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Tuning the Geometrical Structures and Optical Properties of Blue-Emitting Iridium(III) Complexes through Dimethylamine Substitutions: A Theoretical Study

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Abstract: The geometrical structures and photophysical properties of Ir(4,6-dFppy)₂(pic) (Flrpic) and its derivative (*o*-Flr, *m*-Flr, *p*-Flr) with dimethylamine substituted at the picolinic acid (N[^]O) ligand were fully investigated by density functional theory and time-dependent density functional theory. The simulated electronic structure, as well as absorption and emission spectra of Flrpic are in good agreement with the experimental observations. The introduction of dimethylamine at the N[^]O ligand at different positions is beneficial to extend the π -electron delocalization, increase HOMO energy levels, and hence improve the hole injection and transfer ability compared with those of Flrpic. Furthermore, *o*-Flr, *m*-Flr, and *p*-Flr have large absorption intensity and participation of metal-to-ligand charge transfer (MLCT) contribution in the main absorption spectra, which would be useful to improve the intersystem crossing (ISC) from the singlet to triplet excited state. More importantly, the high quantum yield of *o*-Flr (which is explained based on the detailed analysis of triplet energy, E_{T1}), participation of ³MLCT contribution in the phosphorescent spectra, and energy difference between ³MLCT and triplet metal centered (³MC) d-d excited state compared with *m*-Flr and *p*-Flr indicate that *o*-Flr is expected to be an excellent blue phosphorescence emitter with high efficiency.

Keywords: DFT; dimethylamine; phosphorescence; ³MC d-d excited states

1. Introduction

Phosphorescent transition metal materials such as Re(I), Ru(II), and Ir(III) complexes are of great interest [1–4], because these materials have been extensively investigated for optoelectronic and microelectronic applications, such as dye-sensitized solar cells (DSSCs) and organic light-emitting diodes (OLEDs) [5–7]. Among these emitters, Ir(4,6-dFppy)₂(pic) (Flrpic) is one of the most well-known blue phosphorescent emitters, which is widely used as a dopant in the electroluminescent layers of OLEDs [8–10]. However, the quantum yield and equilibrium between the hole and electron transportation of Flrpic are considered to be difficult for its further application. The development of high-performance blue-emitting materials is strongly required for the fabrication of OLEDs. Unfortunately, it is a challenging task to achieve stable and high-efficiency blue phosphorescent emitters because the blue phosphorescence complex has a high emission state compared with other

color complexes. This enables thermal activation to triplet metal centered (3MC) d-d state, which is one of the main nonradiative processes from the emissive excited state for Ir(III) complexes [11,12].

In OLEDs, researchers can control the emission color and photophysical properties of emitters by a modification of the chelating ligands [13–17]. The common approach in tuning the photophysical properties of FIrpic is to modify the chemical structures of 2-(2,4-difluorophenyl) pyridine ($C^{\wedge}N$) ligands. Baranoff et al. [18] replaced the F atoms on FIrpic by H, Cl, and Br. They found that the large halogen substituent contributed to sizable distortions of specific ($C^{\wedge}N$) ligands, which were likely to play a role in the emissive and nonradiative properties when coupled with the heavy-atom effect. Owing to the introduction of the strong electron withdrawing perfluoro carbonyl group on ($C^{\wedge}N$) ligands of FIrpic, the newly designed Ir complexes displayed deepest blue emissions and considerably high external quantum efficiencies (EQEs) [19]. Li et al. [20] investigated a series of Ir(III) complexes by introducing tert-butyl, *n*-heptyl, 3-ethylheptyl, and 2,4,6-trimethylphenyl substituted groups into the ($C^{\wedge}N$) ligands of FIrpic. They revealed that both the substituted group and position had an effect on tuning the emission energies and quantum yield. However, the important chelating picolinic acid ($N^{\wedge}O$) ligand of FIrpic has not received as much attention, to the best of our knowledge. Herein, we will theoretically design and investigate a type of Ir complexes *o*-FIr, *m*-FIr, *p*-FIr (Figure 1) by introducing the dimethylamine substituted at the picolinic acid ($N^{\wedge}O$) ligand, because the dimethylamine is one of the most popular functional groups used as hole transfer materials. We perform detailed theoretical calculations on the electronic structure, absorption and emission properties, quantum efficiency, as well as charge transport properties, aiming at revealing the relationship between the structures and photophysical properties. This work is expected to give some insights into the design and synthesis of blue emitters with high quantum efficiency.

