

# Theoretical Study of Efficient Photon–Phonon Resonance Absorption in the Tungsten-Related Vibrational Mode of Scheelite

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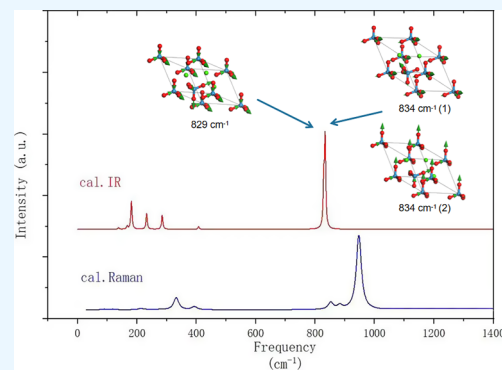


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**ABSTRACT:** Tungsten (W) is an extremely rare and vital metal extensively used in metallurgy, the chemical industry, optoelectronic devices, and machinery manufacturing. In this work, an environmentally friendly and efficient physical method based on photon–phonon resonance absorption (PPRA) is proposed for separating W from scheelite. We calculated the vibrational spectrum of calcium tungstate ( $\text{CaWO}_4$ ) and assigned the infrared (IR) absorption and Raman scattering peaks through a dynamic analysis of the normal modes. We focused on the strong IR absorption peaks related to W and identified three high-intensity IR-active modes at around  $830\text{ cm}^{-1}$ , corresponding to the stretching of the W–O bonds. Therefore, we propose the use of high-power terahertz ( $\sim 25\text{ THz}$ ) laser radiation to facilitate W extraction from compounds, leveraging the high efficiency of PPRA. Experimental testing is required to determine the precise absorption frequency under industrial production conditions.



## 1. INTRODUCTION

Tungsten (W) is a rare and essential element and has the highest melting point among the refractory metals. Moreover, it possesses excellent high-temperature strength and good corrosion resistance to molten alkali.<sup>1</sup> W is extensively applied in several fields, including metallurgy, the chemical industry, optoelectronic devices, and the machinery industry.<sup>2</sup> In particular, W has vital applications in quantum electronics, medical treatments, and military product manufacturing.<sup>3–6</sup> With the development of machinery, automobile, national defense, and electronics industries in China and the continuous progress of hard-alloy and W processing technology, the consumption of raw W in China has grown rapidly. Between 1990 and 2019, consumption increased from 10,000 tons to 48,000 tons, representing an average annual growth rate of 5.5%. The global demand for W is expected to reach 110,000 tons by 2025.<sup>7,8</sup> However, the W content in the earth's crust is relatively low, 0.001%. In 2016, the global reserve of W was 3.1 million tons with a basic reserve at 6.2 million tons, making W an extremely rare resource. China's W reserve was 1.9 million tons with a basic reserve of 4.2 million tons. China possesses 61% of the world's W reserves, establishing it as the country with the largest W reserve globally.

Tungsten is predominantly found in various types of rocks in the forms of wolframite and scheelite, making it challenging to extract. The primary component of scheelite is calcium tungstate ( $\text{CaWO}_4$ ), a critical and commonly used material in W production. Consequently, we conducted a spectral analysis of  $\text{CaWO}_4$  to identify an efficient method for W separation. Typically, W can be separated and extracted

through organic solvent extraction and ion exchange methods.<sup>9,10</sup> In China, the primary methods for smelting wolframite involve NaOH leaching and  $\text{Na}_2\text{CO}_3$  sintering following water leaching, while alkali dissolution following hydrochloric acid decomposition is commonly employed for scheelite smelting. Additional techniques such as hot ball milling,<sup>11</sup> alkali pressure boiling,<sup>12,13</sup> and sulfur and phosphorus mixed acid leaching<sup>14,15</sup> have been developed. W recovery methods include zinc melting<sup>15</sup> and selective electrochemical dissolution.<sup>16</sup> Although numerous W extraction methods exist, they often lead to W pollution in air, water, and soil. Soil contamination with W is a significant concern, as W can interact with the environmental medium, posing a threat to both the ecological environment and human health.<sup>17,18</sup> In China, research on W pollution characteristics and prevention technologies remains relatively limited. Consequently, there is an urgent need for a novel, environmentally friendly technology for extracting W from ore with minimal pollution and higher efficiency suitable for large-scale industrial production.

