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# Surface acidity of tin dioxide nanomaterials revealed with $^{31}\text{P}$ solid-state NMR spectroscopy and DFT calculations†

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Tin dioxide ( $\text{SnO}_2$ ) nanomaterials are important acid catalysts. It is therefore crucial to obtain details about the surface acidic properties in order to develop structure–property relationships. Herein, we apply  $^{31}\text{P}$  solid-state NMR spectroscopy combined with a trimethylphosphine (TMP) probe molecule, to study the facet-dependent acidity of  $\text{SnO}_2$  nanosheets and nanoshuttles. With the help of density functional theory calculations, we show that the tin cations exposed on the surfaces are Lewis acid sites and their acid strengths rely on surface geometries. As a result, the (001), (101), (110), and (100) facets can be differentiated by the  $^{31}\text{P}$  NMR shifts of adsorbed TMP molecules, and their fractions in different nanomaterials can be extracted according to deconvoluted  $^{31}\text{P}$  NMR resonances. The results provide new insights on nanosized oxide acid catalysts.

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## 1. Introduction

With a relatively large band gap of 3.6 eV, tin dioxide ( $\text{SnO}_2$ ) is an intrinsic n-type semiconductor.<sup>1,2</sup>  $\text{SnO}_2$  and  $\text{SnO}_2$  based-materials have been widely used in a variety of fields including gas-sensing,<sup>3</sup> photovoltaics,<sup>4</sup> energy storage,<sup>2,5</sup> and acid catalysis.<sup>6,7</sup> In particular,  $\text{SnO}_2$  nanomaterials, which are associated with a larger surface area and may expose specific facets, exhibit improved catalytic properties compared to their bulk counterparts.<sup>8,9</sup> Therefore, it is important to understand the details about the acidic properties of  $\text{SnO}_2$  nanomaterials in order to develop better  $\text{SnO}_2$ -based catalysts.

The most widely applied methods for determining the acidic properties of solids include  $\text{NH}_3$ -temperature programmed desorption (TPD)<sup>10</sup> and pyridine-infrared spectroscopy (IR).<sup>11</sup>  $\text{NH}_3$ -TPD data gives general information on the relative concentration and strength of different acid sites, however, they cannot differentiate Lewis and Brønsted acid sites.<sup>12</sup> Although pyridine-IR can be used to identify the type of acid sites, it is impossible to obtain quantitative results (*i.e.*, the

concentrations of Lewis and Brønsted acid sites), owing to the overlapping band if both acid sites are present.<sup>13</sup> Therefore, a method that is able to provide both qualitative and quantitative information on Lewis and Brønsted acid sites is preferable. Solid-state nuclear magnetic resonance (NMR) spectroscopy is a powerful method with which to study detailed local structure and properties of solids.<sup>14–21</sup> Combined with appropriate probe molecules such as trimethylphosphine (TMP), the smallest alkyl phosphine that can be bound the acid sites of oxide surfaces, solid-state NMR data can provide rich information on the location, type, strength as well as concentration of acid sites.<sup>22–24</sup> This approach has been extensively used to study metal oxide based catalysts and porous catalytic materials, such as zeolites and mesoporous materials.<sup>25,26</sup>

Despite the importance of exposed facets on the catalytic properties,<sup>27</sup> only very recently, this method was extended to study the facet-dependent acidity of oxide nanomaterials, including  $\text{TiO}_2$ ,<sup>28</sup>  $\text{ZnO}$ ,<sup>29</sup> and  $\text{CeO}_2$ .<sup>13</sup> Based on  $^{31}\text{P}$  chemical shift of TMP molecules adsorbed, different surface species at each facet can be identified and their fractions can be extracted according to spectral deconvolution.<sup>29,30</sup> Since the facet dependent acidity was not studied for  $\text{SnO}_2$  nanostructures, in this work, we adopt NMR spectroscopy combined with DFT calculation to study the surface acid sites of two  $\text{SnO}_2$  nanostructures, *i.e.*, nanosheets and nanoshuttles, and obtain facet-dependent acidity information.  $^{31}\text{P}$  NMR spectra show that  $\text{SnO}_2$  nanosheets possess stronger Lewis acidity than nanoshuttles, which is supported with DFT calculation results on adsorption energies.