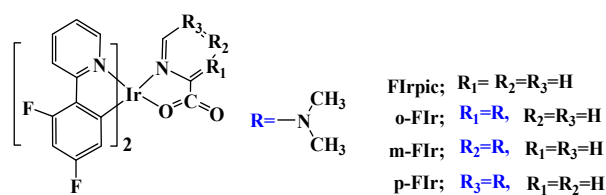


Figure 1. Schematic structures of the studied complexes.

2. Results and Discussion

2.1. The Optimized Geometries in the Ground and Lowest Lying Triplet Excited States

The schematic structures of FIrpic, *o*-FIr, *m*-FIr, and *p*-FIr are depicted in Figure 1. The ground-state (S_0) geometry of FIrpic was optimized by B3LYP, CAMB3LYP, M06L, MPW1PW91 functionals with the same basis set. The detailed geometries and available crystal data of FIrpic [21] are collected in Table S1 in supporting information. It is found that M06L functional predicts more satisfied geometrical structures by compared with the experimental data. Therefore, the S_0 geometries of *o*-FIr, *m*-FIr, *p*-FIr are calculated by M06L functional, and the optimized geometries as well as the numbering of some key atoms are depicted in Figure 2. The optimized primary metal–ligand bond lengths of these designed complexes in the ground-state (S_0) geometry are collected in Table 1.

As listed in Table 1, the N_1 -Ir- N_2 and O_1 -Ir- N_2 of FIrpic are ca. 90° , and the dihedral angle of N_1 - O_1 - C_2 - C_1 of FIrpic is 6.8° . Obviously, FIrpic has distorted octahedral geometry around the Ir(III) atom. Table 1 shows that the N_1 -Ir- N_2 , O_1 -Ir- N_2 , and N_1 - O_1 - C_2 - C_1 of *o*-FIr, *m*-FIr, and *p*-FIr are similar to those of FIrpic. Clearly, the introduction of a dimethylamine moiety on the pyridine moiety of picolinic acid ($N^{\wedge}O$) does not affect the distorted octahedral geometry of the Ir(III) complex. For *o*-FIr and *p*-FIr, because of the introduction of dimethylamine group on the pyridine, the Ir- N_1 and Ir- O_1 of *o*-FIr and *p*-FIr are shortened (0.005 Å and 0.01 Å for *o*-FIr, while 0.05 Å and 0.016 Å for *p*-FIr) in comparison to those of FIrpic. These calculated results indicate that there is a strong interaction

between the Ir(III) and (N[^]O) ligand. Furthermore, the Ir-N₁ bond length of *o*-FIr, *m*-FIr, and *p*-FIr are significantly shortened compared with that of FIrpic; therefore, the possible of dissociation of Ir-N₁ bond in the ³MC state is expected to be decreased by the incorporation of N(CH₃)₂ groups.

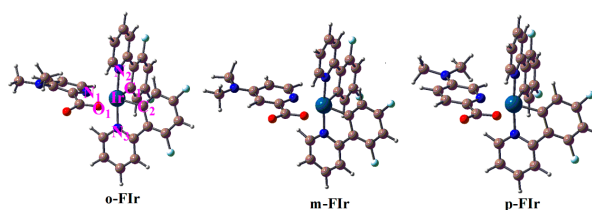


Figure 2. Optimized structures of *o*-FIr, *m*-FIr, and *p*-FIr in the ground state, together with the number of some key atoms.

Table 1. Calculated metal–ligand bond lengths (Å), bond angles (°), and dihedral angles (°) for studied complexes in the ground state (S₀) and first excited triplet state (T₁).

Molecule	State	Ir-N ₁	Ir-O ₁	Ir-N ₂	Ir-C ₁	Ir-N ₃	Ir-C ₂	N ₁ -Ir-N ₂	O ₁ -Ir-N ₂	N ₁ -O ₁ -C ₂ -C ₁
FIrpic	S ₀	2.215	2.190	2.066	2.001	2.058	1.997	88.4	93.5	6.8
	T ₁	2.313	2.172	2.069	1.979	2.060	1.975	97.1	88.7	−3.4
<i>o</i> -FIr	S ₀	2.210	2.180	2.066	2.002	2.054	1.997	87.3	94.0	5.0
	T ₁	2.316	2.154	2.057	1.981	2.067	1.977	84.4	93.7	3.1
<i>m</i> -FIr	S ₀	2.206	2.193	2.064	2.002	2.056	1.995	88.6	93.0	7.2
	T ₁	2.285	2.171	2.069	1.979	2.057	1.978	95.9	89.0	−4.6
<i>p</i> -FIr	S ₀	2.165	2.177	2.055	2.010	2.045	2.001	87.6	92.9	6.3
	T ₁	2.305	2.160	2.068	1.982	2.058	1.977	97.4	88.7	−3.0

