

Theoretical and Experimental Study of CaMgSi Thermoelectric Properties

Jesus L. A. Ponce-Ruiz,* Seiya Ishizuka, Yoshikazu Todaka, Yuki Yamada, Armando Reyes Serrato, and J. M. Herrera-Ramirez*



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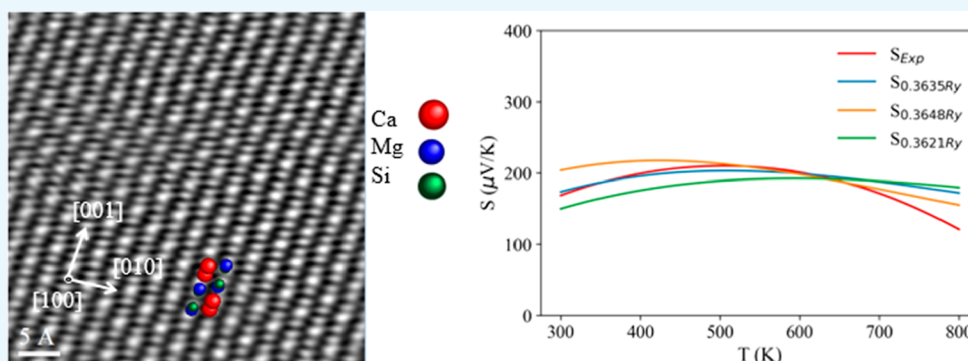


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ABSTRACT: Pure CaMgSi was successfully synthesized by mechanical milling, followed by spark plasma sintering. Rietveld refinement was used to calculate the structural parameters, where a crystallite size (D_{XRD}) of 79 nm was estimated. This value was confirmed by the Williamson–Hall analysis. Transmission electron microscopy was used to analyze the microstructure, revealing the presence of extensive interfaces, nanoparticles, and a high crystallinity. First-principles calculations were performed with the WIEN2k package, finding a band gap of 0.27 eV. The thermoelectric properties were determined combining experimental measurements and theoretical results from the BoltzTraP code. The highest value of the electronic figure of merit (ZT_e) was 1.67 at 415 K. However, when the lattice thermal contribution (k_L) is considered, the highest value of the figure of merit (ZT) was 0.144 at 644 K.

1. INTRODUCTION

Thermoelectricity is the phenomenon in which heat is directly converted into electrical energy making use of the Seebeck effect. This makes it possible to utilize the wasted heat from industrial processes or transport. Thermoelectric materials can be used to form thermoelectric generators stacking p–n semiconductors without the need of moving parts and low maintenance is required. However, despite their benefits, the mass application of these devices has been limited by their low efficiency, the cost of the materials, and the toxicity of some of their constituent elements. The conversion efficiency of thermoelectric (TE) materials is related to a quantity called figure of merit (ZT), which is defined in eq 1.

$$ZT = \frac{S^2 \sigma T}{k} = \frac{\text{PFT}}{k} = \frac{S^2 T}{k \rho} = \frac{S^2 T}{(k_e + k_L) \rho} \quad (1)$$

where S is the Seebeck coefficient, σ is the electrical conductivity, T is the absolute temperature, PF is the power factor equivalent to $S^2 \sigma$, ρ is the electrical resistivity, and k is the thermal conductivity, which have contributions from electrons (k_e) and the crystal lattice (k_L). On the other

hand, eq 2 defines the electronic TE figure of merit (ZT_e). Note that ZT_e is always greater than ZT because k_L is not considered. ZT_e can be used to estimate how good a promising candidate is for TE applications.¹ In semiconductors, the main source of thermal conductivity is the phonon contribution of k_L , so it must be considered.

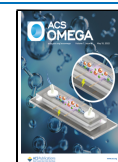
$$ZT_e = \frac{S^2 T}{k_e \rho} \quad (2)$$

In order to find new TE materials, two approaches have been used: first, to explore materials with an intrinsic high ZT and second, to optimize TE properties of a known material by a physical modification.^{2,3} Kagada et al. used the first

