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# Control of Polarization Orientation Angle of Scattered Light Based on Metasurfaces: $\mathbf{- 9 0}$ to $+90^{\circ}$ Linear Variation 

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#### Abstract

Metasurfaces can be used to precisely control polarization state of the scattered light. Here, we present a metasurface-based terahertz device. On the one hand, it serves as a high performance linear polarization converter in transmission of over $80 \%$ power with weak reflection. It is capable of rotating linear polarization orientation angle with respect to $x$-axis continuously from $-90^{\circ}$ to $+90^{\circ}$ at 0.84 THz . On the other hand, it serves as a circular polarizer. It can transform a linear polarized wave into a circular polarized wave at 2.49 THz . The transmitted and reflected field are both circular polarized with $50 \%$ power. The proposed device with dual functionalities can be applied to modulate the polarization state of the signal in THz wireless communication.


Keywords: linear to linar polarization; linear to circular polarization; polarizer; polarization manipulation

## 1. Introduction

The terahertz frequency range $(0.1-10 \mathrm{THz})$ could form the basis of the next generation of high-speed wireless communication, as well as a range of spectroscopy, sensing, and imaging applications [1-4]. Spatiotemporal control of the far- and near-field distribution of terahertz signals, and the manipulation of terahertz fields will be crucial in enabling these applications.

Metasurfaces are two-dimensional surfaces with precisely designed scatterers that create controlled field transformation of incident wavefronts across the properties of amplitude [5,6], phase [7-9], frequency [10,11], and polarization [12-15]. Polarization is one of the basic properties of electromagnetic waves conveying valuable information in signal transmission and sensitive measurements. Cross-polarization conversion has been demonstrated using 2D vertical-shape metasurface [16,17]. V-shape metasurfaces with different sizes are arranged together to transform linearly polarized waves into circularly polarized waves and work for any orientation of the incident linear polarization [18]. Rod-shape metasurface and a metal ground plane form a Fabry-Pérot-like cavity to result in cross-polarized scattering in reflection [19]. Stacked rod-shaped arrays with a tailored rotational twist can manipulate circular polarization in broadband spectrum [9,20]. However, the currently available designs suffer from lack of continuous modulation and low effective power.

In this article, we report a metasurface comprised of circular-shape split-ring resonator (SRR) that offer continuous adjustment of polarization orientation angle and realize linear to circular polarization. Circular-shape SRR, as shown in Figure 1a, is simultaneous control of phase and amplitude of the incident light, which are separately controlled by opening angle $\alpha$ and rotation angle, $\gamma$. All models and simulated calculations are carried out in CST Microwave Studio 2020, and circular-shape SRR is made of Au. For a linearly polarized incident beam, the polarization orientation angle $\theta$ is continuously varied from $-90^{\circ}$ to $+90^{\circ}$ by controlling the rotation angle $\gamma$ at 0.84 THz and the incident linear polarization is transformed into circular polarization by controlling the opening angle $\alpha$ at 2.49 THz . High effective power of $80 \%$ is achieved by means of stacking circular-shape SRR.


Figure 1. (a) Schematic of circular-shaped SRR with opening angle $\alpha$, radius $r$, width $w$ and rotation angle $\theta$ with respect to $x$-axis. (b) Simulated transmission amplitude and (c) phase of an array of circular-shape SRR with $(r, \alpha, w)=\left(34 \mu \mathrm{~m}, 45^{\circ}, 5 \mu \mathrm{~m}\right)$ and $\gamma$ varying from $-45^{\circ}$ to $45^{\circ}$ at 0.84 THz .
(d) Simulated reflection amplitude. The period of the array is $80 \mu \mathrm{~m}$ in both $x$ - and $y$-direction.

