

# Effect of Geometric Parameters on the Performance of P-Type Junctionless Lateral Gate Transistors



Farhad Larki<sup>1\*</sup>, Arash Dehzangi<sup>1</sup>, Sawal Hamid Md Ali<sup>1</sup>, Azman Jalar<sup>1</sup>, Md. Shabiul Islam<sup>1</sup>, Mohd Nizar Hamidon<sup>2</sup>, Burhanuddin Yeop Majlis<sup>1</sup>

1 Institute of Microengineering and Nanoelectronics (IMEN), Universiti Kebangsaan Malaysia, Bangi, Selangor, Malaysia, 2 Functional Devices Laboratory, Institute of Advanced Technology, Universiti Putra Malaysia, Serdang, Selangor, Malaysia

#### **Abstract**

This paper examines the impact of two important geometrical parameters, namely the thickness and source/drain extensions on the performance of low doped p-type double lateral gate junctionless transistors (DGJLTs). The three dimensional Technology Computer-Aided Design simulation is implemented to calculate the characteristics of the devices with different thickness and source/drain extension and based on that, the parameters such as threshold voltage, transconductance and resistance in saturation region are analyzed. In addition, simulation results provide a physical explanation for the variation of device characteristics given by the variation of geometric parameters, mainly based on investigation of the electric field components and the carries density variation. It is shown that, the variation of the carrier density is the main factor which affects the characteristics of the device when the device's thickness is varied. However, the electric field is mainly responsible for variation of the characteristics when the source/drain extension is changed.

Citation: Larki F, Dehzangi A, Ali SHM, Jalar A, Islam MS, et al. (2014) Effect of Geometric Parameters on the Performance of P-Type Junctionless Lateral Gate Transistors. PLoS ONE 9(4): e95182. doi:10.1371/journal.pone.0095182

Editor: Sefer Bora Lisesivdin, Gazi University, Turkey

Received December 12, 2013; Accepted March 25, 2014; Published April 17, 2014

**Copyright:** © 2014 Larki et al. This is an open-access article distributed under the terms of the Creative Commons Attribution License, which permits unrestricted use, distribution, and reproduction in any medium, provided the original author and source are credited.

**Funding:** Funding provided by "DANA PEMBANGUNAN PENYELIDIKAN, DPP-2013-061" and "DANA IMPAC PERDANA, DIP-2012-16". The funders had no role in study design, data collection and analysis, decision to publish, or preparation of the manuscript.

Competing Interests: The authors have declared that no competing interests exist.

\* E-mail: farhad.larki@gmail.com

## Introduction

As the conventional planar metal-oxide-semiconductor fieldeffect transistors (MOSFETs) dimensions scale down to tens of nanometers, the formation of sharp source/channel and channel/ drain junctions becomes crucial for suppressing the issues such as leakage current and short-channel effect (SCE). In nano scale devices, on the other hand, the formation of the ultra sharp junctions imposes severe challenges on doping techniques due to the difficulty to control the distribution of dopants at the metallurgical junctions and the intrinsic discreteness of the dopant itself. Recently, a novel device called junctionless transistor (JLT) has gain significant attention due to the advantages associated to the specific design of the device and simplification of the fabrication process [1,2,3]. From the structural point of view, the JLTs are heavily doped gated resistors with narrow silicon body in multiple gate architecture. Operationally, the JLTs are volume depleted in the off state (at zero gate bias) due to the effective work-function difference between the gate and the channel [4]. Conduction mechanism in JLTs only occurs in the bulk of the channel. Besides, no surface accumulation layer is formed below the threshold voltage [5]. The freedom of JLTs from different types of doping and associated gradients is considered as the main advantage of this design which simplifies the fabrication process. In addition, in subthreshold region device can obtain enhanced immunity from short channel effects in comparison to conventional inversion mode devices [6]. Performance comparison between bulk and SOI JLTs [7,8], temperature dependence [9], scalability [10,11], ballistic transport [12], and analog and digital

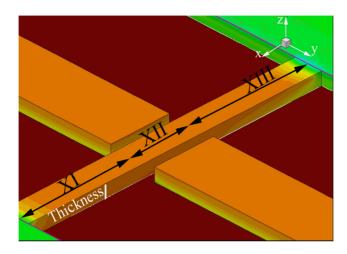
applications of JLTs [13,14] have been investigated in the literature.

Recently, the researchers of the present study have reported the fabrication and experimental characterization of low doped p-type junctionless transistors with single and double lateral gate(s) through the unconventional method of atomic force microscopy (AFM) nanolithography [15,16,17]. Simulations of carrier's transport and pinch off mechanism of the device have been presented in Refs. [18,19]. Computational studies are absolutely critical to predict the ultimate performance of the nanoscale devices and to impel forward future experiments or technological developments. As a result, in this work, the impact of the two important geometric parameters to the operations of the p-type double lateral gate junctionless transistors (DGJLT) structure is comprehensively explored. First, the device simulation structures and models implemented in the present research are presented. Then, the operation and performance of the device in on and off state is briefly explained. Finally, the impact of geometric parameters of the device and particularly active layer thickness and source/drain extensions on characteristics of the device, are investigated based on the variation of hole density and electric field components.

