^*}, \quad \text{Sc} = \frac{\nu}{D}, \quad Bi_1 = \frac{h_f}{k} \left(\frac{\nu}{c}\right)^{1/2}, \\ C_r = \frac{K_1}{c}, \quad Bi_2 = \frac{h_s}{D} \left(\frac{\nu}{c}\right)^{1/2}, \quad \gamma = P \left(\frac{c}{\nu}\right)^{1/2} (> 0), \quad \delta = Q \frac{c}{\nu} (< 0), \quad \lambda = \frac{d}{c}. \end{aligned} \tag{16}$$

The dimensionless form of the skin friction, couple stress, and Nusselt and Sherwood numbers are defined as:

$$\begin{aligned} \sqrt{\text{Re}_x} C_f = 2(\alpha(1 - Mr) + 1)f''(0), \quad C_s = \left(\frac{\alpha}{2} + 1\right)g'(0), \\ \frac{1}{\sqrt{\text{Re}_x}} Nu = -(1 + Rd(\theta_w)^3)\theta'(0), \quad \frac{1}{\sqrt{\text{Re}_x}} Sh = -\phi'(0), \end{aligned} \tag{17}$$

In which $\text{Re}_x = \frac{c x^2}{\nu}$ is Reynolds number.

Solution by HAM

In understanding of (11–14) with (15), HAM is used with the following procedure.

Preliminary assumptions:

$$f_0(\xi) = \frac{1}{(1 + \gamma - \delta)}(1 - e^{-\xi}), \quad g_0(\xi) = -\frac{M_r}{(1 + \gamma - \delta)}e^{-\xi}, \quad \theta_0(\xi) = \frac{Bi_1}{(1 + Bi_1)}e^{-\xi}, \quad \phi_0(\xi) = \frac{Bi_2}{(1 + Bi_2)}e^{-\xi}. \tag{18}$$

Linear operators:

$$L_f(f) = \frac{d^3 f}{d\xi^3} - \frac{df}{d\xi}, \quad L_g(g) = \frac{d^2 g}{d\xi^2} - g, \quad L_\theta(\theta) = \frac{d^2 \theta}{d\xi^2} - \theta, \quad L_\phi(\phi) = \frac{d^2 \phi}{d\xi^2} - \phi, \tag{19}$$

with

$$L_f(a_1 + a_2 e^{-\xi} + a_3 e^\xi) = 0, \quad L_g(a_4 e^{-\xi} + a_5 e^\xi) = 0, \quad L_\theta(a_6 e^{-\xi} + a_7 e^\xi) = 0, \quad L_\phi(a_8 e^{-\xi} + a_9 e^\xi) = 0. \tag{20}$$

where $a_i (i = 1 - 9)$ are constants in general solution.

Convergence analysis by HAM

It is very well understood that the definition of homotopy ensures excellent versatility in interpreting the auxiliary factors ($\tilde{h}_f, \tilde{h}_g, \tilde{h}_\theta, \tilde{h}_\phi$) for regulating and modifying the series solutions convergence. In Figs. 2, 3, 4, \tilde{h} -curves are displayed to interpret the appropriate values of $\tilde{h}_f, \tilde{h}_g, \tilde{h}_\theta$ and \tilde{h}_ϕ . The acceptable ranges for the modeled problem are $-0.54 \leq \tilde{h}_f \leq -0.08, -0.58 \leq \tilde{h}_g \leq -0.04, -0.8 \leq \tilde{h}_\theta \leq 0.0$ and $-0.75 \leq \tilde{h}_\phi \leq 0.0$.

Results and discussion

The impact of relevance parameters on velocity, micro-rotation, thermal, and concentration fields for first order slip γ and second order slip δ parameters are offered in Figs. 5, 6, 7, 8, 9, 10, 11, 12, 13, 14, 15, 16, 17, 18 and 19. Figures 5 and 6 represent the effect of micropolar parameter on velocity and micro-rotation fields when $Rd = 0.3, \theta_w = 0.1, Bi_1 = Bi_2 = 0.5, \text{Pr} = 7.0, \gamma = 0.5, \delta = 1.0, M = 1.0, C_r = 0.5, \text{Sc} = 1.0, \lambda = 0.2, A^* = B^* = 0.2$ and $M_r = 0.5$. The heightening in micropolar factor increases the velocity field however reduces the micro-rotation field. This findings indicate that the momentum exchange layer-by-layer is enhanced positively by the increase in viscosity induced by the mutual micro-rotation of the molecules, i.e. the significant estimates of α ; across the other side, the thermal diffusions are marginally weakened. Figures 7 and 8 present the influence of micro-rotation factor on velocity and micro-rotation fields when $Rd = 0.3, \alpha = 2.0, \theta_w = 0.1, Bi_1 = Bi_2 = 0.5, \text{Pr} = 7.0, \gamma = 0.5, \delta = 1.0, M = 1.0, C_r = 0.5, \text{Sc} = 1.0, \lambda = 0.2$ and $A^* = B^* = 0.2$. The micro-rotation parameter has

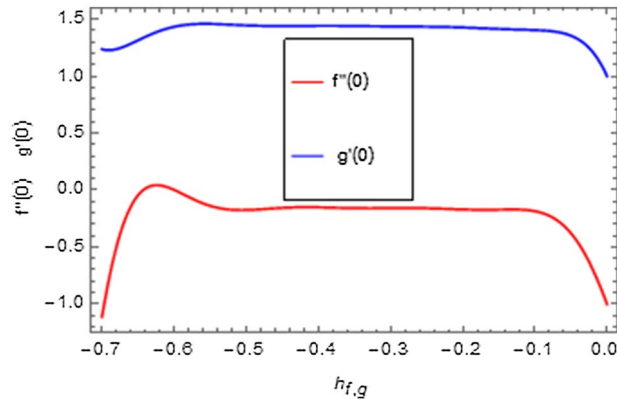


Figure 2. \bar{h} -curves for $f'(\xi)$ and $g(\xi)$.

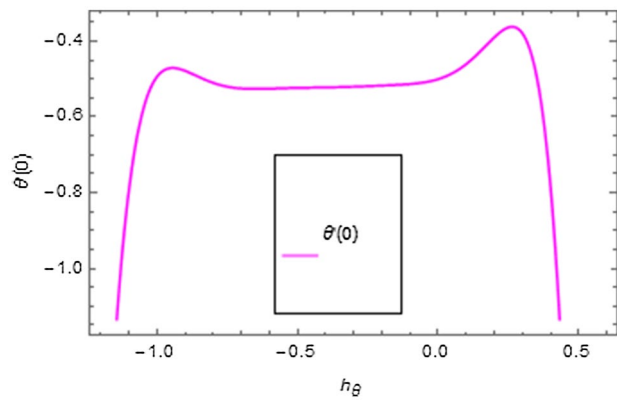


Figure 3. \bar{h} -curve for $\theta(\xi)$.

