

SCIENTIFIC REPORTS



OPEN

The new competitive mechanism of hydrogen bonding interactions and transition process for the hydroxyphenyl imidazo [1, 2-a] pyridine in mixed liquid solution

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The new competitive mechanism of intermolecular and intramolecular hydrogen bond can be proposed with an improved mixed model. Upon the photoinduced process, the twisting intramolecular charge transfer (TICT) structure of the hydroxyphenyl imidazo [1, 2-a] pyridine (HPIP) can be obtained. TICT character prompts the fluorescent inactivation via non-radiative decay process. For exploring the photochemical and photophysical properties, the electronic spectra and the infrared (IR) vibrational spectra of titled compounds have been detailedly investigated. In addition, the frontier molecular orbitals (MOs) analysis visually reveals that the unbalanced electron population can give rise to the torsion of molecular structure. To further give an attractive insight into the non-radiative decay process, the potential energy curves have been depicted on the ground state (S_0), the first excited state (S_1) and the triple excited state (T_1). Minimum energy crossing point (MECP) has been found in the S_1 and T_1 state. On the MECP, the intersystem crossing (ISC) might be dominant channel. The density functional theory (DFT) and the time-dependent density functional theory (TDDFT) methods have been throughout employed in the S_0 state, T_1 state and S_1 state, respectively. The theoretical results are consistent with experiment in mixed and PCM model.

Hydrogen bond is a non-covalent interaction between electronegative atom Y (acceptor) and covalent bond group X-H (donor)¹. Generally, the nitrogen, oxygen, or fluorine usually act as electronegative atoms (X, Y). The hydrogen bond X-H...Y plays an important role on the photophysical and photochemical properties, which are correlated with electronegative atoms of hydrogen bond². The human body is comprised of the most abundant several elements including carbon, hydrogen, oxygen and nitrogen, so we will further research the hydrogen bond type that involves the above several elements. To date, an increasing number of researchers have reported that the hydrogen bond plays a critical role in the biochemistry, organic chemistry, photochemistry and physical chemistry, etc.³⁻⁷. Especially the hydrogen bonding interaction has been widely found in the biomolecules such as proteins, nucleic acids, and so on ref. 8. The hydrogen bonding interaction not only be essential to establish stable building blocks of the life, but also act as the sites for the catalyzing reactions of a variety of enzymes. For example, the researchers have reported that the chemical property can be improved *via* changing one single atom in the InsP6 inhibitor, which can strengthen its hydrogen bonding capabilities with toxin molecules. The change has strengthened InsP6 binding to the allosteric modulator by 26-fold⁹. Moreover, hydrogen bond can modulate the metabolism process, such as the hydrogen bond can facilitate the ubiquitous ultra-weak photon emission mode^{9,10}. In addition, the fluorescent phenomena originated from de-excitation process are very widespread in all sorts of scale of biological systems¹¹. Therefore, the dynamic mechanism of hydrogen bond will be extensively studied in the photo-excitation process by contemporary investigators¹²⁻¹⁴. Especially, the excited state hydrogen bond strengthening mechanism has been put forward for the first time by Han *et al.*^{3,10,15-26}. Hydrogen bond strength

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depends on its bond length, the bond angle, the local dielectric constant, the electronegativity of the donor and acceptor groups, temperature, and pressure, *etc.*^{27, 28}. The hydrogen bonding interaction and corresponding dynamical behaviour have been interpreted *via* a variety of photophysical and photochemical phenomena^{16, 29–33}. For example, the ESIPT reaction, the intramolecular charge transfer (ICT) and the TICT, *etc.*^{19, 28}.

The tautomer structures can be obtained by the ESIPT reaction^{28, 34–36}. The ultrafast hydrogen bond strengthening can offer the driving forces for the ESIPT process. Upon the photo-induced process, the electron population has an obvious change for the molecule, which converts from the π character to the π^* character. The electron density of proton donor group and acceptor group will reduce and increase, respectively²⁹. Therefore, the proton transfer process will be facilitated in the S_1 state^{11, 37}. However, the ESIPT process can be understood as a four-level circular loop model and it is a great important process in the photochemistry, biochemistry and so on refs 38, 39.

As discussed by Toshiki Mutai *et al.* isomer of HPIP originated from ESIPT process had shown an extremely weak fluorescence in the polar solvent tetrahydrofuran (THF). They have drawn a conclusion that the fluorescence yield can be widely enhanced by changing the surroundings from the polar liquid state to the solid state⁴⁰. Actually, the non-radiative decay process plays a prominent role in the biological systems^{41, 42}. The long lifetime deactivation process of HPIP will occur *via* the TICT character in the polar solvent, and then the ISC between the S_1 state and the T_1 state might be dominant channel for the non-radiative decay. On the MECP, the molecular structures have almost identical energy in the S_1 and T_1 state⁴¹. Subsequently, the radiationless decay will jump from the T_1 state to the S_0 state. However, the torsion of HPIP isomer can be prohibited in the solid state surroundings, the nearly co-planar form will emit a drastic fluorescence⁴⁰. In this study, we will carefully investigate the fluorescence quenching mechanism of deactivation process. As Zhao and Liu *et al.* discussed^{33, 43}.

In the present work, we utilize a mixed solvent model⁴⁴ to investigate the reaction where a THF molecule bounds to imino group of HPIP, the other solvent can be substituted by polarizable continuum model (PCM). The complex can be established by hydrogen bond interaction. As Fileti *et al.* have discussed that for alcohol-water complex the stability of complex constitution depends on not only different donor and acceptor molecules, but binding energy of hydrogen bond also is influenced by spatial configurations of complex⁴⁵. Especially, Fileti *et al.*⁴⁶ and Malaspina *et al.*⁴⁷ have accurately provided with the most stable complex of pyridine and water molecule *via* combining Monte Carlo computer simulation and first-principles quantum mechanical calculations method in an aqueous environment⁴⁸. Therefore, in this study HPIP complexes have been certified as most stable structures when considering different spatial configurations of complex. The TICT character can be clearly revealed when we consider the interaction between THF molecule and imine group of HPIP molecule. As Wang *et al.* discussed multiple proton transfer *via* an intermolecular hydrogen-bonded water wire can clearly exhibited the effect of hydrogen bonding dynamics for 3-hydroxypyridine²⁹. Otherwise it cannot be primely explained in the conventional PCM model.

Results and Discussion

In this study, we primarily investigate the photophysics and photochemistry properties of HPIP. The HPIP in the THF solvent phase has a dual emitting in the PCM solvation, but the fluorescence of keto form isomer will be totally quenched via the TICT character in the mixed liquid model. However, we are interested in the fluorescence quenching process, and we have speculated that the TICT structure might go through an ISC process from S_1 state to T_1 state on the MECP^{41, 49–51}. Then, the nonradiative decay will be dominant in the $T_1 \rightarrow S_0$ state as shown in Fig. 1. Chu *et al.* has introduced the competition mechanism between the intermolecular hydrogen bonding interaction and the ESIPT process⁵². However, in this study a new competitive mechanism of intramolecular and intermolecular hydrogen bond will be investigated in the torsional process.