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approach and calculated the TE properties of GeTe with the first-principles calculation plus the BoltzTraP code, obtaining $ZT = 0.7$ at 1300 K, while Reyes et al. obtained $ZT = 0.8$ for ReCN at 1200 K.^{4,5} On the other hand, using the second approach, Haque and Rahaman explored the TE behavior of BaGaSnH, replacing Ba by Sr using first-principles calculations and the BoltzTraP code, predicting a $ZT \sim 1.0$.⁶ Hong et al. were able to maximize the ZT value of GeTe by doping with Sb and Se, reaching a value of $ZT = 2.20$ at 780 K.⁷ Hicks and Dresselhaus proposed increasing the ZT by preparing multilayered superlattice materials.⁸ For the case of SiGe compounds, p-type and n-type and the influence of nanostructuring over the TE properties have been reported.^{9–11} Thin films made of $\text{AgPb}_{18}\text{SbTe}_{20}$ were synthesized by molecular beam epitaxy, obtaining $ZT = 2.1$ at 800 K; however, this synthesis process is slow and expensive, making it difficult to manufacture on a large scale.¹² Bi_2Te_3 -based materials have been the most widely studied for TE applications. Poudel et al. made nanocomposites with the addition of Sb by mechanical milling (MM) and hot pressing (HP), reaching a value of $ZT = 1.4$ at 373 K, which is a good value for low-temperature applications, but it could still be improved.¹³ Consolidating milled powders by the spark plasma sintering (SPS) technique has several advantages over HP since SPS was originally designed to inhibit the grain size and provide a better densification.¹⁴ The phase diagram of Ca–Mg–Si was calculated by Gröbner et al. using the Calphad method.¹⁵ The potential applications of CaMgSi have motivated studies in the field of hydrogen storage and biodegradable implants.^{16–18} Besides, the electronic structure information presented by Whalen et al. suggests the CaMgSi compound as a possible candidate for TE applications.¹⁶ However, obtaining a pure phase is a difficult task by conventional synthesis methods. TE properties of Ca–Mg–Si alloys were presented by Niwa et al., obtaining a maximum of 37% CaMgSi, 53% Mg_2Si , and 10% $\text{Ca}_7\text{Mg}_{7.25}\text{Si}_{14}$.¹⁹ CaMgSi, synthesized by MM and SPS, is a promising candidate for TE applications, and even with several secondary phases, it obtained a value of $\text{PF} = 0.42 \text{ mW m}^{-1} \text{ K}^{-2}$ at 433 K; however, thermal conductivity and ZT were not determined.²⁰ Theoretical simulation using first-principles and Boltzmann transport theory have been used to predict the ZT of CaMgSi by tuning the carrier concentration, reaching a value of $ZT = 1.75$ at 800 K.²¹ The present work aims to determine the structure, microstructure, and TE properties of an experimental sample of CaMgSi compound by combining theoretical simulation and experimental measurements.

2. RESULTS AND DISCUSSION

2.1. X-ray Diffraction. Figure 1 shows the refined X-ray diffraction (XRD) pattern of CaMgSi synthesized by MM and SPS, where the signals of pure CaMgSi with the space group $Pnma$ (PDF #43-1399) were observed. It is worth mentioning that pure CaMgSi has not been reported previously. In this case, this finding can be attributed to the rotational speed of 180 rpm under the MM conditions, which guaranteed a better homogeneity than that obtained by Miyazaki et al., who used 125 rpm and obtained six secondary phases.²² Obtaining pure phases makes it possible to compare the experimental measurements with the simulation. Rietveld refinement was performed to determine the lattice parameters (a , b , c , α , β , γ), the crystallite size (D_{XRD}), atomic positions (x , y , z),

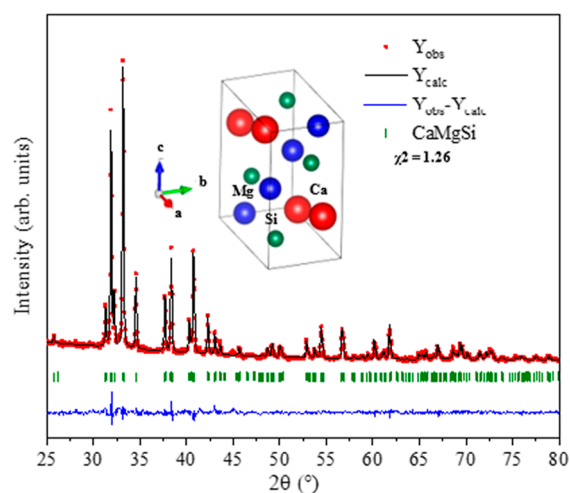


Figure 1. XRD pattern of CaMgSi synthesized by MM and SPS.

occupation factor (g), volume cell (V), and density (ρ). Table 1 summarizes the values calculated.

Figure 2 shows the Williamson–Hall analysis, with which microstrains (ϵ) and crystallite size were calculated.²³ A crystallite size of 79 nm was determined, which is practically the same as the one found by the Rietveld refinement. The crystallite size is small enough to contribute to the reduction of the lattice thermal conductivity (k_L) by increasing phonon scattering and acting as a nanostructured bulk material.¹³ The density of dislocations (δ) was determined by the Williamson–Smallman relation $\delta = 1/D^2$.²⁴ Table 1 presents these values.

2.2. Transmission Electron Microscopy. Figure 3a shows a transmission electron microscopy (TEM) micrograph, where several grain boundaries are observed. Figure 3b shows a high-resolution TEM image in which the presence of a spherical nanoparticle is identified. A similar morphology was reported by Minnich et al. for nanostructured TE materials.⁹ The high density of interfaces and the presence of nanoparticles reduces the thermal conductivity and increases the ZT according to eq 1.^{12,13} Figure 3c shows a region with high crystallinity, where CaMgSi exhibits an orthorhombic arrangement, which is consistent with the $Pnma$ structure oriented in the $[100]$ direction. Figure 3d presents a selected area electron diffraction (SAED) pattern showing sharp and bright diffraction spots, which were assigned to the well-crystallized orthorhombic phase of CaMgSi, and a $[100]$ zone axis was confirmed.