The lowest lying triplet excited states (T₁) of studied complexes are also listed in Table 1. As listed in Table 1, the Ir-N₁ bond length of these complexes is dramatically elongated in the T₁ state in comparison with that in S₀ state. This suggests that the interaction between the N[^]O ligand and the metal is weakened in the T₁ state. From the S₀ to the T₁ state, the Ir-C_{1/2} and Ir-N_{2/3} bond lengths of 2-(2,4-difluorophenyl)pyridine (C[^]N) ligands in *o*-FIr, *m*-FIr, and *p*-FIr are contracted by the substitutions of the dimethylamine groups on the pyridine moiety. The strengthened metal–ligand bonds will strength the interaction between the metal and C[^]N ligands and further benefit the charge transfer from the metal to ligand, which may improve the quantum efficiency.

2.2. Frontier Molecular Orbitals

Considering that frontier molecular orbitals are useful to provide the nature of the charge transfer and photophysical properties in the excited state, the HOMO and LUMO energies, HOMO–LUMO energy gaps, and the contour plots of these orbitals of studied complexes are drawn in Figure 3. Detailed descriptions of the frontier molecular compositions are collected in Table S2 in the supporting information.

Normally, when the interfacial vacuum energy shift can be neglected, the hole (or electron) injection energy barrier from the electrode to semiconductor (Φ_h) (or Φ_e) can be calculated by Φ_h = |HOMO| − Φ_m (or Φ_e = Φ_m − |LUMO|), where Φ_m is the work function of the electrode [22]. Consider that Au (Φ_m = 5.1 eV) is widely used as electrode in OLEDs; therefore, it is used as reference electrode [22]. As depicted in Figure 3, the HOMO and LUMO energies of *o*-FIr, *p*-FIr, and *m*-FIr increase remarkably (−4.65 eV and −2.15 eV for *o*-FIr, −4.59 eV and −2.07 eV for *m*-FIr, −4.65 eV and −2.16 eV for *p*-FIr) compared with FIrpic (−4.80 eV and −2.30 eV). Therefore, the calculated Φ_h of these complexes are (in the region of 0.30–0.51 eV) much smaller than those of the calculated Φ_h values (in the region of 2.80–3.03 eV). It is obvious that these emitters are suitable for use as p-type organic semiconductors [23]. Furthermore, the HOMO energy level of *o*-FIr (0.15 eV), *m*-FIr (0.21 eV), and

p-FIr (0.15 eV) is significantly higher than that of FIrpic. The higher HOMO energies will benefit the hole injection abilities. Together with the increasing HOMO and LUMO energy, the HOMO–LUMO energy gaps of our designed molecules are similar to that of FIrpic, which means that absorption and phosphorescent spectra of these designed complexes might be similar to that of FIrpic when these spectra are predominantly contributed by the HOMO→LUMO transition.

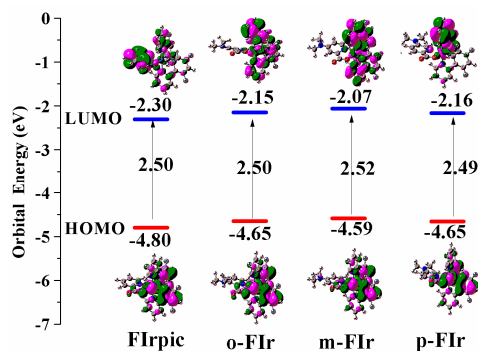


Figure 3. The orbital energies and contour plot of the HOMO and LUMO of studied complexes.

For FIrpic, *o*-FIr, *p*-FIr, and *m*-FIr, the electron densities of HOMOs are mainly located on 50% d(Ir) orbital and 38% $\pi(\text{C}^{\wedge}\text{N}$ ligands) orbital, as shown in Figure 3 and Table S2. Owing to the introduction of $\text{N}(\text{CH}_3)_2$ groups, the LUMOs in *o*-FIr, *m*-FIr, and *p*-FIr are mainly focused on the $\text{C}^{\wedge}\text{N}$ ligands with about 90% compositions, while the LUMO in FIrpic is contributed by $\text{N}^{\wedge}\text{O}$ ligand. The HOMO-1 of these studied complexes is mainly focused on the ca. 50% d orbital the ca. 30% π orbitals of $\text{N}^{\wedge}\text{O}$ ligand, while their HOMO-2 are mainly contributed by $\text{N}^{\wedge}\text{O}$ ligand. The contributions of d(Ir) orbital on HOMO-3 of *o*-FIr (67.5%) and *m*-FIr (56.9%) are similar to that of FIrpic (67.1%). This means that the $\text{N}(\text{CH}_3)_2$ substitutions have a minor effect on the composition of the frontier molecular orbital. Therefore, the nature of electron transition upon excitations of *o*-FIr, *p*-FIr, and *m*-FIr might be similar with those of FIrpic, which will be further investigated in the following section.