Geometric structures and spectra analysis. Four stable structures have been found in the S_0 state, the HPIP (a), the keto form of HPIP (k-HPIP) (b), the trans- keto form HPIP (trans-k-HPIP) (c) and the open-ended intramolecular hydrogen bond HPIP (o-HPIP) (d) have been shown in the Fig. 2. It should be noted that the k-HPIP has not occurred torsion. However, Four stable structures in the S_1 state, the cis-HPIP (a), in PCM the k-HPIP (ik-HPIP) (b), the nearly vertical isomer form of HPIP (v-HPIP) (c) and the trans-k-HPIP (d) have been shown in the Fig. 3. The k-HPIP form cannot exist in the S_1 state. After finishing the ESIPT process of HPIP molecule the isomer form has been directly optimized into v-HPIP form. The ik-HPIP and the v-HPIP forms present the rotations to $\sim 33^\circ$ and $\sim 80^\circ$ between benzothiazole and phenyl group, respectively. For illustrating the reliability of our computation, the absorption and emission spectra have been calculated in the mixed liquid model. The absorption peak values of the cis-HPIP are located in 331 nm. The emission peak values of the cis-HPIP are located in 377 nm. In addition, the fluorescence peak values of the ik-HPIP are located in 602 nm. The above those peak values have coincided with the computational results of Toshiki Mutai *et al.*, which have been shown in Fig. 4. Herein, we have proved the availability of the computing method. The fluorescence of the tautomer form in the S_1 state exists a Stokes' shift beyond 200 nm, which indicates that the ik-HPIP form has a drastic change compared with cis-HPIP form. However, the stable structures have not been investigated exclusively in the mixed liquid model. In the Fig. 5, we have calculated the absorption and fluorescence spectra of the cis-HPIP (a) and trans-k-HPIP (b). As shown in Fig. 5(a), the absorption peak values of the cis-HPIP are located in 328 nm. The emission peak values of the cis-HPIP are located in 374 nm. The emission peak value of v-HPIP form is nearly nonluminous located in the about 847 nm. As shown in Fig. 5(b), the absorption peak is about 497 nm and the fluorescence peak is about 595 nm for the trans-k-HPIP. Although the trans-k-HPIP and k-HPIP cannot be directly gained in the S_0 state, the two structures can be received in the radiationless decay way from the S_1 state to the S_0 state. The reaction mechanism has been exhibited in the Fig. 1. Subsequently, for above structures we have dissected carefully the changes of crucial bond parameters in the S_0 and S_1 state. As shown in the Tables 1–4, the intriguing bond parameters have been revealed. In Table 1, bond lengths O_1-H_1 , $H_1 \cdots N_2$ of the cis-HPIP are 0.959 Å, 1.832 Å in S_0 state and are 1.020 Å, 1.632 Å in the S_1 state, moreover the corresponding

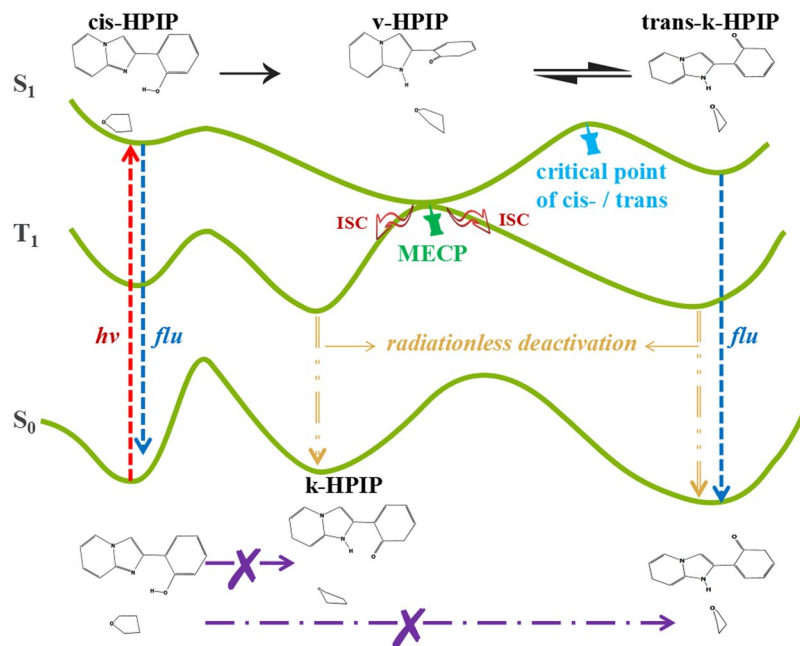


Figure 1. The detailed non-radiative deactivation process.

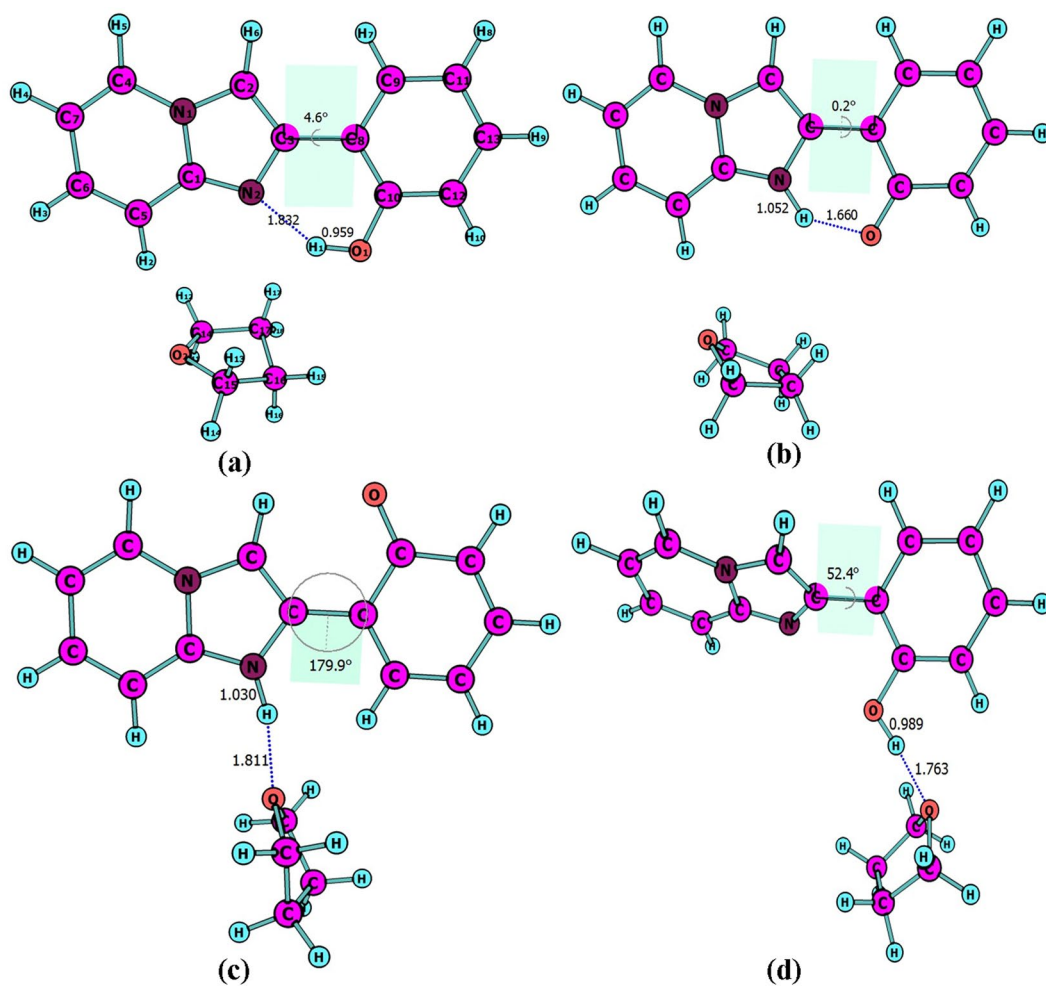


Figure 2. The optimized structures of (a) cis-HPIP, (b) k-HPIP, (c) trans-k-HPIP and (d) o-HPIP in the S_0 state. Some crucial bond parameters have been shown.