## **Methods**

1

The structure of a typical DGJLT used for the simulations is shown in Fig. 1. Parameters used for the device simulations are listed in Table 1. The simulated structures have a uniform low doping concentration (Boron 10<sup>15</sup>) throughout all active regions of



**Figure 1. Isometric view of simulated DGJLT.** doi:10.1371/journal.pone.0095182.g001

the device. The complete structure sits on a 145 nm ideal oxide. Different zones  $(X_{\rm I},~X_{\rm II},~{\rm and}~X_{\rm III})$  correspond to the source extension, the channel and the drain extension, respectively are indicated for enhancement of the discussion and will be referred throughout the discussion.

The models used for the simulations are calibrated against a long-channel experimentally demonstrated DGJLT, as described in [18,19]. We use a hydrodynamic model in Sentaurus software D-2010.03 [20] as the platform for the 3-D TCAD simulation presented in this study. Besides the fundamental equations, the doping-dependent Masetti mobility model and the default Shockley–Read–Hall (SRH) recombination-generation model is activated. SRH recombination-generation is applied in order to consider the leakage current and recombination through deep defect levels in the gap [21].

## **Results and Discussion**

## On and off state behavior of the device

Fig. 2 shows the simulated transfer characteristics for the nominal DGJLT with 100 nm thickness and width for  $V_D = -0.05 \text{ V}$  and -1.0 V. Zones  $X_I$ ,  $X_{II}$ , and  $X_{III}$  have the doping density of  $10^{15}$  cm<sup>-3</sup>, and the dimensions are considered as 2  $\mu$ m, 200 nm, and 2  $\mu$ m, respectively. The nominal simulated device is considered analogous to the fabricated device in order to give the basic view of the device operation. It can be seen that the device is in on state at zero gate voltage with  $I_{on}/I_{off}$  ratio of  $10^7$ . The device is driven into off state by increasing the positive gate voltage. However, negative gate voltage is not able to increase the current

Table 1. Simulated devices parameters.

Parameters	Value	
Device layer thickness	20, 40, 100 nm	
Zones X <sub>I</sub> and X <sub>III</sub> length	500 nm, 1 μm and 2 μm	
Zone X <sub>II</sub> length	200 nm	
Contact work function	5.12 eV	
Gate voltage	-2 V to +2 V	
Drain voltage	−0.05 V to −1 V	

doi:10.1371/journal.pone.0095182.t001

significantly. Since the device is in on state and already near the flatband condition at zero gate voltage, it can be claimed that the device is principally working similar to the pinch-off transistors [22,23]. In the on state, the majority carriers moves in the volume of the channel and the depletion mechanism of the device due to the positive gate voltage applied to the lateral gates starts from the bottom corners of the channel at the Si/BOX interface, face to the lateral gates, and expands toward the center and top surface of the channel [18]. The electrostatic squeezing of the channel creates a large barrier along the holes transmission path. Increasing the drain voltage creates a depletion layer due to the sweeping of the carriers and at very high drain voltage, this depletion area along the electric field works as a barrier and prevents the electrical field from the propagating into the channel. This creates the current saturation, even for  $V_G = 0 \text{ V}$ .

### Influence of thickness

In the p-type DGILTs, the conduction mechanism is through the current of majority carriers (holes) flowing in the bulk of the channel when the device is in the on state, and this volume as a conduction path would be gradually shrunk from bottom surface at Si/BOX interface to the top of the channel by increasing the positive gate voltage to turn the device off. This implies that the thickness has a strong influence on the characteristics and switch behavior of the DGILTs. In order to investigation the effect of thickness on the behavior of the DGILTs, the structures with thickness of 20 and 40 nm were simulated and the results compared to the nominal simulated device with 100 nm thickness. To avoid any quantum mechanical confinement effects, the thickness lower than 20 nm has not been investigated in this study. The dependence of transfer characteristics of DGILT on thickness variation, when the drain is biased by  $V_D = -1.0 \text{ V}$  is shown in Fig. 3. The inset shows the variation of on and off current with the thickness. As it is expected, the device with 20 nm thickness provides the highest  $I_{on}/I_{off}$  ratio ( $\approx 5 \times 10^8$ ) and a small decrease of the on state drain current is obtained for thinner devices.

The threshold voltage variation over the change of the thickness  $(dV_{th}/dT_{Si}\approx 5~mV/nm)$  was negligible, mainly due to the suppression of random impurity fluctuation in these low doped devices [24]. The low doped channel in thin SOI devices has been implemented experimentally to minimize threshold voltage variations caused by random impurity fluctuation effects [25]. Another source of threshold voltage fluctuation is the scattering of source/drain doping atoms into the channel. This source of

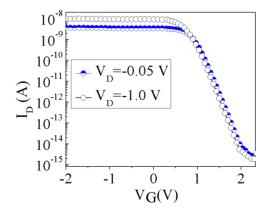


Figure 2. Transfer characteristics simulations of DGJLT with 100 nm width and thickness, nanowire length of 4.2  $\mu\text{m}$ , and channel length of 200 nm.