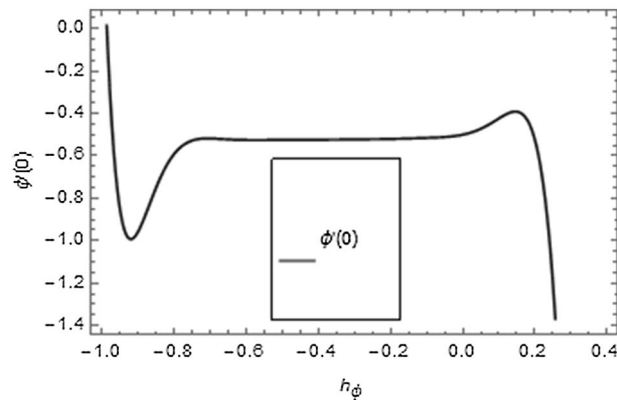


Figure 4. \bar{h} -curve for $\phi(\xi)$.

direct relation with velocity field, while has reverse relation with micro-rotation field. The increasing micro-rotation parameter heightens the velocity field and has reducing impact on micro-rotation field. In addition, it is also determined that there is no micro-rotation velocity when $Mr = 0$. Figure 9 indicates the relation between magnetic parameter and velocity field when $Rd = 0.3$, $\alpha = 2.0$, $\theta_w = 0.1$, $Bi_1 = Bi_2 = 0.5$, $Pr = 7.0$, $C_r = 0.5$, $Sc = 1.0$, $\gamma = 0.5$, $\delta = 1.0$, $\lambda = 0.2$, $A^* = B^* = 0.2$ and $Mr = 0.5$. It contests the physical understanding of the contraction of the magnetic force to the electrically conductive fluid, yet this result in the increase of a drag force which resulted in the slowing down force on the velocity. Figure 10 designates the variation in velocity

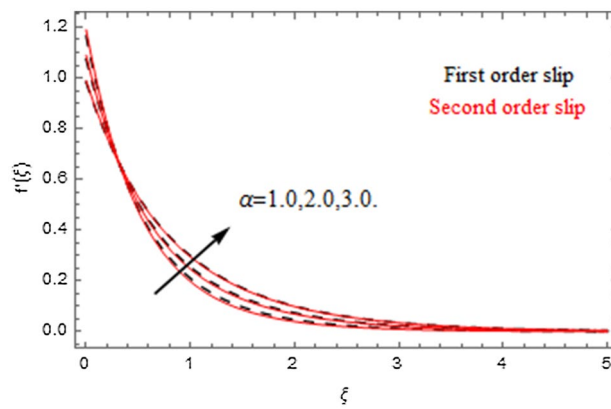


Figure 5. α versus $f'(\xi)$.

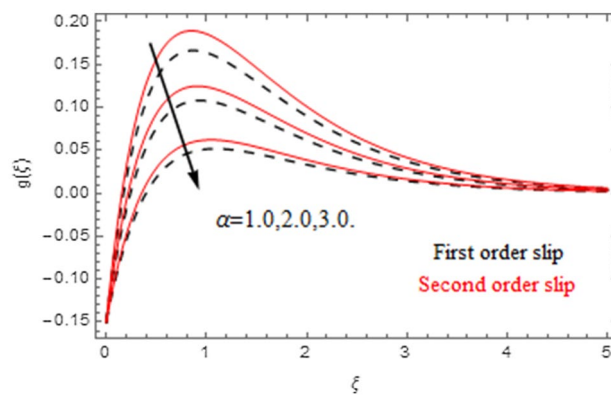


Figure 6. α versus $g(\xi)$.

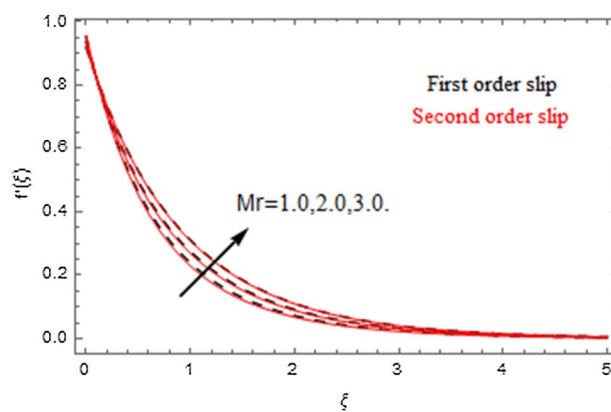


Figure 7. Mr versus $f'(\xi)$.

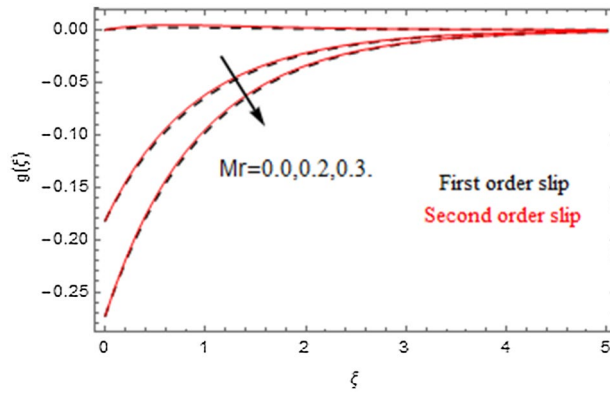


Figure 8. Mr versus $g(\xi)$.

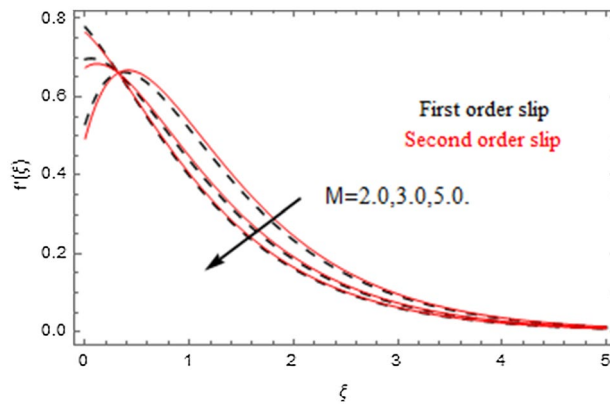


Figure 9. M versus $f'(\xi)$.