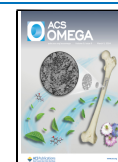
Numerous spectroscopic studies have focused on  $\text{CaWO}_4$ , a crucial and common material in W production.<sup>19–21</sup> Basiev et al. studied  $\text{CaWO}_4$  via Raman spectroscopy.<sup>22</sup> Ryu et al.

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measured the Raman spectrum and ultraviolet–visible spectrum of  $\text{CaWO}_4$  nanoparticles collected from a colloidal suspension.<sup>3</sup> Additionally, Golubovic et al. grew  $\text{CaWO}_4$  and  $\text{Nd}:\text{CaWO}_4$  (0.8% Nd) single crystals from melts through the Czochralski technique and measured the infrared (IR) spectra of  $\text{Nd}:\text{CaWO}_4$  and  $\text{CaWO}_4$  single crystals at both room and liquid nitrogen temperatures.<sup>23</sup>

Despite the aforementioned experimental and spectroscopic studies of  $\text{CaWO}_4$ , there remains a lack of theoretical research on its lattice dynamic processes and the assignment of IR and Raman spectra peaks for  $\text{CaWO}_4$ . In this study, we simulated the vibrational spectrum of  $\text{CaWO}_4$  and compared it with the experimental spectrum to accurately assign the Raman and IR peaks. It is worth noting that the Raman spectrum we calculated was in the IR region so it was nonresonant Raman spectrum. Based on the IR-active peaks related to W, we propose a new method to enhance the separation process of W from oxide, namely photon–phonon resonance absorption (PPRA).<sup>24,25</sup>

## 2. SIMULATION METHOD

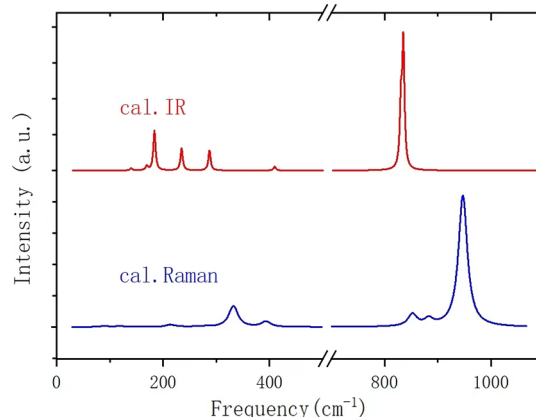
$\text{CaWO}_4$  has a crystal structure of the scheelite type with  $C_{4h}$  point-group symmetry and  $I4_1/a$  space-group symmetry. In this structure,  $\text{Ca}^{2+}$  ions occupy 4b Wyckoff positions (0, 1/4, and 5/8) with local symmetry, while  $\text{W}^{6+}$  ions occupy 4a Wyckoff positions (0, 1/4, and 1/8) with the same local symmetry.  $\text{W}^{6+}$  ions are tetrahedrally coordinated to  $\text{O}^{2-}$  ions, with a coordination number of 4, while  $\text{Ca}^{2+}$  ions are in distorted hexahedral coordination with  $\text{O}^{2-}$  ions, with a coordination number of 8.<sup>26</sup> Vibrational spectral simulations of the  $\text{CaWO}_4$  crystal were performed using a primitive cell that contained two  $\text{CaWO}_4$  molecules.