doi:10.1371/journal.pone.0095182.g002

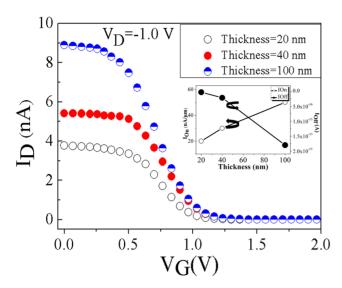


Figure 3. Transfer characteristics simulations of DGJLTs with three different thickness (20, 40, 100 nm). The inset plots the  $\emph{on}$  (V $_G$ =0 V) and  $\emph{off}$  (V $_G$ =2.0 V) current variation with thickness, at V $_D$ =-1.0 V.

doi:10.1371/journal.pone.0095182.g003

threshold voltage fluctuation is eliminated in DGJLTs because of the absence of junctions and doping gradients [26]. Based on the characteristics of the devices, we try to extract some of the important parameters in order to describe the variation of the device operation with respect to the thickness.

The transconductance  $(g_m)$  of the DGJLTs with three different thicknesses as a function of gate voltage is shown in Fig. 4a. The inset shows the variation of the  $g_m$  as a function of thickness at three different gate voltages. The overall trend demonstrates the variation of all devices from on to off state. Sharp decrease of transconductance shows the effect of positive gate voltage to deplete the channel according to the pinch off mechanism. Moreover, the lower transconductance in thinner device is attributed to the scarceness of the carriers flowing through the body of channel. Accordingly, the on state current is decreased in the thinner device.

A useful function for the extraction of the parameters in the conventional devices is the combination of the drain current and transconductance. This parameter  $(I_D/g_m^{\phantom{m}0.5})$  known as  $Y(V_g)$  is conventionally used to extract parameters such as, the threshold voltage and low field mobility [27]. The  $I_D/g_m^{0.5}$  variation versus gate voltage for devices with different thickness is plotted in Fig. 4b. Before investigating this function, it worth noting that, the operational principles of p-type DGJLTs, is different with that of gated resistors JLTs. In the gated resistor JLTs, below threshold the device is fully depleted. The conversion of the channel from full depletion which no conduction occurs to partial depletion which the bulk conduction starts can be achieved by increasing the gate voltage. Further increase in the gate voltage to the flatband voltage (V<sub>fb</sub>) makes channel to be completely neutral and above flatband an accumulation layer can be created at the interface with the gate [1,5]. On the other hand, the DGJLTs are normally on device which forced to the off state by increasing the gate voltage and depleting the channel from the majority carriers. The variation of majority carriers density in the cross section of channel by increasing the gate voltage is shown in [18]. It is shown that, by small increasing of the gate voltage, the gates can just influence on the corner of the channel and it cannot affect the

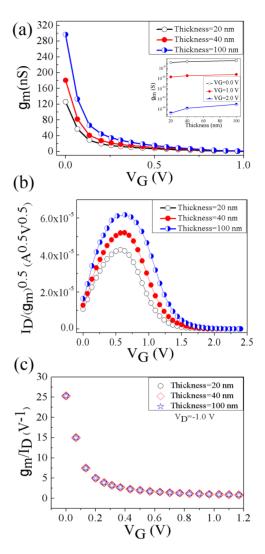


Figure 4. Variation of transconductance  $(g_m)$  (a), plot of the ratio of the drain current to the square root of the transconductance  $(I_D/g_m^{0.5})$  (b) and  $g_m/I_D$  versus gate voltage (c) for three different thicknesses (20, 40, and 100 nm) in saturation  $(V_D=-1.0\ V)$ .

doi:10.1371/journal.pone.0095182.g004

current significantly, mainly due to the lateral gates design and lack of the gate oxide. As a result, in  $I_D/g_m^{\phantom{m}0.5}$  ratio, the numerator cannot be changed significantly; however the decrement of  $g_m$  in denominator is very sharp which is appeared as an increase of  $I_D/g_m^{\phantom{m}0.5}$  in Fig. 4b. On the other hand, by increasing the gate voltage, the lateral gate influence on the current is more significant, while the transconductance variation is negligible. This behavior is shown with a graduate decrease in  $I_D/g_m^{\phantom{m}0.5}$  for higher gate voltage. It should be noticed that, the lower value of  $I_D/g_m^{\phantom{m}0.5}$  in thinner devices might be due the lower carrier concentration and lower value of transconductance; however in thinner devices the current increment is dominant.

A plot of  $g_m/I_D$  as a function of gate voltage when the devices with different thickness are biased in the saturation region ( $V_D = -1.0 \text{ V}$ ), is shown in Fig. 4c. This ratio indeed is a universal characteristic of all the transistors belonging to a same process and known as a measure of the transconductance generation efficiency, since it is a measure of the efficiency of the device to transform current into transconductance [28]. When compared to the