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## 2. Experimental section

### 2.1 Synthesis of SnO<sub>2</sub> nanosheets

In a typical synthesis, 225.6 mg SnCl<sub>2</sub>·2H<sub>2</sub>O was added into 80 mL ethylenediamine, and the mixture was vigorously stirred for 30 minutes before it was transferred into a 100 mL Teflon-lined autoclave. The autoclave was sealed and heated at 180 °C in an oven for 24 h. After the autoclave was cooled down to room temperature, the yellow precipitate was washed with distilled water and ethanol for five times, and then dried in a vacuum oven at 60 °C overnight to obtain SnO<sub>2</sub> nanosheets.<sup>31</sup>

### 2.2 Synthesis of SnO<sub>2</sub> nanoshuttles

Typically, 285.2 mg SnCl<sub>2</sub>·2H<sub>2</sub>O was added into a mixture of 40 mL ethanol and 40 mL ammonia solution (pH ≈ 11). The solution was stirred vigorously for 1 h before it was transferred into a 100 mL Teflon-lined autoclave. The autoclave was sealed tightly and then heated at 120 °C in an oven for 6 h. SnO<sub>2</sub> nanoshuttles were obtained after washing and drying the yellow solids with similar procedures used for SnO<sub>2</sub> nanosheets.<sup>32</sup>

### 2.3 Characterization

X-ray diffraction (XRD) was operated on a Philips X'pert Pro diffractometer with Ni filter and Cu K $\alpha$  irradiation ( $\lambda = 0.15418$  nm) at 40 mA and 40 kV. The  $2\theta$  scanning range was 10–80°. High resolution transmission electron microscopy (HRTEM) images were recorded at 200 kV on a JEOL JEM-2010 instrument. Fourier transform infrared (FT-IR) spectra were obtained with a Thermo Nicolet 8700 FT-IR spectrometer with diamond ATR accessory. X-ray photoelectron spectra (XPS) were collected on an ESCALAB 250 instrument with Al K $\alpha$  irradiation ( $h\nu = 1486.6$  eV) as the excitation source. The Brunauer–Emmett–Teller (BET) surface areas of the samples were determined from nitrogen adsorption and desorption data at 77 K using a Micromeritics ASAP 2000 instrument. The contents of Na and N/C were determined with an Optima 5300DV inductively coupled plasma mass spectrometer (ICP-MS) and a Heraeus CHN-0-Rapid elemental analyser, respectively.

### 2.4 TMP adsorption

In a typical process of TMP adsorption, 200 mg sample was placed in a glass tube. The tube was sealed and then connected to a vacuum line. The sample was first activated under  $1 \times 10^{-3}$  torr at 100 °C for 3 h before 40 mbar of TMP was introduced. The tube was kept at room temperature for 0.5 h, before it was evacuated for another 0.5 h to remove physisorbed TMP molecules.<sup>24</sup> The sample was then transferred into a N<sub>2</sub>-filled glove box and packed zirconia rotors for NMR tests.

### 2.5 <sup>31</sup>P solid-state NMR spectroscopy

<sup>31</sup>P solid state NMR experiments were performed on a Bruker Avance III 400 MHz spectrometer equipped with a 9.4 T superconducting magnet and a double-tuned 4.0 mm magic angle spinning (MAS) probe at room temperature (25 °C). The Larmor

frequencies of <sup>1</sup>H and <sup>31</sup>P nuclei were 400.13 and 161.98 MHz, respectively. A single pulse sequence with <sup>1</sup>H decoupling was used for data acquisition with a spinning speed of 14 kHz. A short excitation pulse of 2.6  $\mu$ s, corresponding to 90 degree flip angle and an optimized recycle delay of 6.0 s were used. <sup>31</sup>P chemical shifts were referenced to ammonium dihydrogen phosphate (ADP) at 0.81 ppm.

### 2.6 Computational details

Density functional theory (DFT) calculations were performed with the Vienna *ab initio* simulation package (VASP).<sup>33,34</sup> The Perdew–Burke–Ernzerhof (PBE)<sup>35</sup> functional was applied to treat the electron exchange and correlation effects by the generalized gradient approximation (GGA).<sup>36</sup> The electron–ion interaction was evaluated by the projector augmented plane wave (PAW) method with a frozen-core approximation.<sup>37</sup> The Brillouin zone integration was approximated by a sum over special  $k$ -points using the Monkhorst–Pack grids.<sup>38</sup> The Gaussian smearing method with a value of 0.1 eV and a cutoff energy of 450 eV was used for all calculations. The Kohn–Sham equations were solved self-consistently. The energy tolerance was set to  $1.0 \times 10^{-4}$  eV and the maximum Hellmann Feynman force tolerance was set to 0.05 eV  $\text{\AA}^{-1}$  for structural optimization.