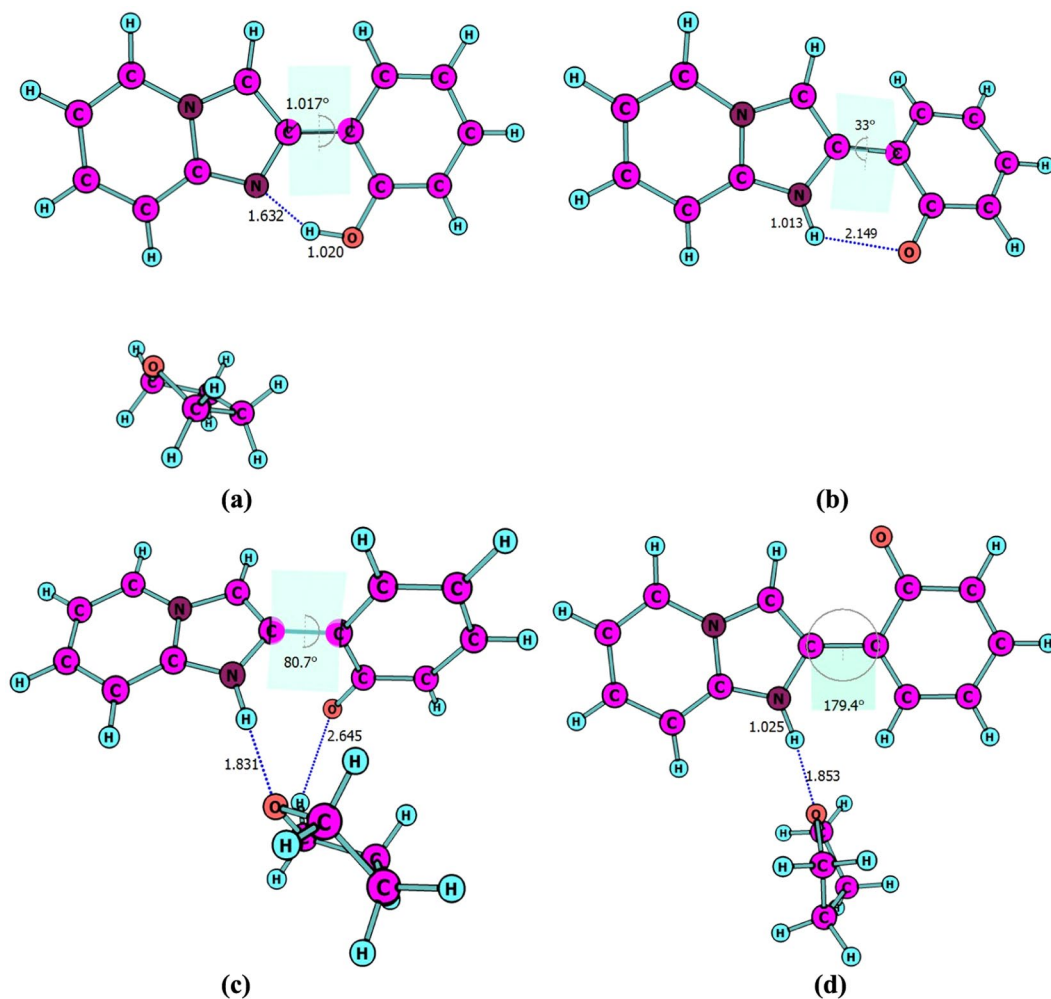


Figure 3. The optimized structures of (a) cis-HPIP, (b) ik-HPIP, (c) v-HPIP, (d) trans-k-HPIP in the S_1 state. Some crucial bond parameters have been shown.

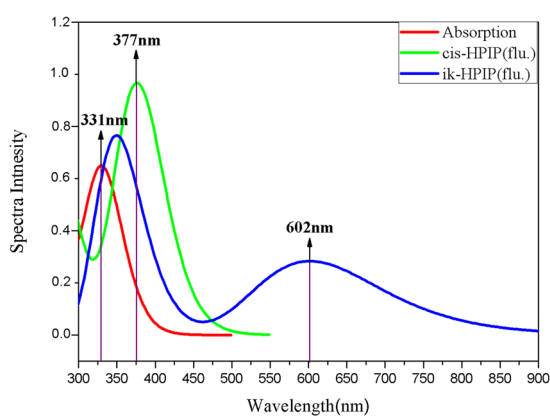


Figure 4. Theoretically simulating the absorption and fluorescence spectra of the HPIP in the PCM solvation, the violet vertical lines stand for the corresponding peak values in the theoretical calculation discussed by Toshiki Mutai *et al.* The detail explanations of curves can be given by the legend on the top right corner.

bond angle $\delta(O_1-H_1-N_2)$ increases from 145.4° to 149.4° . There is a definite conclusion that intramolecular hydrogen bond $O_1-H_1\cdots N_2$ of the cis-HPIP can be strengthened in the S_1 state. The dihedral angle $\delta(N_2-C_3C_8-C_9)$ is 4.6° in the S_0 state, but it reduces to 1.0° in the S_1 state. The structure of cis-HPIP is near planar in the S_1 state.

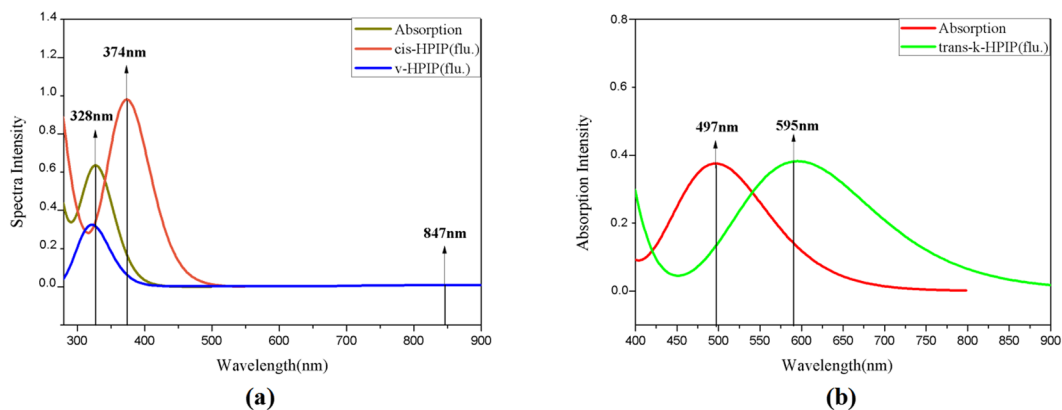


Figure 5. Theoretically simulating the absorption and fluorescence spectra of the HPIP in the mixed liquid model. **(a)** The electron spectra for the cis-HPIP and v-HPIP, **(b)** the electron spectra for the trans-k-HPIP. The detail explanations of curves can be given by the legend on the top right corner.

Parameters of bonds Electronic state	cis-HPIP	
	S ₀	S ₁
O ₁ -H ₁	0.959	1.020
H ₁ -N ₂	1.832	1.632
δ(O ₁ -H ₁ -N ₂)	145.4°	149.4°
δ(N ₂ -C ₃ -C ₈ -C ₉)	4.6°	1.0°

Table 1. The bond parameter (bond length (Å), length angle (°) and dihedral angle (°)) of crucial moiety for the cis-HPIP in the S₀ and S₁ state.

Parameters of bonds Electronic state	ik-HPIP	
	S ₀	S ₁
O ₁ -H ₁	1.644	2.150
H ₁ -N ₂	1.055	1.013
δ(O ₁ -H ₁ -N ₂)	137.9°	118.5°
δ(N ₂ -C ₃ -C ₈ -C ₉)	0.01°	33°

Table 2. The bond parameter (bond length (Å), length angle (°) and dihedral angle (°)) of crucial moiety for the ik-HPIP in the S₀ and S₁ state.

