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Breakup of a long-period comet as the origin of the dinosaur extinction

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The origin of the Chicxulub impactor, which is attributed as the cause of the K/T mass extinction event, is an unsolved puzzle. The background impact rates of main-belt asteroids and long-period comets have been previously dismissed as being too low to explain the Chicxulub impact event. Here, we show that a fraction of long-period comets are tidally disrupted after passing close to the Sun, each producing a collection of smaller fragments that cross the orbit of Earth. This population could increase the impact rate of long-period comets capable of producing Chicxulub impact events by an order of magnitude. This new rate would be consistent with the age of the Chicxulub impact crater, thereby providing a satisfactory explanation for the origin of the impactor. Our hypothesis explains the composition of the largest confirmed impact crater in Earth's history as well as the largest one within the last million years. It predicts a larger proportion of impactors with carbonaceous chondritic compositions than would be expected from meteorite falls of main-belt asteroids.

Strong evidence suggests that the Chicxulub impact led to the K/T mass extinction event, which was the largest in the past ~ 250 Myr and brought about the demise of the dinosaurs^{1,2}. However, the nature of the Chicxulub impactor is poorly understood. The latest scenario suggested postulated that the breakup of the Baptistina asteroid family could have led to the formation of the Chicxulub impactor³. However, spectroscopic follow-up indicated that the Baptistina family has an S-type, rather than an Xc-type composition, making it an unlikely source of the Chicxulub impactor, which had a carbonaceous chondritic composition^{4–6}, although not ruling out entirely the possibility due to the stochastic nature of asteroid collisions and the subsequent disruptive processes⁷. Observations of the Baptistina family also suggested that the breakup age may be ~ 80 Myr⁸ rather than ~ 160 Myr³, further reducing the likelihood that the Baptistina breakup formed the Chicxulub impactor.

The Chicxulub impactor could have originated from the background populations of asteroids or of comets. Main-belt asteroids (MBAs) with diameters $D \gtrsim 10$ km, capable of producing Chicxulub impact events, strike the Earth once per ~ 350 Myr^{9,10}. Based on meteorite fall statistics¹¹, one such object with a carbonaceous chondritic composition impacts the Earth over a characteristic timescale of ~ 3.5 Gyr, too rare to account for the K/T event³. Long-period comets (LPCs) capable of producing Chicxulub-scale impacts strike Earth also too rarely, once per $\sim 3.8–11$ Gyr³, based on the rate of Earth-crossing LPCs and the impact probability per perihelion passage^{12,13}, and adopting a cumulative power-law index within the range -2.0 to -2.7 ^{14–16}. The only cometary sample-return mission to date, *Stardust*, found that Comet 81P/Wild 2 had carbonaceous chondritic composition, suggesting that such a composition could potentially be widespread in comets^{17–20}. As a result, the rate of LPC impacts with carbonaceous chondritic composition could be similar to the overall LPC impact rate. Within a timescale of ~ 100 Myr, stellar encounters could boost the impactor flux by an order of magnitude for a Myr timescale²¹, which are insufficient in magnitude to explain a Chicxulub impact event. We note that comets are typically more fragile and porous than asteroids^{22,23}.

To find the fraction of LPCs with orbital behavior that could affect the impact flux at Earth, we simulated gravitational interactions between LPCs and the Jupiter–Earth–Sun system using a semi-analytic approach. Initially, there are N Jupiter-crossing LPCs (initial pericenter distance $q \lesssim 5.2$ AU) with semi-major axis $a \sim 10^4$ AU and the distribution of pericenter distances scaling as q^2 , the corresponding cross-sectional area^{21,24}. The initial inclination distribution is taken as uniform^{21,24}. We then follow the orbital perturbation prescription for a restricted three-body scattering²⁵. At the initial closest approach to Jupiter, calculated by selecting a random phase angle in Jupiter's orbit and computing the minimum distance between Jupiter and the LPC's orbit b_j , the change in semi-major axis a resulting from the three-body interaction is computed