2.3. First-Principles Calculation. The structural parameters reported in Table 1 were used to calculate the optimization of CaMgSi; $a = 7.4460 \text{ \AA}$, $b = 4.4224 \text{ \AA}$, $c = 8.2950 \text{ \AA}$, and a volumetric reduction of 0.29% were obtained. Figure 4a shows the calculated density of states (DOS) determined by the WIEN2k package using the Tran–Blaha modified Becke Johnson (TB-mBJ) potential for the exchange–correlation, obtaining a band gap value of 0.27 eV. It is well documented that using the TB-mBJ improves the band gap calculation of different semiconductors compared to other approximations such as local density approximation and general gradient approximation (GGA) that usually underestimate the band gap.²⁵ Figure 4b shows a strong hybridization between Ca d–Si p and Mg p–Si d, which is consistent with results previously reported using different potentials.^{20,21} Figure 5 shows the calculated band structure with TB-mBJ that confirms a direct band gap in the gamma direction. Miyasaki et

Table 1. Parameters Calculated by Rietveld, Williamson–Hall, and TEM Analysis

a (Å)	b (Å)	c (Å)	$\alpha = \beta = \gamma$ (deg)	D_{XRD} (nm)	D_{TEM} (nm)	ε	δ (cm ⁻²)
7.4764(18)	4.4264(12)	8.3014(2)	90.0	79.02	110	1.21×10^{-3}	1.601×10^{10}
atom	site	x	y	z	g	B	
Ca	4c	0.0189(5)	0.2500	0.6809(4)	1.00219	1.0573	
Mg	4c	0.1432(7)	0.2500	0.0656(7)	0.96110	0.5407	
Si	4c	0.2690(6)	0.2500	0.3842(5)	0.99767	0.8511	

$$V = 274.730(12) \text{ \AA}^3, \rho = 2.209 \text{ g/cm}^3$$

reliability factors: $R_p = 5.78\%$, $R_{wp} = 7.62\%$, $R_e = 6.79\%$, $R_B = 5.79\%$, $R_F = 7.69\%$

goodness-of-fit indicator: $S = R_{wp}/R_e = 1.122$

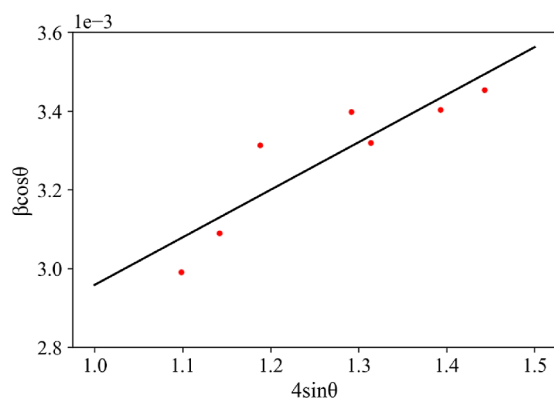


Figure 2. Williamson–Hall plot analysis for the CaMgSi lattice.

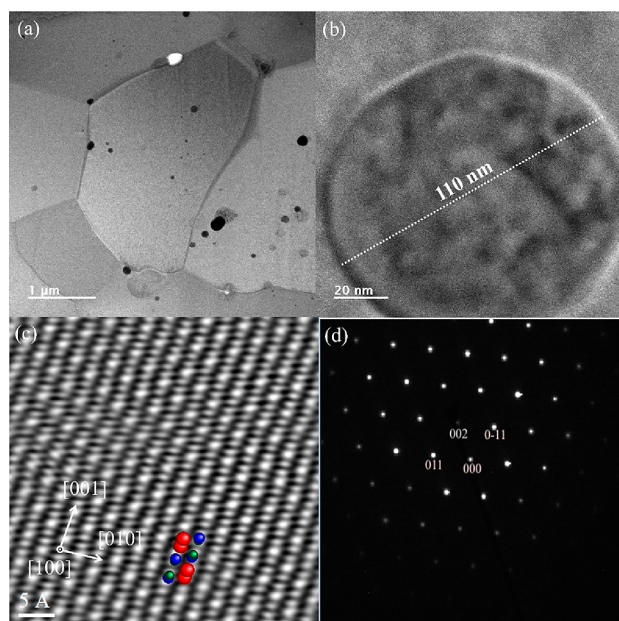


Figure 3. TEM images showing the microstructure of CaMgSi: (a) micrometric grains with clear grain boundaries, (b) high-resolution image showing a nanoparticle, (c) high-crystallinity zone, and (d) indexed SAED pattern with the [100] zone axis.

al. used the WIEN2k package and a GGA, reporting a band gap of 0 eV. Besides, they determined a narrow band gap of 0.26 eV by photoemission spectrum, which showed a clear inconsistency between the first-principles calculations and the experimental results.²² Yang et al. calculated the electronic properties using the Vienna ab initio simulation package (VASP) applying the hybrid functional Heyd–Scuseria–Ernzerhof (HSE06), obtaining a direct band gap of 0.29 eV.²¹