In calculating molecule adsorption, we established the 6 layers-(001), 4 layers-(101), 4 layers-(110) and 7 layers-(100) surface of SnO<sub>2</sub> (Fig. S1†). Except (110) surface ( $p(2 \times 1)$  super cells),  $p(2 \times 2)$  super cells were applied for the other surfaces. For sampling the Brillouin zone,  $3 \times 3 \times 1$ ,  $3 \times 4 \times 1$ ,  $2 \times 3 \times 1$  and  $4 \times 4 \times 1$  Monkhorst–Pack  $k$ -point grid were used for the (001), (100), (101) and (110) surfaces, respectively. To avoid interactions between slabs, all of the above slab models are established with the vacuum space of 12  $\text{\AA}$ .

The adsorption energies of the TMP molecules were calculated as follows:

$$E_{\text{ads}} = E_{\text{TMP/SnO}_2} - (E_{\text{TMP}} + E_{\text{SnO}_2}),$$

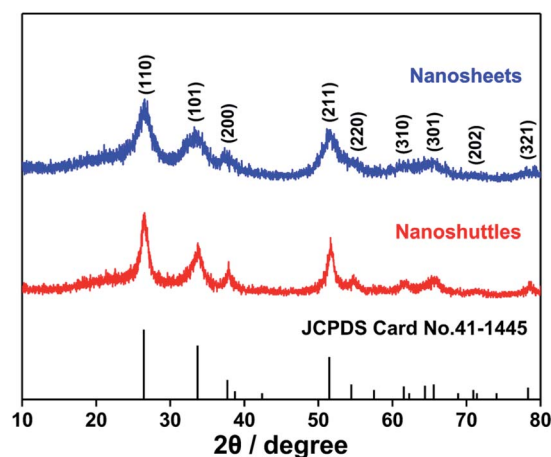


Fig. 1 XRD patterns of SnO<sub>2</sub> nanosheets and nanoshuttles.

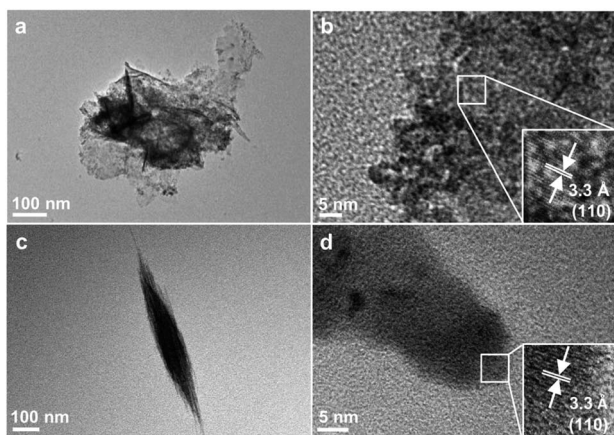


Fig. 2 TEM images of (a and b) SnO<sub>2</sub> nanosheets and (c and d) nanoshuttles.

where  $E_{\text{TMP}/\text{SnO}_2}$  is the total energies of the adsorption model,  $E_{\text{SnO}_2}$  and  $E_{\text{TMP}}$  are the energy of slab and free molecules, respectively.

### 3. Results and discussion

#### 3.1 Basic characterization

The X-ray diffraction (XRD) patterns of both materials prepared hydrothermally (Fig. 1) can be indexed to rutile phase SnO<sub>2</sub> with a tetragonal structure (JCPDS no. 41-1445). The broad widths diffraction peaks indicate their grain sizes are in nanoscale.