## 2. Linear Polarization Orientation Angle Modulation

We choose incident light $\vec{E}_{i n}(\mathrm{z}, t)=\cos (2 \pi f t+k \mathrm{z}) \vec{e}_{x}$ propagating along-z-axis direction, where $f$ is the frequency of the incident light, $t$ is the time, and $k$ is the wavenumber, whereas the scattered light has a component which is polarized orthogonal to that of the incident light and excited by the symmetric and antisymmetric current modes [6,17,21]. The scattered light can be expressed as:

$$
\begin{equation*}
\vec{E}_{\mathrm{scat}}(\mathrm{z}, t)=E_{x}(\mathbf{z}, t) \vec{e}_{x}+E_{y}(\mathbf{z}, t) \stackrel{\rightharpoonup}{e}_{y} \tag{1}
\end{equation*}
$$

where $E_{x}(\mathrm{z}, t)=\left|t_{x x}\right| \cos \left(2 \pi f t+k z+\varphi_{x x}\right), E_{y}(\mathrm{z}, t)=\left|t_{y x}\right| \cos \left(2 \pi f t+k z+\varphi_{y x}\right)$. $t_{x x}=\left|t_{x x}\right| e^{j \varphi_{x x}}$ and $t_{y x}=\left|t_{y x}\right| e^{j \varphi_{y x}}$ are the transmission coefficient for $x$ - and $y$-components, respectively. The scattered wave is linear polarized wave if $\varphi_{x x}-\varphi_{y x}=0^{\circ}, \pm 180^{\circ}$. Polarization orientation angle is $\theta=\arctan \left(\left|t_{y x}\right| /\left|t_{x x}\right|\right)$. The scattered power of interest is determined by the amplitude of the scattered wave, $\left|\vec{E}_{s c a t}\right|^{2}=\left|t_{x x}\right|^{2}+\left|t_{y x}\right|^{2}$. Circularshape SRR metasurface is able to change transmission coefficient continuously and meet the phase condition by rotating $\gamma$. Figure $1 \mathrm{~b}, \mathrm{c}$ show the simulated transmission amplitude and phase variations of an array of circular-shape SRR at different $\gamma$ from $-45^{\circ}$ to $+45^{\circ}$ at $0.84 \mathrm{THz} .\left|t_{y x}\right|$ is approximate to vary as $\sin (2 \theta) / 2$ while the phase remains constant in two separated ranges of $-45^{\circ}$ to $0^{\circ}$ and $0^{\circ}$ to $+45^{\circ}$, with an abrupt change of $180^{\circ}$ at $\gamma=0^{\circ}$. $\left|t_{x x}\right|$ follows a $\cos (4 \theta / 3)$ dependence while the phase remains constant. Thus, the rotation angle $\gamma$ serves as one important parameter to control the amplitude of a scattered wave without resorting to new geometry.

Figure $2 \mathrm{a}-\mathrm{g}$ show the variation of electric field of scattered light in time domain at $z=-10 \mu \mathrm{~m}$. The scattered light is linearly polarized along $x$-axis with polarization orientation angle $\theta$ varying from $-45^{\circ}$ to $+45^{\circ}$. The polarization orientation angle $\theta$ is proportional to the rotation angle $\gamma, \theta / \gamma \approx 1$. The larger polarization orientation angle is, the lower amplitude and effective power of the scattered light is, as is shown in Figure 2 h . The scattered power of interest is higher than $80 \%$ at $\theta=0^{\circ}$ with low reflections (Figure 1d) and less than $50 \%$ at $\theta= \pm 45^{\circ}$ with high reflections (Figure 1d). So, the single-layered metasurface produced weak scattered beams with the increase of $\gamma$ and with increased reflections.


Figure 2. Variation of electric field of scattered light in time domain at $z=-10 \mu \mathrm{~m}$ for (a) $\gamma=+15^{\circ}$ and (b) $\gamma=-15^{\circ}$, (c) $\gamma=+30^{\circ}$ and (d) $\gamma=-30^{\circ}$, (e) $\gamma=+45^{\circ}$ and (f) $\gamma=-45^{\circ}$, (g) $\gamma=0^{\circ}$. (h) Electric field vector of the scattered light with $\gamma$ varying from $-45^{\circ}$ to $+45^{\circ}$ at 0.84 THz .

In order to improve the power of scattered light in transmission with reduced reflections, we use two metasurfaces to result in much improved performance. An example consisting of two-separated circular-shape SRRs is shown in Figure 3a, which are separated by air of thickness $d=50 \mu \mathrm{~m}$. This corresponds to a separation of $\lambda_{0} / 8$ at 0.84 THz . The rotation angle of layer 1 is set as $\gamma=0^{\circ}$.