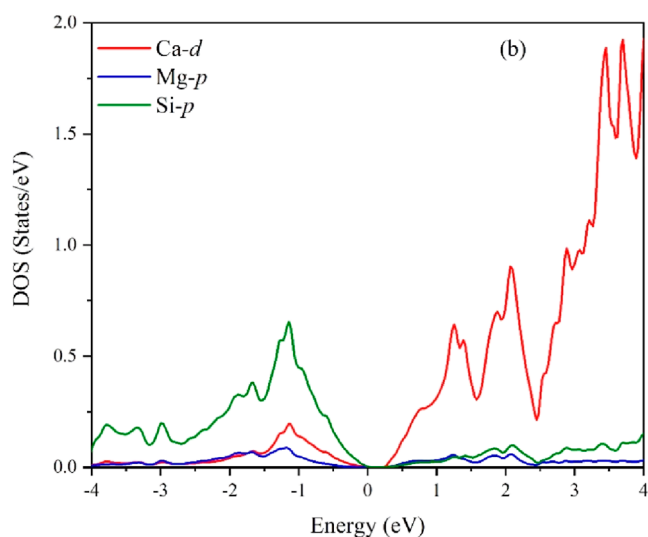
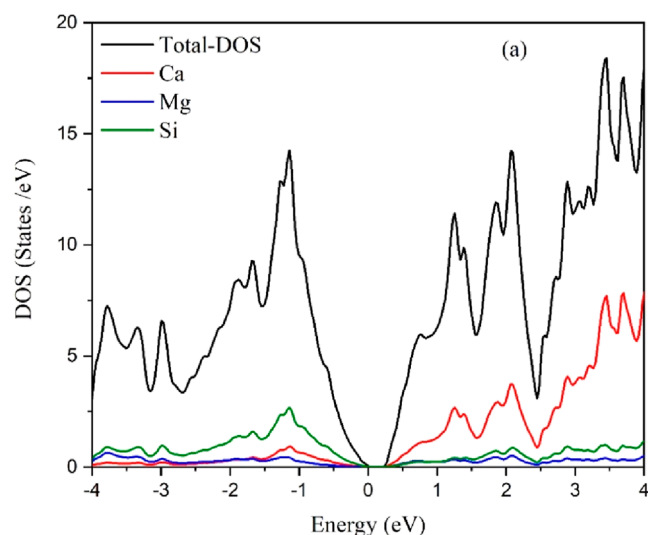


Figure 4. DOS of CaMgSi calculated using WIEN2k with TB-mBJ (a) total and (b) partial.

2.4. TE Properties. TE properties were calculated by BoltzTraP once the theoretical band gap matched with the experimental values.²⁶ Figure 6 shows the simulated value of S as a function of μ at different temperatures (300–800 K) for the CaMgSi compound. In the case of the p-type doping, the maximum S value was 286 μV at 0.3691 Ry, while in n-type doping, it was $-372 \mu\text{V}$ at 0.3759 Ry for a fixed temperature of 300 K. Besides, the transition from p-type to n-type was observed at a μ of 0.3724 Ry. Most of the theoretical reports of TE materials such as PbSe, SnSe, GeSe, and hybrid perovskites

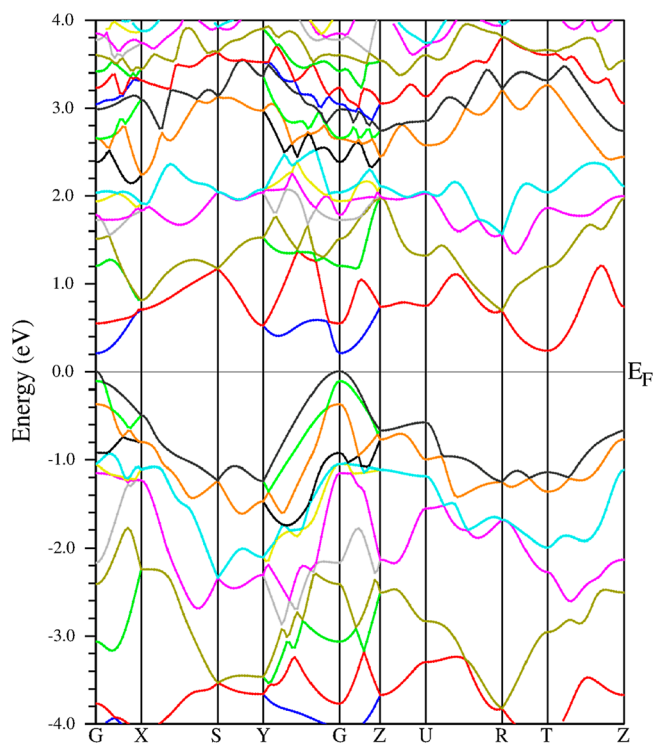


Figure 5. Calculated band structure of CaMgSi using WIEN2k with the TB-mBJ potential.

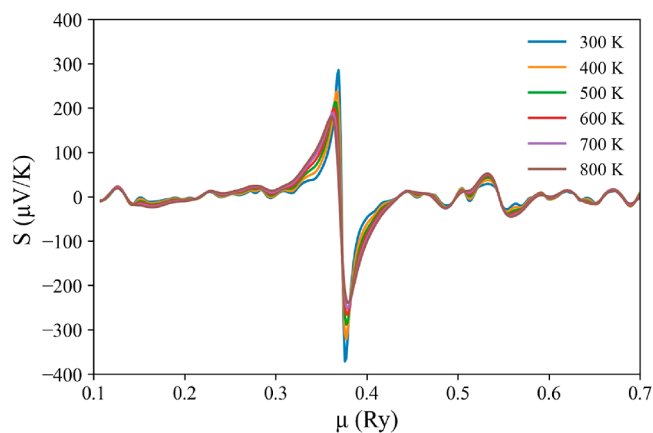


Figure 6. Seebeck coefficient (S) of CaMgSi as a function of the chemical potential (μ) at different temperatures.