Parameters of bonds Electronic state	trans-k-HPIP	
	S ₀	S ₁
O ₁ -H ₁	4.954	4.945
H ₁ -N ₂	1.030	1.025
O ₂ -H ₁	1.811	1.853
δ(O ₁ -H ₁ -N ₂)	42.2°	43.9°
δ(O ₂ -H ₁ -N ₂)	168.0°	174.8°
δ(N ₂ -C ₃ -C ₈ -C ₉)	179.9°	179.4°

Table 3. The bond parameter (bond length (Å), length angle (°) and dihedral angle (°)) of crucial moiety for the trans-k-HPIP in the S₀ and S₁ state.

The bond parameters of ik-HPIP form has been shown in the Table 2, the bond lengths O₁...H₁, H₁-N₂ and δ(O₁-H₁-N₂) of ik-HPIP change from 2.150 Å, 1.013 Å and 118.5° to 1.644 Å, 1.055 Å, and 137.9° in the S₁ → S₀ state, which indicates that the intramolecular hydrogen bonds O₁...H₁-N₂ of ik-HPIP form is weaker in the S₁ state than that in the S₀ state. In addition, we can find that the dihedral angle δ(N₂-C₃-C₈-C₉) of ik-HPIP form decreases from 33° to 0.01° in the S₁ → S₀ state from this Table. In the Fig. 2(d), the o-HPIP molecular structure has been optimized with mixed liquid model. We find that the o-HPIP form has the torsion of about 52.4° in the

Parameters of bonds Electronic state	v-HPIP	ik-HPIP
	S ₁	S ₁
O ₁ -H ₁	3.389	2.150
H ₁ -N ₂	1.026	1.013
O ₂ -H ₁	1.831	—
H ₁₂ -O ₁	2.645	—
δ(O ₁ -H ₁ -N ₂)	80.1°	118.5°
δ(O ₂ -H ₁ -N ₂)	176.5°	—
δ(N ₂ -C ₃ C ₈ -C ₉)	80.7°	33°

Table 4. The bond parameters (bond length (Å), length angle (°) and dihedral angle (°)) of crucial moiety for the v-HPIP and ik-HPIP in the S₁ state.

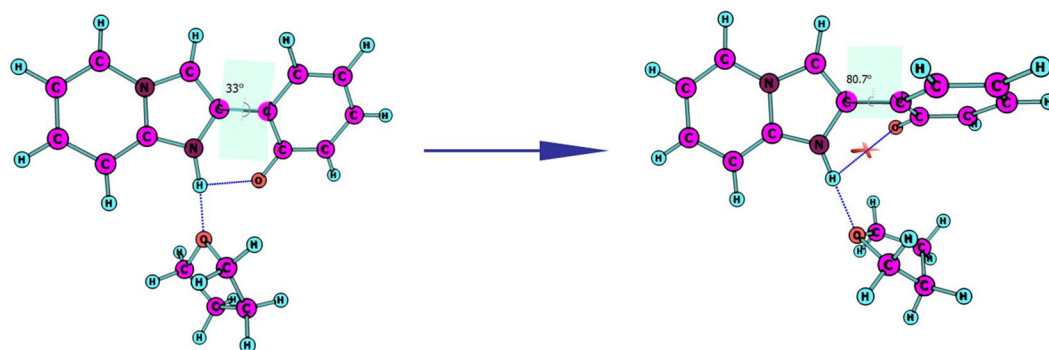


Figure 6. (a) the interaction between THF molecule and ik-HPIP leads to (b) the formation of v-HPIP in the S₁ state.

S₀ state, since the intramolecular hydrogen bond has been destroyed *via* reversing the bond O₁-H₁. Therefore, we can conclude that the intramolecular hydrogen bond O₁-H₁...N₂ can preclude the torsion of HPIP structure. In the Table 3, the bond parameters of trans-k-HPIP form have been shown. The dihedral angle δ (N₂-C₃C₈-C₉) of trans-k-HPIP compared with that of cis-HPIP has twisted about 180°. In addition, a strong intermolecular hydrogen bond N₂-H₁...O₂ has formed. In the S₀ state, the bond lengths O₁-H₁, H₁-N₂ and O₂...H₁ are 4.954 Å, 1.030 Å and 1.811 Å, respectively. The corresponding bond lengths are 4.945 Å, 1.025 Å and 1.853 Å in the S₁ state. Finally, the bond parameters of ik-HPIP and v-HPIP form have been shown in the Table 4, it should be noted that the bond length O₁-H₁ increases from 2.150 Å to 3.389 Å with torsion of the molecule system, which indicates the torsional behavior has given rise to weakening of the intramolecular hydrogen bond N₂-H₁...O₁. In addition, the intermolecular hydrogen bond N₂-H₁...O₂ has been found in the v-HPIP form, and the bond length O₂-H₁ is 1.831 Å. The intermolecular hydrogen bonding interaction can compete with the intramolecular hydrogen bonding interaction of the ik-HPIP form in the mixed liquid model. Yan *et al.* have concluded that the slightly weaker hydrogen bond allows the competition with other type of interaction⁵³. Non-coplanar ik-HPIP form results in the intramolecular hydrogen bond is extremely weak. So when the THF molecule approaches to the imino group (=N-H) of ik-HPIP, the intermolecular hydrogen bond will be strengthened and the intramolecular hydrogen bond will be further weakened, the torsion of HPIP molecule can be facilitated as shown in the Fig. 6. We can propose a viewpoint that the intermolecular hydrogen N₂-H₁...O₂ can give rise to the weakening of intramolecular hydrogen bond N₂-H₁...O₁. Furthermore, we have found that a nonatomic ring structure has been generated in the v-HPIP form linked by two intermolecular hydrogen bonds, which are the strong hydrogen bond N₂-H₁...O₂ and the weak hydrogen bond C₁₄-H₁₂...O₁. The reduced density gradient (RDG) isosurfaces have been shown in the Fig. 7. Herein, the intensity of hydrogen bond N₂-H₁...O₂ can be visibly compared with that of hydrogen bond C₁₄-H₁₂...O₁.

For guaranteeing these structures are the true most stable, the corresponding IR vibrational frequency has been calculated. Meanwhile, the anharmonic effects have been considered in stretching frequencies and ΔZPE correction by means of multiplying correction factors 0.991 and 0.977 that Truhlar *et al.* have fitted⁵⁴. The IR vibrational spectra of hydrogen bond have been shown in the Fig. 8. In the Fig. 8(a), the vibrational frequency of O₁-H₁ is about 3200 cm⁻¹ (the anharmonic frequency is 3171.2 cm⁻¹) in the S₀ state and is about 2729 cm⁻¹ (2704.4 cm⁻¹) in the S₁ state for cis-HPIP form. The 471 cm⁻¹ (466.8 cm⁻¹) red shift indicates that the hydrogen bond O₁-H₁...N₂ is stronger in the S₁ state. In the Fig. 8(b), the vibrational frequency of the N₂-H₁ is tremendously blue-shift 680 cm⁻¹ (673.8 cm⁻¹) from 2926 cm⁻¹ (2899.7 cm⁻¹) to 3606 cm⁻¹ (3573.5 cm⁻¹) in the S₀ → S₁ state for the ik-HPIP form, which has indicated that the hydrogen bond N₂-H₁...O₁ is stronger in the S₀ state. The IR vibrational spectrum of the trans-k-HPIP has been revealed in the Fig. 8(c). The blue-shift 68 cm⁻¹ (67.4 cm⁻¹) from 3302 cm⁻¹ (3272.3 cm⁻¹) in the S₀ state to 3370 cm⁻¹ (3339.7 cm⁻¹) in the S₁ state has indicated that the intermolecular hydrogen bond N₂-H₁...O₂ is stronger in the S₀ state. Upon anharmonic effects the energy of ΔZPE correction for cis-HPIP between the S₀ and S₁ state is 0.16 eV within the error range.