We performed geometry optimization and phonon calculations using first-principle density-functional theory (DFT) and the Cambridge Serial Total Energy Package (CASTEP)<sup>27</sup> code. To account for electron density fluctuations in  $\text{CaWO}_4$ , we adopted the generalized gradient approximation of the Perdew–Burke–Ernzerhof (PBE) exchange correlation functional, which was semilocalized and provided more accurate descriptions of energy and structure. It exhibited better consistency with experimental results compared with other functionals.<sup>28</sup> The convergence tolerance values for energy and the self-consistent field were set to  $1 \times 10^{-9}$  eV/atom. An energy cutoff of 800 eV and a  $3 \times 3 \times 4$   $k$ -point mesh were applied. Furthermore, because of the large number of electrons in 5d metal W, we employed norm-conserving pseudopotential, which packages the core electrons and the nucleus together to balance the accuracy and the time, and finite displacement was used for phonon calculations. Additionally, optical ( $\omega = \infty$ ) and dc ( $\omega = 0$ ) dielectric permittivity, optical ( $\omega = \infty$ ), and static ( $\omega = 0$ ) molecular polarizability were calculated in this work, allowing us to obtain simulated IR and Raman spectra. According to the calculated data, we assigned the vibrational spectra through a dynamic process analysis of each wavenumber. Consequently, we were able to identify W-related peaks in these spectra and determine the highest IR absorption peaks in terms of W-related PPRA.

## 3. RESULTS AND DISCUSSION

There are 12 atoms in one  $\text{CaWO}_4$  primitive cell and 33 (i.e.,  $[12 \times 3] - 3$ , subtracting 3 for the acoustic normal modes) optical normal modes, which can be examined through IR

absorption and Raman scattering via phonon–photon coupling. Owing to the symmetric structure of  $\text{CaWO}_4$ , changes in the dipole moment and polarizability of the primitive cell are mutually repulsive, resulting in entirely complementary IR-active and Raman-active modes. Figure 1 illustrates the simulated IR and Raman spectra.



**Figure 1.** Simulated IR absorption and Raman scattering spectra.  $M_p = 1 \times 1.0$ . The main peaks are almost complementary.

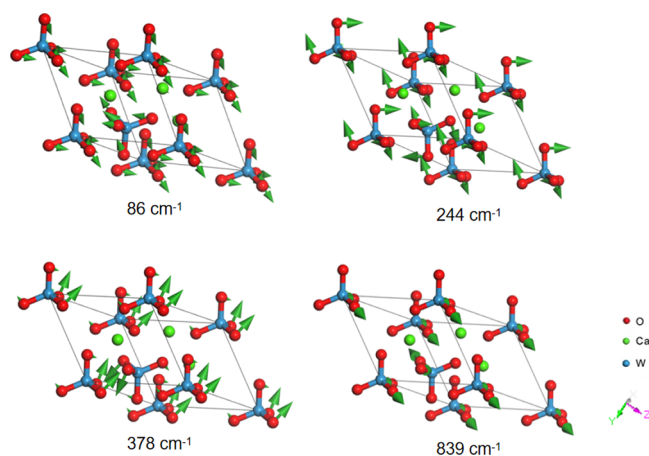
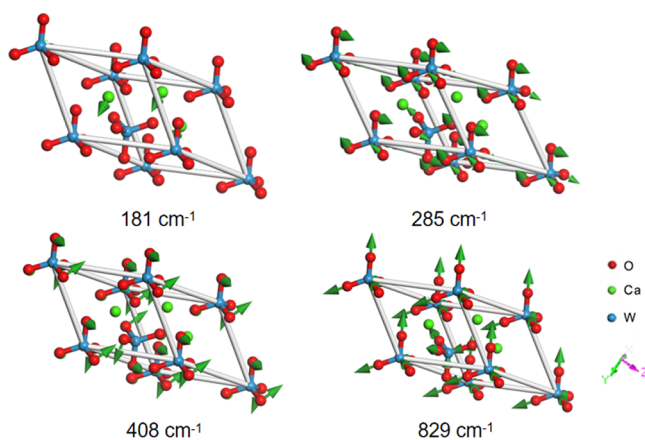
Table 1 compares the calculated wavenumbers and the experimental IR and Raman peaks. The assignments for the 33 vibrational modes are given in the last column. We identify 18 Raman-active vibrational modes, including four wavenumbers (109, 178, 244, and  $839 \text{ cm}^{-1}$ ) that are doubly degenerate, and 12 IR-active vibrational modes, which also feature four wavenumbers (138, 181, 285, and  $834 \text{ cm}^{-1}$ ) that are doubly degenerate. However, the modes at 254, 410, and  $942 \text{ cm}^{-1}$  exhibit neither IR nor Raman activity. According to group theory, the 36 degrees of freedom of a  $\text{CaWO}_4$  primitive cell are divided into 26 species of vibrational modes  $3A_g + 5A_u + 5B_g + 3B_u + 5E_g + 5E_u$ , where the species E vibrations are doubly degenerate. Among these vibrational modes,  $3A_g$ ,  $5B_g$ , and  $5E_g$  are Raman-active modes;  $4A_u$  and  $4E_u$  are IR-active modes;  $3B_u$  are silent modes; and  $1A_u$  and  $1E_u$  are acoustic modes.<sup>20</sup> Compared with the reported experimental results, the simulated peaks show a small blue shift (about  $10 \text{ cm}^{-1}$ ). This can validate the theoretical methodology employed in our work.