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as $\Delta(1/a) = (4M_J v_J \sqrt{a} (\cos \gamma + K \cos \delta) / M_\odot^{3/2} b_J \sqrt{G(1 + K^2)})$, where M_J is the mass of Jupiter, M_\odot is the mass of the Sun, v_J is the heliocentric orbital speed of Jupiter, G is the gravitational constant, γ is the angle between the velocity vectors of Jupiter and the LPC, δ is the angle between the normal in the orbital plane to the approach of the LPC at the time of its closest approach to Jupiter and the velocity vector of Jupiter, and $K \equiv (GM_J a / M_\odot b_J)$. The new inclination is approximated by the numerically derived fitting function, $\approx \arccos[\cos i - 0.38 \sin i^2 Q^{-1/2} (b_J/a)]$, where $Q \equiv (q/b_J)$. The updated eccentricity is calculated through conservation of the Tisserand parameter, $T = (1/a) + 2\sqrt{a(1 - e^2)} \cos i$, across the encounter. If the LPC crosses the orbit of Earth, defined as $q \lesssim 1$ AU, the same process of updating the orbital is repeated for the closest encounter with the Earth, for a random Earth phase angle. We consider LPCs with $a > 2 \times 10^5$ AU or $e \geq 1$ to be ejected and remove them from the simulation as well as any that collide with Jupiter, the Sun, or the Earth. Tidal disruption by Jupiter is similar in likelihood to collision with Jupiter, $\sim 10^{-8}$ per Jupiter-crossing orbit.

We find that for $N = 10^5$ particles, $\sim 20\%$ of Earth-crossing events, defined as perihelia within the orbital radius of the Earth $q \lesssim 1$ AU), were immediately preceded by perihelia within the Roche radius of the Sun, $q \lesssim r_\odot (2\rho_\odot / \rho_{obj})^{1/3}$, where r_\odot is the radius of the Sun, ρ_\odot is the mean mass density of the Sun, and $\rho_{obj} \sim 0.7 \text{ g cm}^{-3}$ is the mean density²⁶ of the LPC, since they were captured into highly eccentric orbits by interacting with the Sun-Jupiter system. This is consistent with previous estimates of the sungrazing LPC population²⁷. If the LPC is solely bound by gravity, then it is tidally disrupted. This is consistent with comets being the most fragile bodies in the Solar system, being mostly formed by weakly bound aggregates^{22,28–30}. Some comets may be highly heterogeneous rubble piles as a result of impact gardening and collisional processes^{23,31}, with some pieces having relatively higher strengths, as was proposed to explain the origin of rare H/L chondrites^{29,32}. The characteristic change in v_∞ for the fragments is, $\Delta v_\infty \sim \sqrt{v \Delta v}$, where $v \sim \sqrt{GM_\odot/d_{\odot,R}}$ and $\Delta v \sim \sqrt{Gm/R}$, where $d_{\odot,R}$ is the Sun's Roche radius, m is the mass of the progenitor, and R is the radius of the progenitor. The change in v_∞ , Δv_∞ , is comparable to the original v_∞ for an LPC. The time between disruption and crossing the Earth's orbit is $\sim (d_\oplus / \sqrt{GM_\odot/d_\oplus}) \sim 10^3 \tau$, where $d_\oplus \sim 1$ AU is the distance of the Earth and τ is the tidal disruption encounter timescale, $\tau \equiv \sqrt{d_\oplus^3 / GM_\odot}$. This is consistent with the conversion³³ of $R \sim 30$ km LPCs into fragments with effective radii of $R \sim 3.5$ km, as required for the Chicxulub impactor, using a framework consistent with the Shoemaker-Levy 9 event³⁴ as well as the formation of the Gomul and Gipul crater chains. Data from Gomul and Gipul, as well other crater chains on Callisto and Ganymede, indicate that the fragments typically vary in size only by a factor of order unity³⁵, due to the gravitationally bound rubble pile fragmentation model, although some second-order disruption effects are possible. We note that the canonical equation³⁶ $z_b = z_* - 2H [\ln 1 + (l/2H) \sqrt{f_p^2 - 1}]$ for the parameters considered here is only consistent with $z_b < 0$, implying that despite experiencing disruption during atmospheric entry³⁷, the comet fragment does not suffer an airburst, which was the fate of the Tunguska impactor^{38,39}, but instead forms a crater, as observed. In the equation above, z_b is the altitude at which the airburst occurs, z_* is the altitude at which the comet begins to disrupt, H is the scale height of the atmosphere, $l = L_0 \sin(\theta) \sqrt{\rho_{obj} / (C_D \rho_a(z_*))}$ is the dispersion length scale, $f_p = (L(z)/L_0)$ is the pancake factor, $L = 2R$ is the impactor diameter, ρ_{obj} is the impactor density, ρ_a is the atmospheric density, θ is the impact angle, and C_D is a drag coefficient.