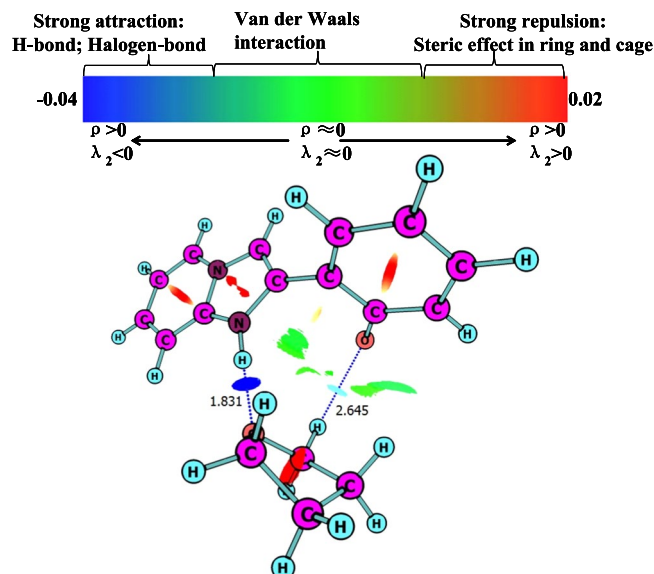


Figure 7. The visual diagram of RDG isosurface and the color gradient axis.

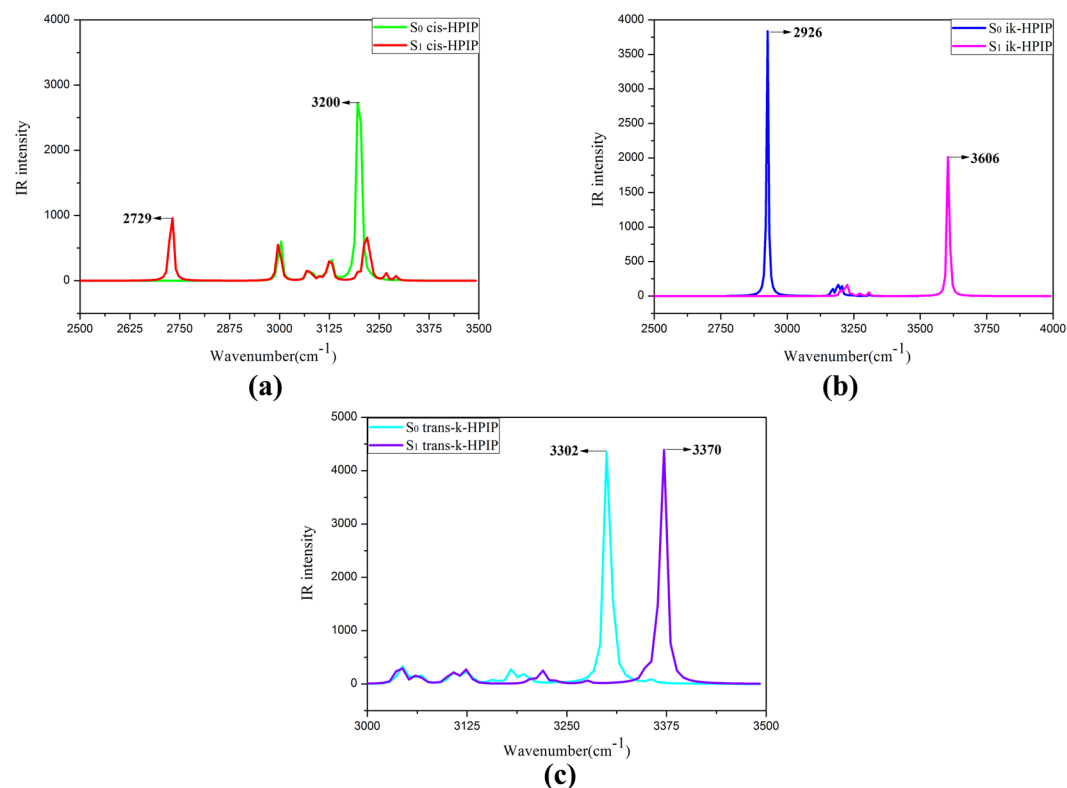


Figure 8. The calculated IR vibrational spectra for (a) cis-HPIP in the spectra region of O₁-H₁, (b) ik-HPIP in the spectra region of N₂-H₁, and (c) trans-k-HPIP in the spectra region of N₂-H₁. The stretching bands of these bonds in the S₀ and S₁ state have been shown. The legend can give reader the detail explanations.

The ESP and the frontier molecular orbitals (MOs) analysis. We guess the rotation is connected with relative displacement of the electron clouds and the nucleus in molecule under the photo-induced. The centre of gravity of positive and negative charges is tipped, which results in the drastic change of dipole moment. In order to prove our conjectures, the ESP values of HPIP have been calculated by the Multiwfn program⁵⁵ and the ESP surface has been shown in the Fig. 9. The corresponding maxima and minima have been exhibited on the figure, we can clearly found that the negative and positive electrostatic potential exist a drastic polarization distribution on the surface. Moreover, in the Fig. 10, the distribution ratio of the different ESP regions has been

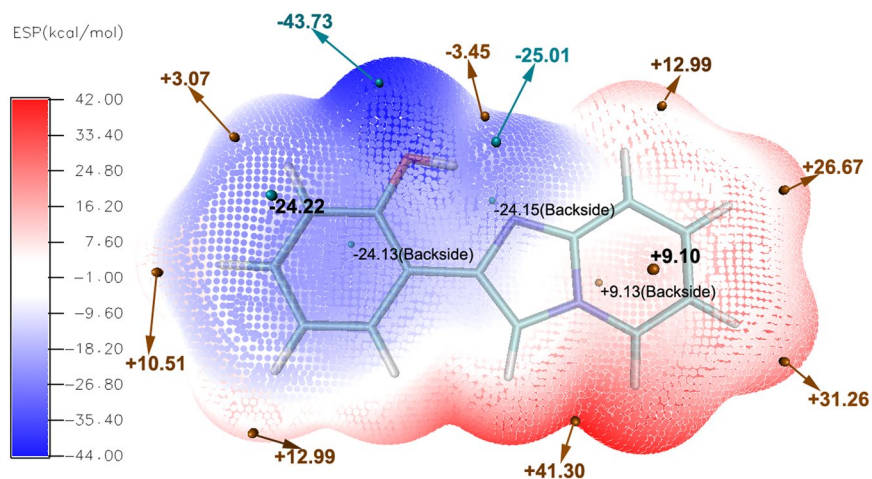


Figure 9. The ESP surface of HPIP molecule structure. The corresponding maxima and minima have been labelled in the surface.