$\text{CH}_3\text{CH}_2\text{NH}_3\text{GeI}_3$ use the μ that describes the best properties.^{27,28}

In our case, the μ determined is the one that best describes the S experimental value (S_{Exp}) since it will represent the TE properties on the synthesized CaMgSi. Figure 7 shows the Seebeck values as a function of the temperature. The red line represents the S_{Exp} obtained from ZEM-1, while the other lines represent the three closest theoretical Seebeck (S_{T}) values obtained at a fixed μ . A python code was developed to make a semiempirical adjustment by the least squares method, in which the S_{T} , S_{Exp} , theoretical slope, and experimental slope at fixed μ were considered. The S_{T} value at a μ of 0.3635 Ry is the one that best describes the S_{Exp} (blue line). This fitting method to find a μ that describes the experimental values has been used by Hayashi et al. for single-wall carbon nanotubes.²⁹ This semi-

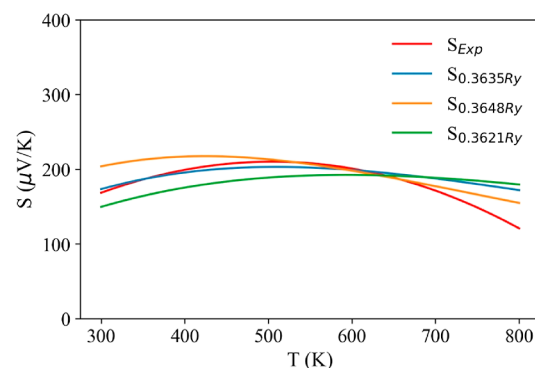


Figure 7. Semi-empirical adjustment from the Seebeck coefficient of CaMgSi to determine the chemical potential.

empirical method has been previously described by Prashun et al.¹

Figure 8a shows the simulated properties σ/τ , and Figure 8b shows k_e/τ as a function of the temperature for the determined

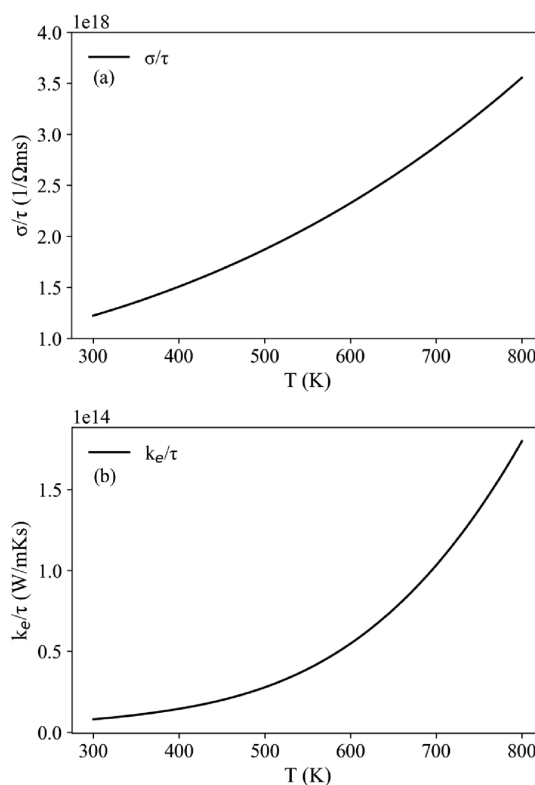


Figure 8. Obtained conductivities of CaMgSi: (a) electrical conductivity σ/τ and (b) thermal conductivity k_e/τ .

$\mu = 0.3635$ Ry. An increase is observed in both properties as the temperature increases.

Figure 9a shows the experimental conductivity (σ_{Exp}), and Figure 9b shows the relaxation time (τ) as a function of the temperature, which were determined by $\tau = \sigma_{\text{Exp}}/(\sigma/\tau)$. As can be seen, τ decreases as the temperature increases until stabilizing at a value of 0.44×10^{-14} s. The determination of the relaxation time using this methodology has been reported by Kumar et al.³⁰

In order to determine the nature of the lattice thermal conductivity, the phonon band structure and phonon DOS were determined (Figure 10). It can be noted that all bands

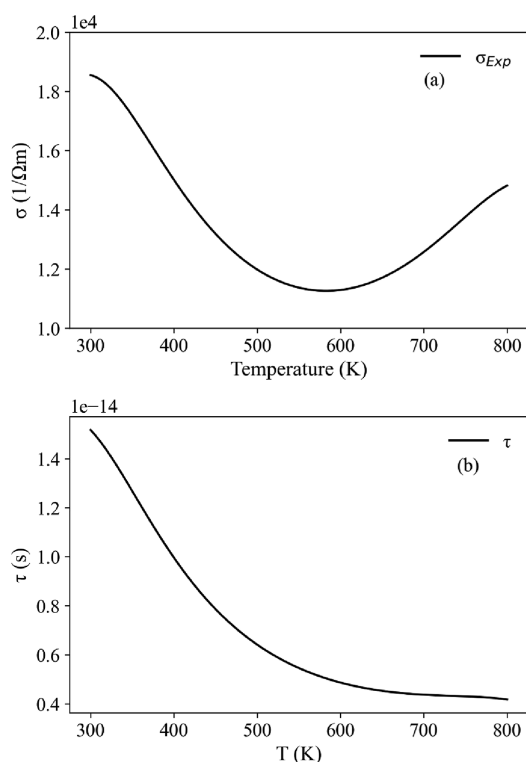


Figure 9. (a) Experimental electrical conductivity σ_{Exp} of CaMgSi and (b) relaxation time τ obtained from the experimental and theoretical electrical conductivities.