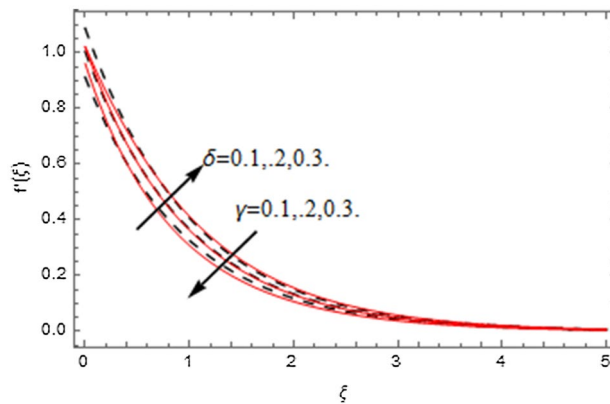


Figure 10. δ and γ versus $f'(\xi)$.

field against γ and δ factors when $Rd = 0.3$, $\alpha = 2.0$, $C_r = 0.5$, $Sc = 1.0$, $\theta_w = 0.1$, $Bi_1 = Bi_2 = 0.5$, $Pr = 7.0$, $M = 1.0$, $\lambda = 0.2$, $A^* = B^* = 0.2$ and $Mr = 0.5$. γ and velocity field has inverse relation to each other. That is, the heightening in γ the reduction in velocity field is detected. δ and velocity field has direct relation to each other. That is, the heightening in δ the rise in velocity field is detected. Figure 11 indicates the similar impact of γ and δ against micro-rotation field when $Rd = 0.3$, $\alpha = 2.0$, $\theta_w = 0.1$, $Bi_1 = Bi_2 = 0.5$, $Pr = 7.0$, $M = 1.0$,

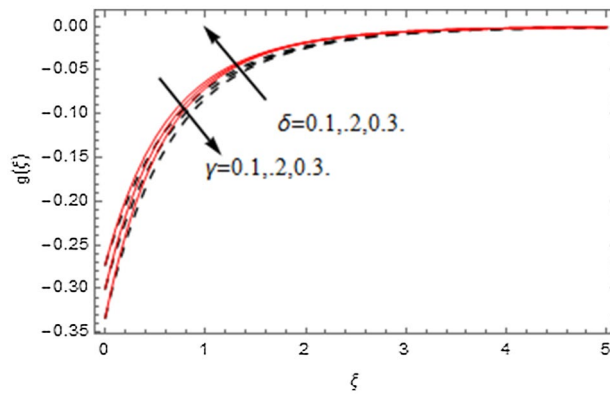


Figure 11. δ and γ versus $g(\xi)$.

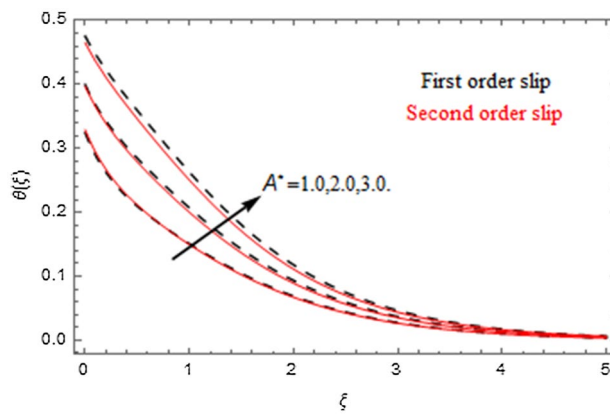


Figure 12. A^* versus $\theta(\xi)$.

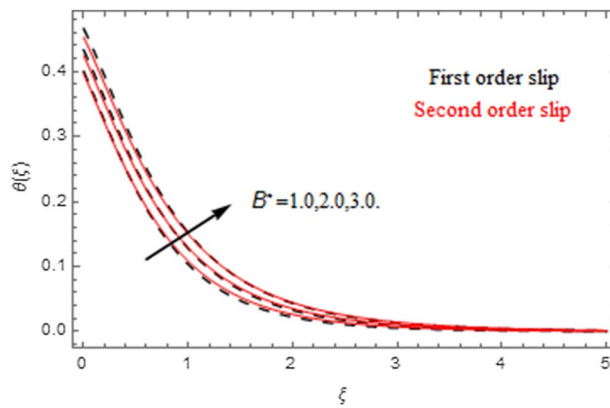


Figure 13. B^* versus $\theta(\xi)$.

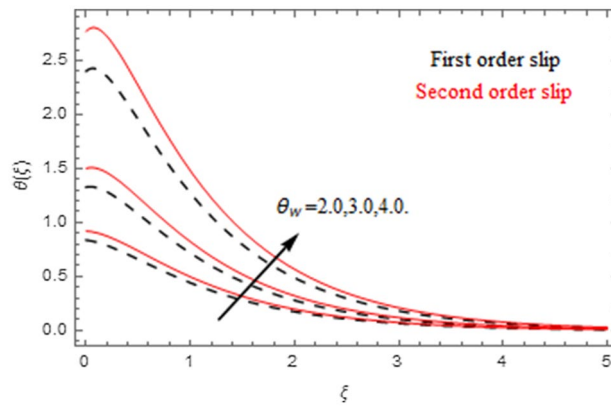


Figure 14. θ_w versus $\theta(\xi)$.

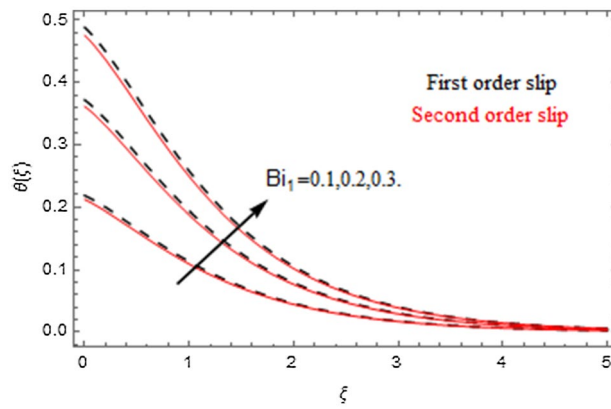


Figure 15. Bi_1 versus $\theta(\xi)$.

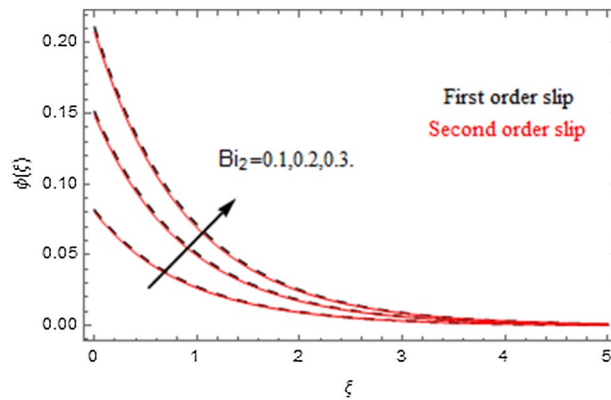


Figure 16. Bi_2 versus $\phi(\xi)$.

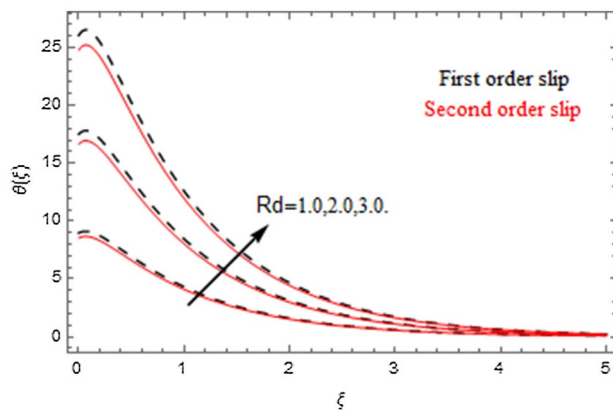


Figure 17. Rd versus $\theta(\xi)$.

