

Transition-Metal Chemistry of Alkaline-Earth Elements: The Trisbenzene Complexes $M(\text{Bz})_3$ ($M = \text{Sr}, \text{Ba}$)

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Dedicated to Professor Christian Reichardt on the occasion of his 85th birthday

Abstract: We report the synthesis and spectroscopic identification of the trisbenzene complexes of strontium and barium $M(\text{Bz})_3$ ($M = \text{Sr}, \text{Ba}$) in low-temperature Ne matrix. Both complexes are characterized by a D_3 symmetric structure involving three equivalent η^6 -bound benzene ligands and a closed-shell singlet electronic ground state. The analysis of the electronic structure shows that the complexes exhibit metal–ligand bonds that are typical for transition metal compounds. The chemical bonds can be explained in terms of weak donation from the π MOs of benzene ligands into the vacant $(n-1)d$ AOs of M and strong backdonation from the occupied $(n-1)d$ AO of M into vacant π^* MOs of benzene ligands. The metals in these 20-electron complexes have 18 effective valence electrons, and, thus, fulfill the 18-electron rule if only the metal–ligand bonding electrons are counted. The results suggest that the heavier alkaline earth atoms exhibit the full bonding scenario of transition metals.

Introduction

The alkali metals and the alkaline earth elements with a $(n)s^1$ or $(n)s^2$ valence shell configuration, respectively, are the most electropositive elements in the periodic table. These classical s-block main group elements were for a long time considered to be among the rather uninteresting elements in terms of chemical bonding due to preconceived notions about their valence orbitals, their reactivity and handling.^[1] They readily lose the outermost $(n)s$ electron(s) to form ionic compounds in the +1 and +2 oxidation states. But recent reports showed that the alkali and alkaline earth elements have a much richer chemistry than hitherto thought.^[2–11] In particular, the heavier alkaline earth atoms $M = \text{Ca}, \text{Sr}, \text{Ba}$ may use their $(n-1)d$ functions as genuine valence orbitals in

the octa-coordinated complexes $M(\text{CO})_8$ ^[7] and $M(\text{N}_2)_8$.^[11] The cubic (O_h) adducts possess metal–ligand bonds that are typical for transition-metal complexes. The participation of the $(n-1)d$ atomic orbitals (AOs) of barium in chemical bonding had been suggested already in earlier theoretical studies^[12–17] and barium was even called by Pyykkö an “honorary transition metal”.^[17] Recent investigations demonstrated that $(n-1)d$ AO bonding is not restricted to Ba only and that the heavier atoms Ca, Sr, Ba exhibit a transition metal chemistry to a much larger extent than ever known so far.^[7–11] Here, we report the isolation and spectroscopic identification of the trisbenzene complexes of heavier alkaline earth atoms $M(\text{Bz})_3$ ($M = \text{Sr}, \text{Ba}$). The synthesis of the first transition metal benzene complex $\text{Cr}(\text{Bz})_2$ by Fischer and Hafner in 1955 was a milestone in transition metal chemistry.^[18] The isolation of $\text{Sr}(\text{Bz})_3$ and $\text{Ba}(\text{Bz})_3$ now opens the door to another chapter of metal complexes.

Results and Discussion

The strontium and barium–benzene complexes were prepared by co-deposition of laser-evaporated strontium and barium atoms with benzene molecules in excess neon at 4 K.^[19] The experiments were performed using relatively low laser energy to avoid the formation of multinuclear species. A series of experiments were performed using different benzene concentrations ranging from 0.025–0.5%. The IR spectra in the 1650–600 cm^{-1} region from co-deposition of laser-evaporated strontium atoms with 0.1% C_6H_6 are represented in Figure 1.

A group of metal-dependent product absorptions are observed after sample deposition at 4 K, remain almost

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<https://doi.org/10.1002/anie.201908572>.

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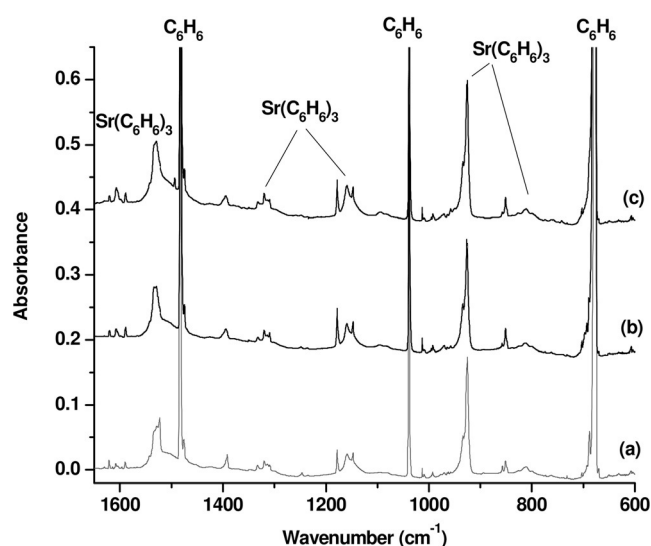


Figure 1. Infrared absorption spectra of strontium–benzene complexes in the 1650–600 cm^{-1} region from co-deposition of laser-evaporated strontium atoms with 0.1% C_6H_6 in neon. a) 30 min of sample deposition at 4 K, b) after annealing at 12 K, c) after 20 min of visible light irradiation.

unchanged on sample annealing to 12 K and increase under 20 min of visible light excitation using a high pressure mercury arc lamp with a 495 nm long-wavelength pass filter ($\lambda = 495\text{--}600\text{ nm}$). Since some bands are partially overlapped by the benzene absorptions, a difference spectrum taken from the spectrum after visible light excitation minus the spectrum after 12 K annealing is obtained, as shown in Figure 2. The upward bands represent the product absorptions, while the downward bands are due to benzene absorptions. It clearly shows the formation of the product absorptions with the consumption of the benzene absorptions. The bands are quite

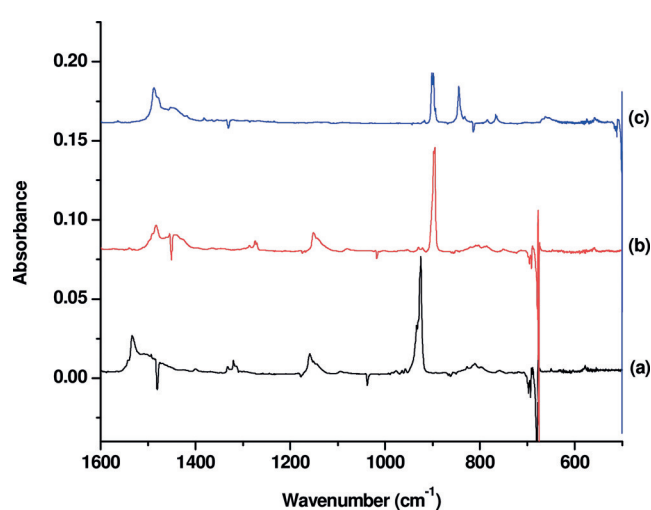


Figure 2. Difference spectra of strontium–benzene complexes in the 1600–500 cm^{-1} region from co-deposition of laser-evaporated strontium atoms with isotopic-labeled benzene in excess neon (the spectrum taken after 20 min of visible light irradiation minus the spectrum taken after 12 K annealing). a) 0.1% C_6H_6 , b) 0.2% $^{13}\text{C}_6\text{H}_6$, and c) 0.2% C_6D_6 .

broad and some bands show site absorptions. Laser evaporation of alkaline earth metal targets might easily lead to formation of cations. In order to verify whether the observed species is a cation or a neutral, additional experiments were performed by adding a trace of CCl_4 in the reactant gas to serve as an electron trap.^[20] CCl_4 captures most of electrons produced by laser ablation during sample deposition, and thus would facilitate the survival of more cation products. The spectra from the experiment with 0.01% CCl_4 added to the neon matrix gas are compared to the spectra from the experiment without CCl_4 doping in Figure S1 of Supporting Information. The results show that the observed product absorptions are reduced relative to other absorptions, indicating that these product absorptions are due to a neutral species rather than a cation.