2.3. Absorption Spectra

The calculated absorption spectra of studied complexes associated with oscillator strengths (f) and the corresponding compositions, as well as the experimental data [24] are given in Table 2. Simulated absorption spectra of these complexes with the data calculated at the TD-M06L method are drawn in Figure 4.

As listed in Table 2, the main calculated values of these bands are localized at 453.00, 380.95, and 265.86 nm for FIrpic, respectively, which are in good agreement with the experimental results of 458, 379, and 256 nm, respectively [24]. For the low-lying absorption spectra of FIrpic, the excitations of HOMO→LUMO + 1 and HOMO – 1→LUMO + 2 are dominantly responsible for the absorption peaks at 453.00 and 380.95 nm, respectively. As listed in Table S2, the HOMO and HOMO – 1 are delocalized on the Ir(d), $\text{N}^{\wedge}\text{O}$, and $\text{N}^{\wedge}\text{C}$ ligand, while the LUMO + 1 and LUMO + 2 are mainly contributed by the $\text{N}^{\wedge}\text{C}$ ligand. Therefore, from the analysis of FMOs, the absorption peaks at 453.00 and 380.95 nm can be characterized as metal-to-ligand charge transfer (MLCT), inter-ligand charge transfer (ILCT), and ligand–ligand charge transfer (LLCT). The absorption peak with largest oscillator strength ($f = 0.2104$) of FIrpic is located at 265.86 nm, which is arising from the HOMO – 8→LUMO + 2 and HOMO→LUMO + 7 transitions. Because of the large amount of Ir(d) in HOMO (50.7%) and HOMO – 8 (20.8%), the absorption peak at 265.86 nm can mainly be assigned to the MLCT/LLCT/ILCT character. From the analysis of frontier molecule orbital compositions, the amount of MLCT of the 453.00, 380.95, and 265.86 nm for FIrpic is 21.07%, 21.07%, and 7.16%, respectively. The significant MLCT contributions on the lower-lying absorptions of FIrpic are in good agreement with the experimental results [24].

Table 2. The calculated maximum absorption wavelength (nm), oscillator strengths (f), and major contribution transition composition for studied complexes obtained by TD-M06L method. FIrpic: Ir(4,6-dFppy)₂(pic).

Molecule	Transition	λ (nm)	f	Composition	CI	Exp. [24] (nm)
FIrpic	S ₀ →S ₂	453.00	0.0250	HOMO→LUMO + 1	0.68971	455
	S ₀ →S ₇	380.95	0.0614	HOMO – 1→LUMO + 2	0.55936	379
	S ₀ →S ₅₀	265.86	0.2104	HOMO – 8→LUMO + 2 HOMO→LUMO + 7	0.38300 0.24359	256
o-FIr	S ₀ →S ₁	463.87	0.0258	HOMO→LUMO	0.70168	
	S ₀ →S ₇	395.15	0.0266	HOMO – 2→LUMO + 2 HOMO – 1→LUMO + 2	0.56244 –0.40890	
	S ₀ →S ₅₂	269.04	0.1135	HOMO – 9→LUMO + 1 HOMO – 5→LUMO + 5 HOMO – 1→LUMO + 7	0.25480 0.42176 0.23321	
	S ₀ →S ₅₃	268.66	0.1642	HOMO – 7→LUMO + 4 HOMO – 1→LUMO + 7	–0.35667 0.36355	
	S ₀ →S ₅₄	266.76	0.1318	HOMO – 9→LUMO + 2 HOMO – 2→LUMO + 7 HOMO – 1→LUMO + 7 HOMO→LUMO + 7	0.21920 0.26962 –0.21838 0.41903	
	S ₀ →S ₅₅	266.32	0.0836	HOMO – 5→LUMO + 5	0.42068	
	S ₀ →S ₅₆	264.54	0.1094	HOMO – 9→LUMO + 2 HOMO – 8→LUMO + 2 HOMO – 1→LUMO + 5	0.48856 –0.19242 0.27865	
	S ₀ →S ₅₇	262.33	0.1716	HOMO – 9→LUMO + 2 HOMO – 5→LUMO + 5 HOMO – 4→LUMO + 5 HOMO – 2→LUMO + 5 HOMO – 2→LUMO + 7 HOMO – 1→LUMO + 5	–0.24371 –0.24091 –0.25407 0.20626 0.30782 0.32002	
	S ₀ →S ₁	462.95	0.0263	HOMO→LUMO	0.70166	
	S ₀ →S ₆	393.55	0.0423	HOMO – 2→LUMO HOMO – 2→LUMO + 1 HOMO – 1→LUMO + 1	–0.41398 0.42182 0.31148	
S ₀ →S ₅₃	267.06	0.3060	HOMO – 8→LUMO + 2 HOMO→LUMO+7	–0.24031 0.25154		
p-FIr	S ₀ →S ₁	463.29	0.0286	HOMO→LUMO	0.70148	
	S ₀ →S ₈	388.02	0.0581	HOMO – 1→LUMO + 1 HOMO→LUMO + 3	–0.43685 0.47613	
	S ₀ →S ₅₂	271.08	0.2512	HOMO – 7→LUMO + 4 HOMO – 3→LUMO + 5 HOMO – 2→LUMO + 5 HOMO→LUMO + 7	–0.27219 0.21898 0.29024 –0.21723	