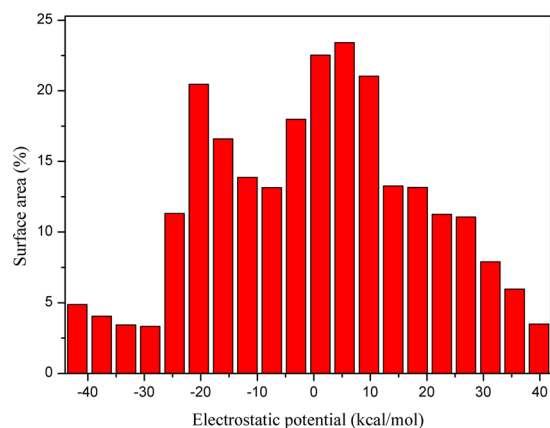


Figure 10. The ESP quantitative distribution column diagram: the X-axis serves as the Surface Area ratio (%), the Y-axis serves as the ESP regions (Kcal/mol).

quantificationally calculated. The ESP values ranged from -20 Kcal/mol to 20 Kcal/mol widely distribute on the surface area. Thereinto, the negative ESP values mainly origin from the π -electron cloud of aromatic nucleus, the aromatic nucleus of C-H hydrogens mainly contribute to the positive areas. In the Fig. 11, the highest occupied molecular orbital (HOMO) and the lowest unoccupied molecular orbital (LUMO) have been revealed. Upon the photo-induced process, the electron distribution mainly changes from the π character on the HOMO to the π^* character on the LUMO. Herein, we only analyze the HOMO and LUMO, because transient excitation mainly stems from the contribution of HOMO \rightarrow LUMO transition. The transition contribution of $\pi\pi^*$ character and corresponding oscillator strength have been listed in the Table 5. In this table, Oscillator strength (f) of ik-HPIP and v-HPIP is 0.1334 and 0.0026, respectively. Their fluorescence compared with the fluorescence of cis-HPIP has been quenched. In the Fig. 11(a), electron redistribution in the PCM is nearly identical to that in the mixed solvation. For further studying the relationship between torsion of molecule structure and electron redistribution, we have obtained non-coplanar HPIP (np-HPIP) form by factitiously twisting the dihedral angle $\delta(N_2-C_3C_8-C_9)$ in the theoretical calculation, in which the np-HPIP molecule cannot be obtained actually in the S_0 state. The MOs graphs of np-HPIP form have been shown in the Fig. 11(b). Its electron redistribution of the HOMO \rightarrow LUMO compared with that of cis-HPIP still is not significant change. Therefore, we draw a conclusion that factitious torsion of molecule structure cannot result in the drastic electron redistribution on the np-HPIP. However, the HOMO and LUMO of the ik-HPIP, v-HPIP and trans-k-HPIP have been depicted in the Fig. 11(c,d) and (e). The electron redistribution of these isomeric forms is great different from that of cis-HPIP. It should be noted that electron population of the isomeric forms is unbalanced for HOMO \rightarrow LUMO. These isomeric forms are zwitterions that the imino group carries positive electric charges and the ketonic oxygen atom carries negative electric charges. ICT character of these zwitterions can lead to the increasing of the dipole moments in the S_1 state, so the corresponding electron population will be unbalanced in the molecule. Moreover, the unbalanced electron population can further result in the torsion of isomeric forms. It is worth noting that TICT character of

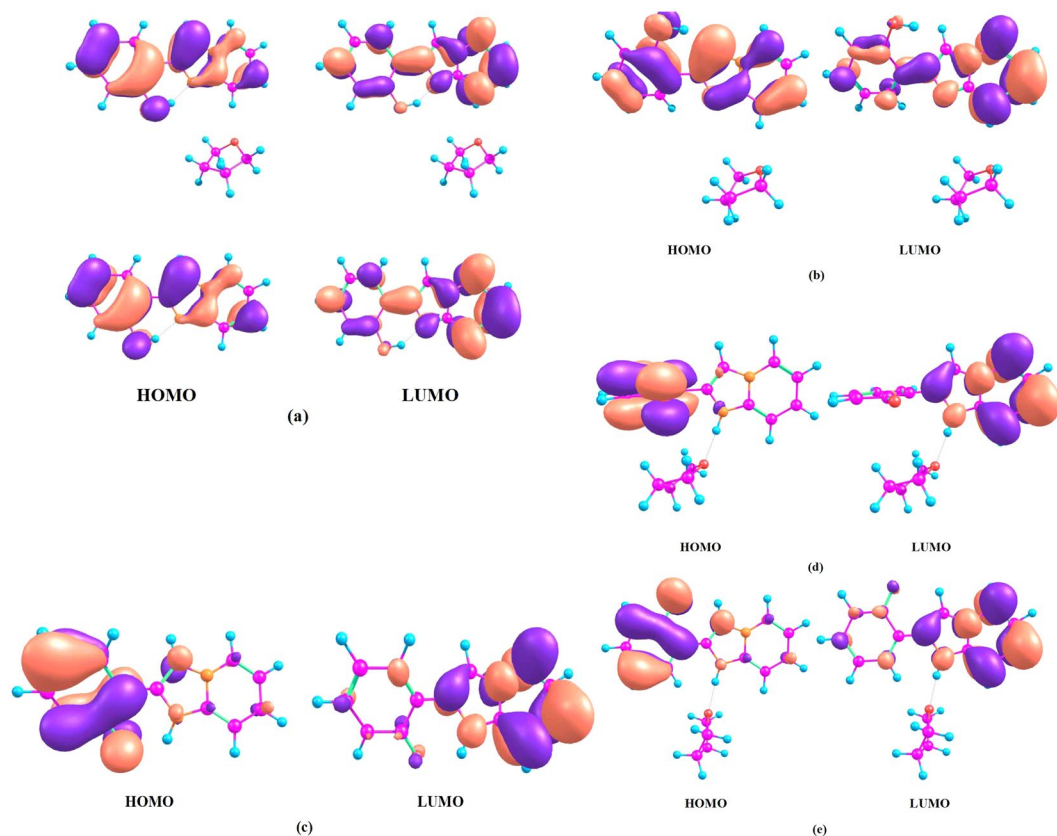


Figure 11. The visual electron population of the frontier molecular orbitals HOMO and LUMO for (a) the cis-HPIP in the mixed liquid and PCM solvation, (b) np-HPIP, (c) ik-HPIP, (d) v-HPIP and (e) trans-k-HPIP.

Molecule name	Oscillator strength (f)	CI (%)	Orbital transition
cis-HPIP (PCM)	0.3945	96.27	55 → 56
cis-HPIP (complex)	0.3839	96.08	75 → 76
ik-HPIP	0.1334	99.67	55 → 56
v-HPIP	0.0026	99.69	75 → 76
trans-k-HPIP	0.2318	99.58	75 → 76
np-HPIP	0.2661	94.03	75 → 76

Table 5. The corresponding oscillator strength (f), contribution index (CI) and orbital transition for the different molecular structures in the $S_0 \rightarrow S_1$ state.

the v-HPIP is shown in the Fig. 11(d), the electron population almost entirely transfers from the phenyl group to the benzothiazole group in the HOMO → LUMO.

The analysis of potential energy curves and MECP. For further studying the non-radiative decay process, comprehending the optical properties and pathways of electronic transition, the potential energy curves of the HPIP in the S_0 and S_1 state have been carefully investigated. The energy of structure is a function of bond length O_1-H_1 and dihedral angle $\delta(N_2-C_3-C_8-C_9)$, respectively. The potential energy curves will be drawn with the bond length and dihedral angle increased by the fixed step sizes, respectively. The reactive potential barriers and stable structures have been obtained in the potential curves. Firstly, we study the potential energy curves in the S_0 state. As shown in the Fig. 12, the two potential energy curves are the reactive pathways of (a) intramolecular proton transfer and (b) twisting dihedral angle, respectively. In the Figure the k-HPIP form is a stable structure and it has been used as the original structure in dihedral angle scan. Proton transfer cannot spontaneously occur in the S_0 state, because the process needs to cross a potential barrier 7.81 Kcal/mol in the Fig. 12(a). Similarly, the dihedral angle torsion cannot also spontaneously occur, and the process needs to cross an energy barrier 9.92 Kcal/mol in the Fig. 12(b). However, the energy instantly reduces about 3.17 Kcal/mol in the torsion process from the Fig. 12(b), because the intermolecular hydrogen bond $N_2-H_1 \cdots O_2$ takes shape in torsion process. Moreover, when the dihedral angle turns to about 180°, the trans-k-HPIP will be a stable structure. Therefore, cis-HPIP, k-HPIP