Isotopic substitution experiments ($^{13}\text{C}_6\text{H}_6$, C_6D_6 , and $^{12}\text{C}_6\text{H}_6 + ^{13}\text{C}_6\text{H}_6$) were performed for product identification based on isotopic shifts and absorption splitting. As shown in Figure 2, all the absorptions are shifted with the $^{13}\text{C}_6\text{H}_6$ and C_6D_6 samples. The assignment of these product absorptions to the tris(benzene)strontium complex is based on the isotopic splitting of the most intense band at 924.9 cm^{-1} , which is due to the benzene ring breathing mode based on the isotopic data. The spectrum from the experiment using an equimolar mixture of $^{12}\text{C}_6\text{H}_6$ and $^{13}\text{C}_6\text{H}_6$ shows that a quartet with two weak intermediates is produced (The low frequency intermediate band is partially overlapped by the strong absorption of $\text{Sr}(^{13}\text{C}_6\text{H}_6)_3$, See Figure 3). This mixed isotopic spectral feature indicates that the 924.9 cm^{-1} band is due to a doubly degenerate benzene ring breathing mode of a tris(benzene) complex involving three equivalent benzene ligands. The positions and mode assignment of the observed bands are listed in Table 1. Besides the most intense ring breathing mode, the second most intense band at 1534.8 cm^{-1} is assigned to the C–C stretching mode. The third strongest band at

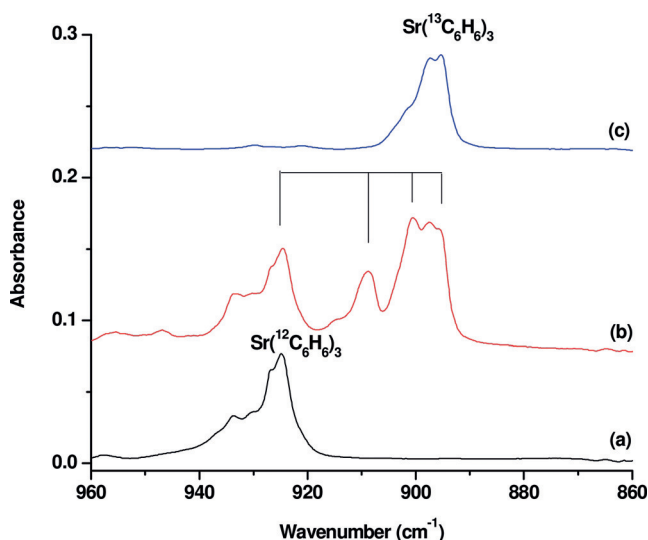


Figure 3. Infrared absorption spectra of strontium–benzene complexes in the 960–860 cm^{-1} region from co-deposition of laser-evaporated strontium atoms with isotopic-labeled benzene in excess neon (the spectra were taken 20 min of visible light irradiation). a) 0.1% C_6H_6 , b) 0.15% $^{12}\text{C}_6\text{H}_6 + 0.15\%$ $^{13}\text{C}_6\text{H}_6$, and c) 0.2% $^{13}\text{C}_6\text{H}_6$.

Table 1: Experimental infrared absorptions (cm^{-1}) of $\text{Sr}(\text{C}_6\text{H}_6)_3$, $\text{Sr}({}^{13}\text{C}_6\text{H}_6)_3$ and $\text{Sr}(\text{C}_6\text{D}_6)_3$ in solid neon and calculated values at the M06-2X(D3)/def2-TZVPP level (D_3 symmetry). Calculated IR intensities (km mol^{-1}) are given in parentheses.

Mode	Experimental			${}^{12}\text{C}_6\text{D}_6$		Calculated			${}^{12}\text{C}_6\text{D}_6$	
	${}^{12}\text{C}_6\text{H}_6$	${}^{13}\text{C}_6\text{H}_6$	$\Delta^{[a]}$		$\Delta^{[a]}$	${}^{12}\text{C}_6\text{H}_6$	${}^{13}\text{C}_6\text{H}_6$	$\Delta^{[a]}$		$\Delta^{[a]}$
C=C stretch	1534.8	1483.0	-51.8	1487.6	-47.2	1556.8 (700)	1506.8 (608)	-50.0	1504.1 (812)	-52.7
C=C stretch and C-D in-plane bend				1382.1					1399.6 (54)	
C=C stretch and C-D in-plane bend				1366.0					1372.2 (21)	
C=C stretch and C-H in-plane bend	1332.6	1286.4	-46.2			1365.3 (49)	1317.5 (47)	-47.8		
C=C stretch and C-H in-plane bend	1320.1	1274.2	-45.9			1346.4 (34)	1299.7 (32)	-46.7		
C-H(D) in-plane bend	1159.1	1151.6	-7.5	844.6	-314.5	1173.3 (119)	1163.4 (155)	-9.9	860.8 (50)	-312.5
C-H(D) in-plane bend	1094.6	1080.6	-14.0	832.6	-262.0	1141.6 (15)	1129.0 (17)	-12.6	848.5 (65)	-293.1
C-H(D) out-of-plane bend	933.8			766.8	-167.0	993.6 (5)			809.0 (53)	-184.6
Ring breath	924.9	895.5	-29.4	898.1	-26.8	968.0 (415)	935.4 (374)	-32.6	936.8 (203)	-31.2
C-H(D) out-of-plane bend	810.7	803.1	-7.6	662.0	-148.7	853.1 (58)	845.5 (53)	-7.6	690.2 (50)	-162.9

[a] Shift with respect to ${}^{12}\text{C}_6\text{H}_6$.

1159.1 cm^{-1} is attributed to the C-H in-plane bending mode. The remaining bands are weak. The bands at 1332.6 and 1320.1 cm^{-1} are assigned to mixed modes involving C-C stretching and in-plane C-H bending. The bands at 933.8 and 810.7 cm^{-1} are attributed to the out-of-plane C-H bending modes.

The $\text{Sr}(\text{C}_6\text{H}_6)_3$ absorptions are the only product absorptions observed in the experiments using different benzene concentrations as well as laser energy. The spectra from experiments using different laser energy are shown in Figure S2. There is no any experimental evidence of formation of the $\text{Sr}(\text{C}_6\text{H}_6)$ and $\text{Sr}(\text{C}_6\text{H}_6)_2$ complexes in these experiments.