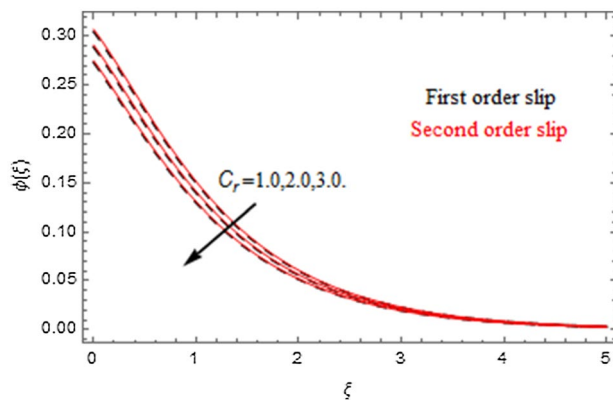


Figure 18. C_r versus $\phi(\xi)$.

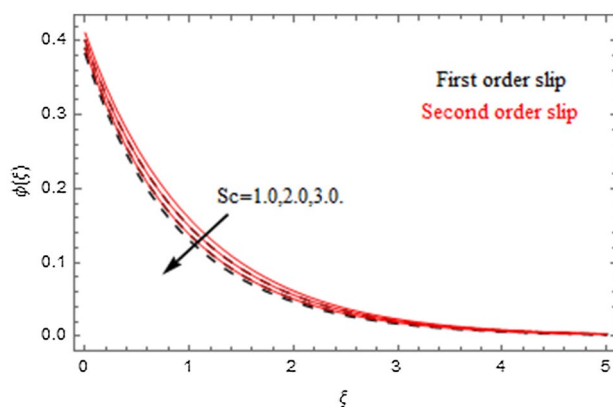


Figure 19. Sc versus $\phi(\xi)$.

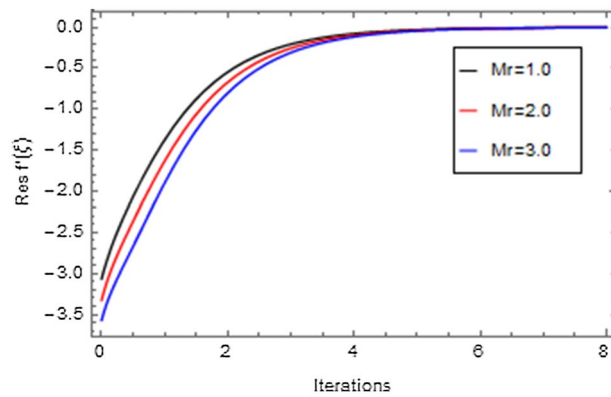


Figure 20. Residual error for Mr on $f'(\xi)$.

$\lambda = 0.2$, $C_r = 0.5$, $Sc = 1.0$, $A^* = B^* = 0.2$ and $M_r = 0.5$. Figures 12 and 13 display the impact of A^* and B^* on thermal field when $Rd = 0.3$, $\alpha = 2.0$, $\theta_w = 0.1$, $C_r = 0.5$, $Sc = 1.0$, $Bi_1 = Bi_2 = 0.5$, $Pr = 7.0$, $\gamma = 0.5$, $\delta = 1.0$, $M = 1.0$, $\lambda = 0.2$ and $M_r = 0.5$. It is detected that the rising values A^* and B^* heighten the thermal field of the nanofluid flow. Actually, A^* and B^* act like a heat originator. The heightening estimations of A^* and B^* escalate the temperature of the fluid flow. Therefore, the increasing values of A^* and B^* intensifies the thermal field. Figure 14 depicts the impact of θ_w on thermal field when $Rd = 0.3$, $\alpha = 2.0$, $Bi_1 = Bi_2 = 0.5$, $Pr = 7.0$, $C_r = 0.5$, $Sc = 1.0$, $\gamma = 0.5$, $\delta = 1.0$, $M = 1.0$, $\lambda = 0.2$, $A^* = B^* = 0.2$ and $M_r = 0.5$. With nonlinear radiation, the value of the temperature ratio parameter must be greater than 1. Also, a boost throughout the temperature ratio parameter allows the temperature to ruse through the layer. It is important to mention that as the temperature ratio tends to 1, the linear and non-linear Rosseland estimations have the same impact on the fluid flow temperature. The increasing influence of θ_w on temperature field is detected here. Figure 15 displays the impact of Biot number Bi_1 on thermal field when $Rd = 0.3$, $\alpha = 2.0$, $\theta_w = 0.1$, $Bi_2 = 0.5$, $A^* = B^* = 0.2$, $C_r = 0.5$, $Sc = 1.0$, $Pr = 7.0$, $\gamma = 0.5$, $\delta = 1.0$, $M = 1.0$, $\lambda = 0.2$ and $M_r = 0.5$. Biot number takes place in the special counsel investigation due to the implication of the convective boundary condition and tends to mean the proportion of both the diffuse opposition throughout the sheet to the convective opposition at the layer of the piece of paper. The lower estimations of the Biot number, therefore, provide an elevated convective opposition at the layer, and one that tends to lead to a medium heat transfer rate from layer to that of the liquid. So, the higher estimations of Biot number Bi_1 raise the temperature field of the fluid flow. A similar impact of Biot number Bi_2 against the concentration field when $Rd = 0.3$, $\alpha = 2.0$, $C_r = 0.5$, $Sc = 1.0$, $\theta_w = 0.1$, $Bi_1 = 0.5$, $A^* = B^* = 0.2$, $Pr = 7.0$, $\gamma = 0.5$, $\delta = 1.0$, $M = 1.0$, $\lambda = 0.2$ and $M_r = 0.5$ is shown in Fig. 16. Figure 17 signifies the non-linear thermal radiation Rd impact on temperature field when $Rd = 0.3$, $\alpha = 2.0$, $\theta_w = 0.1$, $Bi_1 = Bi_2 = 0.5$, $A^* = B^* = 0.2$, $Pr = 7.0$, $C_r = 0.5$, $Sc = 1.0$, $\gamma = 0.5$, $\delta = 1.0$, $M = 1.0$, $\lambda = 0.2$ and $M_r = 0.5$. The rise in Rd escalates the temperature field. It is well known that the function of radiation and thermal expansion is indeed a phenomenon that creates heat through fluid particles in such a manner that certain extra heat is created all through the flow. Figure 18 denotes the influence of C_r on concentration field when $Rd = 0.3$, $\alpha = 2.0$, $\theta_w = 0.1$, $Bi_1 = Bi_2 = 0.5$, $Sc = 1.0$, $A^* = B^* = 0.2$, $Pr = 7.0$, $\gamma = 0.5$, $\delta = 1.0$, $M = 1.0$, $\lambda = 0.2$ and $M_r = 0.5$. Higher values of the chemical reaction parameter lead to a greater rate of pernicious chemical change that breaks down or halts the liquid species more efficiently and effectively. Thus a decreasing influence in concentration profile is depicted. Figure 19 signifies the association of Sc with concentration field when $Rd = 0.3$, $\alpha = 2.0$, $\theta_w = 0.1$, $Bi_1 = Bi_2 = 0.5$, $A^* = B^* = 0.2$, $Pr = 7.0$, $\gamma = 0.5$, $\delta = 1.0$, $C_r = 0.5$, $M = 1.0$, $\lambda = 0.2$ and $M_r = 0.5$. The Schmidt number is inversely related with concentration field. An increase in Sc deescalates the concentration field. It is evidently understood that concentration, in addition to its related boundary layer thickness, are diminishing functions of Sc . The improvement in Sc , thus leads to a lower coefficient of diffusion. Such a lower coefficient of diffusion results a significant decrease throughout the concentration field.