In the lowest-vibrational-energy region, we assigned the three vibrational modes at 86 and  $109 \text{ cm}^{-1}$  to  $\text{WO}_4^{2-}$  translations in which the two anions vibrate relatively along three orthogonal directions. The two vibrational modes at  $109 \text{ cm}^{-1}$  are degenerate. The mode at  $86 \text{ cm}^{-1}$  is shown in Figure 2 as an example of the  $\text{WO}_4^{2-}$  translation. The anions at the eight apex angles represent one anion in a primitive cell and the other one is in the primitive cell. The two degenerate modes at  $138 \text{ cm}^{-1}$  belong to the  $(\text{WO}_4)^{2-}$  rotation. Lim detected peaks at  $115\text{--}211 \text{ cm}^{-1}$  and assigned them to interior vibrations, in which the two ionic groups vibrated together.<sup>29</sup> Porto and Scott reported that the  $E_u$  mode at  $125\text{--}150 \text{ cm}^{-1}$  was related to interior vibration, and the  $E_u$  mode at  $150\text{--}200 \text{ cm}^{-1}$  was associated with rotation.<sup>20</sup> The tentative assignments of the modes were based on group theory, whereas our assignments are based on first-principle calculations.

We assigned the six modes at 167, 178, 181, and  $204 \text{ cm}^{-1}$  to  $\text{Ca}^{2+}$  translations. We present the mode at  $181 \text{ cm}^{-1}$  as a typical  $\text{Ca}^{2+}$  translation of the IR activity (Figure 3). Ryu et al.

**Table 1.** Comparisons of 33 Calculated Wavenumbers (IR and Raman Activities Are Listed) of CaWO<sub>4</sub> with Experimental Data (Unit cm<sup>-1</sup>)