Similar absorptions are also observed in the barium experiments. The IR spectra are shown in Figure 4. All the absorptions are observed on sample deposition and increase on sample annealing to 10 K and 12 K, and remain almost unchanged under visible light irradiation ($\lambda = 495\text{--}600 \text{ nm}$), but decrease upon UV-visible light irradiation using the high pressure mercury arc lamp without filters ($\lambda = 250\text{--}600 \text{ nm}$). The band positions are only slightly blue- or red-shifted with respect to the corresponding modes of tris(benzene)strontium (within 10 cm^{-1} difference). The spectra from the isotopic-labeled experiments are shown in Figures S3 and S4. The band positions are listed in Table 2. No similar product absorptions are observed in the experiments with calcium.

We calculated the trisbenzene complexes $\text{M}(\text{Bz})_3$ ($\text{M} = \text{Ca}, \text{Sr}, \text{Ba}$) and analyzed their electronic structures with quantum chemical methods using density functional theory (DFT) at the M06-2X(D3)/def2-TZVPP level of theory, which considers also dispersion interactions. Details of the methods are given in Supporting Information. The calcium complex was also included in the theoretical work in order to provide a comparison with the experimentally observed strontium and barium adducts.

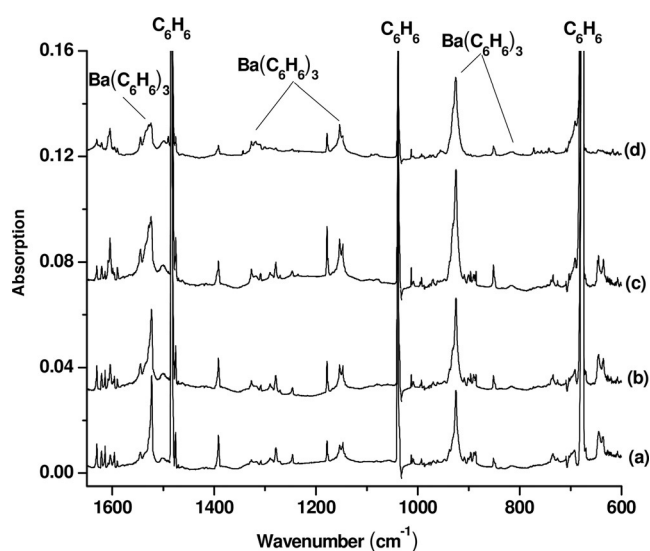


Figure 4. Infrared absorption spectra of barium–benzene complexes in the $1650\text{--}600 \text{ cm}^{-1}$ region from co-deposition of laser-evaporated barium atoms with 0.1% C_6H_6 in neon. a) 30 min of sample deposition at 4 K, b) after annealing at 10 K, c) after annealing to 12 K, and d) after 20 min of UV-visible light irradiation.

The geometry optimizations give D_3 symmetric $\text{M}(\text{Bz})_3$ equilibrium structures with a singlet (1A_1) electronic state as energy minima (Figure 5 a). The benzene ligands are distorted from their D_{6h} geometry in the free molecules to C_2 symmetric fragments with two C–C distances being significantly longer ($1.420\text{--}1.423 \text{ \AA}$) than in free benzene (1.388 \AA), while the other C–C bond lengths are only slightly altered. Calculations with enforced D_{3h} symmetry give structures being marginally higher in energy than the D_3 form (Figure 5 b). The D_{3h}

Table 2: Experimental infrared absorptions (cm^{-1}) of $\text{Ba}(\text{C}_6\text{H}_6)_3$, $\text{Ba}({}^{13}\text{C}_6\text{H}_6)_3$ and $\text{Ba}(\text{C}_6\text{D}_6)_3$ in solid neon and calculated values at the M06-2X(D3)/def2-TZVPP level (D_3 symmetry). Calculated IR intensities (km mol^{-1}) are given in parentheses.

Mode	Experimental					Calculated				
	${}^{12}\text{C}_6\text{H}_6$	${}^{13}\text{C}_6\text{H}_6$	$\Delta^{[a]}$	${}^{12}\text{C}_6\text{D}_6$	$\Delta^{[a]}$	${}^{12}\text{C}_6\text{H}_6$	${}^{13}\text{C}_6\text{H}_6$	$\Delta^{[a]}$	${}^{12}\text{C}_6\text{D}_6$	$\Delta^{[a]}$
C=C stretch	1529.0	1479.5	-49.5	1486.7	-42.3	1575.0 (651)	1523.7 (562)	-51.3	1523.1 (769)	-51.9
C=C stretch and C-D in-plane bend				1382.2					1438.9 (124)	
C=C stretch and C-D in-plane bend				1360.7					1386.8 (50.2)	
C=C stretch and C-H in-plane bend	1327.1	1281.9	-45.2			1360.3 (76)	1312.0 (70)	-48.3		
C=C stretch and C-H in-plane bend	1318.4	1272.2	-46.2			1343.6 (18)	1296.9 (17)	-46.7		
C-H(D) in-plane bend	1153.6	1145.0	-8.6	845.2	-308.4	1172.7 (152)	1163.1 (180)	-9.6	858.4 (40)	-314.3
C-H(D) in-plane bend	1084.6	1073.0	-11.6	838.7	-245.9	1144.1 (2)	1130.5 (3)	-13.6	853.8 (98)	-290.3
C-H(D) out-of-plane bend	932.2			766.9	-165.3	986.5 (8)			801.5 (70)	-185.0
Ring breath	925.8	894.8	-31	899.2	-26.6	969.4 (472)	937.0 (430)	-32.4	937.8 (224)	-31.6
C-H(D) out-of-plane bend	813.8	806.8	-7.0	665.1	-148.7	831.2 (11)	824.3 (11)	-6.9	659.4 (4)	-171.8

[a] Shift with respect to ${}^{12}\text{C}_6\text{H}_6$.

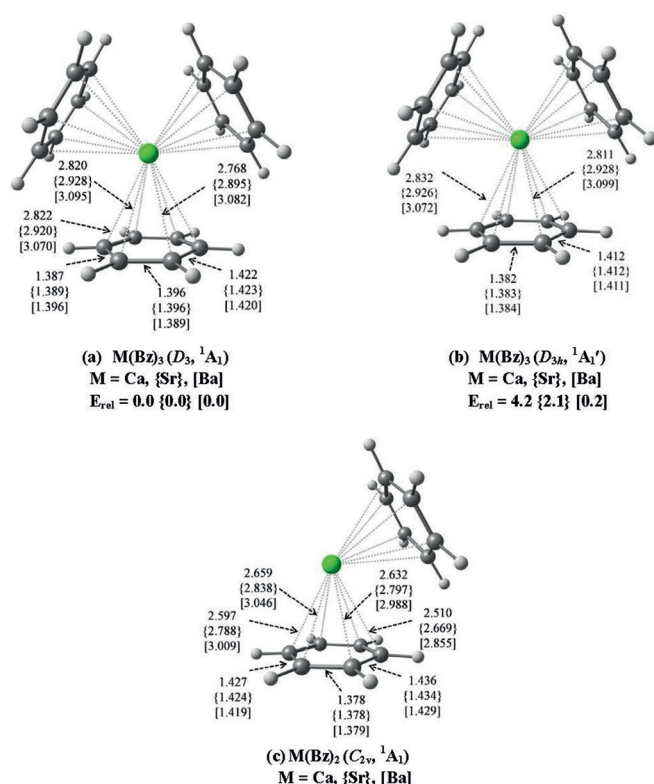


Figure 5. Calculated geometries of $\text{M}(\text{Bz})_3$ and $\text{M}(\text{Bz})_2$ complexes (M = Ca, Sr, Ba) at M06-2X-D3/def2-TZVPP. Bond distances are in Å. Relative energies are in kcal mol^{-1} .