Figures 20, 21, 22, 23, 24, 25, 26, 27, 28, 29, 30, 31, 32, 33 and 34 are displayed to examine the fitting deviation or residual error for velocity, micro-rotation, temperature, and concentration profiles via influential parameters for the persistence of defining the correctness and determining the convergence of solution method. The residual errors for velocity profile $f'(\xi)$ via Mr , M , γ , α and δ are displayed in Figs. 20, 21, 22, 23, and 24. The residual error for Mr converges quickly after $0.0 < \text{iteration} < 5.0$ which is shown in Fig. 20. This effect shows the correctness and convergence of the activated technique (HAM). Figure 21 displays the residual error for M on $f'(\xi)$. The residual error for M converges speedily after $0.0 < \text{iteration} < 6.0$. Figures 22, 23, and 24 indicate the residual errors for γ , α and δ on $f'(\xi)$, respectively. The residual errors for γ , α and δ rapidly converges after $0.0 < \text{iteration} < 5.5$, $0.0 < \text{iteration} < 5.5$, $0.0 < \text{iteration} < 6.0$ and $0.0 < \text{iteration} < 4.0$, respectively. The residual errors for micro-rotation profile $g(\xi)$ via α and Mr are displayed in Figs. 25 and 26 respectively. The residual error for α converges quickly after $0.0 < \text{iteration} < 7.5$. Also the residual error for Mr converges quickly for $0.0 < \text{iteration} < 2.8$. The residual errors for thermal profile $\theta(\xi)$ via Rd , Bi_1 , A^* , B^* and θ_w are displayed in Figs. 27, 28, 29, 30 and 31. The residual error for Rd converges quickly after $0.0 < \text{iteration} < 6.0$ which is shown in Fig. 27. The

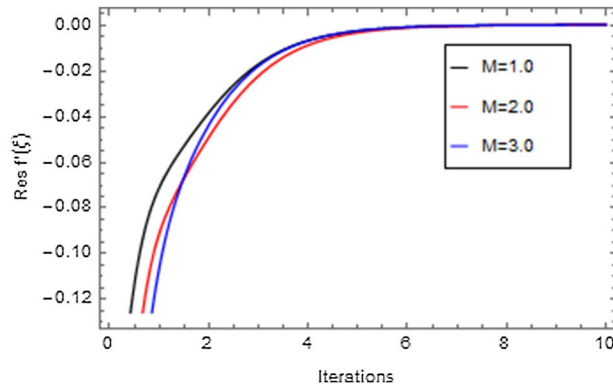


Figure 21. Residual error for M on $f'(\xi)$.

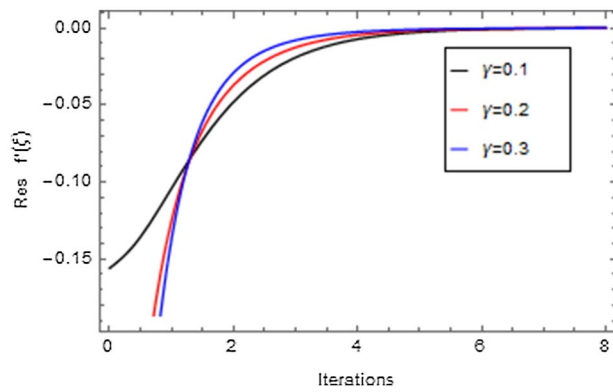


Figure 22. Residual error for γ on $f'(\xi)$.

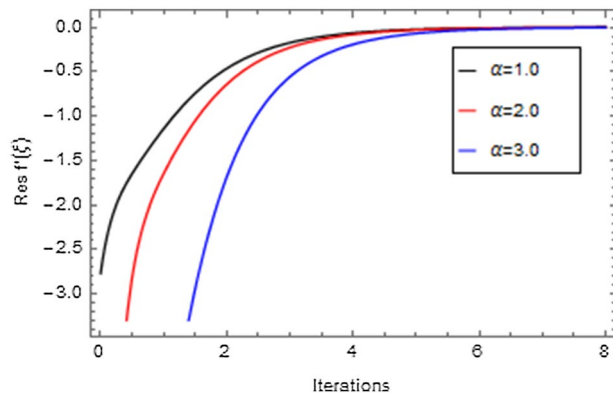


Figure 23. Residual error for α on $f'(\xi)$.

similar effect of Bi_1 is observed in Fig. 28. The residual errors for $\theta(\xi)$ via A^* and B^* are displayed in Figs. 29 and 30 respectively. The residual errors converge quickly after $0.0 < \text{iteration} < 5.0$. The residual error for θ_w converges quickly for $0.0 < \text{iteration} < 4.5$ which is displayed in Fig. 31. The residual errors for concentration profile $\phi(\xi)$ via Bi_2 , C_r and Sc are displayed in Figs. 32, 33, and 34. The residual errors of Bi_2 , C_r and Sc are quickly converge after $0.0 < \text{iteration} < 7.0$, $0.0 < \text{iteration} < 6.0$ and $0.0 < \text{iteration} < 5.0$, respectively.

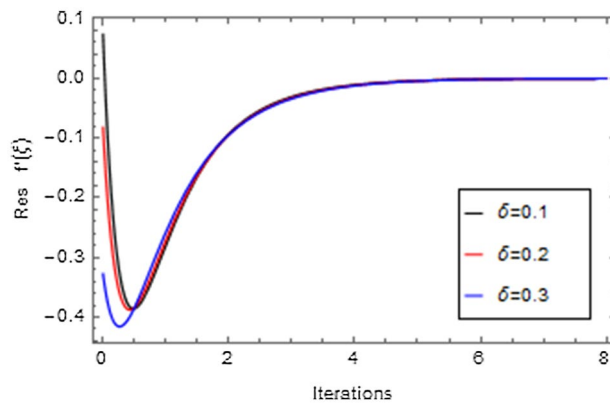


Figure 24. Residual error for δ on $f'(\xi)$.