wavenumbers	activity	IR exp.	Raman exp.	assignment	wavenumbers	activity	IR exp.	Raman exp.	assignment
86	Raman		84 <sup>a</sup> , 86 <sup>b</sup> , 84 <sup>c</sup>	WO <sub>4</sub> <sup>2-</sup> translation	285	IR			WO <sub>4</sub> <sup>2-</sup> rotation
109	Raman		117 <sup>a</sup> , 118 <sup>b</sup> , 117 <sup>c</sup>	WO <sub>4</sub> <sup>2-</sup> translation	315	Raman			WO <sub>4</sub> <sup>2-</sup> bending
109	Raman			WO <sub>4</sub> <sup>2-</sup> translation	316	Raman		336 <sup>a</sup> , 334 <sup>b</sup> , 334 <sup>c</sup>	WO <sub>4</sub> <sup>2-</sup> bending
138	IR	143 <sup>a</sup> , 143 <sup>b</sup>		WO <sub>4</sub> <sup>2-</sup> rotation	374	Raman			WO <sub>4</sub> <sup>2-</sup> bending
138	IR			WO <sub>4</sub> <sup>2-</sup> rotation	378	Raman		403 <sup>b</sup> , 401 <sup>c</sup>	WO <sub>4</sub> <sup>2-</sup> bending
167	IR			Ca <sup>2+</sup> translation	378	Raman			WO <sub>4</sub> <sup>2-</sup> bending
178	Raman		180 <sup>b</sup> , 185 <sup>d</sup>	Ca <sup>2+</sup> translation	408	IR	409 <sup>a</sup> , 435 <sup>b</sup> , 439 <sup>c</sup> , 433 <sup>d</sup>		WO <sub>4</sub> <sup>2-</sup> bending
178	Raman			Ca <sup>2+</sup> translation	410				WO <sub>4</sub> <sup>2-</sup> bending
181	IR	180 <sup>a</sup> , 180 <sup>b</sup>		Ca <sup>2+</sup> translation	829	IR	797 <sup>a</sup> , 793 <sup>b</sup> , 754 <sup>c</sup>		WO <sub>4</sub> <sup>2-</sup> stretching
181	IR			Ca <sup>2+</sup> translation	834	IR	822 <sup>c</sup> , 862 <sup>d</sup>		WO <sub>4</sub> <sup>2-</sup> stretching
195	Raman		195 <sup>a</sup> , 196 <sup>b</sup>	WO <sub>4</sub> <sup>2-</sup> rotation	834	IR			WO <sub>4</sub> <sup>2-</sup> stretching
204	Raman		210 <sup>a</sup> , 210 <sup>b</sup> , 212 <sup>c</sup> , 210 <sup>d</sup>	Ca <sup>2+</sup> translation	839	Raman		838 <sup>a</sup> , 838 <sup>b</sup> , 839 <sup>c</sup>	WO <sub>4</sub> <sup>2-</sup> stretching
232	IR	237 <sup>a</sup> , 237 <sup>b</sup>		WO <sub>4</sub> <sup>2-</sup> rotation	839	Raman			WO <sub>4</sub> <sup>2-</sup> stretching
244	Raman			WO <sub>4</sub> <sup>2-</sup> rotation	873	Raman			WO <sub>4</sub> <sup>2-</sup> stretching
244	Raman			WO <sub>4</sub> <sup>2-</sup> rotation	939	Raman		912 <sup>a</sup> , 922 <sup>b</sup> , 912 <sup>c</sup>	WO <sub>4</sub> <sup>2-</sup> stretching
254				WO <sub>4</sub> <sup>2-</sup> rotation	942				WO <sub>4</sub> <sup>2-</sup> stretching
285	IR	275 <sup>a</sup> , 309 <sup>b</sup>		WO <sub>4</sub> <sup>2-</sup> rotation					

<sup>a</sup>Ref 20. <sup>b</sup>Ref 30. <sup>c</sup>Ref 34. <sup>d</sup>Ref 26.**Figure 2.** Four typical Raman-active vibrational modes. The mode of wavenumber at 86 cm<sup>-1</sup> represents anion translation, in which the two anion groups in one primitive cell vibrate in opposite directions. The mode at 244 cm<sup>-1</sup> represents WO<sub>4</sub><sup>2-</sup> rotation, the mode at 378 cm<sup>-1</sup> represents WO<sub>4</sub><sup>2-</sup> bending, and the mode at 839 cm<sup>-1</sup> represents O–W stretching.**Figure 3.** Four typical IR-active vibrational modes. The mode at 181 cm<sup>-1</sup> corresponds to Ca<sup>2+</sup> translation, the mode at 285 cm<sup>-1</sup> corresponds to WO<sub>4</sub><sup>2-</sup> rotation, the mode at 408 cm<sup>-1</sup> corresponds to WO<sub>4</sub><sup>2-</sup> bending, and the mode at 829 cm<sup>-1</sup> corresponds to O–W stretching.