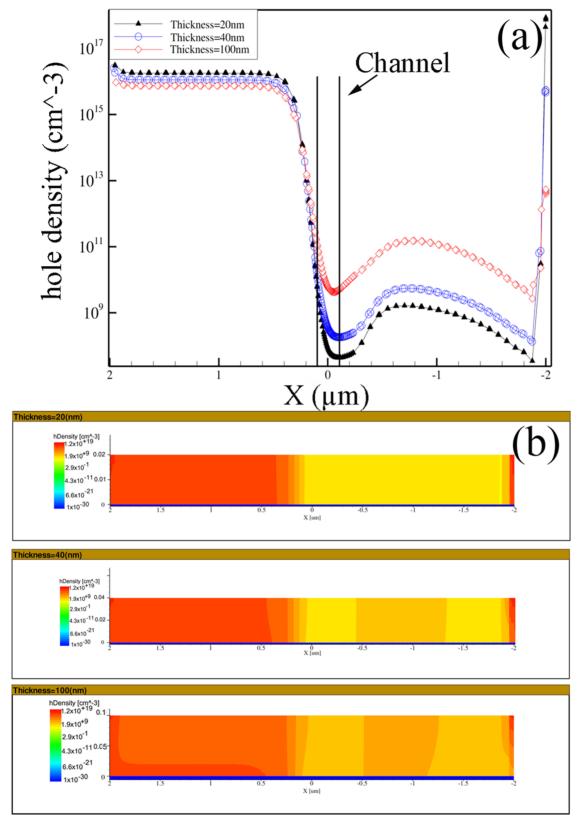


Figure 5. Hole density as a function of position along a horizontal cut line (a) and hole density along a vertical cut at y = 0 (b), for devices with different thicknesses of 20, 40, 100 nm.  $V_D = -1.0 \text{ V}$ ,  $V_G = +2 \text{ V}$ . doi:10.1371/journal.pone.0095182.g005

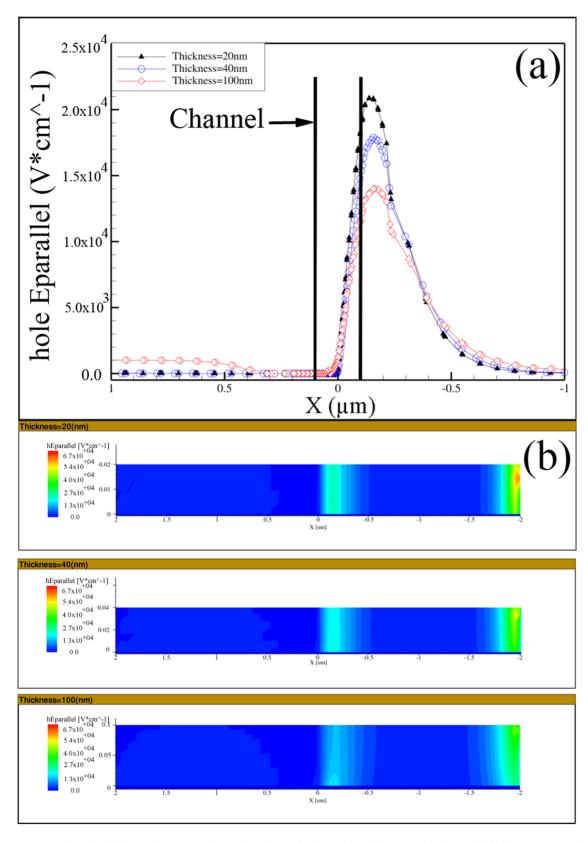


Figure 6. Electric field as a function of position along a horizontal cut line (a) and Electric field along a vertical cut at y = 0 (b), for devices with different thicknesses of 20, 40, 100 nm.  $V_D = -1.0 \text{ V}$ ,  $V_G = +2 \text{ V}$ . doi:10.1371/journal.pone.0095182.g006

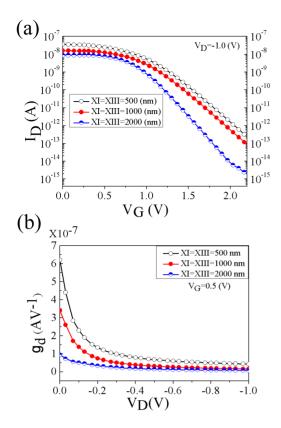


Figure 7. Transfer characteristics ( $I_D$ – $V_G$ ) (a) and drain conductance as a function of drain voltage (b), for different zones  $X_I$ /  $X_{III}$  length at  $V_G$ =0.5 V. Zones  $X_I$  and  $X_{III}$  have the length of 500, 1000, and 2000 nm. The length of zone  $X_{II}$  is constant at 200 nm. doi:10.1371/journal.pone.0095182.q007

degradation of drive current and transconductance in thinner devices, the degradation in  $g_{\rm m}/I_{\rm D}$  is seen to be smaller and  $g_{\rm m}/I_{\rm D}$  is nearly identical for the different thickness. This trend is attributed to the fact that all devices have similar efficiency to transform current into transconductance. In addition,  $g_{\rm m}/I_{\rm D}$  in JLTs is principally controlled by the body factor of the device which is equal to unity in these devices [14]. In the next section, in order to give a deeper physical perception about the effect of thickness variation on the device performance, we investigate the parameters which directly influence the output characteristics.

