

## Erratum: “Ultrafast demagnetization by hot electrons: Diffusion or super-diffusion?” [Struct. Dyn. 3, 055101 (2016)]

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The total heat of the electron gas per volume was taken as  $\gamma(T_e^2 - T_0^2)$ , whereas the correct value is  $\frac{1}{2}\gamma(T_e^2 - T_0^2)$ . The corrected diffusion equation reads

$$\gamma T_e \partial_t T_e = k(x) \Delta T_e - G(x)(T_e - T_l) + P(t, x). \quad (3)$$

The simulation was calculated correctly in the paper.<sup>1</sup> In Section V (Analytical Model), the following equations need to be adapted accordingly:

$$Q_{\text{dep},A} = \frac{1}{2} \int_0^{\bar{d}} \gamma(x) (T_e^2(x, t=0) - T_0^2) dx, \quad (7)$$

$$\partial_t Q_{\text{el},A} \approx -\frac{T_0}{\tau} \bar{\gamma} \bar{d} (T_e(t) - T_0), \quad (10)$$

$$\tau = T_0 \frac{\gamma_{\text{Al}} d_{\text{Al}} + \gamma_{\text{Ni}} d_{\text{Ni}}}{G_{\text{Al}} d_{\text{Al}} + G_{\text{Ni}} d_{\text{Ni}}}, \quad (11)$$

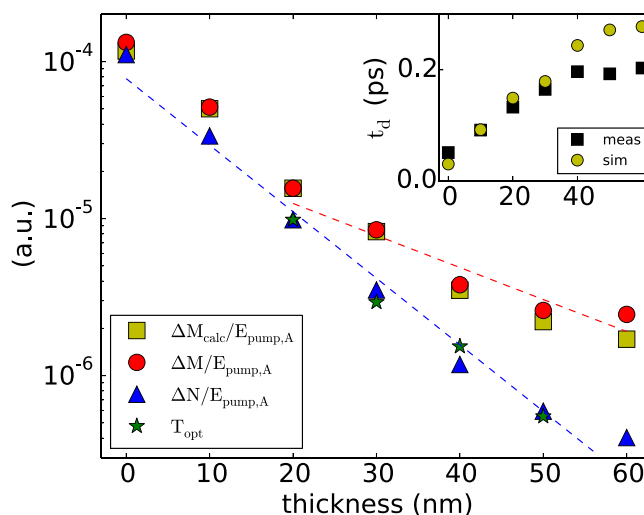


FIG. 3. Amplitude dependence of the ultrafast demagnetization (red) and the non-magnetic contrast (blue) as a function of the absorber film thickness  $d_{\text{Al}}$ , scaled by the pump pulse energy. The non-magnetic contribution follows the optical transmission. The demagnetization initially follows the non-magnetic signal but decays on a longer length scale of 23.5 nm for  $d_{\text{Al}} > 30$  nm. The inset shows the demagnetization time as a function of  $d_{\text{Al}}$  for the measurement and the simulation.

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$$T_e \approx \left( \sqrt{\frac{2Q_{\text{dep},A}}{\bar{\gamma}\bar{d}} + T_0^2} - T_0 \right) e^{-t_d/\tau} + T_0. \quad (12)$$

These changes affect the result of the analytical model presented in Figure 3. The correct model is in better agreement with the experimental data.

<sup>1</sup>G. Salvatella, R. Gort, K. Bühlmann, S. Däster, A. Vaterlaus, and Y. Acremann, [Struct. Dyn.](#) 3, 055101 (2016).