We now consider the effect that tidal disruption of a fraction of LPCs has on the impact rate of cometary bodies capable of producing Chicxulub. We first note that $D \gtrsim 10$ km progenitors, as considered here, are not thermally disrupted at large distances like smaller comets⁴⁰. We adopt the size distribution of Kuiper belt objects (KBOs) as a proxy for large LPCs or Oort cloud objects, due to their shared histories^{41–43}. KBOs with radii ranging from $R \sim 5 - 10$ km and $R \sim 30$ km can be described with a power-law index of $q \sim 2$ for a cumulative size distribution of the form^{44,45}, $N(> R) \propto R^{1-q}$. The size distribution for LPCs, which have been observed up to radii of $R \sim 10$ km, is consistent with the extrapolation of the $q \sim 2$ power law down to a size of a cometary Chicxulub impactor^{40,46}, $R \sim 3.5$ km. KBOs with $R \sim 30$ km are primarily bound by gravity, as indicated by modeling consistent with the observed size-density relationship^{47,48} and as implied by the location of the break in the size distribution^{44,45,49}. Most asteroids with sizes of $D \gtrsim 10$ km are not considered strengthless, meaning that if they passed within the Sun's Roche limit, they most likely would not produce fragments of the necessary size to explain Chicxulub³⁴.

Since the mass of an LPC scales as R^3 and the abundance of LPCs scales as R^{1-q} , the overall enhancement of the time-averaged flux of cometary impactors capable of producing Chicxulub impact events resulting from the breakup and immediate crossing of the ~ 1 AU sphere following perihelion of objects larger than an intact LPC capable of producing a Chicxulub impact event by a factor of ~ 10 in radius is, $\sim 0.2 \times (30 \text{ km}/3.5 \text{ km})^{3+(1-q)} \approx 15$, since 20% of progenitors are tidally disrupted. This results in an impact rate for LPC fragments capable of producing Chicxulub impact events of once per $\sim 250 - 730$ Myr. Irrespective of composition, the total impact rate of LPC fragments that could cause Chicxulub impact events is comparable to the total impact rate of MBAs that trigger events. We note that in order to be in agreement with the lack of an observed increase in the Earth's dust accretion rate across the K/T event over timescales of ~ 1 Myr, the power-law index of the differential size distribution at the time of the tidally disrupted LPC's encounter with Earth must have been $q \gtrsim -3$, which can be tested through detailed modeling of such tidal disruption events.

The carbonaceous chondritic composition fraction of LPCs might be comparable to unity, since the first cometary target of a sample return mission Comet 81P/Wild 2 indicated a carbonaceous chondritic composition. However, the tiny aggregate particles collected had very low tensile strengths, potentially complicating the understanding of cometary structure in general²⁸. Adopting the assumption that the carbonaceous chondritic composition fraction of LPCs might be comparable to unity, the impact rate of tidally-disrupted LPCs is consistent with the Chicxulub impact event being the largest mass extinction event in the last ~ 250 Myr, and is significantly larger than the impact rate of MBAs that could cause Chicxulub impact events. In particular, the

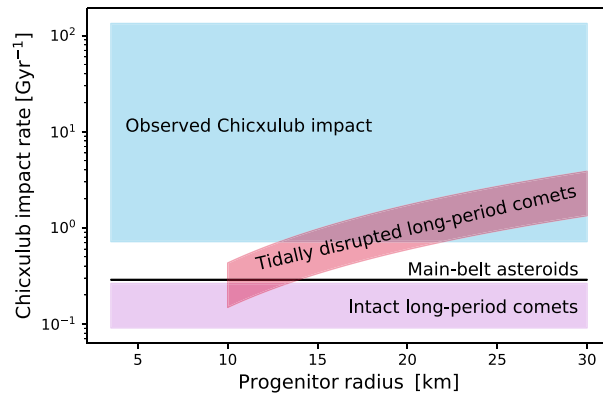


Figure 1. The impact rate of tidally disrupted LPCs with energies comparable to that of the Chicxulub impactor, with the impact rates of intact LPCs and MBAs for reference, in addition to the range of rates that would explain the observed Chicxulub impact, including 95% Poisson errors. Most LPCs and $\sim 10\%$ of MBAs are assumed to have a carbonaceous chondritic composition (see text for details).

probability that the Chicxulub impactor was an LPC fragment is larger than the probability that it was an MBA if the carbonaceous chondritic composition fraction of the LPC progenitors is $\gtrsim 7\text{--}20\%$.