Figure 3. (a) Two-layer circular-shape SRR. (b) Electric field vector variation of the scattered light at 0.84 THz with rotation angle $\gamma_{1}$ varying from $-45^{\circ}$ to $+45^{\circ}$, where the red and blue line represent the single-layer and two-layer metasurface, respectively.

The scattered light from two-layered circular-shape SRRs is linearly polarized with respect $x$-axis. Polarization orientation angle is controlled by rotation angle $\gamma_{1}$ of layer 2 . The electric field vector of the scattered light is shown in Figure 3b with $\gamma_{1}$ varying from $-45^{\circ}$ to $+45^{\circ}$. The polarization orientation angle of two-layered circular-shape SRRs (blue solid line in Figure 3b) is continuously varied from $-45^{\circ}$ to $+45^{\circ}$ with the same scattered power of more than $80 \%$ at 0.84 THz . Compared with the single-layer circular-shape SRR (red solid line in Figure 3b), two-layer circular-shape SRRs have better impedance matching to reduce reflections with the increase of $\gamma_{1}$, so the scattered power of interest is enhanced. So far, continuous polarization orientation angle modulation from $-45^{\circ}$ to $+45^{\circ}$ with high effective power of more than $80 \%$ is achieved by stacking two-layer circular-shaped SRRs. However, two-layered structure cannot result in polarization orientation angle of $-90^{\circ}$ to $-45^{\circ}$ and $+45^{\circ}$ to $+90^{\circ}$ with effective power of more than $80 \%$. When $\gamma_{1}$ exceeds $\pm 45^{\circ}$, the scattered power of interest decreases significantly. One way to overcome this problem is to use a three-layer structure.

Figure $4 b, c$ show two different styles for three-layer structures. Style I is used to continuously control polarization orientation angle from $+45^{\circ}$ to $+90^{\circ}$ by rotating $\gamma_{2}$ of layer 3 from $+45^{\circ}$ to $+90^{\circ}$ (black solid line in Figure 4a). Style II is used for control of polarization orientation angle from $-45^{\circ}$ to $-90^{\circ}$ by varying $\gamma_{2}$ of layer 3 from $-45^{\circ}$ to $-90^{\circ}$ (red solid line in Figure 4a). The scattered power of interest remains constant of more than $80 \%$. Figure 5 summarizes the simulated results of linear polarization orientation angle modulation by using circular-shape SRR. Linear polarization orientation angle with respect to $x$-axis is modulated continuously from $-90^{\circ}$ to $+90^{\circ}$ at 0.84 THz and high effective power of more than $80 \%$ are obtained for different polarization orientation angle.


Figure 4. (a) Electric field vector of the scattered light at 0.84 THz with rotation angle $\gamma_{2}$ varying from $+45^{\circ}$ to $+90^{\circ}$ (red solid line) and $-45^{\circ}$ to $-90^{\circ}$ (black solid line). Schematic of the three-layer structure: (b) Style I with $\gamma_{1}=+45^{\circ}$ and (c) Style II with $\gamma_{1}=-45^{\circ}$.


Figure 5. Linear polarization orientation angle modulation and its corresponding circular-shape SRR structure.

## 3. Circular Polarization Modulation

In the above section, linear polarization orientation angle modulation is demonstrated by the amplitude transformation ability through circular-shape SRR metasurface. $\left|t_{x x}\right|$ and $\left|t_{y x}\right|$ are continuously controlled by rotation angle $\gamma$ while maintaining almost the same phase or out of phase. In this section, the phase transformation ability through circularshape SRR metasurface is demonstrated by condsidering circular polarization. We choose the incident polarization to be at $45^{\circ}$ with respect to $x$-axis and $\vec{E}_{\text {in }}$ can be decomposed into two components, a vertical component $\vec{E}_{\perp}$ and a horizontal component $\vec{E}_{\|}$, as is shown in Figure $6 \mathrm{a} . t_{\perp}=\left|t_{\perp}\right| e^{j \varphi_{\perp}}$ and $t_{\|}=\left|t_{\|}\right| e^{j \varphi_{\|}}$are transmission coefficients for vertical and horizontal component respectively. $\varphi_{\perp}$ and $\varphi_{\|}$are continuously controlled by opening angle $\alpha$ while maintaining similar amplitude, as is shown in Figure $6 \mathrm{~b}, \mathrm{c}$. When $\alpha=55.5^{\circ}$, $\varphi_{\perp}-\varphi_{\|}$will be $-90^{\circ}$ at 2.49 THz . The scattered light is a circular polarization wave with effective power of $50 \%$, as is shown in Figure 6d. Another $50 \%$ of the power is reflected in circularly polarized waves [22]. So the transmitted and reflected wave share the same power and wave form.