TEM images in Fig. 2a and c show that the two SnO<sub>2</sub> samples are nano-sized sheets and shuttles, respectively. The

SnO<sub>2</sub> nanosheets are either stretched or aggregated (Fig. 2a), and an inter-planar spacing of 0.33 nm is readily observed in the HRTEM image (Fig. 2b), corresponding to the (110) surface of rutile SnO<sub>2</sub>, which is the most stable low-index surface.<sup>32</sup> The shuttle-like SnO<sub>2</sub> (Fig. 2c) is about 700 nm long with a maximum diameter of approx. 100 nm and the lattice spacing of (110) surface can be found at the end of the shuttle. It is also expected that other common facets, such as (001), (101) and (100), may also be present on the two SnO<sub>2</sub> nanostructures.<sup>39,40</sup>

The FT-IR transmission spectra of both SnO<sub>2</sub> nanosheets and nanoshuttles in Fig. 3a show six bands stemming from the vibrations and overtones of Sn–O, Sn–O–Sn and surface hydroxyl groups, respectively.<sup>41</sup> The most intense peak centered at around 590 cm<sup>-1</sup> is owing to the stretching vibration of Sn–O–Sn. Similarly, another stretching vibration peak at approx. 3380 cm<sup>-1</sup> corresponds to surface hydroxyl groups. The intensity of this band, however, is significantly weaker than SnO<sub>2</sub> nanoparticles previously reported,<sup>42</sup> suggesting that the concentration of surface Sn–OH is very low. Quantitative <sup>1</sup>H NMR data (Fig. S2 and Table S2†) confirmed that the coverages of proton on SnO<sub>2</sub> nanosheets and nanoshuttles were 1.84 and 1.44%. The influence of hydrogen species is therefore neglected. The peaks at ~970 and 2420 cm<sup>-1</sup> are assigned to the overtones of the SnO<sub>2</sub> lattice and Sn–OH configuration, respectively. Other bands appear at 1240 and 1640 cm<sup>-1</sup> respectively belong to the bending vibrations of Sn–OH and molecular water. In addition to these FT-IR peaks relevant to SnO<sub>2</sub>, no other signals are observed. X-ray photoelectron spectra (XPS) were also recorded to study the surface states of SnO<sub>2</sub> nanosheets and nanoshuttles. As shown in Fig. 3b, the detected

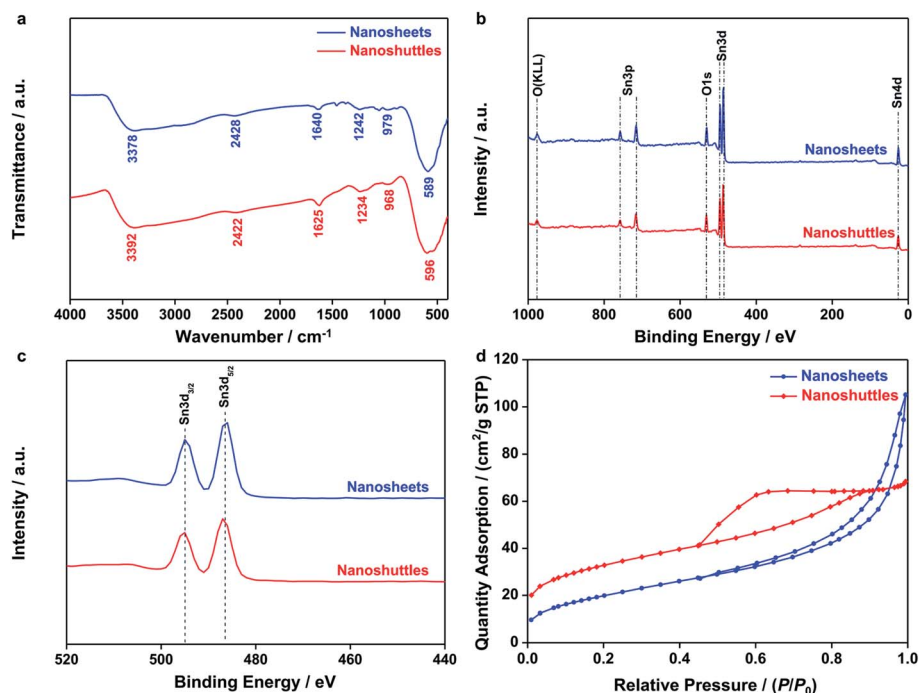


Fig. 3 (a) FT-IR spectra, (b and c) XPS spectra and (d) N<sub>2</sub> adsorption/desorption isotherms of SnO<sub>2</sub> nanosheets and nanoshuttles.