As illustrated in Fig. 1, the LPC fragment hypothesis is consistent with the 95% Poisson limits on the observed Chicxulub impact rate for progenitor carbonaceous chondritic composition fractions of $\gtrsim 20\text{--}50\%$. Future cometary sample-return missions similar to *Stardust* will constrain the fraction of comets with carbonaceous chondritic compositions and thereby serve as important test for our hypothesis. In addition, measurements of the size distribution of Oort cloud objects will improve the precision of our model. Since comets with $D \lesssim 10$ km are thermally disrupted at large distances from the Sun⁴⁰ and also the size distribution of comets with $D \gtrsim 60$ km is described by a power law⁴⁵ with a cumulative power-law index steeper than -3 , our model only applies to the progenitor size range of $10 \text{ km} \lesssim D \lesssim 60 \text{ km}$, thereby not affecting the overall crater size distribution.

Our hypothesis predicts that other Chicxulub-size craters on Earth are more likely to correspond to an impactor with a carbonaceous chondritic composition than expected from the carbonaceous chondritic composition fraction of MBAs. We note that meteorite fall statistics should still reflect the compositions of asteroids, as canonically assumed. For small LPCs that pass within the Sun's Roche radius, the ablated mass is $\sim (R^2 L_{\odot} \tau / 8 d_{\odot,R}^2 Q)$, where L_{\odot} is the luminosity of the Sun, $d_{\odot,R}$ is the Roche radius of the Sun, τ is the encounter timescale, and Q is the energy per unit mass necessary to vaporize the material. Adopting⁵⁰ $Q \sim 3 \times 10^{11} \text{ erg g}^{-1}$, the initial mass is comparable to the ablated mass for object radii of $R \sim 1$ m, resulting in a conservative lower bound on the mass of LPC fragments of $\sim 10^5$ g, which is orders of magnitude above the preatmospheric entry masses of objects that dominate the meteorite flux at the Earth's surface⁵¹. This magnitude of ablation indicates that mass loss is negligible for the progenitor size range considered here. In addition, the heating due to solar irradiation, $\sim 10^3$ K over $\sim 10^3$ s, does not exceed the expected heating from the impact itself³⁶, so no additional signatures of thermal processing would be expected. Shoemaker-Levy 9, 2015 TB145, and the Encke complex are all examples of large fragments resulting from tidal disruption^{30,52,53}. Additionally, the observation that the largest particles in most observed meteoroid streams are cm-sized⁵⁴ is not surprising, since larger particles are naturally more rare than smaller particles.

Indeed, Vredefort, the only confirmed crater on Earth larger than Chicxulub (by a factor of ~ 2 in radius)⁵⁵, may correspond to an impactor with a carbonaceous chondritic composition⁵⁶. Additionally, since LPC fragment Chicxulub impactors should strike Earth once every $\sim 250\text{--}730$ Myr, fragments an order of magnitude smaller in radius, if produced by the same progenitors, would strike Earth no more frequently than once per $\sim 0.25\text{--}0.73$ Myr and if a significant fraction of the progenitors have a carbonaceous chondritic composition, the most recent such crater should reflect such a composition. Indeed, the Zhamanshin crater, the largest confirmed impact crater on Earth formed in the last \sim Myr (an order of magnitude smaller in radius than Chicxulub)⁵⁷, shows evidence that the impactor may have had a carbonaceous chondritic composition⁵⁸, providing support to our model. Additionally, the likely existence of a well-separated reservoir of carbonaceous chondritic material beyond the orbit of Jupiter in the solar protoplanetary disk⁵⁹ lends further support to our model. Our model is in no conflict with the Moon's cratering rate, since it only applies in the size range around Chicxulub-scale impactors. The cross-sectional area of the Moon is an order of magnitude smaller than Earth, implying that a Chicxulub size impactor would be very rare (once per few Gyr), and thereby implying that such an LPC impact event may have not happened for the Moon.

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Author contributions

A.S. and A.L. conceived of the presented idea and developed the theory. A.S. led the writing of the manuscript and A.L. provided critical feedback.

Competing interests

The authors declare no competing interests.

Additional information

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