# Origins of characteristics change by thickness variation

In the DGLJTs, the gradient of carrier concentration along the source/drain and channel is zero, hence no diffusion can take place and the current is dominated by the drift current. For semiconductors with holes as majority carriers, the drift current under an applied field is given by [21]:

$$J = q\mu_P PE \tag{1}$$

hence the current can be written as

$$I = q\mu \int P(A)EdA = q\mu P(x)E(x)$$
 (2)

where q is the electron charge,  $\mu$  is mobility of carriers, P(x) is the carriers density (p(A)) integrated over the cross sectional area (A) of

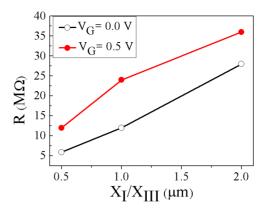


Figure 8. Total resistance versus zone  $X_i/X_{iii}$  length for two different gate voltages of 0.0 and 0.5 V, in the saturation region.

doi:10.1371/journal.pone.0095182.g008

the device, and E is the electric field. It can be seen that, mobility, carrier density and electric field are factors which influence on output characteristics of the device, e.g.  $I_{\rm D}$  and  $g_{\rm m}$ .

In the DGJLT the conduction path locates near the centre of the channel instead of being confined into the surface channel like conventional MOSFETs. This allows for the holes to move through the silicon with bulk mobility. In addition, due to the fabrication method based on the AFM nanolithography the surface and the body of the upper silicon layer in SOI remains intact. This makes the possibility to obtain more bulk properties, such as mobility value close to the bulk mobility even at very thin devices [15]. Accordingly, we can claim that the mobility variation is not the main factor for variation of characteristics of the device when thickness is varied. In this matter, as the most influential factors, the carrier density and electric field are taken into account for further investigation about thickness variation.

The holes density (majority carriers) variation, along a longitudinal cutline and in a cross section of devices with three different thickness (20, 40, 100 nm) at  $V_G$  = +2 V and  $V_D$  = -1 V is presented in Fig. 5 (a–b), respectively. The effect of gate voltage variation on hole density for the nominal device was already investigated in Ref. [18]. In this gate voltage, the lateral gates have the highest impact on the channel with maximum variation of the

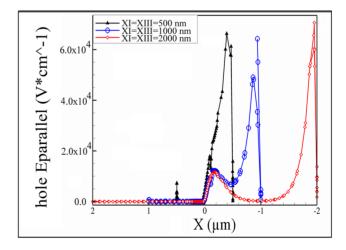


Figure 9. Electric field along the horizontal cut line for different zones  $X_I$  and  $X_{III}$  lengths,  $V_D = -1.0$  V,  $V_G = +2$  V. doi:10.1371/journal.pone.0095182.g009

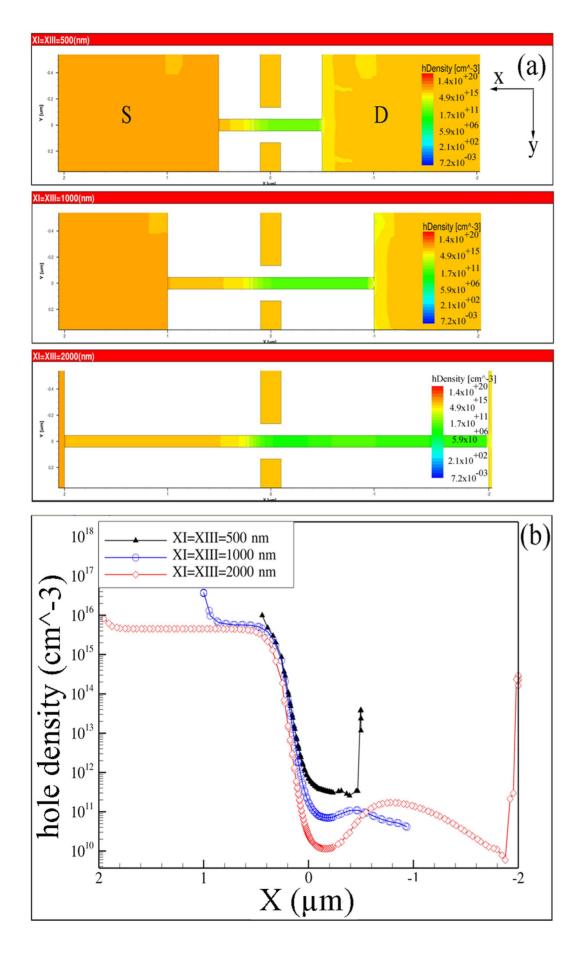


Figure 10. Simulated hole density as a function of position along a cut at Z = 50 nm (a) and hole density as a function of position along a horizontal cut line (b), for different zones  $X_I$  and  $X_{III}$  lengths at  $V_D = -1.0$  V and  $V_G = +2$  V. doi:10.1371/journal.pone.0095182.g010

carrier population. It is observed that, as the channel thickness decreased, the lateral gates were able to vary the carriers more effectively and consequently the current variation is more significant in the thinner devices. This trend confirms the variation of  $g_{\rm m}$  results already presented in Fig. 4a, where lower  $g_{\rm m}$  for smaller thickness can be explained by lower hole density in the channel.