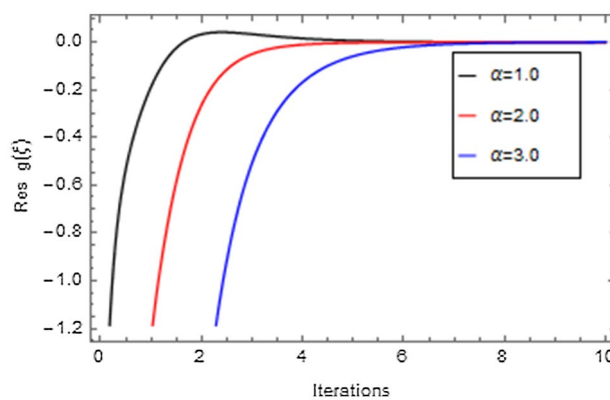


Figure 25. Residual error for α on $g(\xi)$.

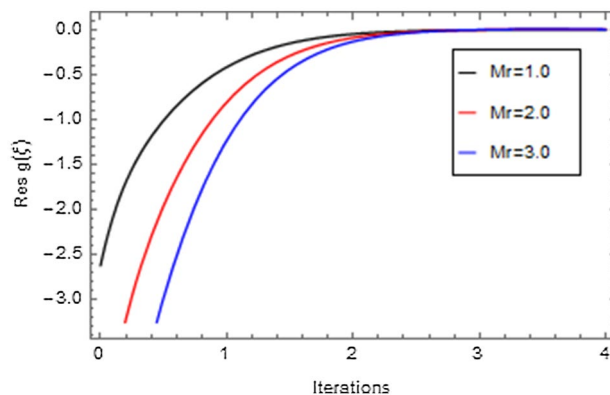


Figure 26. Residual error for Mr on $g(\xi)$.

An analytical solution is the actual solution. An analytical approach with varying properties can be used to analyze the behavior of systems. Unfortunately, there are many functional methods that relate to an analytical solution, and analytical methods are mostly of minimal use. That's why we have used a numerical method to produce response that seems to be closer to realistic outcomes. In the physical world, there are virtually no

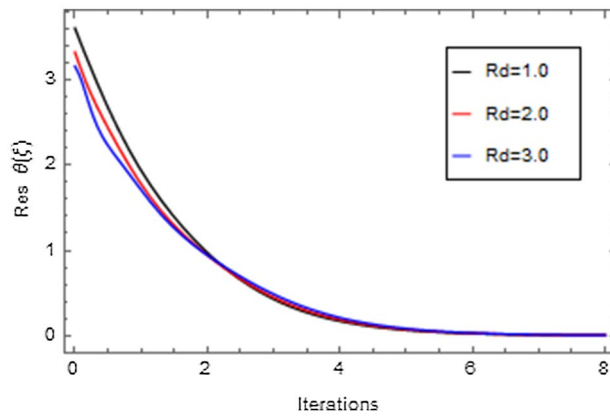


Figure 27. Residual error for Rd on $\theta(\xi)$.

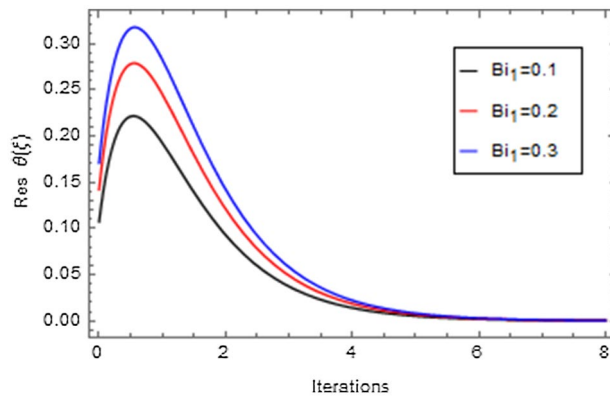


Figure 28. Residual error for Bi_1 on $\theta(\xi)$.

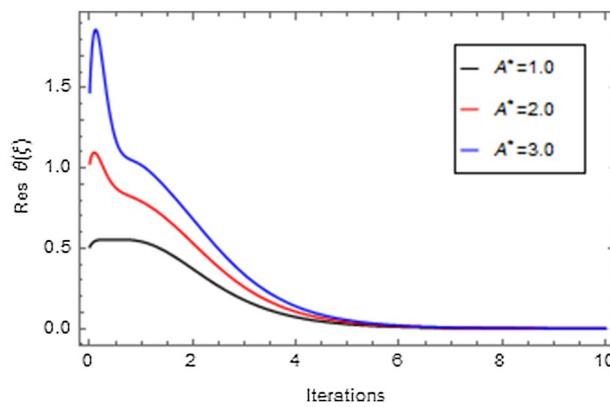


Figure 29. Residual error for A^* on $\theta(\xi)$.

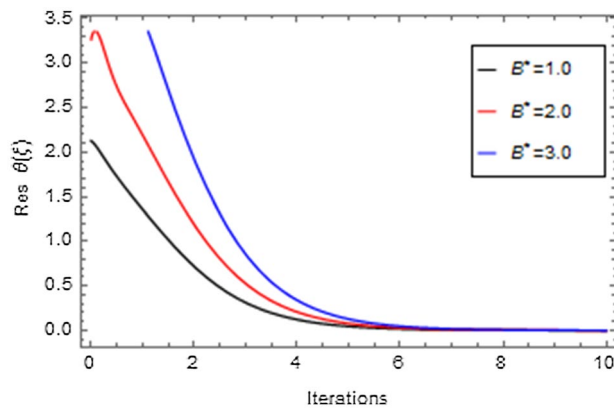


Figure 30. Residual error for B^* on $\theta(\xi)$.

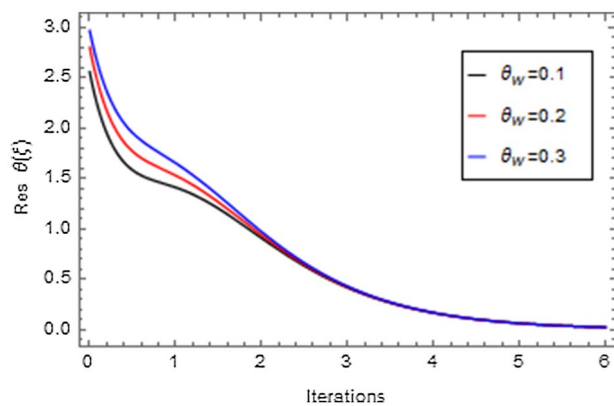


Figure 31. Residual error for θ_w on $\theta(\xi)$.