structures of $\text{Ca}(\text{Bz})_3$ and $\text{Sr}(\text{Bz})_3$ have 3 imaginary frequencies, but the D_{3h} form of $\text{Ba}(\text{Bz})_3$ is calculated as energy minimum, which is only $0.2 \text{ kcal mol}^{-1}$ less stable than the D_3

structure. The D_{3h} structures of $\text{M}(\text{Bz})_3$ have four C-C bonds that are clearly longer (1.411 – 1.412 \AA) than in free benzene while two C-C bonds are a bit shortened (1.382 – 1.384 \AA). The corresponding triplet states of $\text{M}(\text{Bz})_3$ possess C_1 symmetry for M = Ca and C_2 symmetry for M = Sr, Ba and they are 0.6 – $3.8 \text{ kcal mol}^{-1}$ higher in energy than the singlet D_3 structures (see Figure S5).

Table 3 shows the bond dissociation energies (BDEs) of $\text{M}(\text{Bz})_3$ for the loss of one or all three benzene ligands. The calculations suggest that the BDE of one ligand ranges between $D_e = 19.4 \text{ kcal mol}^{-1}$ for M = Ca and $D_e = 26.1 \text{ kcal mol}^{-1}$ for M = Sr. The BDEs at room temperature become smaller after correcting for thermodynamic contributions, having values between $\Delta G^{298} = 6.5 \text{ kcal mol}^{-1}$ (Ca) and $\Delta G^{298} = 14.2 \text{ kcal mol}^{-1}$ (Sr). The total BDEs for fragmentation into three benzene molecules are between $D_e = 38.6 \text{ kcal mol}^{-1}$ (Ca) and $D_e = 46.0 \text{ kcal mol}^{-1}$ (Ba), but the free energies at room temperature are significantly smaller having values between $\Delta G^{298} = 8.7 \text{ kcal mol}^{-1}$ (Ca) and $\Delta G^{298} =$

Table 3: Calculated bond dissociation energies of $\text{M}(\text{Bz})_3$ complexes for the loss of one Bz and three Bz molecules at the M06-2X-D3/def2-TZVPP level. The D_e values give the electronic energies, the D_0 data include the corrections by zero-point vibrational frequencies and the $\Delta G^{298\text{K}}$ values consider the thermodynamic and vibrational corrections at room temperature.

Reaction	M	D_e	D_0	$\Delta G^{298\text{K}}$
$\text{M}(\text{Bz})_3$ (D_3 , ${}^1A_1'$) \rightarrow $\text{M}(\text{Bz})_2$ (C_{2v} , 1A_1) + Bz	Ca	19.4	18.3	6.5
	Sr	26.1	25.3	14.2
	Ba	24.0	23.6	13.4
$\text{M}(\text{Bz})_3$ (D_3 , ${}^1A_1'$) \rightarrow M (1S) + 3 Bz	Ca	38.6	38.3	8.7
	Sr	39.2	39.7	11.0
	Ba	46.0	47.1	19.5

19.5 kcal mol⁻¹ (Ba). Note that the latter process includes the electronic relaxation of the metal atoms to the ¹S electronic ground state. The lower BDE of the calcium complex may be the reason why Ca(Bz)₃ could not be observed experimentally.

The calculated infrared signals of the M(Bz)₃ complexes and their ¹³C and deuterated isotopes in the region > 600 cm⁻¹ as well as the frequency shifts are shown in Tables 1 and 2. The agreement of the computed unscaled values of the harmonic frequencies with the experimental values of the anharmonic mode is quite good. The calculated isotopic frequency shifts are also in very good accord with the recorded spectra. The comparison between the calculated and observed IR spectra leaves no doubt that the experimentally observed species are the neutral trisbenzene complexes M(Bz)₃ (M = Sr, Ba). The radical cations [M(Bz)₃]⁺ (M = Sr, Ba) were also calculated. The optimized structures are shown in Figure S6. Both cations are predicted to have D₃ symmetric structures with a doublet (²A₁) electronic ground state. The calculated vibrational frequencies and intensities are listed in Tables S1 and S2, respectively. Both cations are predicted to have two intense well-separated C=C stretching modes. Therefore, the assignment of the experimentally observed species to the radical cations can clearly be ruled out, as only one C=C stretching band is observed experimentally.

The complexes M(Bz)₃ (M = Ca, Sr, Ba) are formally 20-electron species possessing a singlet (¹A₁) electronic ground state. Our previously reported isoelectronic alkaline earth complexes M(CO)₈^[7] and M(N₂)₈^[11] are 18-electron adducts with cubic (O_h) symmetry that possess a triplet (³A_{1g}) ground state. The latter molecules have two singly occupied degenerate orbitals and one lower lying ligand-based molecular orbital (MO) with a_{2u} symmetry that has zero coefficient at the metal, because it has no valence AO having proper symmetry. We earlier reported the observation of the 20-electron transition metal octa-carbonyl anions [TM(CO)₈]⁻ (TM = Sc, Y, La).^[21] These complexes also have cubic symmetry but a singlet (¹A_{1g}) ground state, where the degenerate HOMO is fully occupied. The 18-electron rule first suggested by Langmuir^[22] remains valid in the above complexes, if only the metal–ligand bonding electrons are counted.

Figure 6a shows the orbital correlation diagram for the splitting of the (n)s(n)p(n-1)d AOs of alkaline earth atoms in the D_{3h} field of three benzene ligands along with plots of the associated occupied MOs of M(Bz)₃. In the D_{3h} symmetry field, the five d orbitals of metal split into three irreducible representations (a₁' + e' + e''). The a₁' orbital, which is the d₂ AO of the metal, provides backdonation to the antibonding π* orbital of the benzene ligands leading to the bonding MO (2a₁'). The remaining eight metal AOs serve as acceptor orbitals for the charge donation from the bonding π orbitals of the benzene ligands. There are nine symmetry-adapted linear combination orbitals from the bonding π orbitals of three benzene ligands. The 1a₂' MO remains as a ligand-only orbital, because there is no valence AO of the metal possessing this symmetry. All of these bonding and non-bonding orbitals are doubly occupied, resulting in a closed-