Figs. 6 (a–b) present the electric field along a horizontal line along the channel and in a cross section for devices with different thicknesses. It can be observed that, the peak of the electric field is occurred in the drain extension and not into the channel which confirms the typical behavior of the JLTs for all thicknesses [29]. The maximum electric field along the channel appeared in the device with the smallest thickness and for all devices the electric field reach to the lowest value at the centre of the channel (x = 0).

In fact, the overall electric field behavior along the nanowire was consistent with the hole density. Wherever the electric field was stronger the higher rate of hole depletion occurred. This trend for device with the lower thickness was even more intensified. The stronger electric field and the highest hole depletion were achieved at the lowest thickness (20 nm).

According to the definition of the transconductance  $\left(\frac{\partial I_D}{\partial V_G}\right)$  and considering that the value of the mobility is almost constant, from the differentiation of Equation. 2, we can obtain:

$$\frac{\partial I_D}{\partial V_G} = q\mu \left( \frac{\partial P(x)}{\partial V_G} E(x) + \frac{\partial E(x)}{\partial V_G} P(x) \right)$$
(3)

This shows that two factors can directly cause the transconductance variation. The first factor is the variation of hole density with gate voltage multiplied by electric field along the channel  $(\frac{\partial P(x)}{\partial V_G}E(x))$  and the second factor is the variation of electric field with gate voltage multiplied by hole density along the channel  $(\frac{\partial E(x)}{\partial V_G}P(x))$ . As it is presented in Figs 6a and 7a the hole density is significantly reduced with decreasing the thickness, however the magnitude of electric field is increased with decreasing the thickness. As the value of  $g_m$  is lowered with decreasing the thickness (Fig. 4a), then it can be concluded that  $(\frac{\partial P(x)}{\partial V_G}E(x))$  is the dominant factor contributing to the reduction of  $g_m$  in thinner devices.

## Influence of zones $X_I/X_{III}$ length, Fixed Zone $X_{II}$

According to the specific design of the device, the variation of source/drain extension lengths  $(X_I/X_{III})$  is another important geometric parameter which can significantly influence on the output characteristics. Since the nominal device was experimentally characterized with the channel length  $(X_{II})$  of 200 nm, the same  $X_{II}$  length is taken for all devices to purely investigate the effect of zones  $X_I/X_{III}$  length variation. The transfer characteristics  $(I_D-V_G)$  and drain conductance variation as a function of drain voltage  $(g_d-V_D)$  of the devices with different zones  $X_I/X_{III}$  length, in comparison with the nominal device, are presented in Figs. 7 (a–b). The device display typical characteristics over the investigated range of zones  $X_I$  and  $X_{III}$  lengths. On and off currents were increased by decreasing the zones  $X_I$  and  $X_{III}$  lengths. According to the model proposed in [30], in addition to the channel cross section, the output current in the flatband condition

depends on the effective drain voltage at the channel  $(\ensuremath{V_{ch}})$  and doping concentration.

The transfer characteristics presented in Fig. 7aindicates that as the zones  $X_I/X_{III}$  length scaled down to 500 nm, the subthreshold leakage increased by about 2 orders of magnitude resulting in degradation of  $I_{on}/I_{off}$  ratio. However, the comparison of drain conductance as a function of drain voltage for different zones  $X_I/X_{III}$  length presented in Fig. 7b indicates that by decreasing zones  $X_I/X_{III}$  length an increasing trend of the drain conductance is observed for all drain voltages. In order to understand the reason for increase of the *on* and *off* state current when the zones  $X_I/X_{III}$  length decreases, some of the factors which influence on the characteristics of the device are analyzed.

Total resistance  $(R_{tot})$  versus zones  $X_I/X_{III}$  length obtained using the slope of the output characteristics data for devices with zones  $X_I/X_{III}$  length of 0.5  $\mu m$ , 1  $\mu m$ , and 2  $\mu m$  at two different gate voltages are shown in Fig. 8. The total resistance is varied from 6  $M\Omega$  for device with 0.5  $\mu m$   $X_I/X_{III}$  length to 28  $M\Omega$  for device with 2  $\mu m$   $X_I/X_{III}$  length, at  $V_G$ =0 V. The resistance obtained includes the nanowire resistance  $(R_{NW})$ , contacts resistance  $(R_c)$  and the resistance of the external circuit  $(R_{ext})$  [31],

$$R_{tot} = R_{NW} + R_{ext} + R_{C1} + R_{C2} \tag{4}$$

The  $R_{c1}$  and  $R_{c2}$  could be estimated by simply extrapolating the curves to L=0 since  $R_{\rm NW}$  decreases to zero and  $R_{\rm tot}$  asymptotically approaches the value of  $R_c$  [32]. By decreasing the zone  $X_{\rm I}/X_{\rm III}$  the only part which varies is the  $R_{\rm NW}$  due to the variation of length. Therefore, the lower  $R_{\rm NW}$  in the device with 0.5  $\mu$ m result in the enhanced on -state performance. However, with decreasing length the off state current ( $I_{\rm off}$ ) is also increased, which is not favorable. The reason for  $I_{\rm off}$  value increment can be understood by investigating the electric field and hole density variation along the channel.