suggested that the bands at 185 cm<sup>-1</sup> are related to the translational mode of the Ca–O bond in CaWO<sub>4</sub>.<sup>3</sup> Russell assigned modes with frequencies below 435 cm<sup>-1</sup> to vibrations between the loosely bound Ca<sup>2+</sup> and WO<sub>4</sub><sup>2-</sup> ions.<sup>30</sup> Zhang et al. reported that the Raman peaks in the range of 100–250

cm<sup>-1</sup> corresponded to the contraction and bending vibrations of the CaO<sub>8</sub> groups.<sup>31</sup>

We propose that the six modes at 232–285 cm<sup>-1</sup> correspond to WO<sub>4</sub><sup>2-</sup> rotations. Additionally, seven modes at 315–410 cm<sup>-1</sup> correspond to WO<sub>4</sub><sup>2-</sup> bending vibrations. The modes at 244 and 378 cm<sup>-1</sup> are demonstrated in Figure 2, while the modes at 285 and 408 cm<sup>-1</sup> are shown in Figure 3. Ryu et al. assigned the bands at 210 cm<sup>-1</sup> to the translational



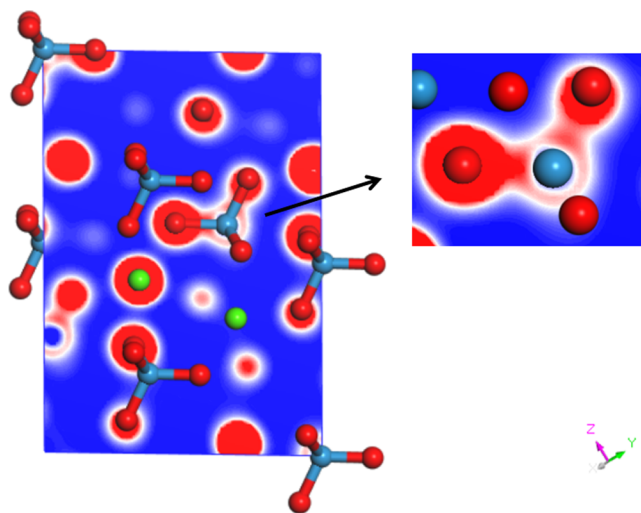
mode of the anion group.<sup>3</sup> However, we disagree with their assignment because the anion group has a greater mass than  $\text{Ca}^{2+}$ . Lv et al. suggested that vibrational peaks within the range of  $350\text{--}600\text{ cm}^{-1}$  were related to W–O stretching vibrations.<sup>32</sup> However, as shown in Table 1, the lowest wavenumber for W–O stretching is  $829\text{ cm}^{-1}$ . Zhang et al. reported that Raman peaks at  $334$  and  $405\text{ cm}^{-1}$  corresponded to the symmetric and nonsymmetric bending vibrations of  $\text{WO}_4^{2-}$ , respectively, which aligns with our calculations.<sup>31</sup> Furthermore, Abozaid et al. examined the IR spectroscopy of single crystals of calcium tungstate doped with neodymium ( $\text{CaWO}_4:\text{Nd}^{3+}$ ). They reported a weak vibration band at  $436\text{ cm}^{-1}$ , related to W–O bending vibration, and a weak but sharp band at  $433\text{ cm}^{-1}$ , related to the Ca–O bond.<sup>26</sup> Similarly, Phuruangrat et al. detected a weak W–O vibration at  $433\text{ cm}^{-1}$ .<sup>33</sup>