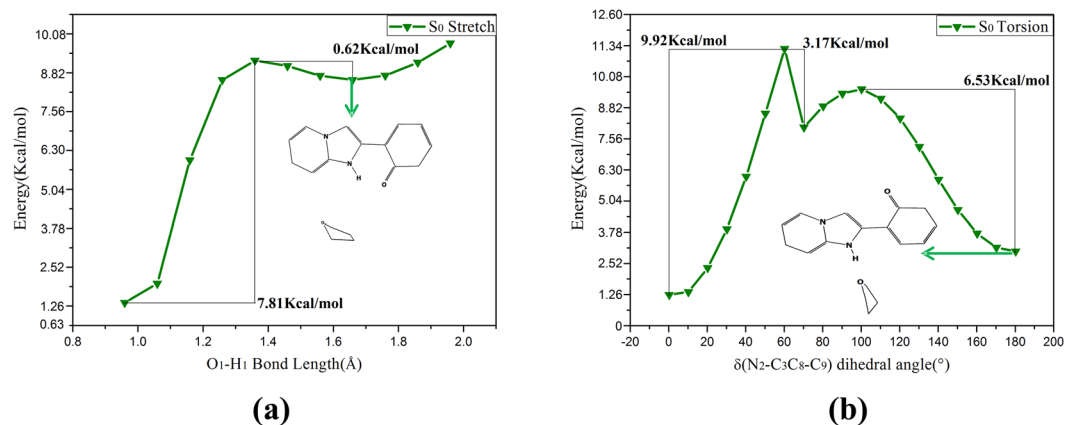


Figure 12. Constructed the potential energy curves of the HPIP: (a) the energies of different versus the O_1-H_1 bond lengths in the S_0 state, (b) the energies of different structures versus the dihedral angles $\delta(N_2-C_3C_8-C_9)$ in the S_0 state. The numerical values in the graphs stand for the energy barriers of the reactions.

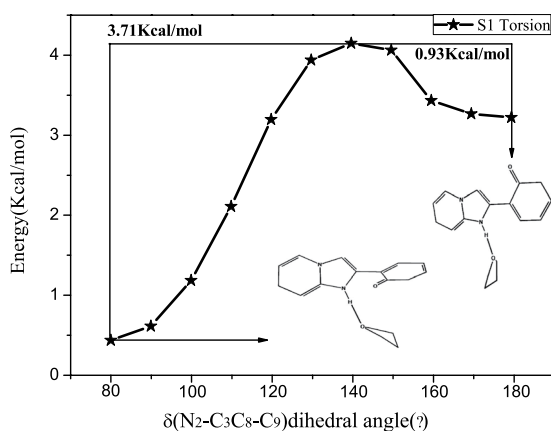


Figure 13. Constructed the potential energy curves of the HPIP: the energies of different structures versus the dihedral angles $\delta(N_2-C_3C_8-C_9)$ in the S_1 state. The numerical values in the graphs stand for the energy barriers of the reactions.

and trans-k-HPIP forms can exist in the S_0 state, but the energy barriers of the proton transfer and structural torsion process are so high that the two processes cannot spontaneously occur in the S_0 state. Secondly, we study the potential energy curves in the S_1 state. In the mixed liquid model the keto form will directly be optimized into the v-HPIP form, so the v-HPIP form has been used as the initial structure in dihedral angle scan. In the Fig. 13 the potential energy curve of twisting dihedral angle has indicated that the v-HPIP form located in about 80° is stable structure, and the stable tran-k-HPIP form located in about 180° . Moreover, the torsion from the v-HPIP to the trans-k-HPIP must get over an energy barrier 3.71 Kcal/mol. The reversed torsion process needs to get over a small energy barrier 0.93 Kcal/mol, which indicates that the trans-k-HPIP can spontaneously revert to v-HPIP form in the S_1 state. The corresponding potential barrier is the critical point of cis- and trans-isomer.

To sum up, we have fully explained the reactive pathways for the HPIP complexes *via* analyzing the corresponding potential energy curves in the S_0 and S_1 state. However, the non-radiative transition process is still ambiguous, so we have further calculated the potential energy curves of the T_1 state. Because the oscillator strength of v-HPIP we have calculated is 0.0026, fluorescence of the v-HPIP has been totally quenched. We speculate that the v-HPIP structure might undergo a non-radiative transition from the S_1 state to the S_0 state. Therefore, herein searching the MECP has become mainly work for us. As shown in the Fig. 14, the potential energy curves of the torsional dihedral angle in the T_1 state and S_1 state have been exhibited. It should be noted that the MECP locates in about 90° in the figure. On this point, the ISC process might be dominant channel in the $S_1 \rightarrow T_1$ process. In the potential energy curves of the T_1 state, the MECP structure is extremely unstable, so this structure will fast slide to the k-HPIP form located in about 0° or will get over a negligible energy barrier 0.33 Kcal/mol and then fast slide to the trans-k-HPIP form located in about 180° along with the potential energy curve of T_1 state. The k-HPIP and trans-k-HPIP existed in the S_0 state could be obtained from the radiationless decay process. As Harvey *et al.* reported that the reliability of hybrid method has been testified about searching the MECP⁵⁶. Therefore, for prove the reliability of MECP, the sobMECP suite⁵⁷ has been used in this study. It should be noted

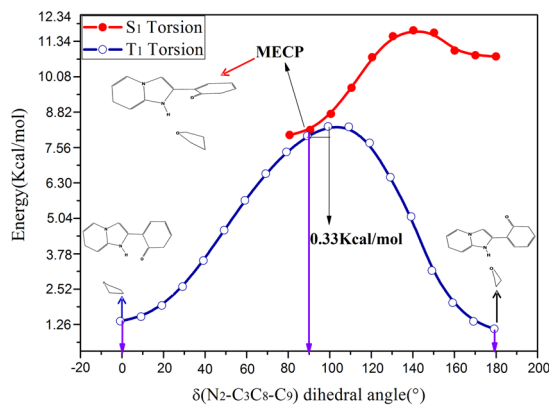


Figure 14. Constructed the potential energy curves of the HPIP: red line: the energies of different structures versus the dihedral angles $\delta(\text{N}_2\text{-C}_3\text{C}_8\text{-C}_9)$ in the S_1 state; blue line: the energies of different structures versus the dihedral angles $\delta(\text{N}_2\text{-C}_3\text{C}_8\text{-C}_9)$ in the T_1 state. The numerical value in the graph stands for the energy barrier of the reaction. The MECP and the corresponding structure have been shown.

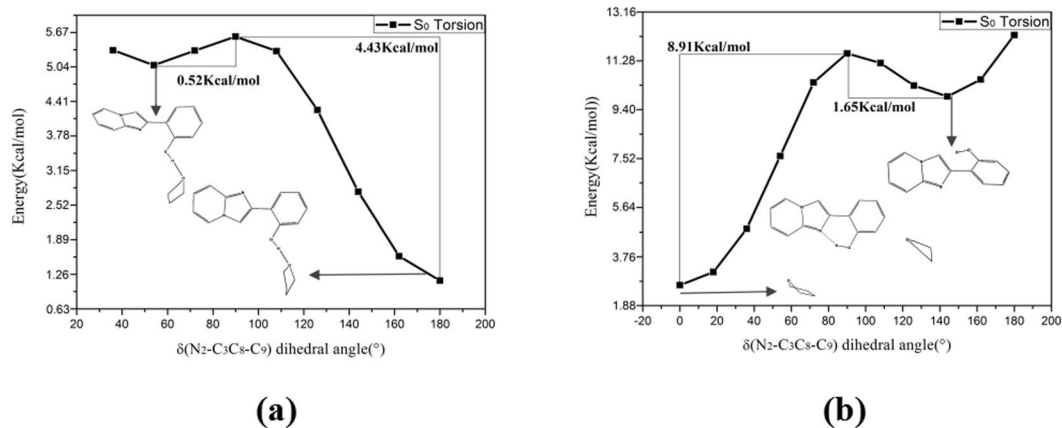


Figure 15. Constructed the potential energy curves of the o-HPIP (a) and np-HPIP (b): the energies of different structures versus the dihedral angles $\delta(\text{N}_2\text{-C}_3\text{C}_8\text{-C}_9)$ in the S_0 state. The numerical values in the graph stand for the energy barriers of the torsion and reversed torsion process.

that the point sobMECP has sought out is about 0.34 Kcal/mol higher than MECP we have confirmed, and their geometries are nearly consistent.