Table 2 also shows the lowest-lying absorption peak of *o*-FIr, *m*-FIr, and *p*-FIr at 463.87, 462.95, and 463.29 nm, respectively, which is dominantly responsible for the excitation of HOMO→LUMO. Table S2 shows that the HOMO of *o*-FIr, *m*-FIr, and *p*-FIr is composed of 50.1% Ir(d) + 47.7% $\pi(\text{N}^{\wedge}\text{O} + \text{N}^{\wedge}\text{C})$, 50.3% Ir(d) + 47.9% $\pi(\text{N}^{\wedge}\text{O} + \text{N}^{\wedge}\text{C})$, 48.6% Ir(d) + 49.3% $\pi(\text{N}^{\wedge}\text{O} + \text{N}^{\wedge}\text{C})$, respectively, while their LUMOs are contributed by ca. 80% $\pi^*(\text{N}^{\wedge}\text{C})$. Thus, the lowest-lying absorptions of these complexes originate from the $[\text{Ir}(\text{d}) + \pi(\text{N}^{\wedge}\text{O} + \text{N}^{\wedge}\text{C})] \rightarrow \pi^*(\text{N}^{\wedge}\text{C})$ excited state with MLCT/ILCT/LLCT character. The amount of MLCT of the lowest-lying absorption of *o*-FIr, *m*-FIr, and *p*-FIr is ca.22%, which is similar to that of FIrpic. Furthermore, as shown in Figure 4, the absorption intensity of N(CH₃)₂-substituted Ir complexes (*m*-FIr, *o*-FIr, *p*-FIr) in the highest energy absorption spectra have been effectively enhanced, which may increase the intersystem crossing rates (ISC). For *o*-FIr, the highest energy absorption peak

(262.33 nm) is contributed by the mixture excitations of HOMO – 9 → LUMO + 2, HOMO – 5 → LUMO + 5, HOMO – 4 → LUMO + 5, HOMO – 2 → LUMO + 5, HOMO – 1 → LUMO + 5. Since the Ir(d) in the HOMO – 9, HOMO – 4, HOMO is 22.7%, 67.5%, 50.1%, respectively, the MLCT contributed 8.9% to the 262.33 nm of *o*-FIr, which is slightly larger than that of FIrpic. For *m*-FIr and *p*-FIr, the absorption peak with large oscillator strength is located at 267.06 (attributed by HOMO – 8 → LUMO + 2 and HOMO → LUMO + 7) and 271.08 nm (attributed by HOMO – 7 → LUMO + 4, HOMO – 3 → LUMO + 5, HOMO – 2 → LUMO + 5, HOMO → LUMO + 7), respectively. Combining the transition configuration with the orbital composition, the amount of MLCT of the high energy absorption peaks for *m*-FIr and *p*-FIr are only 4.4% and 5.2%, respectively. Clearly, the contribution of MLCT in the whole absorption spectra of *o*-FIr are remarkably larger than those of FIrpic, *m*-FIr, and *p*-FIr. Normally, the participation of MLCT in the absorption spectra is a critical factor affecting the intensity of the singlet–triplet transitions, and thus improves the quantum efficiency. Thus, the calculated results indicate that the introduction of N(CH₃)₂ groups at the ortho-position of N[^]O ligand is useful to enhance the intersystem crossing and singlet–triplet transitions, and therefore the possibility of the phosphorescent quantum yield might be enhanced.