In the high-frequency vibrational region, eight modes at  $829\text{--}942\text{ cm}^{-1}$  correspond to intra-anion  $\text{WO}_4^{2-}$  stretching. These kinds of vibrations are related to the O–W bond stretching. Two examples of modes at  $829$  (IR-active) and  $839\text{ cm}^{-1}$  (Raman-active) are presented in Figures 3 and 2, respectively. Ryu et al. reported that the  $\nu_3$  ( $B_g$ ) vibration was located around  $840\text{ cm}^{-1}$ , and the  $\nu_3$  ( $E_g$ ) vibration was found at  $800\text{ cm}^{-1}$ .<sup>3</sup> Russell assigned modes with wavenumbers of  $778\text{--}922\text{ cm}^{-1}$  to the internal vibrations of tightly bound  $\text{WO}_4^{2-}$  ions.<sup>30</sup> Furthermore, Zhang et al. reported that Raman peaks at  $910$ ,  $836$ , and  $797\text{ cm}^{-1}$  corresponded to the stretching vibration of  $\text{WO}_4^{2-}$ , and the mode at  $910\text{ cm}^{-1}$  was related to the symmetric stretching vibration of  $\text{WO}_4^{2-}$ .<sup>31</sup> Abozaid et al. reported an IR absorption band in  $\text{CaWO}_4:\text{Nd}^{3+}$  at around  $742\text{ cm}^{-1}$  and a strong band at  $862\text{ cm}^{-1}$ , attributing them to the O–W–O stretching.<sup>26</sup> Similarly, Phuruangrat et al. detected strong antisymmetric W–O stretching in  $\text{WO}_4^{2-}$  tetrahedrons in the range of  $711\text{--}933\text{ cm}^{-1}$ .<sup>33</sup> Although the region above  $800\text{ cm}^{-1}$  is associated with W–O stretching, the intensities of IR and Raman photon–phonon coupling are different owing to different physical mechanisms.

According to the calculation results, three IR-active modes at  $829$  and  $833\text{ cm}^{-1}$  exhibit the strongest intensities, as illustrated in Figure 1, while the largest Raman peak is at  $939\text{ cm}^{-1}$ . For the dynamic processes of the three IR-active modes, see the Supporting Information Movies. Considering that IR absorption is due to the PPRA effect, the experimental peaks at around  $830\text{ cm}^{-1}$  feature the most efficient photon–phonon absorption. Two experimental peaks occur at approximately  $822$  and  $862\text{ cm}^{-1}$ .<sup>26,34</sup> Accurate IR absorption experiments in this region will be necessary in the future. Applying a high-power terahertz laser to a  $\text{CaWO}_4$  crystal will endow the phonons of the W–O vibrational mode with high energy-absorption efficiency, which may facilitate chemical bond breakage and W separation from  $\text{CaWO}_4$ . This could offer a novel physical method for W separation without chemical pollution. Figure 4 presents the charge distribution between W and O. The oxygen shows a quite electronegative with a small charge distribution around W. This shows that the W–O bond possesses a high property of ionic bond, which reduces the strength of the W–O bond. Therefore, it is possible to achieve a W–O bond break through PPRA.

#### 4. CONCLUSIONS

According to first-principle DFT simulations, we analyzed the dynamic processes of the vibrational modes of  $\text{CaWO}_4$  and identified Raman-active and IR-active vibrational modes. New



**Figure 4.** Electron density at a certain slice. The red color represents a large electron density, and the blue color represents a small electron density.

assignments that were not present in experimental data are presented.

In our analysis of vibrational modes, we focused on the W-related vibrational modes. Experiments and simulations revealed that three W–O stretching vibrational modes with wavenumbers around  $830\text{ cm}^{-1}$  corresponded to high-intensity experimental peaks, which indicates efficient absorption in IR spectroscopy. Given that the physical mechanism of IR spectroscopy is PPRA, we propose the use of a high-power terahertz laser radiation (at  $\sim 25\text{ THz}$ ) on the  $\text{CaWO}_4$  crystals. The resonance absorption is expected to facilitate the breaking of W–O bonds. This may offer an environmentally friendly and effective method for separating W from scheelite. Additional experimental testing is required to determine the precise frequency at which the high-power terahertz laser should operate when it is applied to  $\text{CaWO}_4$  crystals.

#### ■ ASSOCIATED CONTENT

##### Supporting Information

The Supporting Information is available free of charge at <https://pubs.acs.org/doi/10.1021/acsomega.3c08636>.

- IR-active vibrational mode at  $829\text{ cm}^{-1}$  (MP4)
- IR-active vibrational mode-1 at  $834\text{ cm}^{-1}$  (MP4)
- IR-active vibrational mode-2 at  $834\text{ cm}^{-1}$  (MP4)

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## Notes

The authors declare no competing financial interest.

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