In addition, the functions of the potential energy curve versus the dihedral angle $\delta(\text{N}_2\text{-C}_3\text{C}_8\text{-C}_9)$ of o-HPIP and np-HPIP forms have been depicted on the Fig. 15. The negligible energy barrier of o-HPIP is 0.52 Kcal/mol and the energy barrier of reversed twisting is 4.43 Kcal/mol in the Fig. 15(a), it is clearly indicated that twisting of o-HPIP can spontaneously occur in the S_0 state, the energy of o-HPIP form is gradually reduced with the torsional process and the stable structure located in about 180°. On the contrary, in the Fig. 15(b) the energy barrier of np-HPIP is 8.91 Kcal/mol and the energy barrier of reversed twisting is 1.65 Kcal/mol. The stable np-HPIP form locates in about 143.4° and the energy of np-HPIP form is gradually increased with the torsional process. Therefore, the intramolecular hydrogen bond can preclude the torsion of HPIP structure.

Conclusions

For illustrating the reliability of our computation, we have compared the electronic spectra with that of Toshiki Mutai *et al.* Both of the computed results are tremendous coincidental. The hydrogen bonding strengthening mechanism has been proved *via* analyzing the bond parameters of hydrogen bond in the S_0 and S_1 state. In addition, the analysis of IR vibrational spectra can further illustrate the above strengthening mechanism. In this study, competitive mechanism that intermolecular hydrogen bond $\text{N}_2\text{-H}_1\cdots\text{O}_2$ can weaken interaction of intramolecular hydrogen bond $\text{N}_2\text{-H}_1\cdots\text{O}_1$ has been primely proved. MOs analysis has indicated that the unbalanced electron population of the keto forms can give rise to the torsion of structures. For further studying the model of non-radiative decay process, comprehending the optical properties and pathways of electronic transition, the potential energy curves have been studied. Herein, we draw a conclusion that the k-HPIP and trans-k-HPIP cannot spontaneously occur in the S_0 state resulted from calculating the reactive energy barriers. These two structures may be obtained by the non-radiative decay process from $S_1 \rightarrow T_1 \rightarrow S_0$ state. The MECP has further been sought out *via* sobMECP program⁵⁷. Moreover, we have demonstrated that the intramolecular hydrogen bond $\text{N}_2\text{-H}_1\cdots\text{O}_1$

can preclude the torsion of HPIP form *via* analyzing the o-HPIP form and np-HPIP form and corresponding potential energy curves.

Computational details. For providing with available information of geometrical configurations, such as the minimum energy, potential energy surface, electron spectra and infrared vibrational spectra, *etc.*⁴⁴. The density functional theory (DFT) and the time-dependent density functional theory (TDDFT) have been throughout employed in the S_0 state and the S_1 state, respectively^{58–60}. Both of them have been performed at Gaussian 09 program⁶¹. In this study, we only select to use Becke's three-parameter hybrid exchange function with the Lee-Yang-Parr gradient-corrected correlation functional (B3LYP)^{62–64}, the Pople's 6–31 G (d) and 6–31 + G (d) triple- ξ quality basis sets are used in this level computation⁶⁵. Herein, the functional B3LYP had been extensively applied in the past few decades, which indicates that B3LYP method is greatly reliable for calculating theoretically in the S_1 state^{15, 16, 30, 31, 66–80}. We have investigated the solvent effect based on the integral equation formalism polarizable continuum model (IEF-PCM)^{81–83}. The molecular electrostatic potential (ESP) has been calculated for predicting the nucleophilic and electrophilic sites of target molecule. The ESP surface has been visually portrayed *via* the Visual Molecular Dynamics (VMD) software⁸⁴, and we have also applied the RDG function to investigate the weak interaction types in the Multiwfn program⁵⁵. The visual software Chemcraft has been applied to visualize these figures, such as the molecular structures, the MOs and the RDG⁸⁵. The MECPs were sought out *via* sobMECP, which is a modified version of Harvey's MECP program⁵⁶ by Tian Lu. The sobMECP is a wrapper of Harvey's MECP program to simplify the operation of the MECP program^{55, 86}.

For performing fully optimized in our molecular system, the hybrid functional method has been carried out in computed point energies and geometries and corresponding gradients. The effective gradients \mathbf{f} and \mathbf{g} have been defined as:

$$\mathbf{f} = (E_1 - E_2) \left[\left(\frac{\partial E_1}{\partial q} \right) - \left(\frac{\partial E_2}{\partial q} \right) \right] = (E_1 - E_2) \mathbf{x}_1 \quad (1)$$

$$\mathbf{g} = \left(\frac{\partial E_1}{\partial q} \right) - \frac{\mathbf{x}_1}{|\mathbf{x}_1|} \left[\left(\frac{\partial E_1}{\partial q} \right) \cdot \frac{\mathbf{x}_1}{|\mathbf{x}_1|} \right] \quad (2)$$

where terms E_1 and E_2 are energies on the two potential energy curves, respectively. The $\partial E_n / \partial q$ are corresponding partial derivatives of relative nuclear coordinates q .

The energy difference $E_1 - E_2$ can reduce gradually in the vector \mathbf{f} direction and E_1 can reduce in vector \mathbf{g} direction, which the two vectors \mathbf{f} and \mathbf{g} are orthogonal. The MECP can be sought out by optimizing with effective force $\mathbf{f} + \mathbf{g}$ of molecular system. $E_1 - E_2$ of two states presented differ spin status can be reduced with descending total energy.

As hybrid method will be used to find MECP, the spin-orbit coupling will not be taken consideration. In two 2×2 Hessian matrices H_1 and H_2 electronic coupling matrix elements $H_{12} = H_{21}$ become zero. Therefore, when the diagonal matrix elements of two matrices are equivalent, the MECP will be a real minimum. However, MECP is not a stable point in 3N-6 dimensions when the hybrid method is taken a consideration. The energy of MECP will be corrected via second-order Taylor expansion for two states:

$$\begin{aligned} E &= E_{MECP} + \frac{1}{2} \Delta q^T \left(\frac{\left| \frac{\partial E_2}{\partial q} \right| H_2}{|\mathbf{x}_1|} - \frac{\left| \frac{\partial E_2}{\partial q} \right| H_1}{|\mathbf{x}_1|} \right) \Delta q \\ &= E_{MECP} + \frac{1}{2} \Delta q^T H_{eff} \Delta q \end{aligned} \quad (3)$$

where Δq is the displacement along touching hyperline, which is perpendicular to different gradient \mathbf{x}_1 , the H_{eff} is diagonal matrix elements of effective Hessian. For nonlinear molecular system, the MECP can be testified as indeed minimum point in 3N-7 dimensions. So this program can approximatively research dynamic measures in nonadiabatic surfaces⁵⁶.

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Acknowledgements

This work was supported by the National Natural Science Foundation of China (Grant No. 11474141, 11544015, 11604333), the Program for Liaoning Excellent Talents in University (Grant No. LJQ2015040).

Author Contributions

Y.Q. Li supervised the project, Y.F. Yang, Y.D. and Y.Q. Li performed calculations. Y.F. Yang and Y.Q. Li analyzed data and wrote the paper.

Additional Information

Competing Interests: The authors declare that they have no competing interests.

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