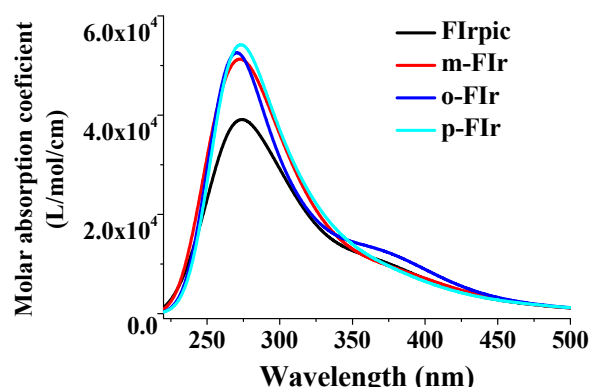


Figure 4. Simulated absorption spectra of studied complexes with the data calculated at the TD-M06L level.

2.4. Phosphorescence Emission Spectra

Based on the optimized lowest triplet excited state (T_1) geometries, vertical T_1 – S_0 transition energies (ΔE_{vert}) and the emission wavelength (nm) obtained by CAMB3LYP functional are collected in Table 3, together with available experimental values [24]. Since the vertical T_1 – S_0 transition energy (ΔE_{vert}) is the electronic energy difference between the T_1 and S_0 state at the optimized T_1 geometry, it is therefore strongly dependent on the functional. The molecular orbital compositions responsible for the emission spectra of the studied complex are given in Table S3. The 0–0 T_1 → S_0 transition energy (ΔE_{0-0}) is calculated by the energy difference between the S_0 and T_1 states at their respective optimized geometries obtained by M06 functional.

As listed in Table 3, the calculated emission peak for FIrpic based on the calculated E_{vert} and E_{0-0} method is located at 490.49 and 480.86 nm, respectively. The slightly longer wavelength derived from E_{vert} is caused by only one single configuration, while real emission is composed of more than one configuration. The phosphorescence at 490.49 nm of FIrpic is contributed by the excitation of HOMO → LUMO. According to the orbital composition analyses, the HOMO is contributed by 42.8% Ir(d) + 54.3% $\pi(\text{C}^{\wedge}\text{N} + \text{N}^{\wedge}\text{O})$, while the LUMO is mainly focused on C[^]N ligand. Accordingly, the emission of FIrpic can be attributed to ³MLCT/³LLCT/³ILCT character. The d(Ir) orbital contributes to 42.8% for HOMO, and thus the emission of FIrpic at 490.49 nm has ³MLCT (20.31%) transition character. For N(CH₃)₂-substituted Ir complexes, the phosphorescent emissions based on E_{vert} method are located at 492.78, 484.32, and 485.83 nm for *o*-FIr, *m*-FIr, and *p*-FIr, respectively, which is dominantly responsible for excitations of HOMO → LUMO and HOMO – 4 → LUMO. This is because the HOMO

and HOMO – 4 of *o*-FIr, *m*-FIr, and *p*-FIr are mainly contributed by d(Ir) and $\pi(\text{C}^{\wedge}\text{N} + \text{N}^{\wedge}\text{O})$ (besides the HOMO – 4 of *m*-FIr), while the LUMOs of *o*-FIr and *m*-FIr are located at $\pi^*(\text{C}^{\wedge}\text{N})$. Table 3 shows that the calculated E_{0-0} for these designed complexes are ca. 480nm. Accordingly, these designed complexes have similar phosphorescent emissions and transition character with that of FIrpic, which might be a potential candidate for blue-emitting materials. Furthermore, the d(Ir) orbital contributions to HOMO and HOMO – 4 (42.8% and 0.97% for *m*-FIr, 30.68% and 36.68% for *p*-FIr) are much smaller than that of HOMO (41.35%) and HOMO – 4 (63.39%) for *o*-FIr. The larger d(Ir) orbital composition results in slightly larger ${}^3\text{MLCT}\%$ in its phosphorescence of *o*-FIr (18.78%) than *m*-FIr (9.11%) and *p*-FIr (11.56%). Therefore, the *o*-FIr may have improved Φ_p compared with *m*-FIr and *p*-FIr.

Table 3. The calculated phosphorescence emission wavelength (nm), along with metal–ligand charge transfer character (${}^3\text{MLCT}\%$, %) and available experimental values.