elements in XPS survey spectrum only included Sn (Sn 3p, Sn 3d, and Sn 4d), O (O KLL and O 1s) and C (C 1s).<sup>43</sup> In Fig. 3c, two splitting peaks attributed to Sn 3d<sub>3/2</sub> (494.96 eV) and Sn 3d<sub>5/2</sub> (486.48 eV) exhibit a difference of 8.48 eV, consistent with the spin-orbit coupling value of Sn<sup>4+</sup>.<sup>44</sup> Their symmetric peaks indicate that there is few surface Sn<sup>2+</sup>, or surface oxygen vacancy.<sup>45</sup> In addition, the tiny C 1s signals are related to carbon contamination according to previous publications. Elemental analysis in Table S1† also confirms that the amounts of the possible impurities in SnO<sub>2</sub> nanosheets and nanoshuttles are small and therefore, their influences can be neglected. The Brunauer-Emmett-Teller (BET) surface areas were measured as 74.1 and 117.0 m<sup>2</sup> g<sup>-1</sup>, respectively, for nanosheets and nanoshuttles (Fig. 3d). The H3-type hysteric loop in the curve of SnO<sub>2</sub> nanosheets verifies the existence of flake particles which

formed slit pores with a wide range of pore-size distribution (Fig. S3†) from 8 to 20 nm. The BET curve of SnO<sub>2</sub> nanoshuttles displays a H2-type hysteric loop and its pore-size distribution (Fig. S3†) ranges from 5 to 25 nm.

### 3.2 <sup>31</sup>P NMR spectroscopy and DFT calculations

The SnO<sub>2</sub> nanosheets and nanoshuttles were then adsorbed with TMP molecules and investigated with <sup>31</sup>P solid state NMR spectroscopy (Fig. 4a). In general, physisorbed TMP molecules lead to a sharp signal at around -60 ppm, while TMP molecules adsorbed at Brønsted and Lewis acid sites generate distinct <sup>31</sup>P peaks at -5 to -2 and -10 to -55 ppm, respectively.<sup>25</sup> No NMR resonance is observed at -60 ppm for both samples, indicating there is no physisorbed TMP molecule on the surface. The