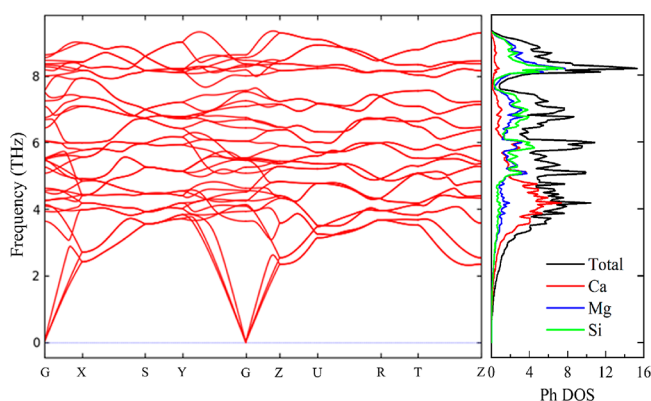


Figure 10. Calculated phonon band structure and phonon DOS of CaMgSi.

start above 0 THz, showing dynamical stability. The CaMgSi unit cell contains 12 atoms, which generate 3 acoustic branches (TA1, TA2, and LA) and 33 optical branches in the first Brillouin zone. The low values of the acoustic branches (0–2.5 THz) suggest a low k_L . The Ph DOS shows that the vibration modes of Mg and Si are strongly coupled due to their similar mass. The high optical frequency modes are presented above 8 THz produced mainly by Mg and Si. The phonon properties of CaMgSi synthesized in this work are consistent with those reported by Yang et al.²¹

Figure 11a shows the Grüneisen parameter (γ) of CaMgSi as a function of the frequency, which represents the degree of anharmonicity.³¹ The acoustic phonon modes (0–2.5 THz) provide γ values of -1.2 to 4.2 . Figure 11b shows the lifetime of CaMgSi obtaining a value of 16.8 ps at acoustic branches. Li et al. reported similar values of the Grüneisen parameter and

lifetime for BaMgSi, Ba₂Mg₃Si₄, and BaMg₂Si₂, identifying that they can be designed as low-thermal-conductivity materials.³²

Figure 12a shows the electronic contribution to the thermal conductivity k_e , which was determined with the BoltzTraP package at a chemical potential of $\mu = 0.3635$ Ry and by taking the obtained relaxation time. Figure 12b shows the lattice contribution to the thermal conductivity k_L , which was obtained combining density functional theory (DFT) calculations implemented in the VASP and phono3py.^{33–35} Figure 12 proves that k_L is several times bigger than k_e ; thus, k_L must be reduced to increase ZT . It is important to consider the influence of the density of dislocations δ in Table 1 because it can reduce k_L reported in Figure 12b through phonon scattering.^{36,37}

Figure 13a shows the ZT_e value as a function of the temperature, reaching a maximum value of $ZT_e = 1.67$ at 415 K. Figure 13b shows the ZT value considering the effect of k_L , for which the values are lower than for ZT_e but reaching a maximum value of $ZT = 0.144$ at 644 K. Table 2 shows a summary of the TE properties obtained in this work. The biggest difference between the values presented in Table 2 and those simulated ($ZT = 1.78$ at 800 K) can be mainly attributed to the fact that the experimental electrical conductivity σ_{Exp} is several times lower than the one predicted by Yang et al.²¹ The presented information suggests that increasing σ will lead to bigger ZT values, which can be achieved by doping the CaMgSi compound.

3. CONCLUSIONS

In the presented work, pure CaMgSi was successfully synthesized using MM followed by SPS, obtaining a crystallite size of 79 nm. The morphology reported confirms the formation of a nanostructured bulk material with a high crystallinity. Through first-principles calculations, a band gap of 0.27 eV was obtained, which is very close to the experimental value of 0.26 eV. The thermal conductivity and relaxation time were obtained combining BoltzTraP results with experimental measurements, making it possible to determine the figure of merit, with a maximum value of $ZT_e = 1.67$ at 415 K. When k_L was included, a maximum value of $ZT = 0.144$ at 644 K was obtained. However, the experimental electrical conductivity σ_{Exp} obtained is several times lower than the one previously simulated, the increment of σ_{Exp} being the key factor to reach bigger values of ZT .