Figure 6. (a) Schematic of circular-shaped SRR to realize linear to circular polarization with rotation angle $\gamma=0^{\circ}$. Simulated transmission (b) phase and (c) amplitude for vertical and horizontal components with opening angle $\alpha$ varying from $1^{\circ}$ to $100^{\circ}$. (d) Variation of electric field of scattered light in time domain at 2.49 THz , at $\mathrm{z}=-10 \mu \mathrm{~m}$.

In order to improve effective power of interest by decreasing the reflected power, the two-layered circular-shape SRR circular polarizer is designed, as is shown in Figure 7a.

Two-layer circular-shape SRRs are the same geometry and spaced by air of thickness $d_{1}=42.3 \mu \mathrm{~m}$. The scattered light is circular polarized at 2.07 THz (blue solid line in Figure 7b). Compared with a single layer polarizer (red solid line in Figure 7b), two-layered polarizer has an effective power of more than $80 \%$ with reduced reflections, which can be seen from Figure 7 b . The amplitude of the transmitted circularly polarized wave is larger than $80 \%$.


Figure 7. (a) Schematic of circular-shape SRR circular polarizer with $(r, \alpha, w, \gamma)=\left(39 \mu \mathrm{~m}, 60^{\circ}\right.$, $1 \mu \mathrm{~m}, 0^{\circ}$ ) and $d_{1}=42.3 \mu \mathrm{~m}$. (b) Variation of electric field of scattered light in time domain at $z=-10 \mu \mathrm{~m}$, where the red and blue line represent the single-layer and two-layered circular-shape SRR polarizers, respectively.

## 4. Conclusions

We have reported a circular-shape SRR with amplitude and phase transformation ability. The rotation angle $\gamma$ can control the scattered amplitude while maintaining almost the same phase. The opening angle $\alpha$ can control the scattered phase while maintaining a similar amplitude. Continuous linear polarization orientation angle modulation from $-90^{\circ}$ to $+90^{\circ}$ is realized by the amplitude transformation ability. Linear to circular polarization modulation is realized by the phase transformation ability. In order to improve the effective power of interest, we use a two- or three-layer structure to result in effective power up to $80 \%$.