Molecule			Evert		E ₀₋₀		Exp [24]
	λ (nm)	E (eV)	Composition	${}^3\text{MLCT}\%$	λ (nm)	E (eV)	λ (nm)
FIrpic	490.49	2.53	0.49 (HOMO→LUMO)	20.31	480.86	2.58	468
<i>o</i> -FIr	492.78	2.52	0.45 (HOMO – 4→LUMO) 0.43 (HOMO→LUMO)	18.78	482.58	2.60	
<i>m</i> -FIr	484.32	2.56	–0.36 (HOMO – 4→LUMO) 0.40 (HOMO→LUMO)	9.11	481.78	2.57	
<i>p</i> -FIr	485.83	2.55	–0.36 (HOMO – 4→LUMO) 0.49 (HOMO→LUMO)	11.56	481.76	2.57	

To give some insight into the nonradiative rate constant (k_{nr}), the high lying metal-centered (${}^3\text{MC}$ d-d) excited states of studied complexes—which is one of the most important deactivation pathways for the phosphorescence emission from T_1 in transition-metal complexes—were optimized by unrestricted M06L method. The energy difference between ${}^3\text{MLCT}/\pi-\pi^*$ and ${}^3\text{MC}$ d-d excited states are drawn in Figure 5, along with the spin density distribution in the ${}^3\text{MC}$ d-d state. As shown in Figure 5, the ${}^3\text{MC}$ d-d state of studied complexes are lying at a relatively higher energy than that of ${}^3\text{MLCT}/\pi-\pi^*$ state, which suggests that the possibility of a radiationless pathway is low. Furthermore, the energy difference between the ${}^3\text{MLCT}/\pi-\pi^*$ and ${}^3\text{MC}$ d-d excited states of *o*-FIr (0.30 eV) and *m*-FIr (0.30 eV) is larger than that of FIrpic (0.24 eV). Therefore, these calculation results show that ortho- and meta- substituted Ir complexes are useful to decrease the K_{nr} . Based on the above viewpoint, it becomes clear that introducing the $\text{N}(\text{CH}_3)_2$ at the ortho-position of $\text{N}^{\wedge}\text{O}$ ligand in FIrpic led to larger ISC and k_r values, as well as the smaller k_{nr} value relative to others. Therefore, *o*-FIr is considered to be used as potential blue-emitting emits with high quantum efficiency.

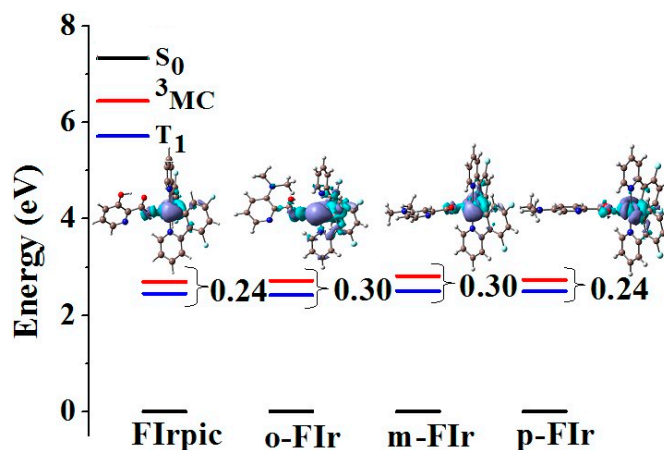


Figure 5. Energy level diagram of the studied complexes in the T_1 and triplet metal centered (${}^3\text{MC}$) excited states, together with the contour plots of the spin density distribution in the ${}^3\text{MC}$ state.

2.5. Performance in OLEDs

To evaluate the performance of OLED devices, it is highly important to investigate the charge injection and transfer ability. The ionization potential (IP) and electron affinity (EA) are usually employed to evaluate the energy barrier for the hole and electron injections. Normally, the lower IP (larger EA) value suggests that it is easier to inject hole (electron). According to the Marcus theory [25], the electron transfer rate constants can be expressed by the following formula:

$$k = \left(\frac{\pi}{\lambda k_B T} \right)^{1/2} \frac{V^2}{\hbar} \exp\left(-\frac{\lambda}{4k_B T}\right) \quad (1)$$

where T and k_B is the temperature and Boltzmann constant, respectively. The λ is the reorganization energy, and V is the coupling matrix element. According to Equation (1), the intermolecular charge transfer rate k is determined by two factors, λ and V. Due to the limited intermolecular charge transfer range in the solid state, the mobility of charges has been demonstrated to be primarily associated with the reorganization energy λ for OLED materials [26]. The detailed description of the calculations of vertical ionization potential (IP(v)), vertical electron affinity (EA(v)), hole extraction potential (HEP), electron extraction potential (EEP) and electron reorganization energy ($\lambda_{\text{electron}}$), as well as hole reorganization energy (λ_{hole}) are listed in our previous report [16]. According to the previous report [16], the IP, EA, and λ of studied complexes are calculated and listed in Table 4.

Table 4. The calculated ionization potential (IP, eV), electron affinities (EA, eV), hole extraction potential (HEP, eV), electron extraction potential (EEP, eV) hole/electron reorganization energy ($\lambda_{\text{hole}}/\lambda_{\text{electron}}$, eV).