The electric field as a function of position along the current direction in devices with three different zones  $X_{\rm I}/X_{\rm III}$  lengths is presented in Fig. 9. It is shown that, the interaction of electric fields from lateral gates and drain contact was diminished in the channel and enhanced in drain extension for all devices. However, for the device with the shortest length ( $X_{\rm I} = X_{\rm III} = 500$  nm), the electric field was stronger at the drain extension compared to the devices with longer  $X_{\rm I}/X_{\rm III}$  length. Moreover, the peak of the electric field is located closer to the channel. For the short  $X_{\rm I}/X_{\rm III}$  length, the strong electric field from the drain could overcome the zone  $X_{\rm II}$  and suck the holes directly from the channel or even zone  $X_{\rm I}$ 

**Table 2.** Dependency of the devices' characteristics to the scaling of thickness and zones  $X_I/X_{III}$  length.

	Thickness scaling	Zone $X_I/X_{III}$ scaling
Threshold voltage (V <sub>th</sub> )	≈Insensitive	≈Insensitive
Leakage current	Decreases	Increases
Hole density in the channel	Decreases	Increases
Electric field	Increases	Increases
The most Influential factor	Hole density	Electric field

doi:10.1371/journal.pone.0095182.t002

(punching effect). The carriers which are dislocated by electric field need to be relocated somewhere in the channel (zone  $X_{\rm II}$ ) or in the zone  $X_{\rm III}$ . This can vary the carrier configuration in different parts of the device.

The holes density variation along the cut at  $Z=50~\mathrm{nm}$  and along a cutline as a function of position for the devices with three different zones  $X_{\mathrm{I}}/X_{\mathrm{III}}$  length are shown in Fig. 10 (a–b), respectively. It can be observed that the lowest depletion in the channel occurred for the shortest length ( $X_{\mathrm{I}}=X_{\mathrm{III}}=500~\mathrm{nm}$ ). For the devices with longer length, the depleted carriers from channel could be spread into the zone  $X_{\mathrm{III}}$ .

However, for the device with the shortest length, there is no area for the depleted holes to be located because of the punching effect and drain influence. This could provide the lower rate of depletion in the channel at off state and explain the higher subthreshold current for the device with the shortest length.

#### Conclusion

Thickness and source/drain extension lengths as two important geometric parameters of the double lateral gate junctionless transistors (DGJLT) have been investigated using 3-D TCAD simulation tools. The device has shown a good scaling ability in the investigated parameters. Unlike the high doped JLTs, the threshold voltage remains approximately insensitive to the

#### References

- 1. Colinge JP, Lee CW, Afzalian A, Akhavan ND, Yan R, et al. (2010) Nanowire transistors without junctions. Nature Nanotechnology 5: 225–229.
- Lee CW, Ferain I, Afzalian A, Yan R, Akhavan ND, et al. (2010) Performance estimation of junctionless multigate transistors. Solid-State Electronics 54: 97– 103.
- Ionescu AM (2010) Electronic devices: nanowire transistors made easy. Nature Nanotechnology 5: 178–179.
- Colinge J, Lee C, Afzalian A, Dehdashti N, Yan R, et al. SOI gated resistor: CMOS without junctions; 2009. IEEE. pp. 1–2.
- Kranti A, Yan R, Lee CW, Ferain I, Yu R, et al. Junctionless nanowire transistor (JNT): Properties and design guidelines; 2010. IEEE. pp. 357–360.
- 6. Lee C, Ferain I, Kranti A, Akhavan ND, Razavi P, et al. Short-channel inctionless papowire transistors: 2010, pp. 1044–1045
- junctionless nanowire transistors; 2010. pp. 1044–1045.

  7. Han M-H, Chang C-Y, Chen H-B, Wu J-J, Cheng Y-C, et al. (2013) Performance comparison between bulk and SOI junctionless transistors.
- Rios R, Cappellani A, Armstrong M, Budrevich A, Gomez H, et al. (2011) Comparison of Junctionless and Conventional Trigate Transistors With Lg Down to 26 nm. Electron Device Letters, IEEE 32: 1170–1172.
- Lee CW, Borne A, Ferain I, Afzalian A, Yan R, et al. (2010) High-temperature performance of silicon junctionless MOSFETs. Electron Devices, IEEE Transactions on 57: 620–625.
- Colinge JP, Lee CW, Ferain I, Akhavan ND, Yan R, et al. (2010) Reduced electric field in junctionless transistors. Applied Physics Letters 96: 073510.
- Gundapaneni S, Ganguly S, Kottantharayil A (2011) Bulk planar junctionless transistor (BPJLT): An attractive device alternative for scaling. Electron Device Letters, IEEE 32: 261–263.
- 12. Sels D, Sorée B, Groeseneken G (2011) Quantum ballistic transport in the junctionless nanowire pinch-off field effect transistor. Journal of computational electronics: 1–6.
- Choi S-J, Moon D-I, Kim S, Ahn J-H, Lee J-S, et al. (2011) Nonvolatile memory by all-around-gate junctionless transistor composed of silicon nanowire on bulk substrate. Electron Device Letters, IEEE 32: 602–604.
- Doria RT, Pavanello MA, Trevisoli RD, de Souza M, Lee C-W, et al. (2011) Junctionless multiple-gate transistors for analog applications. Electron Devices, IEEE Transactions on 58: 2511–2519.
- Dehzangi A, Abdullah AM, Larki F, Hutagalung SD, Saion EB, et al. (2012) Electrical property comparison and charge transmission in p-type double gate and single gate junctionless accumulation transistor fabricated by AFM nanolithography. Nanoscale research letters 7: 1–9.
- Dehzangi A, Larki F, Hutagalung S, Saion E, Abdullah A, et al. (2012) Numerical investigation and comparison with experimental characterisation of side gate p-type junctionless silicon transistor in pinch-off state. Micro & Nano Letters 7: 981–985.
- Larki F, Dehzangi A, Hassan J, Abedini A, Saion E, et al. (2013) Pinch-Off Effect in P-Type Double Gate and Single Gate Junctionless Silicon Nanowire Transistor Fabricated by Atomic Force Microscopy Nanolithography. Nano Hybrids OPEN ACCESS 4: 33–45.