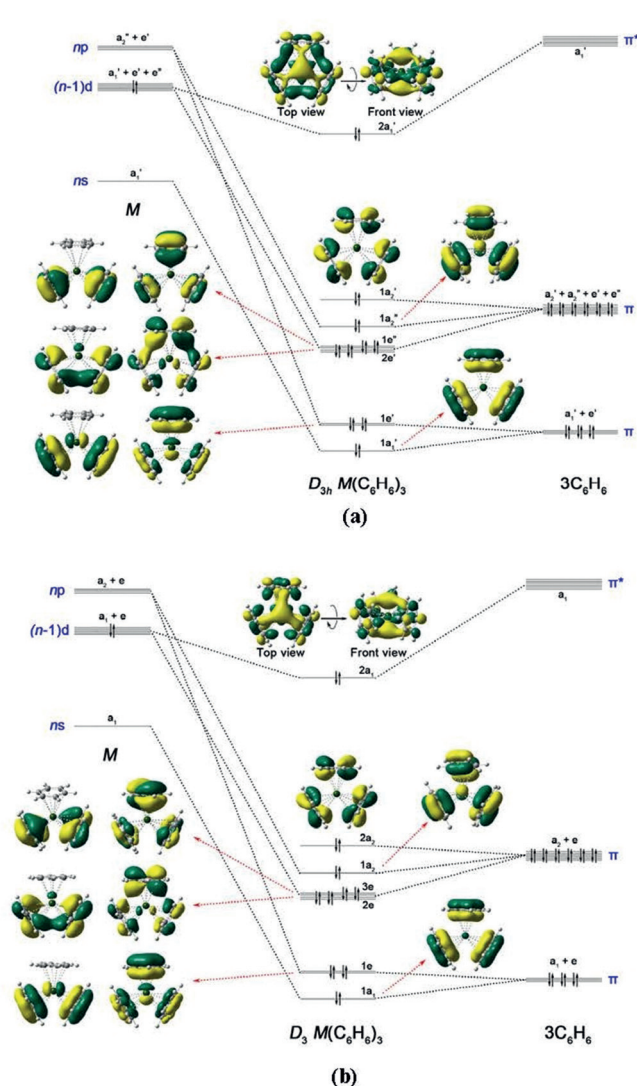


Figure 6. MO correlation diagram of the valence (n)s(n)p(n-1)d AOs of M and (C₆H₆)₃ in a) D_{3h} symmetry and b) D₃ symmetry. Shapes of the MOs of Ba(Bz)₃ are also shown.

shell singlet ground state. Thus, the M(Bz)₃ complexes have a total of 20 valence electrons, but the metal centers have only 18 effective valence electrons, and thus still fulfill the 18-electron rule if only the metal–ligand bonding electrons are considered.

Figure 6b shows the related orbital diagram for M(Bz)₃ in the D₃ field of the benzene ligands at the equilibrium geometries and the shape of the occupied MOs of M(Bz)₃. Comparing the latter MOs with the D_{3h} orbitals shows only negligible differences. The 2a₂ MO of the D₃ form takes the position of the 2a₂' MO of the D_{3h} form and could, in principle, mix with one (n)p AO of the metal. Inspection of the MO coefficients shows that the contribution of the (n)p AO of the metal is nearly zero. Thus, the pattern of the orbital interaction of the D₃ structures is very similar to the D_{3h} form. We think that the deviation from perfect D_{3h} geometry of the M(Bz)₃ complexes is due to weak steric interactions between the benzene ligands, which have a negligible effect on the orbital interactions. This corresponds to the reduction

Table 4: Results of EDA-NOCV calculations at BP86-D3(BJ)/TZ2P for $M(\text{Bz})_3$ (D_3 , 1A_1) complexes using neutral atoms M in the singlet state with $n^0n^0p^0(n-1)d^2$ electron configuration and $(\text{Bz})_3$ (singlet) as interacting fragments.

Term	Interaction	Ca (S, $4s^04p^03d^2$) + $(\text{Bz})_3$ (S)	Sr (S, $5s^05p^04d^2$) + $(\text{Bz})_3$ (S)	Ba (S, $6s^06p^05d^2$) + $(\text{Bz})_3$ (S)
ΔE_{int}		-197.4	-191.5	-135.5
ΔE_{Pauli}		124.5	140.5	156.8
$\Delta E_{\text{disp}}^{[a]}$		-7.1 (2.2%)	-8.4 (2.5%)	-16.8 (5.7%)
$\Delta E_{\text{elstat}}^{[a]}$		-82.7 (25.7%)	-88.8 (26.8%)	-116.4 (39.8%)
$\Delta E_{\text{orb}}^{[a]}$		-232.1 (72.1%)	-234.7 (70.7%)	-159.1 (54.4%)
$\Delta E_{\text{orb}(1)}^{[b]}$ ($2a_1$)	$(\text{Bz})_3 \leftarrow M(d)$ backdonation	-191.8 (82.6%)	-189.9 (80.9%)	-119.9 (75.4%)
$\Delta E_{\text{orb}(2)}^{[b]}$ ($3e$)	$(\text{Bz})_3 \rightarrow M(d)$ donation	-14.0 (6.0%)	-16.8 (7.2%)	-16.8 (10.6%)
$\Delta E_{\text{orb}(3)}^{[b]}$ ($2e$)	$(\text{Bz})_3 \rightarrow M(d)$ donation	-13.6 (5.9%)	-15.0 (6.4%)	-11.8 (7.4%)
$\Delta E_{\text{orb}(4)}^{[b]}$ ($1a_1$)	$(\text{Bz})_3 \rightarrow M(s)$ donation	-2.7 (1.2%)	-2.9 (1.2%)	-2.1 (1.3%)
$\Delta E_{\text{orb}(5)}^{[b]}$ ($1a_2$)	$(\text{Bz})_3 \rightarrow M(p)$ donation	-1.4 (0.6%)	-1.3 (0.6%)	-0.8 (0.5%)
$\Delta E_{\text{orb}(6)}^{[b]}$ ($1e$)	$(\text{Bz})_3 \rightarrow M(p)$ donation	-3.2 (1.4%)	-2.4 (1.0%)	-0.8 (0.5%)
$\Delta E_{\text{orb}(7)}^{[b]}$ ($1a_2$)	Polarization	-1.3 (0.6%)	-1.8 (0.8%)	-2.5 (1.6%)
$\Delta E_{\text{orb}(\text{rest})}^{[b]}$		-4.1 (1.8%)	-4.6 (2.0%)	-4.4 (2.8%)

[a] The values in parentheses give the percentage contribution to the total attractive interactions $\Delta E_{\text{elstat}} + \Delta E_{\text{orb}} + \Delta E_{\text{disp}}$. [b] The values in parentheses give the percentage contribution to the total orbital interactions ΔE_{orb} .

of the energy difference between the D_3 and D_{3h} forms of $M(\text{Bz})_3$ having the order of $\text{Ca} > \text{Sr} > \text{Ba}$, because the metal–ligand bonds and thus the ligand–ligand distances increase (Figure 5).

On the basis of above bonding analysis, it is clear that the $M(\text{Bz})_3$ complexes can be regarded as being formed via the interactions between the metal atoms in the singlet excited state with an $(n-1)d^2ns^0np^0$ electron configuration and three benzene ligands in the singlet state. The strength of the pairwise orbital interactions can be estimated with the EDA-NOCV (energy decomposition analysis in combination with natural orbitals for chemical valence) method.^[23] Details of the method have been described in the literature.^[24–26] The calculations at M06-2X with the Slater basis functions TZ2P of the ADF program for M in the singlet excited state with d^2 configuration did not converge. Therefore, we used the BP86-D3(BJ) functional in conjunction with the TZ2P basis set. Our experience in numerous studies has shown that the EDA-NOCV calculations are not very sensitive to the functional.^[27–37] Table 4 gives the numerical results of the calculations of $M(\text{Bz})_3$ ($M = \text{Ca}, \text{Sr}, \text{Ba}$) at the D_3 equilibrium structures.