Molecule	IP(v)	HEP	EA(v)	EEP	λ_{hole}	$\lambda_{\text{electron}}$
Flrpic	6.56	6.29	0.72	0.98	0.27	0.26
<i>o</i> -Flr	6.32	6.06	0.48	0.69	0.26	0.21
<i>m</i> -Flr	6.39	6.09	0.56	0.75	0.30	0.19
<i>p</i> -Flr	6.34	6.11	0.53	0.73	0.23	0.20

Compared with Flrpic, all these designed complexes have smaller IP(v) and HEP, which is consistent with their higher HOMO energy levels. Thus, the hole injection ability of these designed complexes are improved by introduction of dimethylamine substitutions at the C[^]O ligand. However, *o*-Flr, *m*-Flr, and *p*-Flr have smaller EA values than Flrpic, which will lead to worse electron injection ability compared with Flrpic. According to Equation (1), the λ is required to be small for efficient charge transport. As listed in Table 4, the $\lambda_{\text{electron}}$ of studied complexes are slightly smaller than that of Flrpic, indicating that these complexes have excellent electron transfer ability. The difference between the λ_{hole} and $\lambda_{\text{electron}}$ of Flrpic, *o*-Flr, and *p*-Flr is less than 0.05 eV. Hence, it is expected that the Flrpic, *o*-Flr, and *p*-Flr might have excellent charge transfer balance, which is one of the key factors affecting the performance of OLEDs.

3. Computational Details

The geometrical structures of the ground state (S_0) of Flrpic were fully optimized by B3LYP [27,28], CAMB3LYP [29], M06L [30], and mPW1PW91 [31] methods, respectively, combined with 6-31G(d) basis set for the non-metal atoms (H, C, N, O, and F) and a “double- η ” quality basis set LANL2DZ for Ir atom. The calculated vibrational frequencies with no imaginary frequencies based on the optimized geometries verify that each of the optimized structures was a true minimum on the potential energy surface. The calculated results are collected in Table S1 in the supporting information, together with the crystal data [21]. As depicted in Table S1, the bond lengths obtained for B3LYP and CAMBLYP overestimated the bond lengths, especially Ir-N₁ (0.024 Å for CAMB3LYP and 0.031 Å for B3LYP) and Ir-N₃ (0.049 Å for CAMB3LYP and 0.073 Å for B3LYP), while the bond length differences between the calculated results obtained by M06L and crystal data were small. Therefore, the ground-state

geometries of the studied complexes were calculated by M06L method. Based on the ground-optimized structure, the vertical excitation energies were calculated by TD-M06L method with continuum polarizable continuum model (CPCM). The lowest-lying triplet excited state geometries (T_1) and triplet-centered (3MC) d-d states were optimized by unrestricted M06L method. Based on the T_1 equilibrium geometries, four different functionals—PBE0, B3LYP, CAMB3LYP, and M06L—were performed on the emission spectra calculations (as listed in Table S4). The CAMB3LYP functional was adopted to predict the emission spectra of these complexes, because the calculated results obtained by CAMB3LYP were consistent with the experimental results. All calculations were performed with the Gaussian 09 software package [32].

4. Conclusions

In this work, the geometrical and electronic structures, absorption and emission spectra, as well as phosphorescence efficiencies for a series of dimethylamine-substituted iridium(III) complexes FIrpic have been investigated by density functional theory (DFT) and time-dependent density functional theory (TDDFT) methods. The results indicate that the introduction of dimethylamine is useful to extend the π -electron delocalization, increase HOMO energy levels, and enhance hole injection and transfer ability. Furthermore, the absorption intensity and participation of MLCT of these dimethylamine substitutions are significantly stronger than that of FIrpic, suggesting the possibility of singlet-to-triplet transitions are enhanced, and therefore the phosphorescent quantum yield might be increased by introducing dimethylamine. Importantly, through the analyses of triplet energy (E_{T1}), participation of 3MLCT in phosphorescent spectra, as well as the energy difference between the ${}^3MLCT/\pi$ - π and 3MC d-d state, it can be inferred that o -FIr has larger k_r values and smaller k_{nr} values compared with other molecules. So, the strategy of introducing the dimethylamine in $N^{\wedge}O$ ligand of FIrpic should be useful in the design of more-efficient blue emitters.

Supplementary Materials: Supplementary materials are available on line.

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Author Contributions: Hong-Qu Tang performed the calculation work; Guo-Jun Kang and Xue-Feng Ren analyzed the data and wrote the manuscript.

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