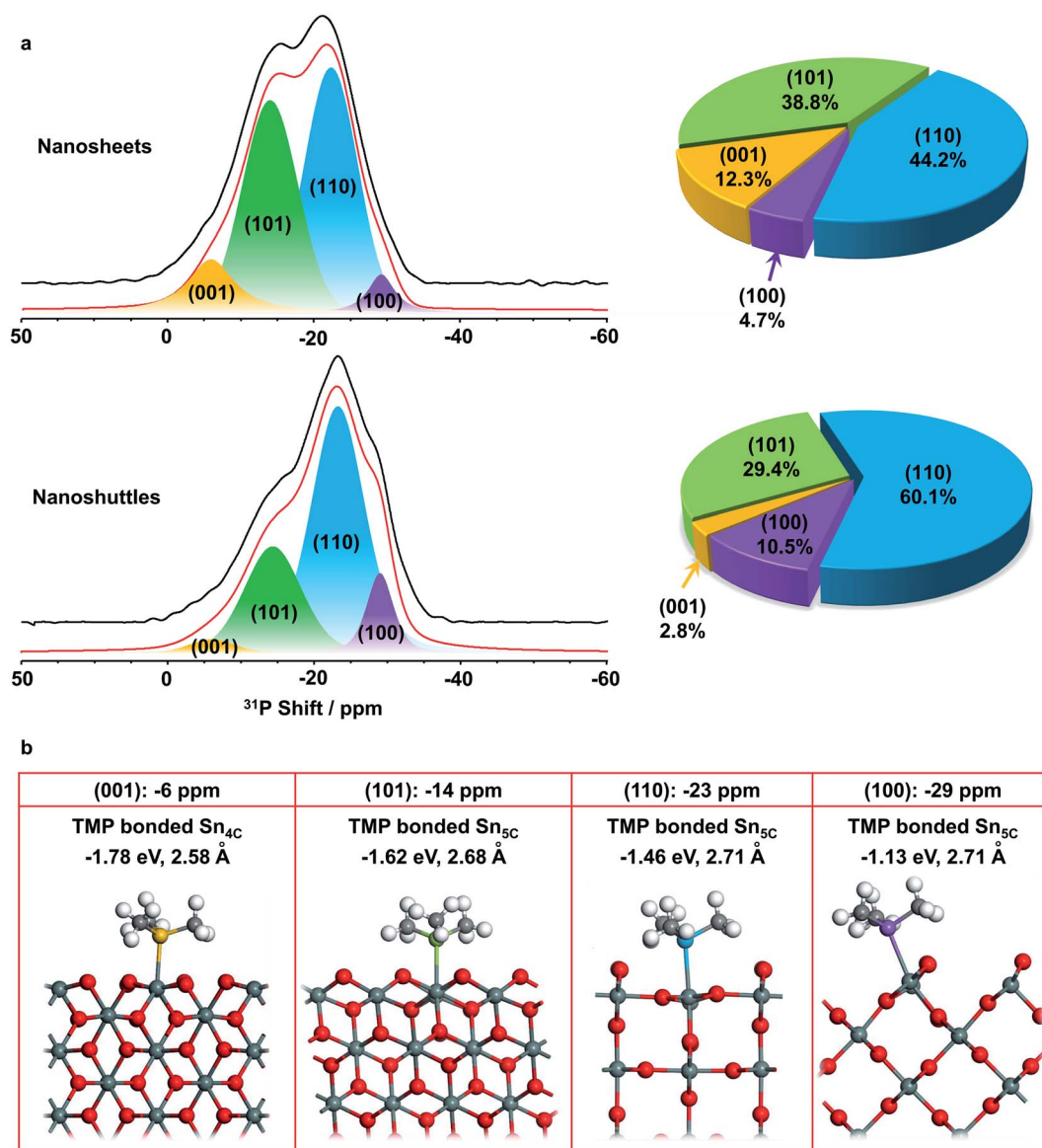


Fig. 4 (a) <sup>31</sup>P solid-state NMR spectra of SnO<sub>2</sub> nanosheets and nanoshuttles adsorbed with TMP (left) and extracted fractions of four low-index facets (right). (b) Surface structures of four facets ((001), (101), (110) and (100)) adsorbed with TMP, corresponding adsorption energy and internuclear distances (P-Sn).



signals are quite broad in both spectra, with the majority of the intensity coming from around  $-6$  to  $-34$  ppm, implying Lewis acid sites dominate the surface of both  $\text{SnO}_2$  nanosheets and nanoshuttles, which should arise from exposed surface Sn cations.<sup>46</sup> The centers of gravity of the spectra of nanosheets and nanoshuttles are at approx.  $-18.1$  and  $-21.6$  ppm, respectively. Since the higher the shift, the stronger the Lewis acidity, the above observation suggests that the Lewis acidity is stronger in  $\text{SnO}_2$  nanosheets than nanoshuttles.

Spectral deconvolution was performed to further explore more structure and properties of the two nanomaterials. In both data, two major peaks are observed at around  $-14$  and  $-23$  ppm, while there are additional weak peaks at approx.  $-6$  and  $-29$  ppm (Fig. 4a and Table S2†), indicating four different acid sites can be distinguished. The resonances at lower frequencies ( $-14$ ,  $-23$  and  $-29$  ppm) can be readily assigned to three different Lewis acid sites, while the peak at  $-6$  ppm is between the frequency range of common Lewis and Brønsted acid sites. Previous pyridine-IR studies show that only Lewis acid sites are present on the surface of  $\text{SnO}_2$ , thus this weak peak is assigned to another Lewis acid site.<sup>11</sup>