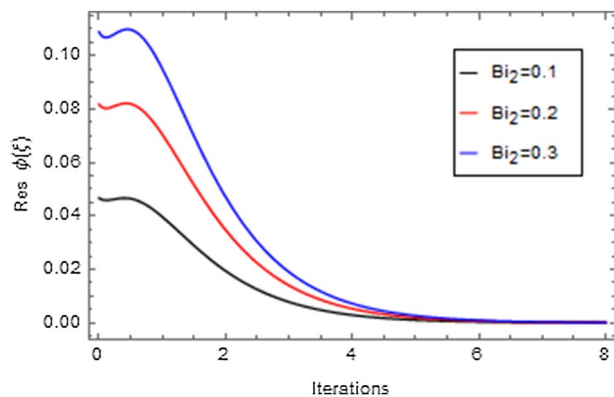


Figure 32. Residual error for Bi_2 on $\phi(\xi)$.

problems that can be perfectly solved, making the problem more complicated than other problems that can be exactly solved. There are three or four of them in existence that have already been resolved, but nevertheless still numerical approaches do not always provide an efficient solution. Numerical methods can be applied to any finalized physical geometry that is sometimes hard to address analytically. Here we have applied both analytical and numerical approaches to solve the modeled system of equations. Both techniques have quite close agreement as shown in Figs. 35, 36, 37 and 38 and Tables 1, 2, 3 and 4.

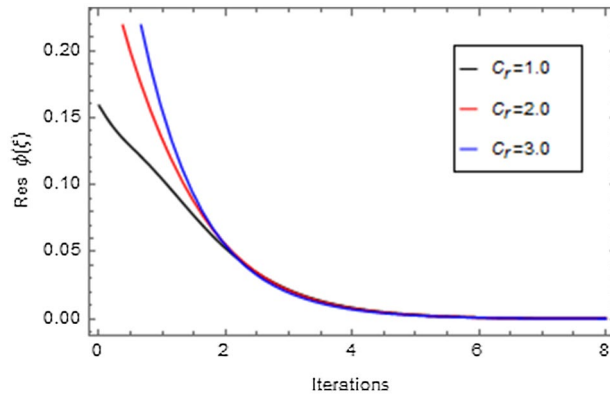


Figure 33. Residual error for C_r on $\phi(\xi)$.

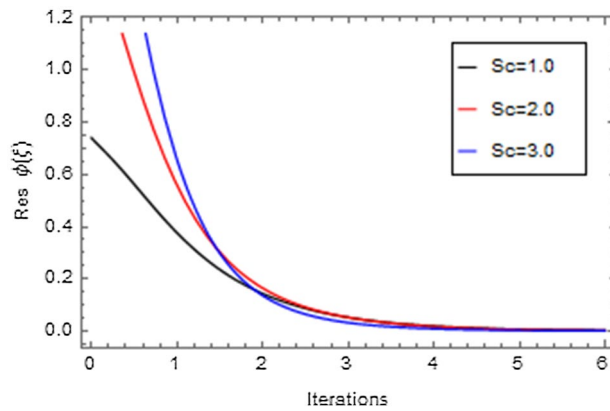


Figure 34. Residual error for Sc on $\phi(\xi)$.

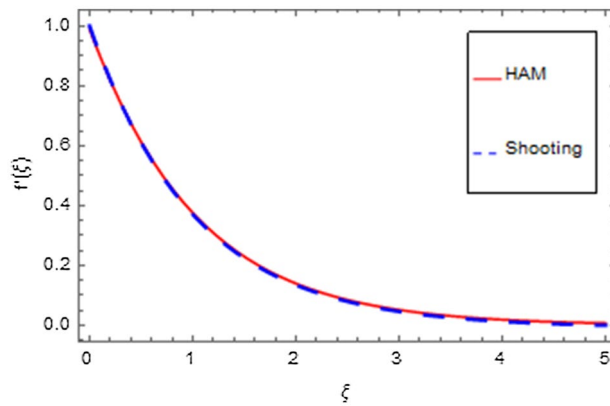


Figure 35. HAM versus shooting for $f'(\xi)$.

The assessment of skin factor, heat, and mass transfer rates are calculated in Tables 5, 6, 7 and 8. Table 5 represents the influence of corresponding parameters on skin factor C_f . The higher estimations of micropolar parameter reduce C_f whereas the larger values of magnetic parameter escalate C_f . Also, the present study in compared with Kumar et al.⁴³ and has agreed with the past analysis. Table 6 expresses the assessments of C_f and C_s against $\delta = 0$ and $\delta = 1$ for different values of the corresponding factors. C_f escalates for higher estimations

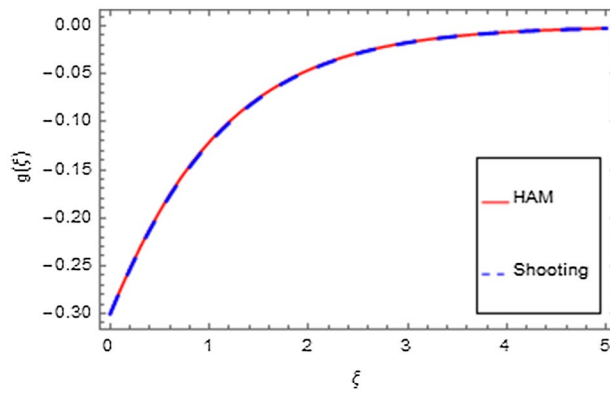


Figure 36. HAM versus shooting for $g(\xi)$.

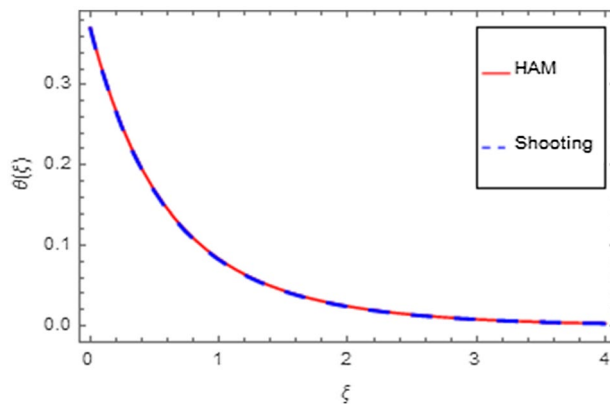


Figure 37. HAM versus shooting for $\theta(\xi)$.

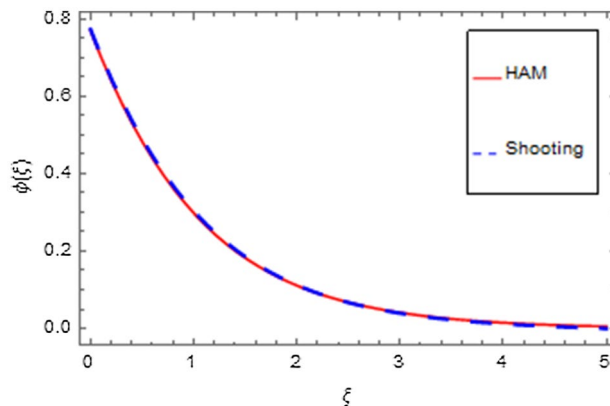


Figure 38. HAM versus shooting for $\phi(\xi)$.

of micropolar and micro-rotation parameters, while this behavior is opposite against magnetic parameter. C_s diminishes for higher values of micropolar, micro-rotation and magnetic field parameters. Table 7 illustrates the assessments of heat transfer rate Nu against $\delta = 0$ and $\delta = 1$ for unlike estimations of the corresponding factors. The higher estimations of heat source and sink, temperature ratio parameter and non-linear thermal energy parameter deescalates Nu , while the Biot number Bi_1 has inverse influence via Nu . Table 8 signifies assessment

ξ	HAM	Shooting
0.0	-2.42861×10^{-17}	0.000000
0.5	0.393406	0.394616
1.0	0.635715	0.634836
1.5	0.784043	0.780444
2.0	0.879508	0.868196
2.5	0.929562	0.920581
3.0	0.963022	0.951360
3.5	0.983341	0.968874
4.0	0.997526	0.978229
4.5	1.003160	0.982535
5.0	1.007700	0.983686

Table 1. HAM versus shooting for $f'(\xi)$.