4. EXPERIMENTAL CONDITIONS AND COMPUTATIONAL METHODS

4.1. Experimental Procedure. CaH₂ (Sigma-Aldrich 99%, <200 μ m), Mg (Mitsuwa chemicals 99.9%, <150 μ m), and Si (Kojundo Chemicals, 99.9%, <5 μ m) were used to synthesize the CaMgSi compound. Initially, 100 g of precursor powders was weighed in a stoichiometric ratio into a glovebox under an Ar atmosphere. Then, the powders were mechanically milled under an Ar atmosphere using a planetary ball mill (Pulverisette 5); a 500 cm³ bowl (SUS304 stainless steel) and 100 bearing balls (SUJ2 bearing steel, 10 mm) were used. The ball-to-powder weight ratio was 4:1, the rotational speed of the main disk was 180 rpm, and the powders were milled for 20 h to obtain a better homogeneity. The milled powders (5 g) were put into a graphite die with an inner diameter of 25 mm and a thickness of 4 mm. Subsequently, an SPS machine (Dr. Sinter SPS, Sumitomo Coal Mining) was utilized under an Ar

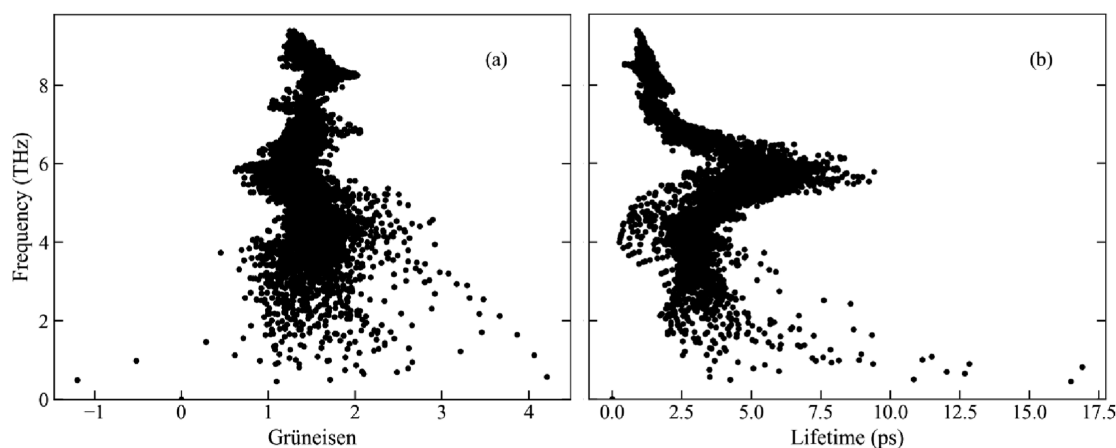


Figure 11. (a) Grüneisen parameter and (b) lifetime of CaMgSi.

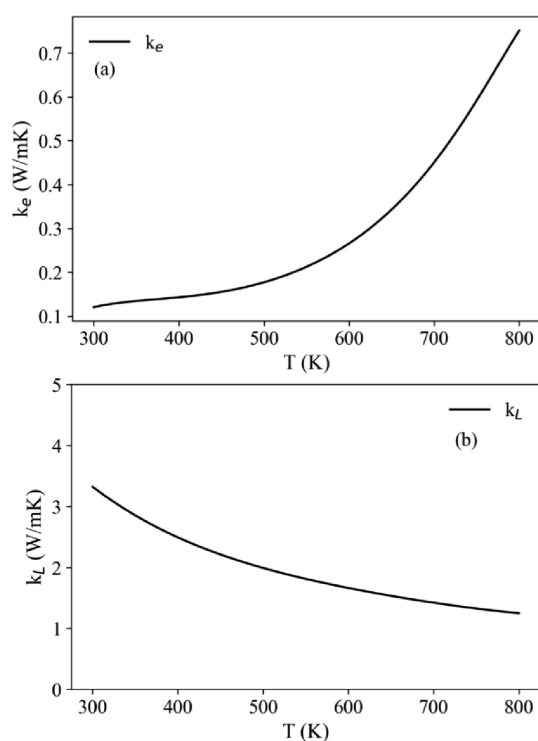


Figure 12. Thermal conductivities of CaMgSi due to (a) electrons obtained from relaxation time and BolzTraP and (b) lattice obtained from phono3py.

atmosphere and a constant pressure of 50 MPa. The SPS conditions used in this case were taken from Miyazaki et al.²⁰ Thus, the milled powders were sintered at 1273 K for 20 min with a heating ramp of 97.5 K/min (dehydrogenation of CaH₂ and formation of CaMgSi). Then, a controlled cooling was carried out with a ramp of 33.3 K/min up to a temperature of 773 K. Later, a natural cooling up to room temperature was done.

4.2. Characterization. Powder XRD patterns of CaMgSi powders were acquired with a PANalytical X'Pert Pro diffractometer equipped with a monochromator and a radiation source of Cu K α 1 ($\lambda_1 = 1.54056 \text{ \AA}$) operating at 40 kV/30 mA. Diffraction patterns were acquired in the 2θ range of 25–80° with a step size of 0.02°. Rietveld refinement was performed using FullProf software, and the crystal structure was generated with the VESTA software.^{38–40} TEM