$^{31}\text{P}$  NMR studies on other metal oxide nanostructures clearly show that  $^{31}\text{P}$  chemical shift is dependent on the facet TMP adsorbed.<sup>28–30</sup> Therefore, the four peaks may correspond to TMP bound to Lewis acid sites, which are expected to be Sn cations, at different facets of  $\text{SnO}_2$  nanostructures. (001), (101), (100), and (110) surfaces are the commonly observed low-index facets for  $\text{SnO}_2$ <sup>39,40</sup> and the DFT calculations on the TMP adsorption energy were performed (Fig. 4b). Fig. S4† provides the charge density difference of acid sites. The calculation results confirm that TMP molecules are adsorbed on the surface Sn cations in all of the four facets. TMP adsorbed on the 4-coordinated  $\text{Sn}^{4+}$  ( $\text{Sn}_{4\text{C}}$ ) on (001) surface is associated with the largest adsorption energy of  $-1.78$  eV, followed by TMP adsorbed on  $\text{Sn}_{5\text{C}}$  on (101), (110), and (100) surfaces with energies of  $-1.62$ ,  $-1.46$ , and  $-1.13$  eV, respectively. The internuclear distance between P and Sn also increases with the same order (Fig. 4b). Since a stronger adsorption for TMP gives rise to a more positive  $^{31}\text{P}$  chemical shift,<sup>16</sup> it is reasonable to attribute the resonances at around  $-6$ ,  $-14$ ,  $-23$  and  $-29$  ppm to TMP adsorbed on the tin cations on (001), (101), (110), and (100) surfaces, respectively. It has also been shown that  $^{31}\text{P}$  chemical shift exhibits a linear correlation with the adsorption energies.<sup>29,47–49</sup> Using physisorbed TMP molecules on other oxides ( $^{31}\text{P}$  chemical shift is  $-62.2$  ppm and adsorption energy is  $0$ )<sup>50</sup> as a reference, the experimentally observed  $^{31}\text{P}$  chemical shifts are plotted against the calculated adsorption energies (Fig. 5). Clearly, there is a strong linear correlation, which can be described by the equation:  $y = -62.9 - 30.3x$ , where  $x$  and  $y$  are adsorption energy (eV) and  $^{31}\text{P}$  chemical shift (ppm), respectively. Therefore, this result further supports our spectral assignments.

In both samples, the (110) facet has the largest fraction, which is in agreement with the understanding that (110) is often the most stable surface in rutile structure.<sup>39</sup> In contrast, the fraction of (001) facet in  $\text{SnO}_2$  nanosheets is much higher than nanoshuttles, which can be related to the sample preparation procedure. Ethylenediamine, which was used in

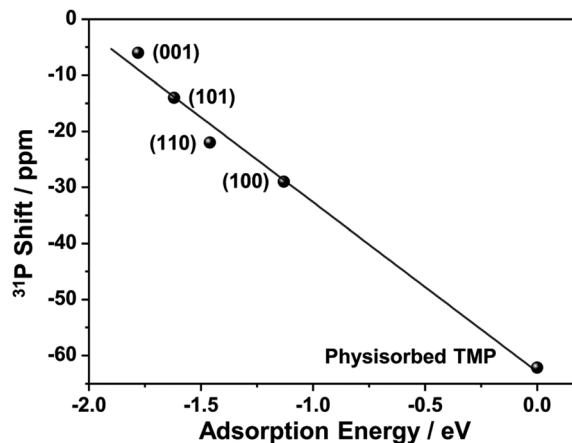


Fig. 5 Correlations of observed  $^{31}\text{P}$  chemical shifts and calculated adsorption energies.

preparing  $\text{SnO}_2$  nanosheets, prefers to adsorb on  $\text{SnO}_2$  (001) surfaces,<sup>31</sup> prompting the formation of more high energy (001) surfaces in  $\text{SnO}_2$  nanosheets. The total fraction of (001) and (101) facets, which have stronger Lewis acidity, is higher than 50% for  $\text{SnO}_2$  nanosheets, while this value is only around 30% for  $\text{SnO}_2$  nanoshuttles. Such difference is expected to make a difference in acid catalysis.

## 4. Conclusion

Facet dependent acidity of  $\text{SnO}_2$  nanosheets and nanoshuttles were investigated with TMP probe molecule assisted  $^{31}\text{P}$  solid-state NMR spectroscopy combined with DFT calculations for the first time. Strong Lewis acid sites dominate the surface of  $\text{SnO}_2$  and four different sites arising from different facets can be distinguished. The fractions of (001) and (101) facets, associated with stronger Lewis acidity, are larger in  $\text{SnO}_2$  nanosheets than  $\text{SnO}_2$  nanoshuttles, which can be ascribed to the different conditions in sample preparation. These results should shed light on the rational design of better nanosized oxide acid catalysts.

## Author contributions

W. Zhang, Z. Lin, J. Chen and L. Peng, contributed the design of the experiment, analysis of data and writing of the paper. Z. Lin, and H. Li carried out synthesis and characterization of the samples. W. Zhang, Y. Wen, M. Xu, Y. Wang, X. Ke and X. Xia performed the solid state NMR experiments. F. Wang carried out the DFT calculation. All the authors have given their approval to the final version of the manuscript.

## Conflicts of interest

Authors declare no competing interests.

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