ξ	HAM	Shooting
0.0	-0.300000	-0.300000
0.5	-0.194235	-0.192929
1.0	-0.122029	-0.120811
1.5	-0.075511	-0.074624
2.0	-0.046340	-0.045751
2.5	-0.028303	-0.027027
3.0	-0.017238	-0.017003
3.5	-0.010482	-0.010337
4.0	-0.006367	-0.006278
4.5	-0.003865	-0.003811
5.0	-0.002345	-0.002312

Table 2. HAM versus shooting for $g(\xi)$.

ξ	HAM	Shooting
0.0	0.188612	0.221368
0.5	-0.016859	0.015829
1.0	-0.049618	-0.026464
1.5	-0.042942	-0.027918
2.0	-0.030446	-0.021020
2.5	-0.020068	-0.014194
3.0	-0.012698	-0.009128
3.5	-0.007903	-0.005725
4.0	-0.004867	-0.003541
4.5	-0.002978	-0.002172
5.0	-0.001816	-0.001327

Table 3. HAM versus shooting for $\theta(\xi)$.

of Sh against $\delta = 0$ and $\delta = 1$ for different values of the corresponding factors. The higher estimations of Biot number Bi_2 and chemical reaction parameter hikes the Sh while the Schmidt number reduces Sh .

Conclusion

The electrically accompanying magnetohydrodynamic micropolar nanofluid flow over an extending sheet with secondary slips conditions and chemical reaction is considered here. The nanofluid flow is considered in two dimensional coordinates system. The proposed model is treated analytically and numerically. In order to improve

ξ	HAM	Shooting
0.0	0.768626	0.772196
0.5	0.482797	0.489327
1.0	0.297773	0.303516
1.5	0.182298	0.185537
2.0	0.111148	0.112200
2.5	0.067617	0.067150
3.0	0.041084	0.039587
3.5	0.024945	0.022635
4.0	0.015139	0.012001
4.5	0.009186	0.005016
5.0	0.005572	-1.4124×10^{-26}

Table 4. HAM versus shooting for $\phi(\xi)$.

α	M	Kumar et al. ⁴³	Present analysis
1.0		0.31709	0.317097
2.0		0.30676	0.306764
3.0		0.29713	0.297135
4.0		0.28841	0.288410
	0.1	0.32196	0.321963
	0.2	0.32623	0.326239
	0.3	0.32933	0.329332

Table 5. Judgment of the skin friction C_f of the present analysis with previous investigation against different estimations of α and M .

	C_f		C_s	
	$\delta = 0$	$\delta = 1$	$\delta = 0$	$\delta = 1$
$\alpha = 1.0$	-0.600453	-0.820725	-0.306835	-0.429662
$\alpha = 2.0$	-0.566971	-0.725193	-0.275236	-0.358507
$\alpha = 3.0$	-0.540782	-0.664839	-0.246085	-0.307275
$M_r = 1.0$	-0.270378	-0.276755	0.595872	0.610186
$M_r = 2.0$	-0.239194	-0.230289	0.632875	0.682285
$M_r = 3.0$	-0.215495	-0.197424	0.732696	0.767432
$M = 2.0$	-0.315575	-0.346335	0.058832	0.064125
$M = 3.0$	-0.342973	-0.392325	0.051585	0.058475
$M = 5.0$	-0.359486	-0.426532	0.045825	0.053846

Table 6. Assessments of C_f and C_s against different values of the corresponding factors.

the validity of the solutions and the method convergence, a descriptive demonstration of residual errors for various factors is presented. The main results are set out below.

- The velocity field heightens with the rise in micropolar factor, micro-rotation factor and primary velocity factor whereas reduces with escalation in magnetic factor, and secondary velocity slip parameter.
- The micro-rotation field rises with the escalation in primary order velocity slip factor while reduces with micro-rotation parameter, secondary order velocity slip parameter, and micropolar factor.
- The thermal field heightens with escalating non-uniform heat sink/source, Biot number, temperature ratio factor, and thermal radiation factor.
- The concentration field escalates with the increasing Biot number, while reduces with heightening chemical reaction and Schmidt number.
- Analytical and numerical approaches have quite close agreement.

	<i>Nu</i>	
	$\delta = 0$	$\delta = 1$
$A^* = 1.0$	0.855429	0.836741
$A^* = 2.0$	0.739836	0.742832
$A^* = 3.0$	0.420749	0.423428
$B^* = 1.0$	0.457476	0.397397
$B^* = 2.0$	0.123974	0.125165
$B^* = 3.0$	0.096440	0.096542
$\theta_w = 1.0$	0.336174	0.339975
$\theta_w = 2.0$	0.328617	0.329640
$\theta_w = 3.0$	0.326418	0.328393
$Bi_1 = 0.1$	0.107880	0.158824
$Bi_1 = 0.2$	0.184497	0.185759
$Bi_1 = 0.3$	0.252874	0.254822
$Rd = 1.0$	0.348153	0.350865
$Rd = 2.0$	0.345652	0.347967
$Rd = 3.0$	0.342979	0.345142

Table 7. Assessment of *Nu* against different values of the corresponding factors.

	<i>Sh</i>	
	$\delta = 0$	$\delta = 1$
$Bi_2 = 0.1$	0.481612	0.481265
$Bi_2 = 0.2$	0.634121	0.634742
$Bi_2 = 0.3$	0.708909	0.709680
$Cr = 1.0$	0.715317	0.716092
$Cr = 2.0$	0.531242	0.732006
$Cr = 3.0$	0.747005	0.747758
$Sc = 1.0$	0.857845	0.860115
$Sc = 2.0$	1.002340	1.006230
$Sc = 3.0$	1.131470	1.136380

Table 8. Assessment of *Sh* against different values of the corresponding factors.

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Author contributions

A.D and Z.S modeled and solved the problem. Z.S, A.D and S.I wrote the manuscript. S.I, W.K P.K and H.A contributed in the numerical computations and plotting the graphical results. N.S, P.K and W.K contributed in the revision. All authors finalized the manuscript after its internal evaluation.

Competing interests

The authors declare no competing interests.

Additional information

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