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## References

1. Nagatsuma, T.; Ducournau, G.; Renaud, C.C. Advances in terahertz communications accelerated by photonics. Nat. Photonics 2016, 10, 371-379. [CrossRef]
2. Sengupta, K.; Nagatsuma, T.; Mittleman, D.M. Terahertz integrated electronic and hybrid electronic-photonic systems. Nat. Electron. 2018, 1, 622-635. [CrossRef]
3. Al Hadi, R.; Sherry, H.; Grzyb, J.; Zhao, Y.; Forster, W.; Keller, H.M.; Cathelin, A.; Kaiser, A. 1 k-Pixel Video Camera for 0.7-1.1 Terahertz Imaging Applications in 65-nm CMOS. IEEE J. Solid-State Circuits 2012, 47, 2999-3012. [CrossRef]
4. Sengupta, K. Integrated Circuits for Terahertz Communication Beyond 100 GHz : Are We There Yet? In Proceedings of the 2019 IEEE International Conference on Communications Workshops (ICC Workshops), Shanghai, China, 20-24 May 2019.
5. Jiang, Z.H.; Yun, S.; Lin, L.; Bossard, J.A.; Werner, D.H.; Mayer, T.S. Tailoring Dispersion for Broadband Low-loss Optical Metamaterials Using Deep-subwavelength Inclusions. Sci. Rep. 2017, 3, 1571. [CrossRef]
6. Liu, L.; Zhang, X.; Kenney, M.; Su, X.; Xu, N.; Ouyang, C.; Shi, Y.; Han, J.; Zhang, W.; Zhang, S. Broadband Metasurfaces with Simultaneous Control of Phase and Amplitude. Adv. Mater. 2014, 26, 5031-5036. [CrossRef]
7. Sun, S.; He, Q.; Xiao, S.; Xu, Q.; Li, X.; Zhou, L. Gradient-index meta-surfaces as a bridge linking propagating waves and surface waves. Nat. Mater. 2012, 11, 426-431. [CrossRef]
8. Sun, S.; Yang, K.Y.; Wang, C.M.; Juan, T.K.; Tsai, D.P. High-Efficiency Broadband Anomalous Reflection by Gradient Meta-Surfaces. Nano Lett. 2012, 12, 6223-6229. [CrossRef]
9. Huang, L.; Chen, X.; Mühlenbernd, H.; Li, G.; Bai, B.; Tan, Q.; Jin, G.; Zentgraf, T.; Zhang, S. Dispersionless Phase Discontinuities for Controlling Light Propagation. Nano Lett. 2012, 12, 5750-5755. [CrossRef]
10. Huang, F.M.; Zheludev, N.I. Super-resolution without evanescent waves. Nano Lett. 2009, 9, 1249-1254. [CrossRef]
11. Biagioni, P.; Savoini, M.; Huang, J.S.; Duò, L.; Finazzi, M.; Hecht, B. Near-field polarization shaping by a near-resonant plasmonic cross antenna. Phys. Rev. B 2009, 80, 153409. [CrossRef]
12. Chen, W.T.; Zhu, A.Y.; Capasso, F. Flat optics with dispersion-engineered metasurfaces. Nat. Rev. Mater. 2020, 5, 604-620. [CrossRef]
13. Chen, X.; Huang, L.; Mühlenbernd, H.; Li, G.; Bai, B.; Tan, Q.; Jin, G.; Qiu, C.W.; Zhang, S.; Zentgraf, T. Dual-polarity plasmonic metalens for visible light. Nat. Commun. 2012, 3, 1198. [CrossRef] [PubMed]
14. Kang, M.; Chen, J.; Wang, X.; Wang, H. Twisted vector field from an inhomogeneous and anisotropic metamaterial. J. Opt. Soc. Am. B Opt. Phys. 2012, 29, 572-576. [CrossRef]
15. Yu, N.; Capasso, F. Flat optics with designer metasurfaces. Nat. Mater. 2014, 13, 139-150. [CrossRef]
16. Ni, X.; Emani, N.K.; Kildishev, A.V.; Boltasseva, A.; Shalaev, V.M. Broadband Light Bending with Plasmonic Nanoantennas. Science 2012, 335, 427. [CrossRef]
17. Yu, N.; Genevet, P.; Kats, M.A.; Aieta, F.; Tetienne, J.-P.; Capasso, F.; Gaburro, Z. Light Propagation with Phase Discontinuities: Generalized Laws of Reflection and Refraction. Science 2011, 334, 333-337. [CrossRef]
18. Yu, N.; Aieta, F.; Genevet, P.; Kats, M.A.; Gaburro, Z.; Capasso, F. A Broadband, Background-Free Quarter-Wave Plate Based on Plasmonic Metasurfaces. Nano Lett. 2012, 12, 6328-6333. [CrossRef]
19. Grady, N.K.; Heyes, J.E.; Chowdhury, D.R.; Zeng, Y.; Reiten, M.T.; Azad, A.K.; Taylor, A.J.; Dalvit DA, R.; Chen, H.-T. Terahertz Metamaterials for Linear Polarization Conversion and Anomalous Refraction. Science 2013, 340, 1304-1307. [CrossRef]
20. Zhao, Y.; Belkin, M.A.; Alù, A. Twisted optical metamaterials for planarized ultrathin broadband circular polarizers. Nat. Commun. 2012, 3, 870. [CrossRef]
21. Zhang, X.; Tian, Z.; Yue, W.; Gu, J.; Zhang, S.; Han, J.; Zhang, W. Broadband Terahertz Wave Deflection Based on Circular-shape Complex Metamaterials with Phase Discontinuities. Adv. Mater. 2013, 25, 4567-4572. [CrossRef]
22. Wu, S.; Zha, D.; He, Y.; Miao, L.; Jiang, J. Design of stable and high-efficiency graphene-based polarizers for oblique angle of incidence in terahertz frequency. Appl. Opt. 2019, 58, 492. [CrossRef] [PubMed]