The data in Table 4 suggest that the $M-(\text{Bz})_3$ attractive interactions come mainly from the covalent term ΔE_{orb} , which provides between 72% ($M = \text{Ca}$) and 54% ($M = \text{Ba}$) to the total attraction. The breakdown of ΔE_{orb} into pairwise orbital interactions shows that the dominant contribution $\Delta E_{\text{orb}(1)}$ comes from the backdonation of the occupied $(n-1)d$ AO of the metal into the vacant π^* orbitals of the ligands denoted as $(\text{Bz})_3 \leftarrow M(d)$. This orbital interaction contributes between 83% ($M = \text{Ca}$) and 75% ($M = \text{Ba}$) to ΔE_{orb} . The following terms $\Delta E_{\text{orb}(2)}$ and $\Delta E_{\text{orb}(3)}$ also involve the $(n-1)d$ AOs of the metals. They are due to the donation from two sets of degenerate occupied π MOs of the ligands into vacant $(n-1)d$ metal AOs denoted as $(\text{Bz})_3 \rightarrow M(d)$. Thus, the valence AOs of the metals that contribute to the metal–ligand bonding are dominated by the $(n-1)d$ orbitals.

The assignment of the $\Delta E_{\text{orb}(1)}$ interaction to $(\text{Bz})_3 \leftarrow M(d)$, which has a_1 symmetry, may be questioned because the MO diagram shows that the metal $(n)s$ AO correlates in a D_3 environment also with a_1 symmetric MOs (Figure 6b).

Inspection of the related deformation density $\Delta\rho_1$ and the associated orbitals reveals that the contribution of the metal $(n)s$ AO to the $2a_1$ MO is negligible. Also, the EDA-NOCV calculation of the D_{3h} structures of $M(\text{Bz})_3$ gives values that are very similar to the D_3 analysis. The numerical results of the EDA-NOCV calculation using the D_{3h} structures are shown in Table S3. Figure 7 displays the shape of the deformation densities $\Delta\rho_1 - \Delta\rho_7$ of $\text{Ba}(\text{Bz})_3$ which nicely illustrate the charge alteration due to the pairwise orbital interactions $\Delta E_{\text{orb}(1)} - \Delta E_{\text{orb}(7)}$. The color code of the charge flow is red \rightarrow blue. The deformation densities $\Delta\rho_1 - \Delta\rho_7$ of $M(\text{Bz})_3$ calculated with D_{3h} symmetry are displayed in Figures S7–S9. The shape of $\Delta\rho_7$ clearly reveals that the small stabilization energy $\Delta E_{\text{orb}(7)}$ is solely due to the polarization of the ligand orbitals. The deformation densities $\Delta\rho_1 - \Delta\rho_7$ of $\text{Ca}(\text{Bz})_3$ and $\text{Sr}(\text{Bz})_3$ closely resemble those of $\text{Ba}(\text{Bz})_3$. They are shown in Figures S10 and S11.

One referee pointed out that the choice of the electronic reference states of the fragments is somewhat arbitrary and that the same state can be obtained from a ground state metal atom and excited benzene molecules, which would yield a different pattern of donation and back-donation. This is a justified objection, which needs to be addressed. The electronic ground state of the metal atoms has the electron configuration $(n)s^2$, which transforms to a_1' in D_{3h} symmetry. But the shape of the occupied $1a_1'$ MO of the complex $\text{Ba}(\text{Bz})_3$ shown in Figure 6a reveals that the $6s$ AO of Ba has a negligible coefficient in the molecular orbital of the complex, which has only slightly distorted ligand orbitals. The same situation is found in the D_3 symmetric structure (Figure 6b). The visual impression is supported by EDA-NOCV calculations of the complexes $M(\text{Bz})_3$ using the atoms M in the electronic ground state configuration $(n)s^2$ and $(\text{Bz})_3$ in the doubly excited singlet state with two electrons in the π^* MO. It has been shown by us in numerous studies that those fragments, which give the smallest absolute values for the orbital term ΔE_{orb} , are the most suitable species to represent the bonding of the complex, since they change the least in the formation of the bond. This method has proven to be very helpful for a large number of chemical compounds whose

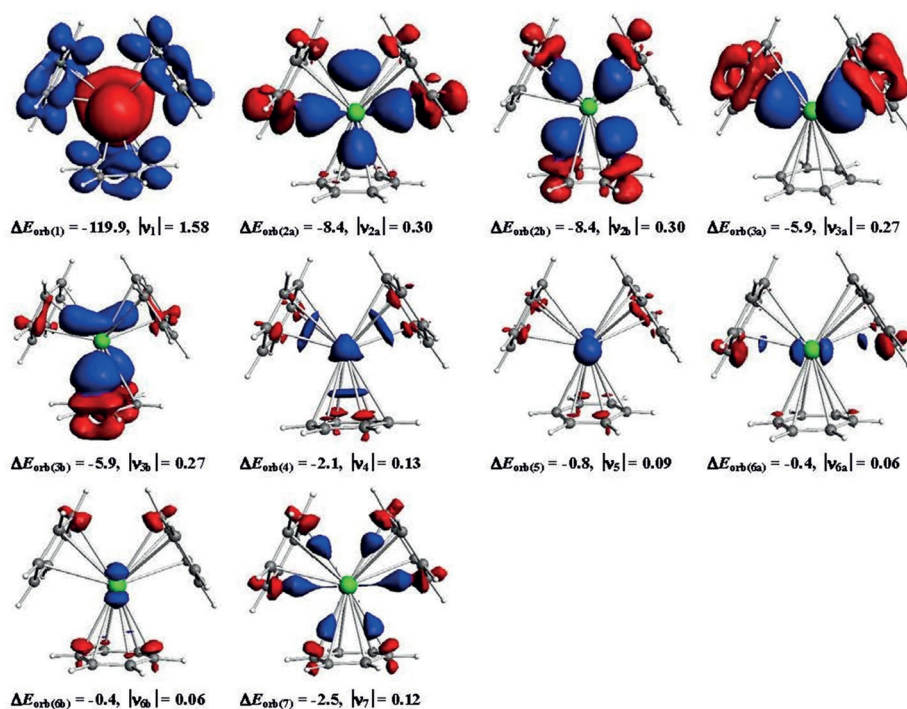


Figure 7. Shape of the deformation densities $\Delta\rho_{(1)-(7)}$, which are associated with the orbital interactions $\Delta E_{orb(1)-(7)}$ between neutral fragments Ba and $(Bz)_3$ in $Ba(Bz)_3$ ($D_3, ^1A_1$) complex, using Ba ($S, 6s^06p^05d^2$) + $(Bz)_3$ (S) as interacting complexes (Table 4), and eigenvalues $|v_n|$ of the charge flow. The color code of the charge flow is red \rightarrow blue. The isosurface values are 0.002 for $\Delta\rho_{(1)}$, 0.0005 for $\Delta\rho_{(2)-(4)}$ and 0.0004 for $\Delta\rho_{(5)-(7)}$.

nature is not directly apparent.^[27–37] Tables S4–S6 show that the ΔE_{orb} values using atom M in the electronic ground state and $(Bz)_3$ in the excited state are significantly larger than those in Table 4.

However, there is an alternative option for the choice of the interacting fragments. The calculated partial charges of $M(Bz)_3$ using the NBO 6.0 method^[38] suggest large positive charges for the metal atoms of 1.47 for Ca, 1.46 for Sr and 1.40 for Ba. The choice of neutral fragments in the electronic references states for $M(Bz)_3$ follows the Dewar–Chatt–Duncanson (DCD) model for transition metal compounds, which considers the metal atom and the ligands prior to bond

formation.^[39] But the formation of chemical bonds can lead to large charge transfers, resulting in an electronic structure of the molecule that differs greatly from the initial components.