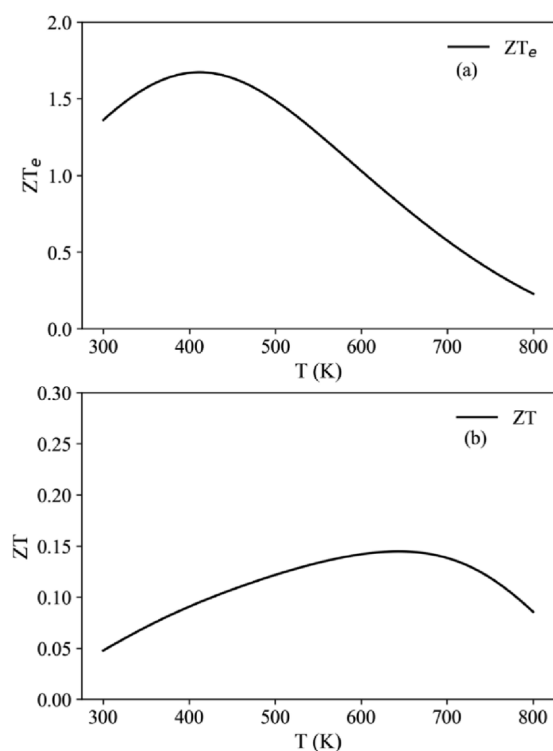


Figure 13. ZT of CaMgSi as a function of the temperature: (a) electronic contribution ZT_e and (b) total ZT.

micrographs were obtained with a JEOL JEM2200FS +CS microscope at an accelerating voltage of 200 kV. The Miller indices' identification on the SAED patterns was made with the Crystallographic Tool Box (CrysTBox) software using the z .⁴¹ The experimental TE measurements were taken with a ZEM-1 Ulvac Sinku-Riko equipment.

4.3. Computational Details. The first-principles calculations were done with the WIEN2k package, which allows us to perform electronic structure calculations using the DFT. It is based on the full-potential linearized augmented plane wave method for solving Kohn–Sham equations.⁴² The exchange–correlation energy was treated with the TB-mBJ potential.²⁵ The separation energy between the core and the valence state was set to -6.0 Ry . Our calculations were carried out in the Brillouin zone with a $13 \times 23 \times 12$ uniform k -point mesh and a convergence criterion of 10^{-4} Ry . The TE properties were

Table 2. TE Properties of CaMgSi for Different Temperatures

T (K)	S ($\mu\text{V/K}$)	$\sigma \times 10^4$ ($1/\Omega \text{ m}$)	PF $\times 10^{-4}$ (W/m K^2)	$\tau \times 10^{-15}$ (s)	ke (W/m K)	kL (W/m K)	ZT _e	ZT
300	168.4766	1.8675	5.2647	15.172	0.1198	3.2906	1.3180	0.0477
400	199.1247	1.5012	5.9531	9.9693	0.1431	2.4931	1.6642	0.0905
500	210.0595	1.1999	5.2632	6.3841	0.1772	1.9933	1.4913	0.1215
600	200.9097	1.1294	4.5579	4.8591	0.2657	1.6606	1.0289	0.1418
700	171.3055	1.2573	3.6840	4.3655	0.4505	1.4232	0.5733	0.1382
800	120.8761	1.4836	2.1649	4.1702	0.7503	1.2452	0.2308	0.0855

calculated using the BoltzTraP code, which uses the approach called constant relaxation time approximation, which combines the electronic structure calculation and Boltzmann statistics. The obtained results were S , σ/τ , and k/τ , where the electrical and thermal conductivities are a function of the relaxation time τ .²⁶ A python script was developed to determine which simulated Seebeck values at fixed chemical potential best describe the experimental Seebeck (S_{Exp}) measurements using the least-squares method. The phonon band structure and phonon DOS were obtained using VASP and the density functional perturbation theory implemented in phonopy with a $2 \times 3 \times 2$ supercell.^{33,43} The phonon lifetime, Grüneisen parameters, and lattice thermal conductivity were obtained combining VASP and phono3py.^{33,35} The VASP parameters are kinetic energy cutoff of 500 eV, a Monkhorst–Pack scheme in the Brillouin zone with $2 \times 2 \times 2$ uniform k -point mesh, and the convergence energy threshold of 10^{-8} eV. The interatomic force constants in phono3py were determined using various positions of the atoms in a supercell made of $2 \times 2 \times 2$ primitive cells.

AUTHOR INFORMATION

Corresponding Authors

Jesus L. A. Ponce-Ruiz – Centro de investigación en Materiales Avanzados (CIMAV), Laboratorio Nacional de Nanotecnología, Chihuahua 31136, Mexico; orcid.org/0000-0002-3123-1245; Email: jesus.ponce@cimav.edu.mx

J. M. Herrera-Ramirez – Centro de investigación en Materiales Avanzados (CIMAV), Laboratorio Nacional de Nanotecnología, Chihuahua 31136, Mexico; orcid.org/0000-0003-4321-6445; Email: martin.herrera@cimav.edu.mx

Authors

Seiya Ishizuka – Toyohashi University of Technology, Department of Mechanical Engineering, Toyohashi, Aichi 441-8580, Japan

Yoshikazu Todaka – Toyohashi University of Technology, Department of Mechanical Engineering, Toyohashi, Aichi 441-8580, Japan

Yuki Yamada – Toyohashi University of Technology, Department of Mechanical Engineering, Toyohashi, Aichi 441-8580, Japan

Armando Reyes Serrato – Universidad Nacional Autónoma de México, Departamento modelación de nanomateriales, Centro de Nanociencias y Nanotecnología, Ensenada, Baja California 22860, Mexico

Complete contact information is available at: <https://pubs.acs.org/10.1021/acsomega.1c07307>

Notes

The authors declare no competing financial interest.

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