variation of investigated parameter. The subthreshold leakage current shows strong sensitivity to both thickness and zone  $X_{\rm I}/X_{\rm III}$  length variation. It is concluded that, when the thickness of the device is scaled down the carrier's density variation is the main factor which influence the output characteristics of the device. However, for further scaling the zone  $X_{\rm I}/X_{\rm III}$  length the electric fields penetration inside the channel is the main source of increases in leakage current. The dependency of the devices' characteristics to the scaling of the investigated parameters is summarized in Table 2. These modeling results clearly encourage what our previous experimental work has proposed: that the junctionless lateral gate transistor can be used to construct extremely simple transistor device.

# **Acknowledgments**

Authors would like to thank Prof. Dr. Elias B. Saion from UPM for his great support and assistance.

# **Author Contributions**

Conceived and designed the experiments: FL AD SA MNH BYM. Performed the experiments: FL AD SA AJ MSI. Analyzed the data: FL AD SA AJ MSI BYM. Contributed reagents/materials/analysis tools: FL AD SA. Wrote the paper: FL AD SA.

- Larki F, Dehzangi A, Abedini A, Abdullah AM, Saion E, et al. (2012) Pinch-off mechanism in double-lateral-gate junctionless transistors fabricated by scanning probe microscope based lithography. Beilstein journal of nanotechnology 3: 817-823.
- Larki F, Dehzangi A, Saion EB, Abedini A, Hutagalung SD, et al. (2013) Simulation of transport in laterally gated junctionless transistors fabricated by local anodization with an atomic force microscope. physica status solidi (a) 210: 1914–1919.
- Sentaurus T (2010) User Guide, Version D-2010.03, March 2010

   Synopsys, 2010. synopsys.
- Sze SM (2008) Semiconductor devices: physics and technology: John Wiley & Sons.
- Soree B, Magnus W (2009) Silicon nanowire pinch-off FET: basic operation and analytical model. IEEE. pp. 245–248.
- Soree B, Magnus W, Vandenberghe W (2011) Low-field mobility in ultrathin silicon nanowire junctionless transistors. Applied Physics Letters 99: 233509-233509-233503.
- Yan R, Kranti A, Ferain I, Lee CW, Yu R, et al. (2011) Investigation of highperformance sub-50nm junctionless nanowire transistors. Microelectronics Reliability 51: 1166–1171.
- Colinge J, Kranti A, Yan R, Lee C, Ferain I, et al. (2011) Junctionless Nanowire Transistor (JNT): Properties and design guidelines. Solid-State Electronics.
- Chiang MH, Lin JN, Kim K, Chuang CT (2007) Random dopant fluctuation in limited-width FinFET technologies. Electron Devices, IEEE Transactions on 54: 2055–2060.
- Ghibaudo G (1997) Critical MOSFETs operation for low voltage/low power IC's: Ideal characteristics, parameter extraction, electrical noise and RTS fluctuations. Microelectronic engineering 39: 31–57.
- Silveira F, Flandre D, Jespers P (1996) A gm/ID based methodology for the design of CMOS analog circuits and its application to the synthesis of a siliconon-insulator micropower OTA. Solid-State Circuits, IEEE Journal of 31: 1314– 1319
- Colinge JP, Ferain I, Kranti A, Lee CW, Akhavan ND, et al. (2011) Junctionless Nanowire Transistor: Complementary Metal-Oxide-Semiconductor Without Junctions. Science of Advanced Materials 3: 477

  –482.
- Dehzangi A, Abdullah AAM, Larki F, Hutagalung SD, Saion EB, et al. (2012)
   Electrical property comparison and charge transmission in p-type double gate and single gate junctionless accumulation transistor fabricated by AFM nanolithography. Nanoscale Research Letters 7: 381.
- Martinez J, Martinez RV, Garcia R (2008) Silicon nanowire transistors with a channel width of 4 nm fabricated by atomic force microscope nanolithography. Nano Letters 8: 3636–3639.
- Su C-J, Tsai T-I, Liou Y-L, Lin Z-M, Lin H-C, et al. (2011) Gate-all-around junctionless transistors with heavily doped polysilicon nanowire channels. Electron Device Letters, IEEE 32: 521–523.