We carried out EDA-NOCV calculations using charged fragments as interacting moieties. We found that the lowest ΔE_{orb} value is given when the singly charged species M^+ in the doublet state with the electron configuration $(n-1)d^1ns^0np^0$ and doublet $(Bz)_3^-$ are employed as interacting fragments. Fragments with other charges or electron configurations give larger ΔE_{orb} values (see the numerical results in Tables S4–S6 in Supporting Information). Table 5 shows the numerical EDA-NOCV results for the $M^+-(Bz)_3^-$ interactions.

Table 5: Results of EDA-NOCV calculations at BP86-D3(BJ)/TZ2P for $M(Bz)_3$ ($D_3, ^1A_1$) complexes using positively charged atoms M^+ in the doublet state with $ns^0np^0(n-1)d^1$ electron configuration and negatively charged $(Bz)_3^-$ (D) as interacting fragments.

Term	Interaction	Ca ⁺ ($D, 4s^04p^03d^1$) + $(Bz)_3^-$ (D)	Sr ⁺ ($D, 5s^05p^04d^1$) + $(Bz)_3^-$ (D)	Ba ⁺ ($D, 6s^06p^05d^1$) + $(Bz)_3^-$ (D)
ΔE_{int}		–232.9	–228.9	–207.0
ΔE_{Pauli}		71.4	85.6	110.1
$\Delta E_{disp}^{[a]}$		–7.1 (2.3%)	–8.4 (2.7%)	–16.8 (5.3%)
$\Delta E_{elstat}^{[a]}$		–155.0 (50.9%)	–159.0 (50.6%)	–176.7 (55.8%)
$\Delta E_{orb}^{[a]}$		–142.2 (46.7%)	–147.0 (46.8%)	–123.4 (38.9%)
$\Delta E_{orb(1)}^{[b]}$ ($2a_1$)	$(Bz)_3 \rightarrow M(d)$ electron-sharing	–70.9 (49.9%)	–76.7 (52.2%)	–53.7 (43.5%)
$\Delta E_{orb(2)}^{[b]}$ ($3e$)	$(Bz)_3 \rightarrow M(d)$ donation	–20.4 (14.3%)	–22.0 (15.0%)	–24.2 (19.6%)
$\Delta E_{orb(3)}^{[b]}$ ($2e$)	$(Bz)_3 \rightarrow M(d)$ donation	–19.4 (13.6%)	–19.4 (13.2%)	–18.4 (14.9%)
$\Delta E_{orb(4)}^{[b]}$ ($1a_1$)	$(Bz)_3 \rightarrow M(s)$ donation	–5.5 (3.9%)	–5.1 (3.5%)	–4.0 (3.2%)
$\Delta E_{orb(5)}^{[b]}$ ($1a_2$)	$(Bz)_3 \rightarrow M(p)$ donation	–3.3 (2.3%)	–2.7 (1.8%)	–2.4 (1.9%)
$\Delta E_{orb(6)}^{[b]}$ ($1e$)	$(Bz)_3 \rightarrow M(p)$ donation	–6.6 (4.6%)	–3.0 (2.0%)	–3.8 (3.1%)
$\Delta E_{orb(7)}^{[b]}$ ($1a_2$)	Polarization	–2.6 (1.8%)	–3.0 (2.0%)	–4.3 (3.5%)
$\Delta E_{orb(rest)}^{[b]}$		–13.5 (9.5%)	–15.1 (10.3%)	–12.6 (10.2%)

[a] The values in parentheses give the percentage contribution to the total attractive interactions $\Delta E_{elstat} + \Delta E_{orb} + \Delta E_{disp}$. [b] The values in parentheses give the percentage contribution to the total orbital interactions ΔE_{orb} .

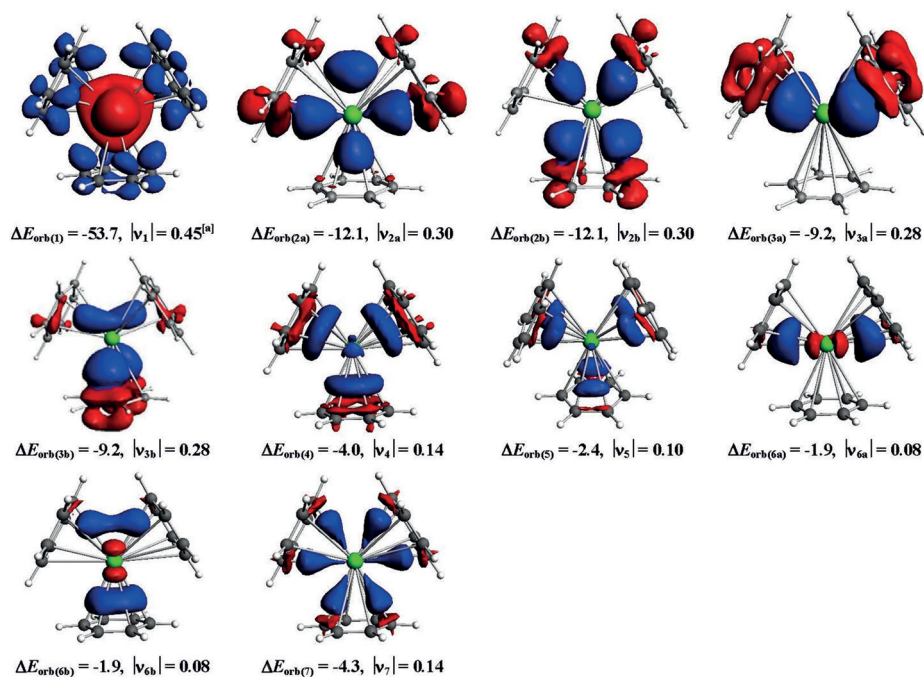


Figure 8. Shape of the deformation densities $\Delta\rho_{(1)-(7)}$, which are associated with the orbital interactions $\Delta E_{\text{orb}(1)-(7)}$ in $\text{Ba}(\text{Bz})_3$ ($D_3, {}^1A_1$) complex, using Ba^+ ($D, 6s^06p^05d^1$) + $(\text{Bz})_3^-$ (D) as interacting complexes (Table 5), and eigenvalues $|v_n|$ of the charge flow. The color code of the charge flow is red \rightarrow blue. The isosurface values are 0.001 for $\Delta\rho_{(1)}$, 0.0005 for $\Delta\rho_{(2)-(4)}$ and 0.0004 for $\Delta\rho_{(5)-(7)}$. ^aNet charge flow of the electron-sharing interactions.

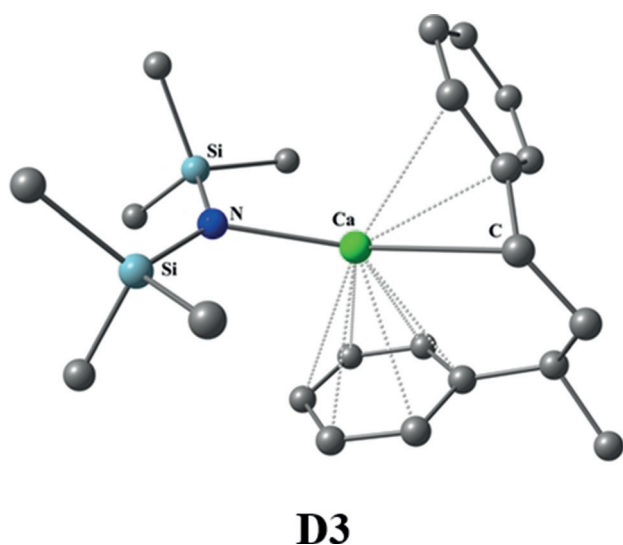
The data in Table 5 show that the covalent term ΔE_{orb} is now a bit smaller than the Coulombic attraction ΔE_{elstat} . The strongest pairwise orbital interaction $\Delta E_{\text{orb}(1)}$ between the fragments in the electronic doublet state comes from the electron-sharing bond formation between the unpaired electrons of M^+ and $(\text{Bz})_3^-$. It is interesting to see that the associated deformation density $\Delta\rho_1$ shown in Figure 8 exhibits a very similar shape as the dative interaction of $\Delta E_{\text{orb}(1)}$ when neutral fragments are used (Figure 7). It means that the electron-sharing interactions between the $(n-1)d$ electron of M^+ and the unpaired electron of $(\text{Bz})_3^-$ have a polarity from the metal towards the ligands. As can be expected, the associated eigenvalue v_1 , which indicates the magnitude of the net charge flow, is much smaller for the charged fragments than for the neutral fragments. The relative contribution of the $2a_1$ orbital interaction between the charged fragments is smaller than the dominating $2a_1$ term of the neutral fragments, whereas the $2e$ and $3e$ orbital terms that come from $(\text{Bz})_3 \rightarrow \text{M}(d)$ π donation are larger when charged fragment are used. But the most relevant results is the finding that the $(n-1)d$ AOs of the metal are still clearly the most important valence orbitals for covalent metal- $(\text{Bz})_3$ binding even when charged fragments are used for the EDA-NOCV calculations.

The two sets of EDA-NOCV results in Tables 4 and 5 address two different questions. The dissociation of $\text{M}(\text{Bz})_3$ gives neutral atoms M and three neutral ligands Bz and the data in Table 4 provide information about the total charge flow taking place during bond formation. The latter results directly refer to the DCD model and to orbital correlation diagrams, which always take the orbitals of the fragments prior to bond formation for explaining molecular structures

and reactivities.^[40] The data in Table 5 give information about the best description of the bonding situation in the eventually formed complexes, which is given in terms of interactions between the metal ion M^+ and the negatively charged ligand cave $(\text{Bz})_3^-$. EDA-NOCV calculations using the doubly charged fragments M^{2+} and $(\text{Bz})_3^{2-}$ give larger ΔE_{orb} values (Tables S4–S6, last columns) than singly charged fragments. The results in Tables 4 and 5 demonstrate the variability of the EDA-NOCV method, which provides quantitative answer to the questions about the binding interactions between fragments prior and after the bond formation. The difference between the two situations is sometimes not recognized, which has led to a misunderstanding of the EDA-NOCV results by some authors.^[41]

Benzene is an excellent ligand in the organometallic chemistry of low-valent transition metals. Since the first synthesis of bis(benzene)chromium $\text{Cr}(\text{Bz})_2$,^[18] a large number of sandwich or half-sandwich transition metal compounds involving one or two benzene ligands have been reported.^[42] Although main group metal cation–benzene complexes^[43] as well as organometallic complexes of the alkaline earth cations such as metallocene complexes that feature ionic bonding with no propensity for π backbonding are quite common,^[44–48] sandwich complexes of neutral alkaline earth metals, in particular with three η^6 -bound benzene ligands, are unprecedented. In the literature, only very few metallocene complexes of actinides involving three η^5 -bound cyclopentadienyl ligands have been reported.^[49] The successful synthesis of trisbenzene complexes of heavier alkaline earth metals may open the door to another chapter of metal complexes.

The observation of the alkaline earth trisbenzene complexes $M(\text{Bz})_3$ is not just a surprising finding of some exotic species. It seems that the bonding situation of the arene complexes plays also an important role in some intermediates of alkaline earth catalyzed reactions. Alonso, Harder and co-workers recently reported the alkene hydrogenation catalyzed by simple compounds of the alkaline earth elements Ca, Sr, Ba.^[4] The calculated reaction profile identified a species **D3** where Ca is sandwiched between two aromatic rings (Figure 9). The structure of **D3** closely resembles the metal–ligand bonds in $M(\text{Bz})_3$. A similar situation was reported by the group for the alkaline earth catalyzed imine hydrogenation.^[5] The involvement of the valence-d orbitals of Ca may also be involved in the recently reported organocalcium-mediated nucleophilic alkylation of benzene.^[3] The transition-metal like chemical reactivity of Ca, Ba, Sr is perhaps more relevant for synthetic chemistry than hitherto thought.



D3

Figure 9. Calculated structure **D3** of a calcium compound sandwiched between two aromatic rings, which is relevant in the hydrogenation of alkene catalyzed by alkaline earth atoms. The Figure is adapted from reference 4.

Conclusion

In summary, we report the synthesis and spectroscopic identification of the unprecedented trisbenzene complexes $M(\text{Bz})_3$ ($M = \text{Sr}, \text{Ba}$) featuring three equivalent η^6 -bound benzene ligands in low-temperature Ne matrix. The analysis of the electronic structure shows that the complexes exhibit metal–ligand bonds that are typical for transition metal compounds. The chemical bonds can be explained in terms of weak donation from the π MOs of benzene ligands into the vacant $(n-1)d$ AOs of M and strong backdonation from the occupied $(n-1)d$ AO of M into vacant π^* MOs of benzene ligands. The metals in these 20-electron complexes have 18 effective valence electrons, and thus still fulfill the 18-electron rule, because two valence electrons occupy a ligand-based MO that has zero coefficient at the metal. The results suggest

that the heavier earth alkaline atoms Ca, Sr, Ba may exhibit the full bonding scenario of transition metals.

Acknowledgements

The experimental work was supported by the National Natural Science Foundation of China (grant numbers 21688102 and 21433005). L.Z. and G.F. acknowledge financial support from Nanjing Tech University (grant numbers 39837123, 39837132) and SICAM Fellowship from Jiangsu National Synergetic Innovation Center for Advanced Materials, Natural Science Foundation of Jiangsu Province for Youth (grant number BK20170964), National Natural Science Foundation of China (grant number 21703099). S.P. thanks Nanjing Tech University for a postdoctoral fellowship and the High Performance Center of Nanjing Tech University for the computational resources. Y.W. thanks National Natural Science Foundation of China (grant number 21273140) for financial support and the High Performance Center of Shanxi University for the computational resources.

Conflict of interest

The authors declare no conflict of interest.

Keywords: 18-electron rule · alkaline earth metals · benzene complexes · bonding analysis · matrix isolation spectroscopy

How to cite: *Angew. Chem. Int. Ed.* **2019**, *58*, 17365–17374
Angew. Chem. **2019**, *131*, 17526–17535

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Manuscript received: July 10, 2019

Revised manuscript received: August 6, 2019

Accepted manuscript online: September 9, 2019

Version of record online: October 17, 2019